



SCIPP

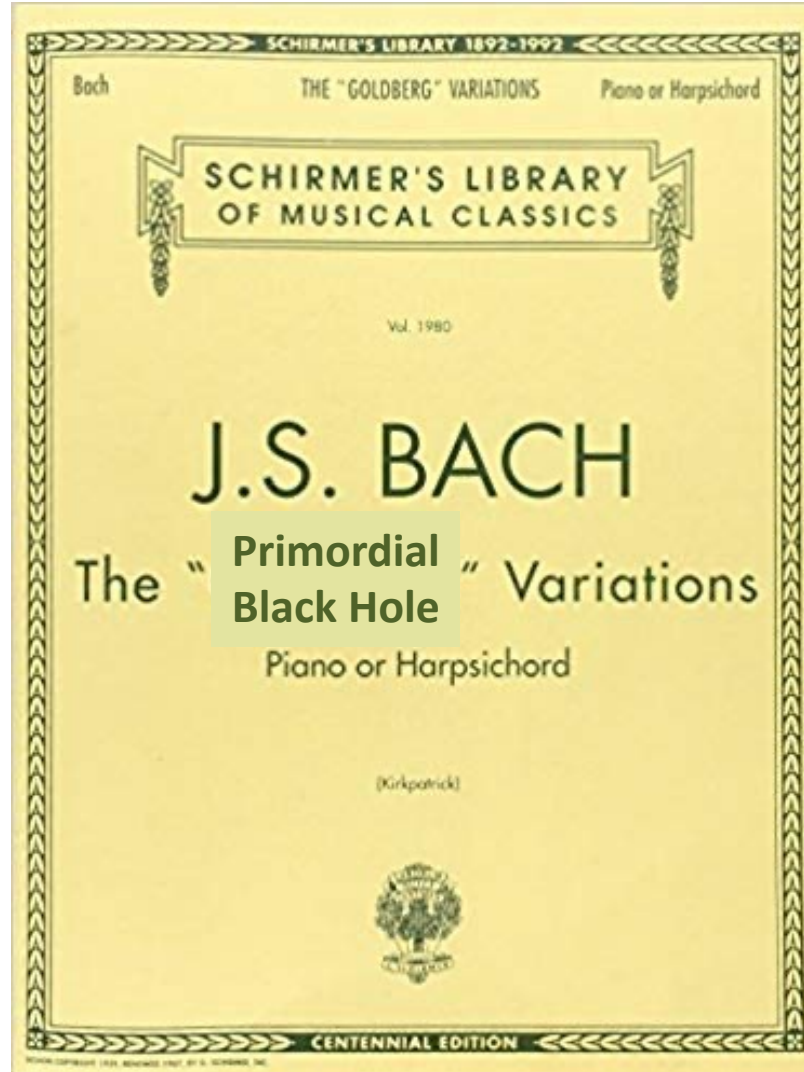
SANTA CRUZ INSTITUTE FOR PARTICLE PHYSICS

# Stefano Profumo

University of California, Santa Cruz



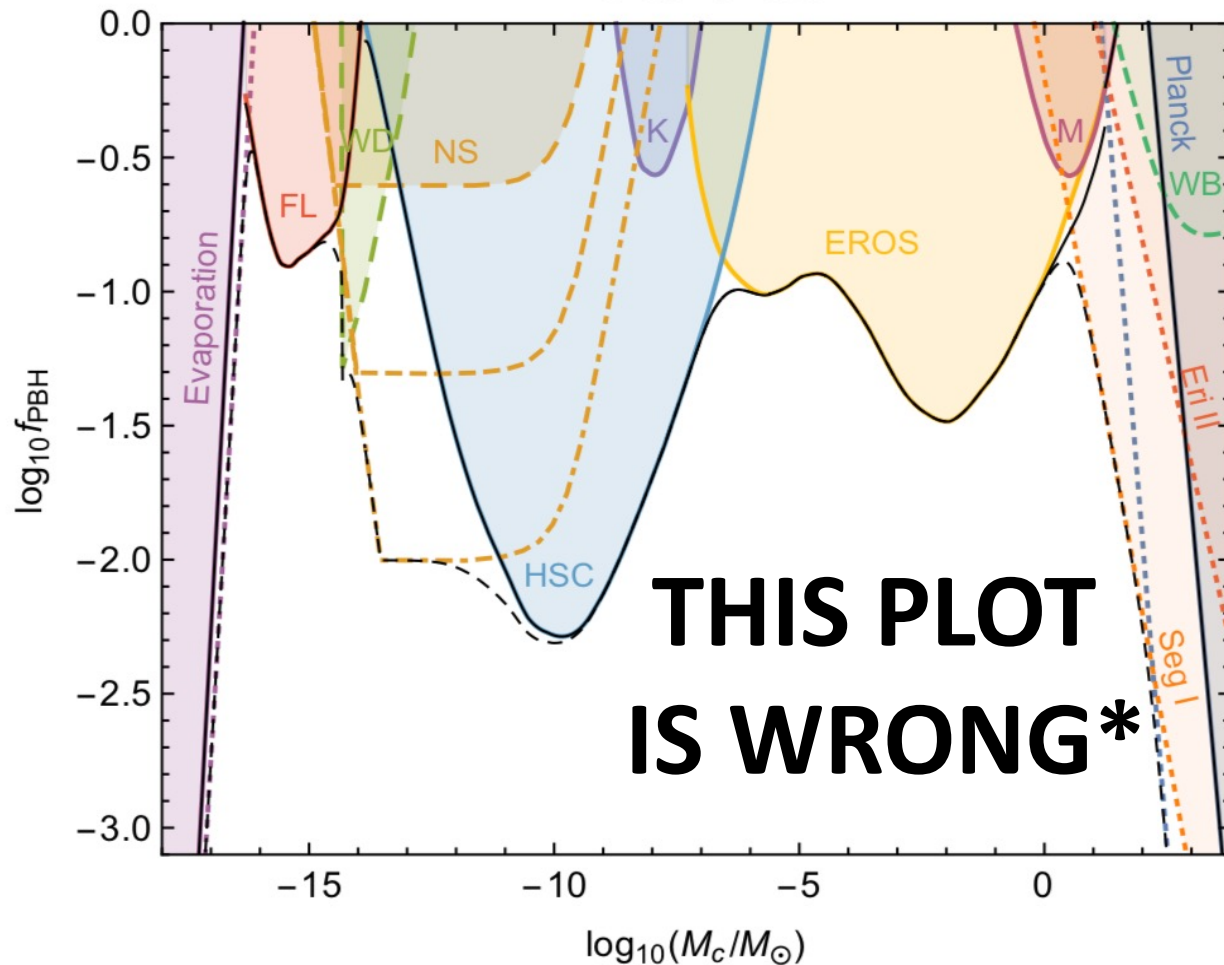
UC SANTA CRUZ



Phenomenology 2021 Symposium – May 26, 2021

Can there be **enough** PBH around to be the **DM**?

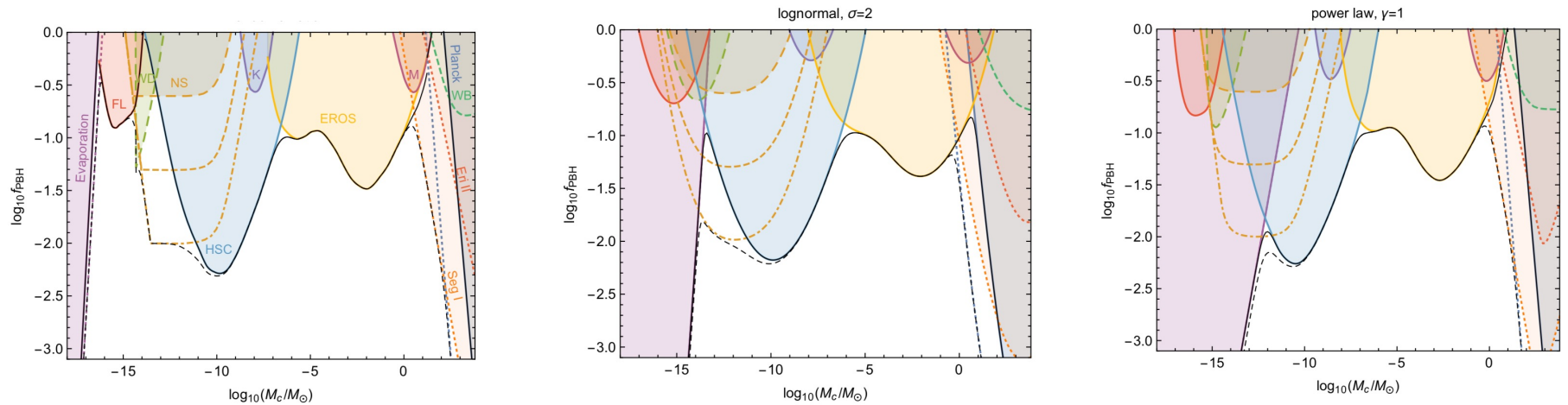
What is the **maximal fraction** of **dark matter** in **PBH**?



\*Carr has since corrected it!

Carr et al, 2017

The **fraction** of PBH that could be the **dark matter** depends on the **mass function**!



...what is the mathematical function that **maximizes** the **mass fraction** of primordial black holes compatibly with **constraints**?

stro-ph.CO] 3 Apr 2018

# The Maximal-Density Mass Function for Primordial Black Hole Dark Matter

**Benjamin V. Lehmann, Stefano Profumo and Jackson Yant**

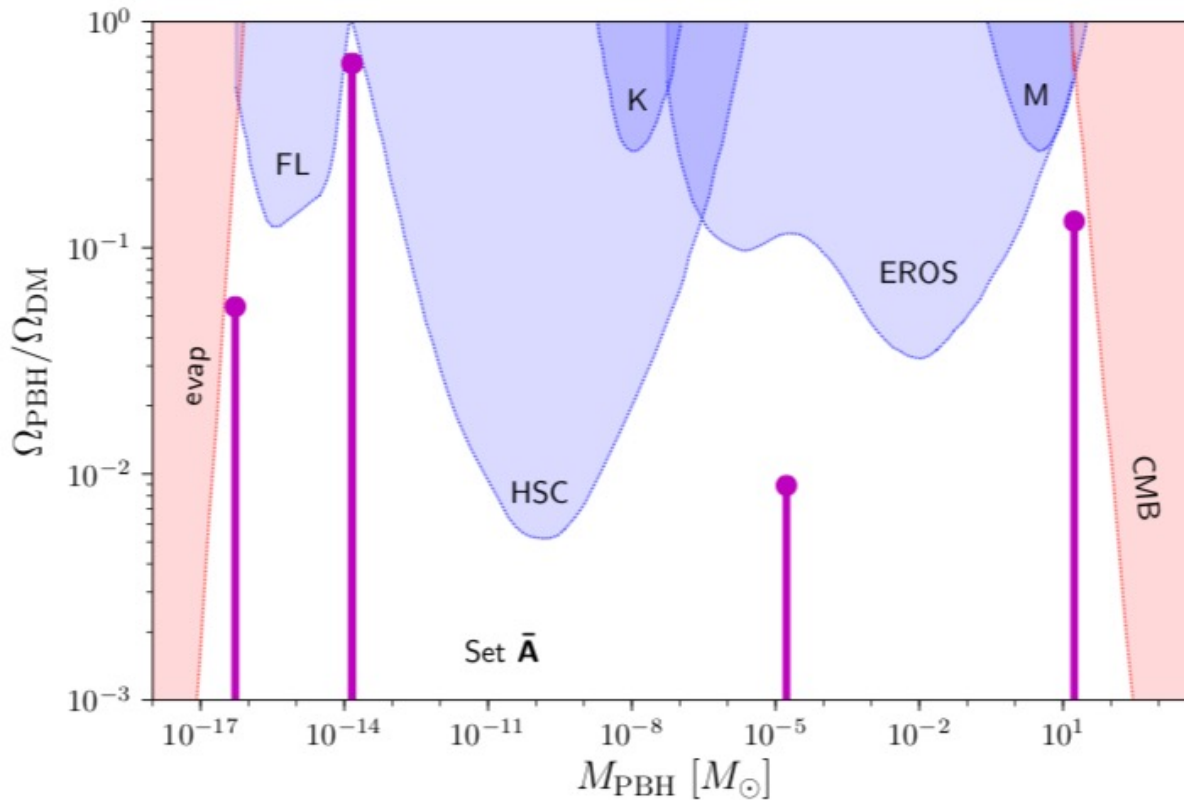
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E-mail: [blehmann@ucsc.edu](mailto:blehmann@ucsc.edu), [profumo@ucsc.edu](mailto:profumo@ucsc.edu), [jyant@ucsc.edu](mailto:jyant@ucsc.edu)

**Abstract.** The advent of gravitational wave astronomy has rekindled interest in primordial black holes (PBH) as a dark matter candidate. As there are many different observational probes of the PBH density across different masses, constraints on PBH models are dependent on the functional form of the PBH mass function. This complicates general statements about

**Answer: with  $N$  independent constraints, the optimal function is a linear combination of  $N$  delta functions with calculable relative weights**

$$\min \{ \|\mathbf{x}\| \mid \mathbf{x} \in \text{conv} \{ \mathbf{g}(M) \mid M \in U \} \}$$



**Answer: with  $N$  independent constraints, the optimal function is a linear combination of  $N$  delta functions with calculable relative weights**

**Planck scale**

**Stellar-mass**

(heavier ~ ruled out by  
dynamical/accretion constraints)

$10^{-3}$  g

$10^7$  g

$10^{17}$  g

$10^{27}$  g

$10^{37}$  g

$10^{30}$  eV

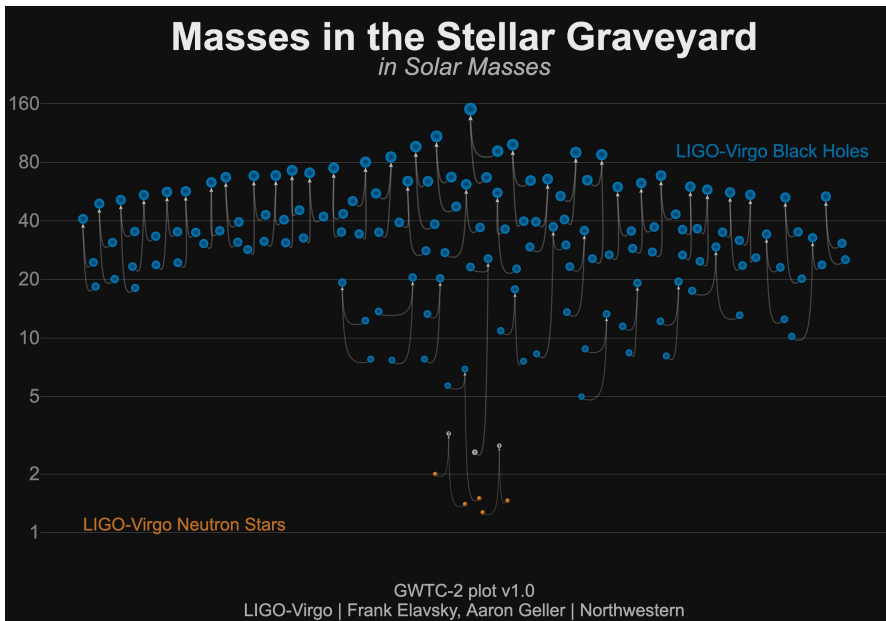
$10^{40}$  eV

$10^{50}$  eV

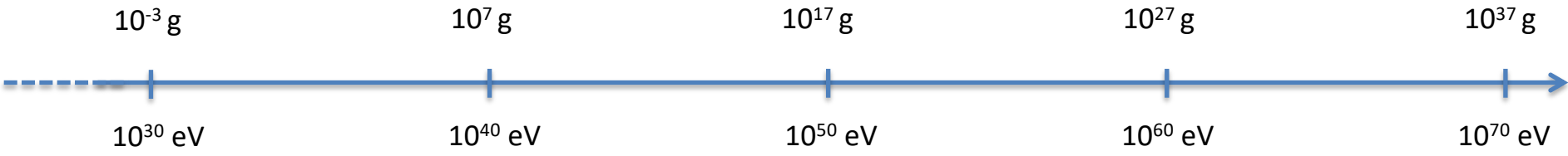
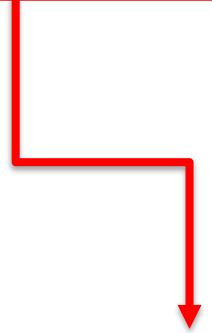
$10^{60}$  eV

$10^{70}$  eV



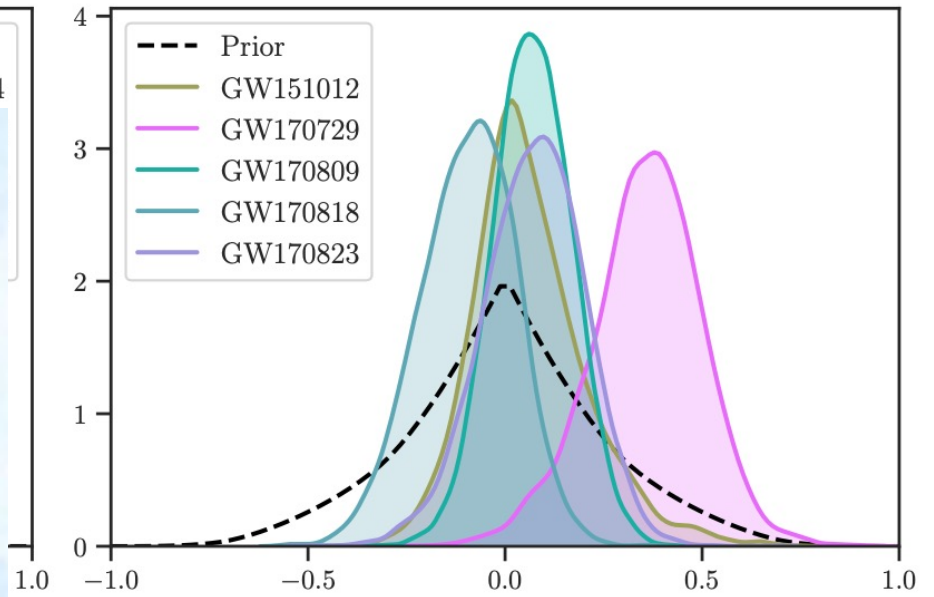


**“Stellar-Mass”  
( $10^{35}$  g)  
Black Holes**



✓ **Spins look a lot like PBH!\***

\* Fernandez and Profumo, 2019

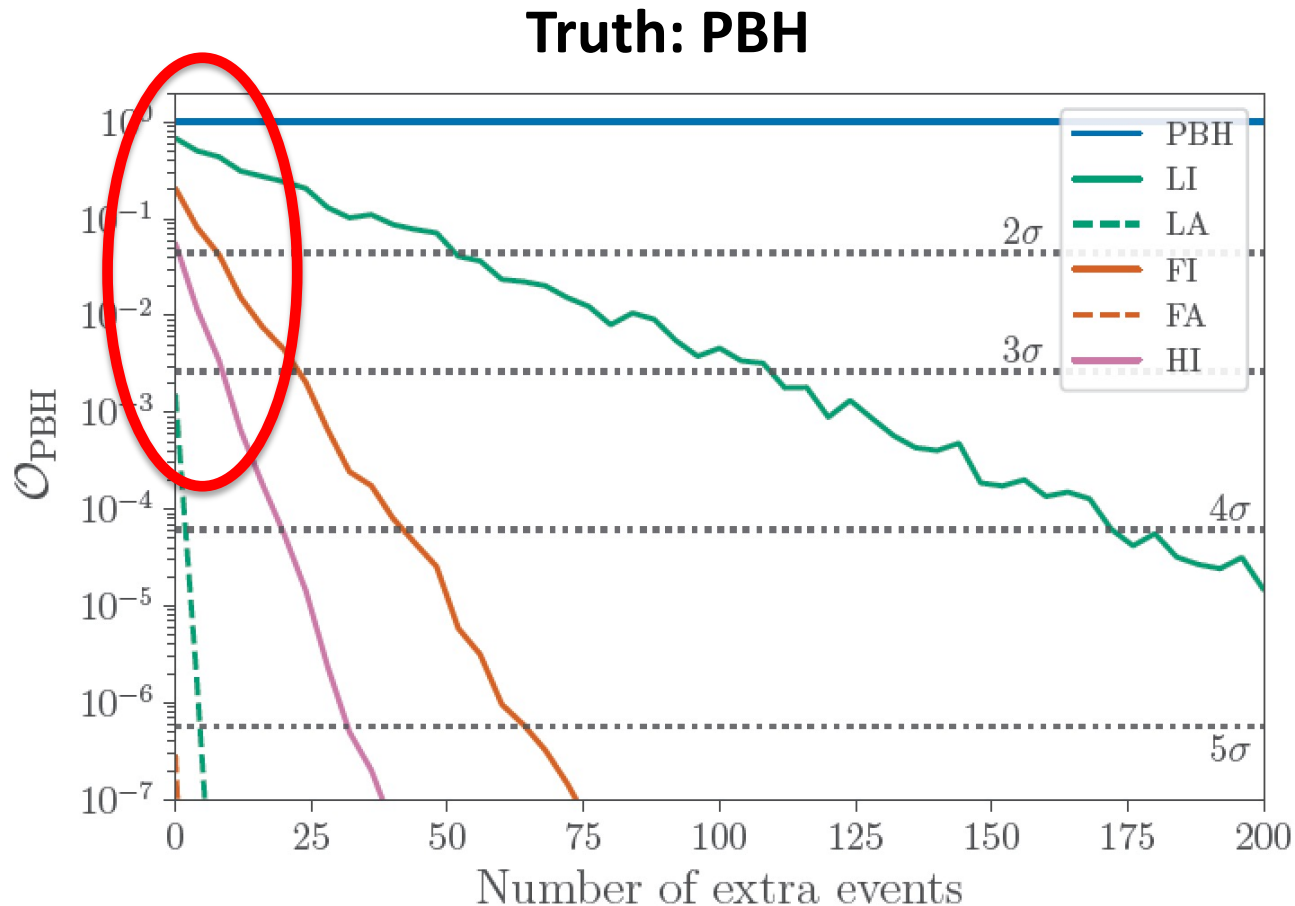


## Odds ratios

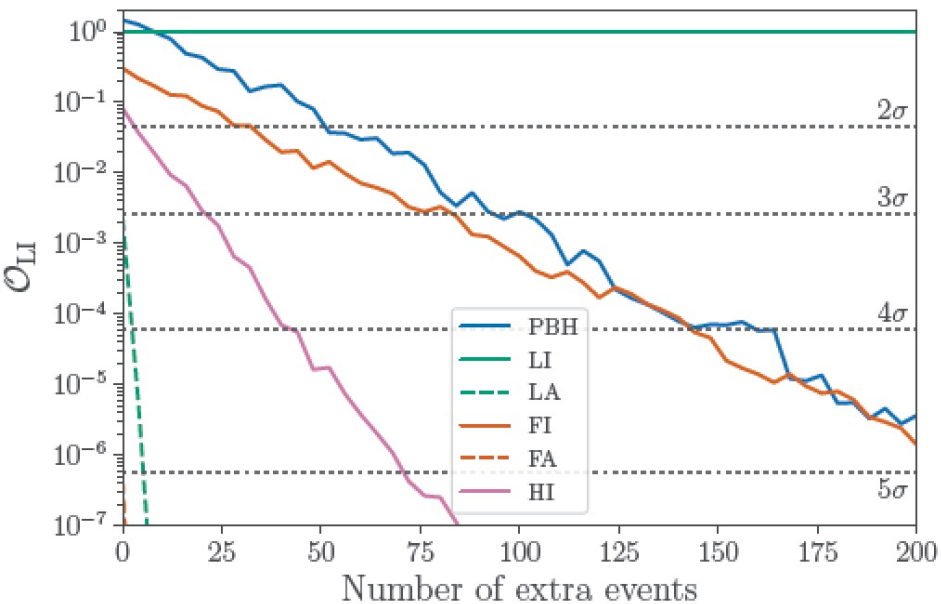
Flat	High	PBH
-1.18	-2.49	0.39
-14.65	-36.41	



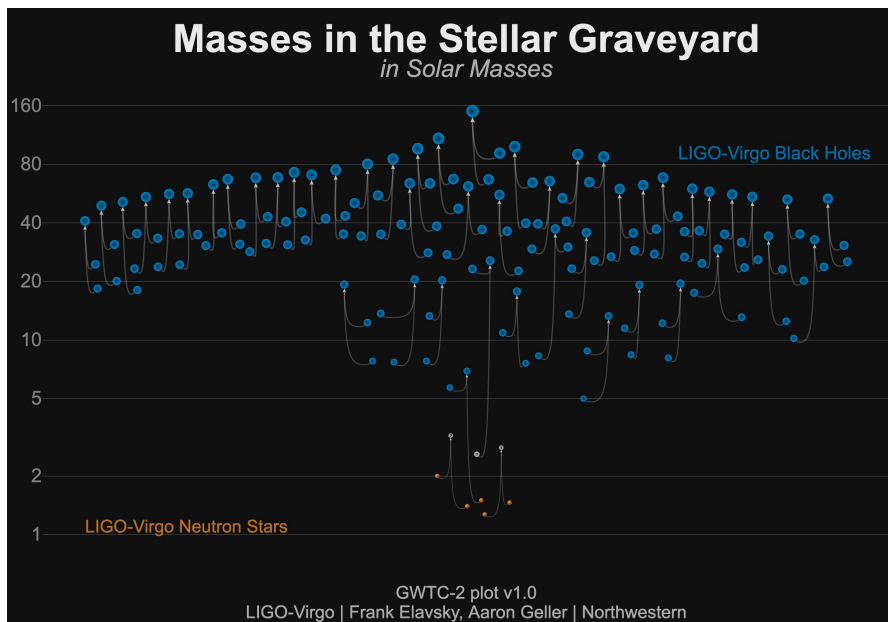
# Evolution of the Odds ratios



# Evolution of the Odds ratios



**Truth: Low-isotropic**



**“Stellar-Mass”  
( $10^{35}$  g)  
Black Holes**

$10^{-3}$  g

$10^7$  g

$10^{17}$  g

$10^{27}$  g

$10^{37}$  g

$10^{30}$  eV

$10^{40}$  eV

$10^{50}$  eV

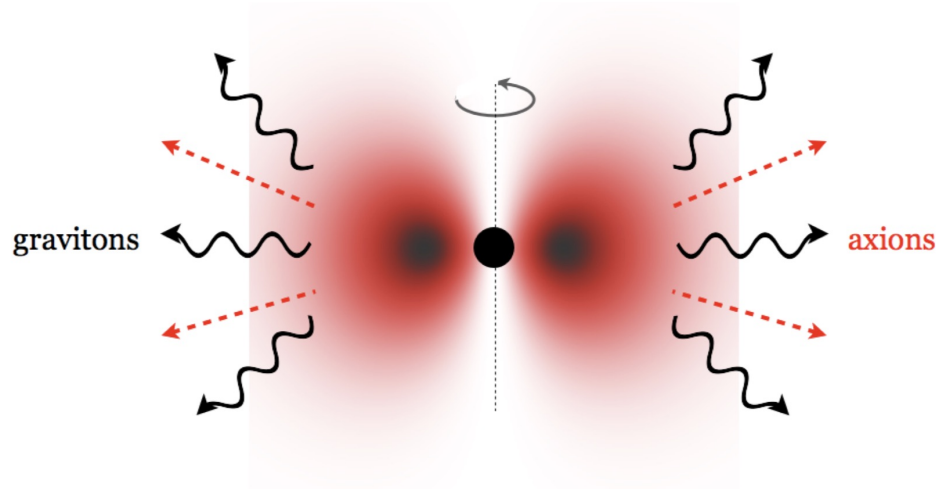
$10^{60}$  eV

$10^{70}$  eV

- ✓ **Spins** look a lot like PBH!
- ✓ ...or maybe they are low because of **superradiance**\*?

\* Fernandez, Ghalsasi, Profumo, 2020

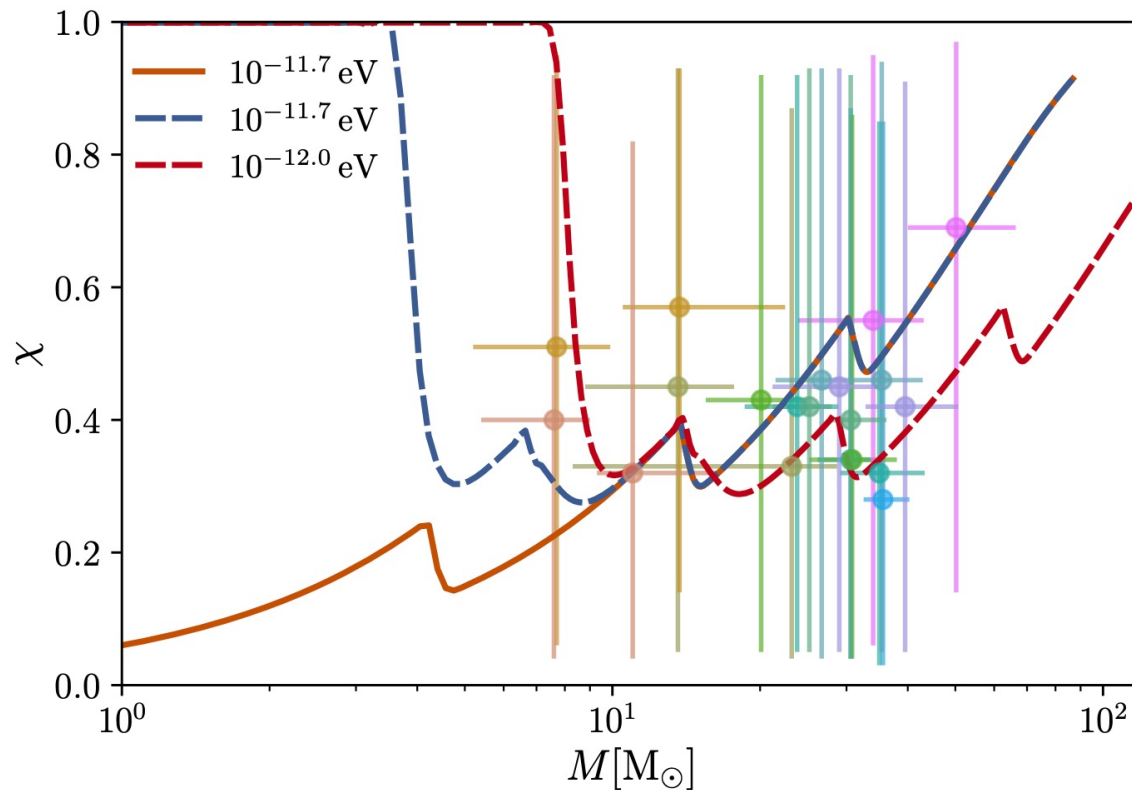
# What else could **fake** a **low-spin PBH**? **Super-radiance!**



Assuming an initial **spin** and **alignment** distribution, one can compute the “**best-fit**” axion mass

Similarly, spin measurements can put **constraints** on axion-like particles

# What else could fake a low-spin PBH? Super-radiance!

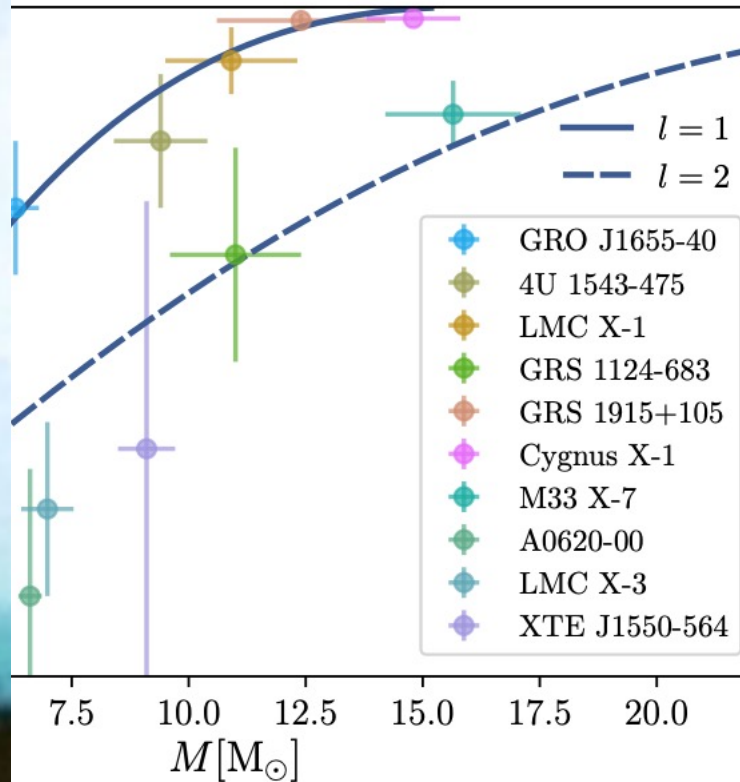


	GW150914		GW151226		GW170608		GW170809		GW170818
	GW151012		GW170104		GW170729		GW170814		GW170823

**Regge** plot (effective spin vs mass) assuming  
**Flat priors** for both mass and spin\*

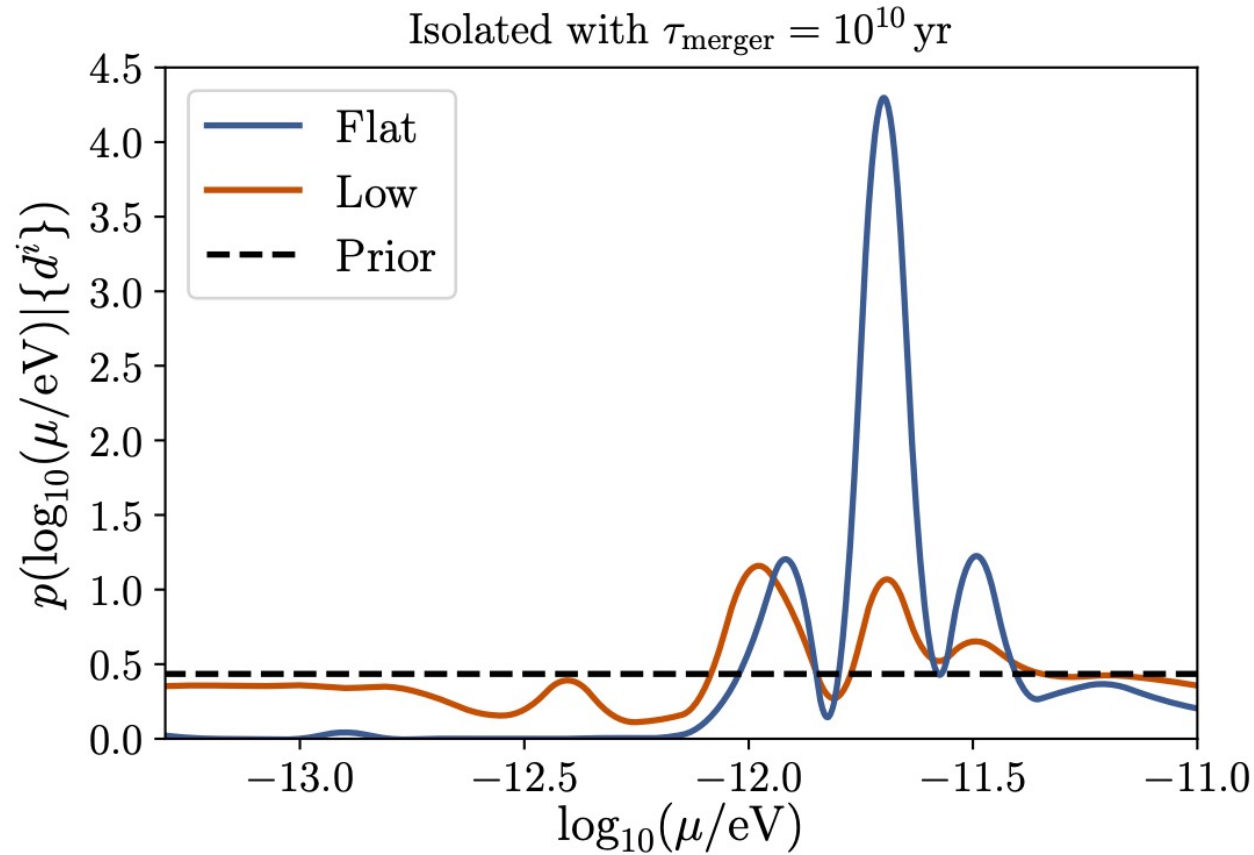
\*Fernandez, Ghalsasy, Profumo, 1911.07862

# What else could fake a low-spin PBH? Super-radiance!

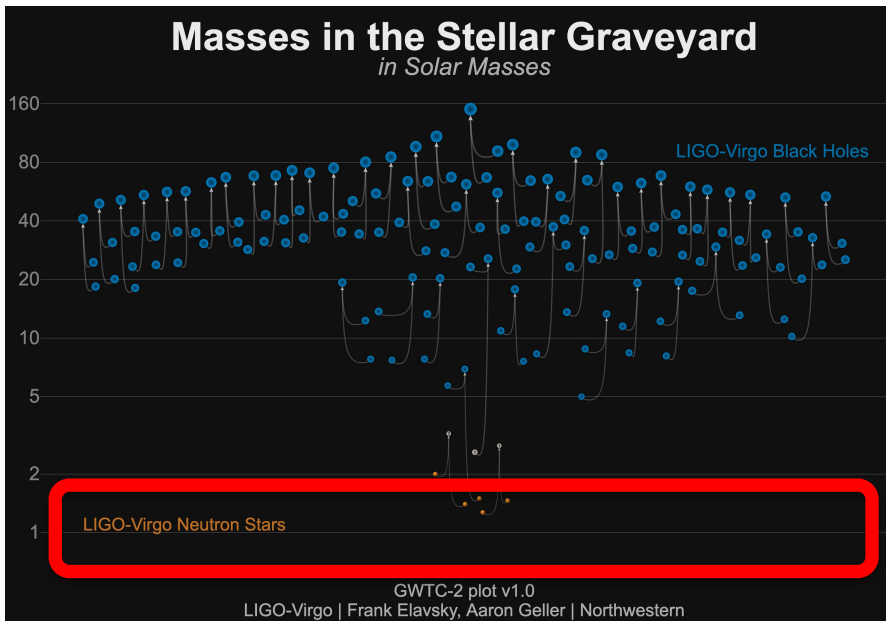


...but these are **massive**, so high- $l$  is **non-super-radiant!**

# What else could fake a low-spin PBH? Super-radiance!



**Posterior Probability** for ALP mass



sub-“Stellar-Mass”  
(<10<sup>33</sup> g)  
Black Holes

10<sup>-3</sup> g

10<sup>7</sup> g

10<sup>17</sup> g

10<sup>27</sup> g

10<sup>37</sup> g

10<sup>30</sup> eV

10<sup>40</sup> eV

10<sup>50</sup> eV

10<sup>60</sup> eV

10<sup>70</sup> eV

✓ Is there an **unmistakable signature** for PBH as DM?



Yes! **BH merger** with a **sub-Chandrasekhar** mass ( $1.4 M_{\text{sun}}$ )

Preliminary LIGO **search results** are out!

Given a **mass function**, one can calculate:

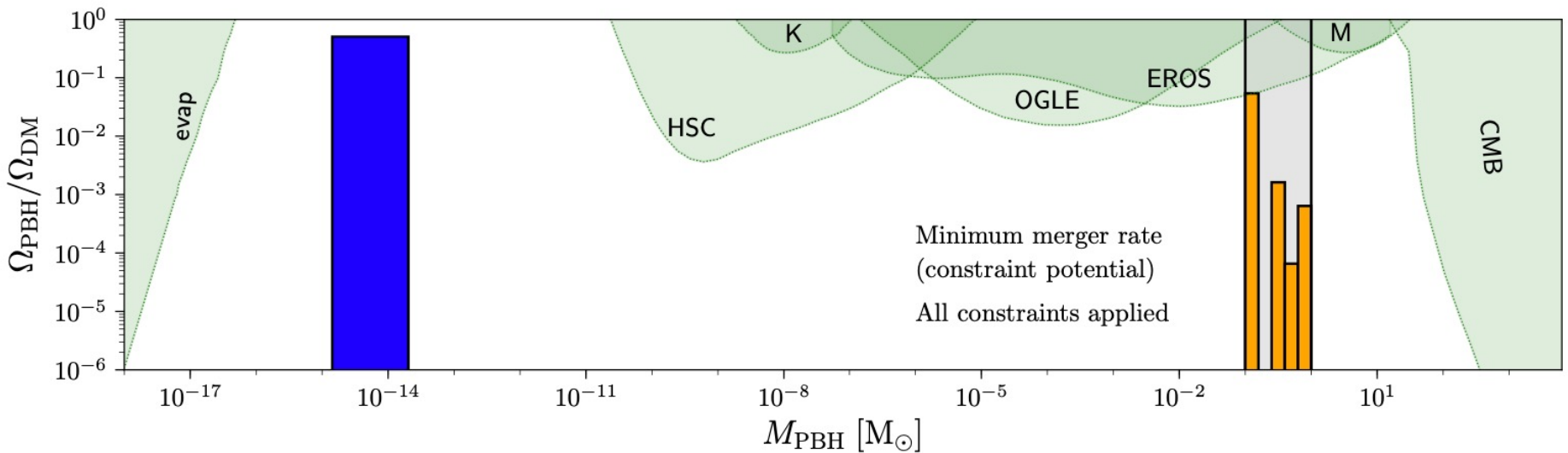
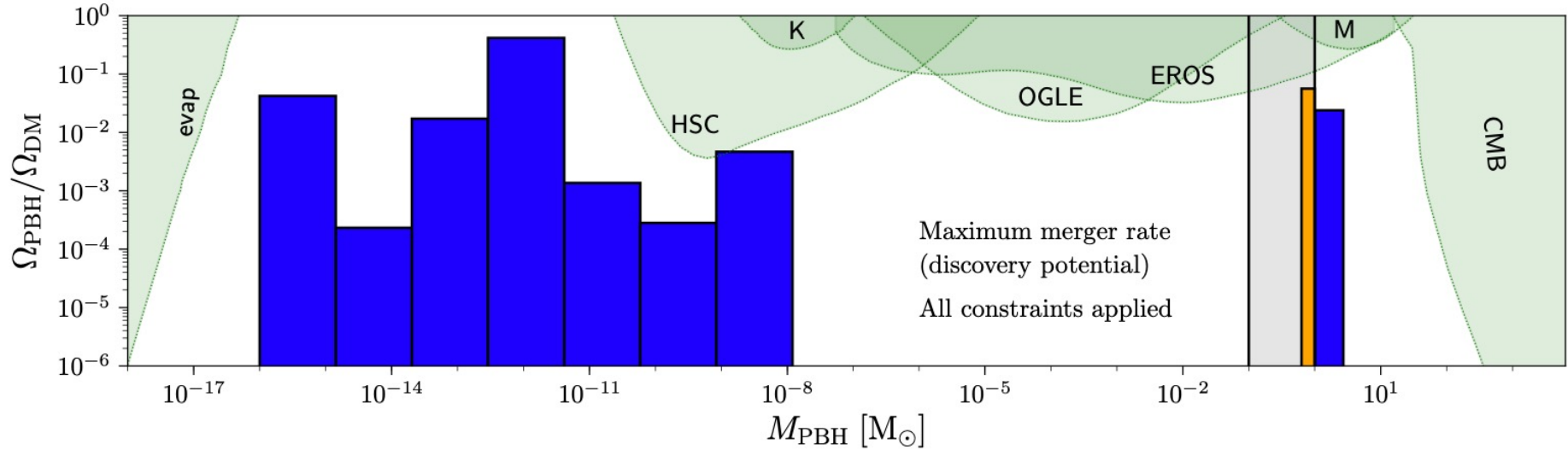
1. **Rate** of “goldilocks events”

$$R_{\text{DP}}(\psi) = \int_{\text{DP}^2} dm_1 dm_2 \mathcal{R}(m_1, m_2) V_{\text{eff}}(m_1, m_2),$$

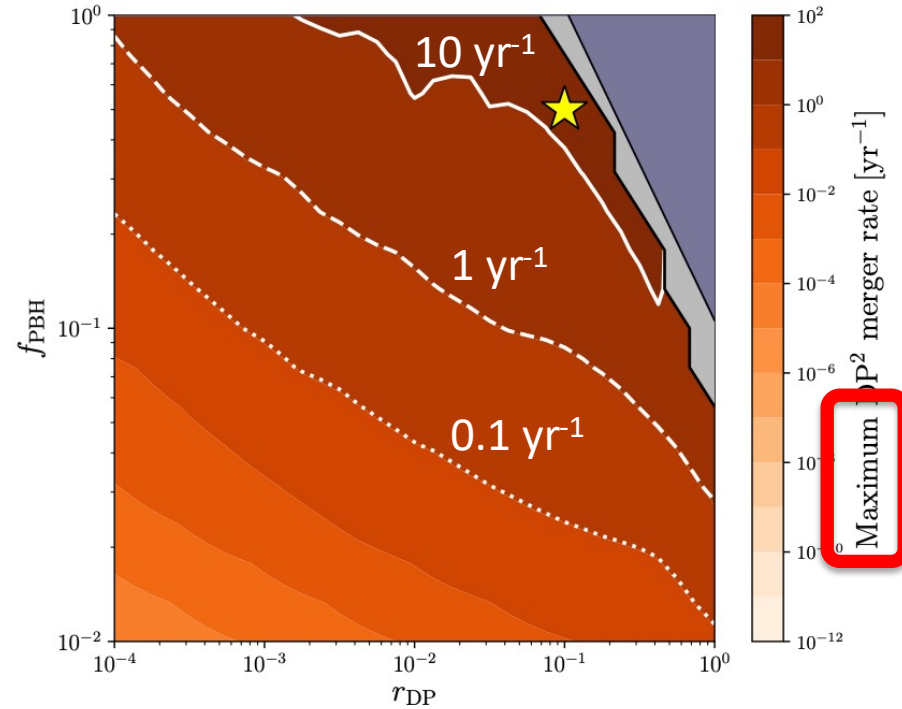
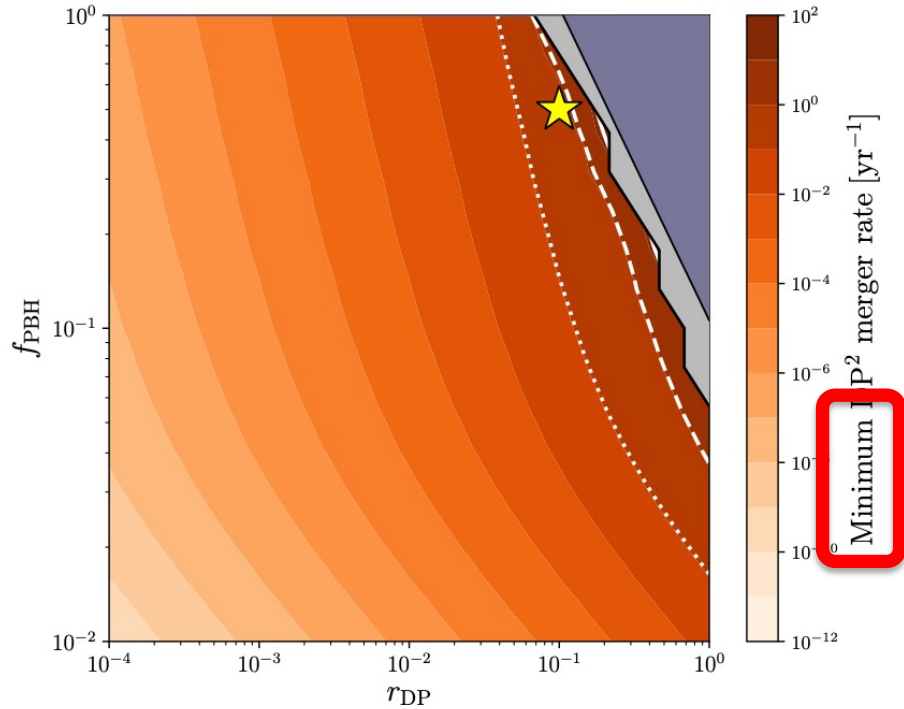
2. **Mass fraction** of **light+detectable** BHs

$$r_{\text{DP}} = \frac{1}{f_{\text{PBH}}} \int_{m_{\text{DP}}^{\text{min}}}^{m_{\text{DP}}^{\text{max}}} dm \psi(m).$$

# We can numerically compute the maximal and minimal possible “goldilocks event rate”



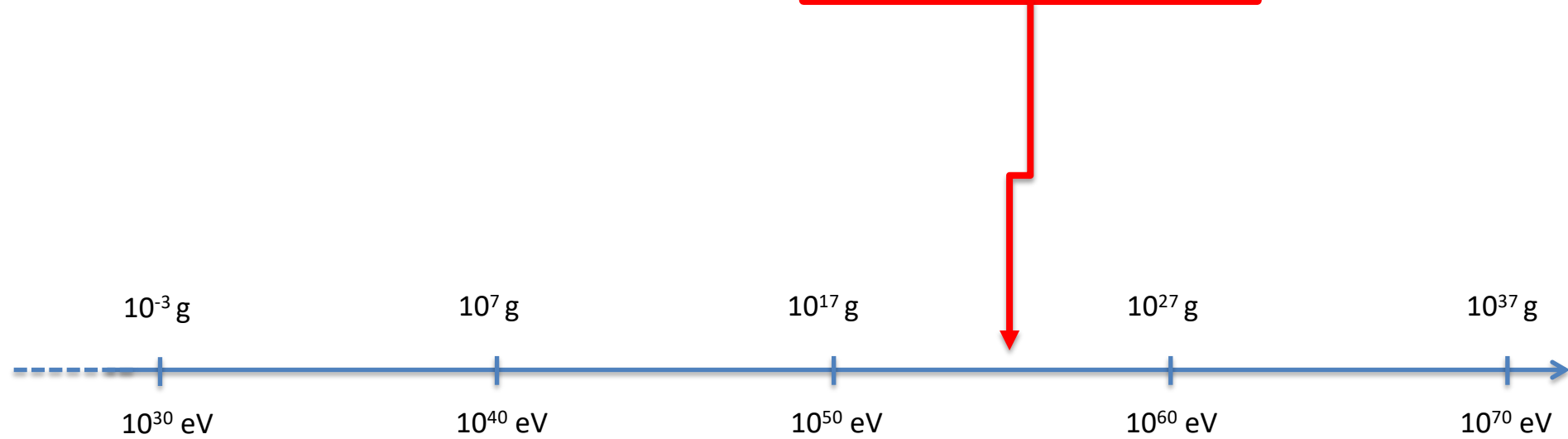
# We can numerically compute the maximal and minimal possible “goldilocks event rate”



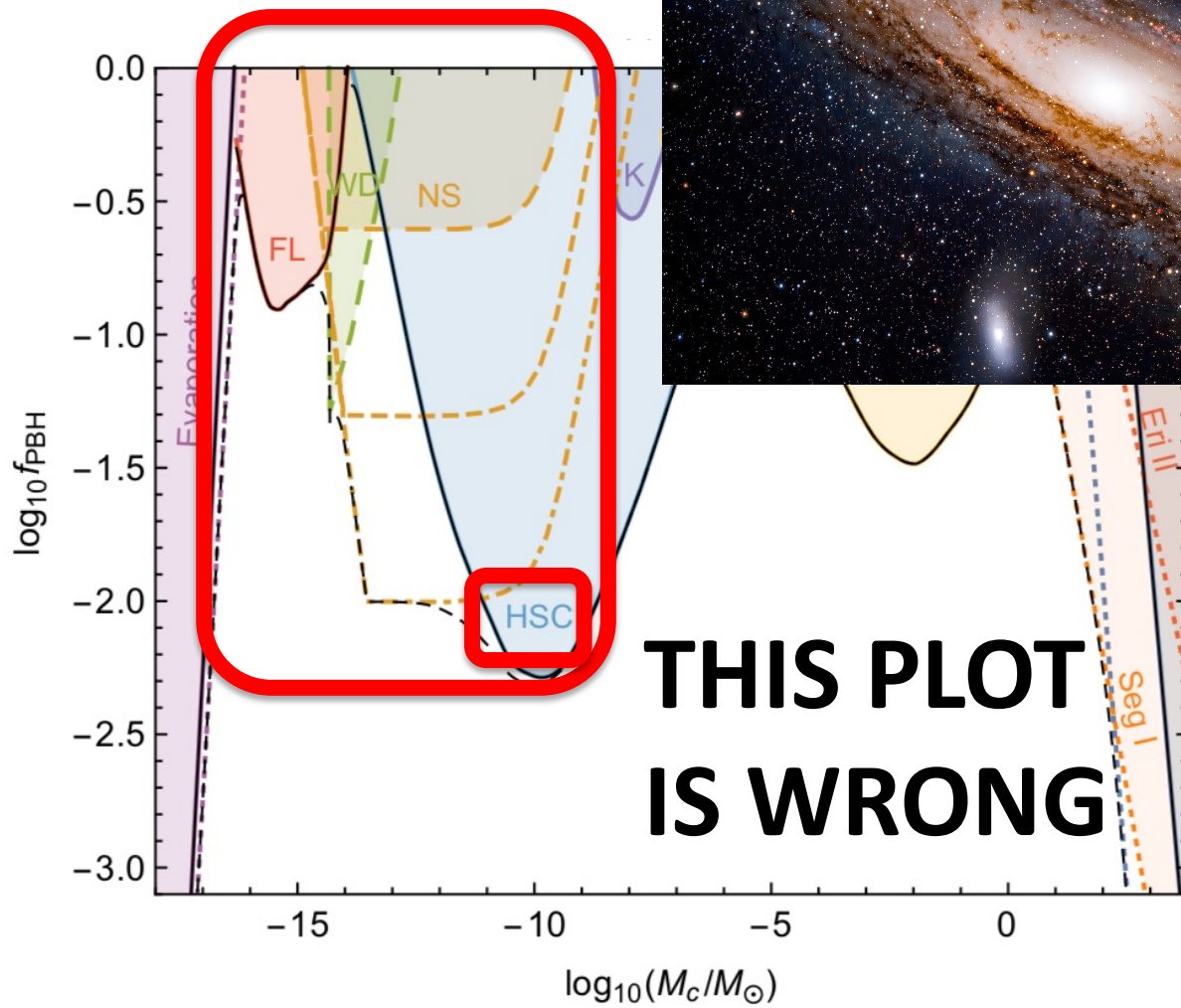
$$\int dm \psi(m) = \frac{\Omega_{\text{PBH}}}{\Omega_{\text{DM}}} \equiv f_{\text{PBH}}.$$

$$r_{\text{DP}} = \frac{1}{f_{\text{PBH}}} \int_{m_{\text{DP}}^{\text{min}}}^{m_{\text{DP}}^{\text{max}}} dm \psi(m).$$

**“Asteroid-Mass”  
( $10^{22}$  g)  
Black Holes**

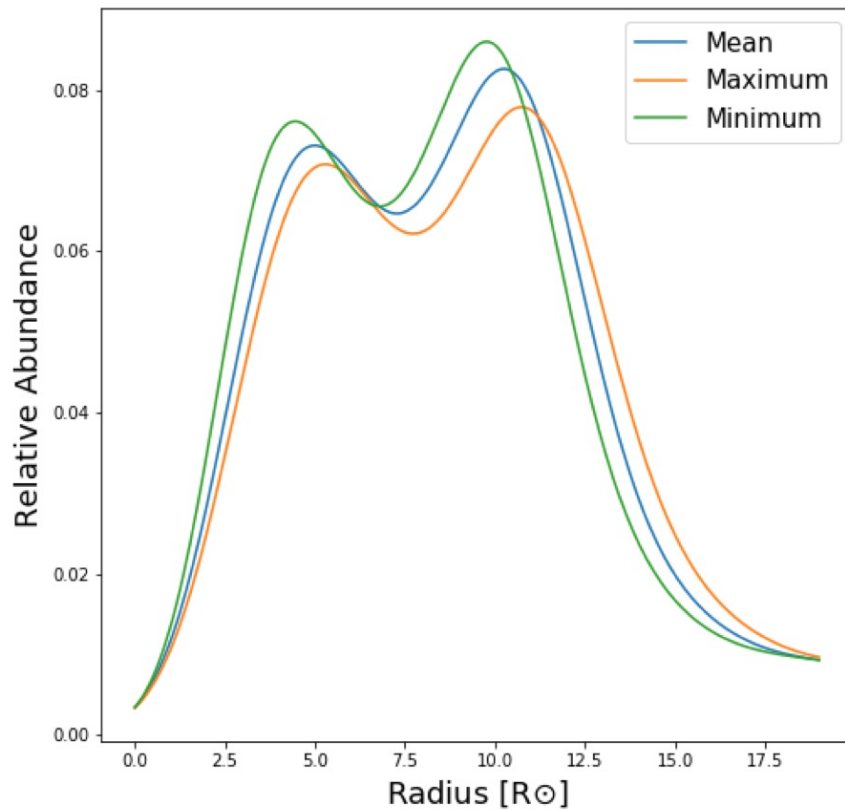


✓ **Microensing** a lot **trickier**  
than previously thought!



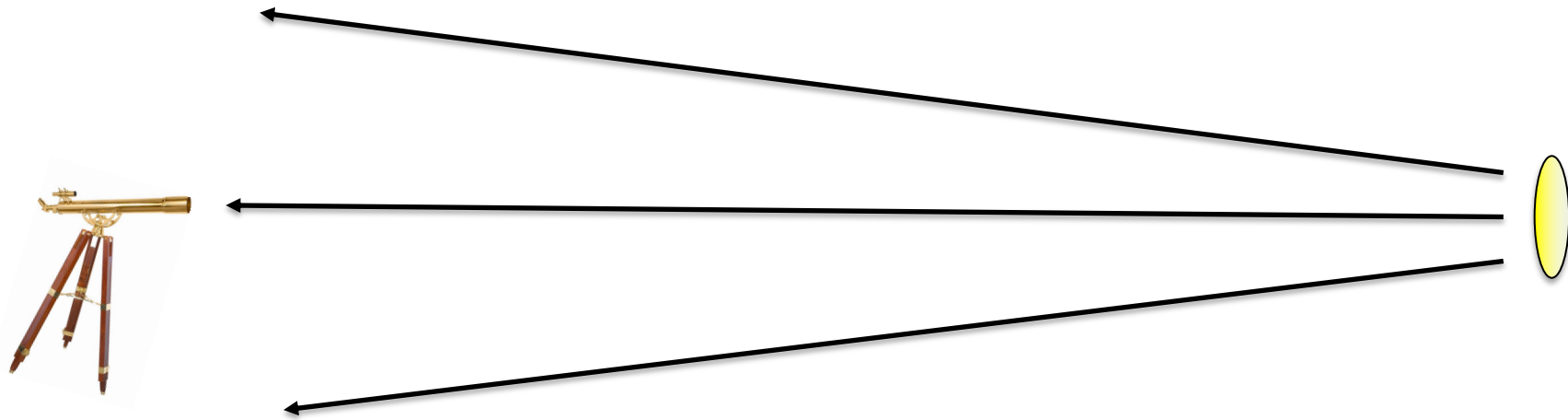
**THIS PLOT  
IS WRONG**

HSC study assumes **all stars** in M31 are Sun-like...  
but Sun-like stars are **too dim** for HSC!

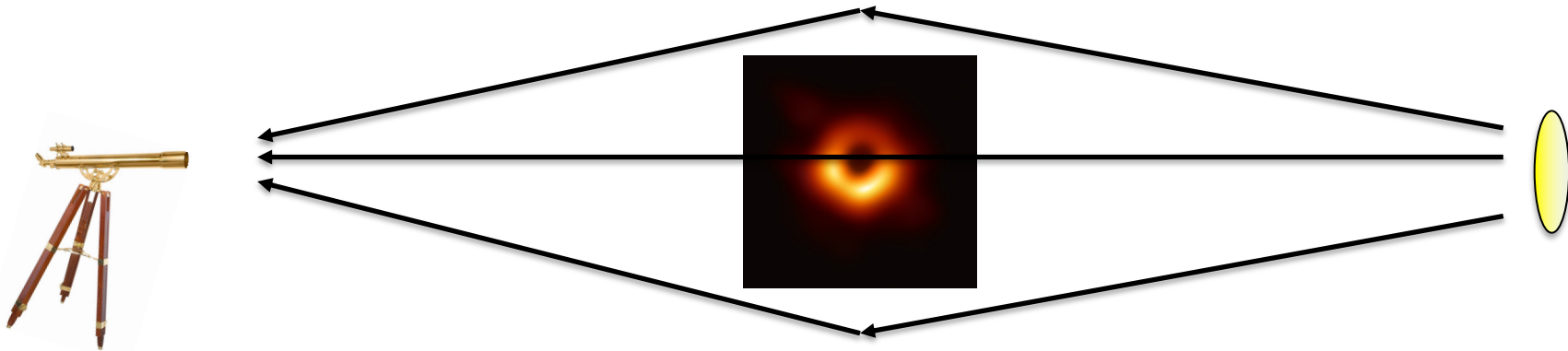


Stars that contribute to the microlensing constraints are **~ 100x larger in the sky** than the Sun!

The bigger the star, the more important  
finite-**source-size** effects!

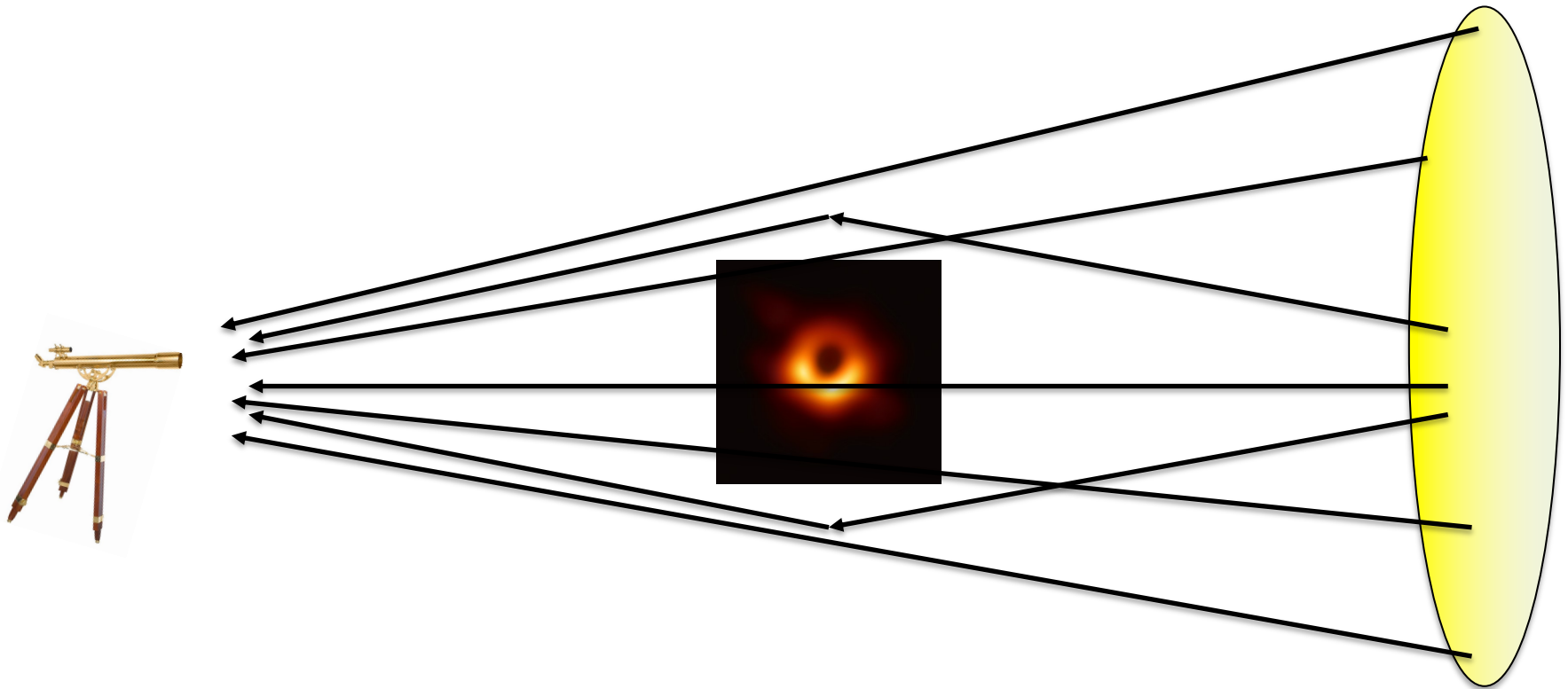


The bigger the star, the more important  
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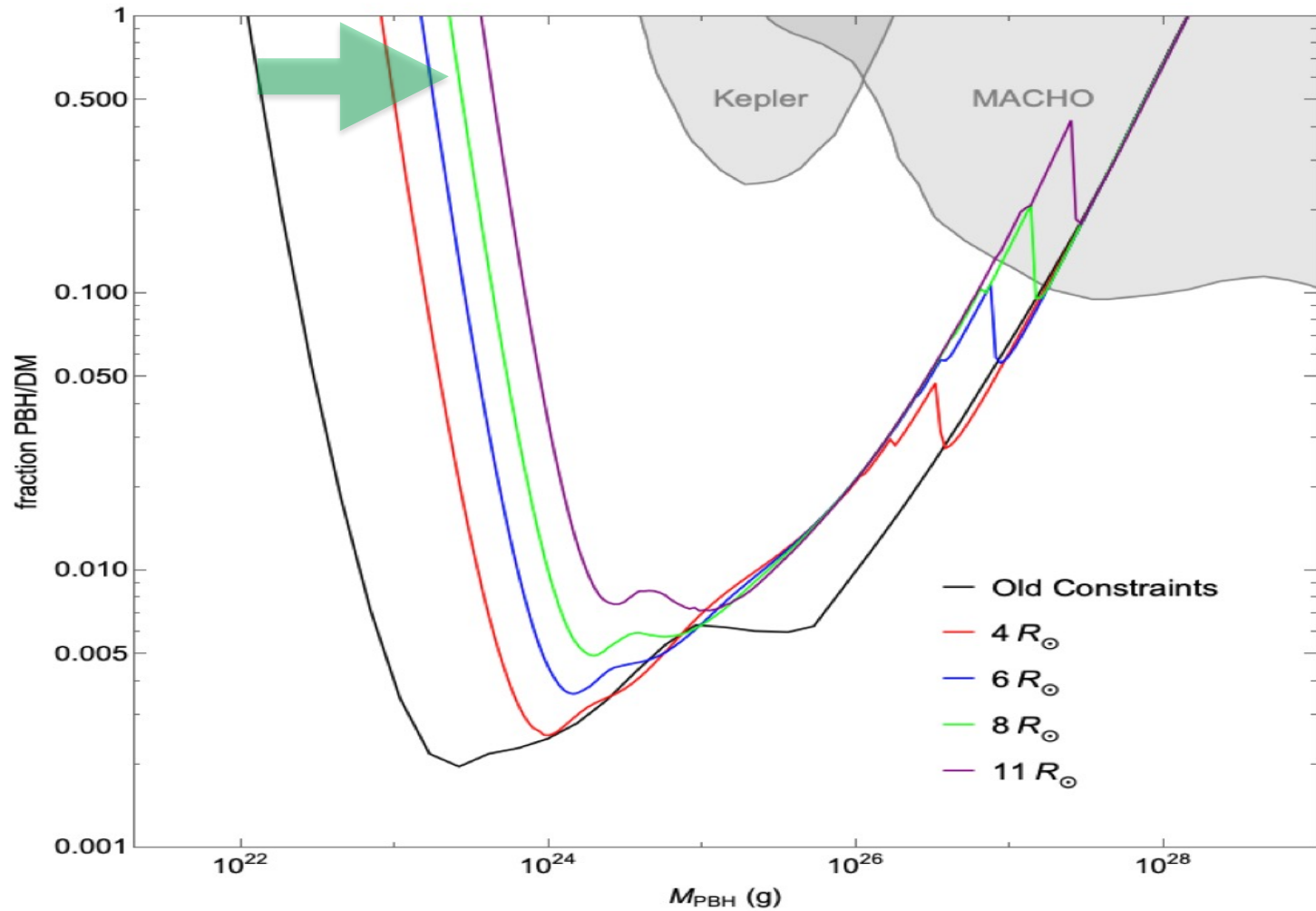


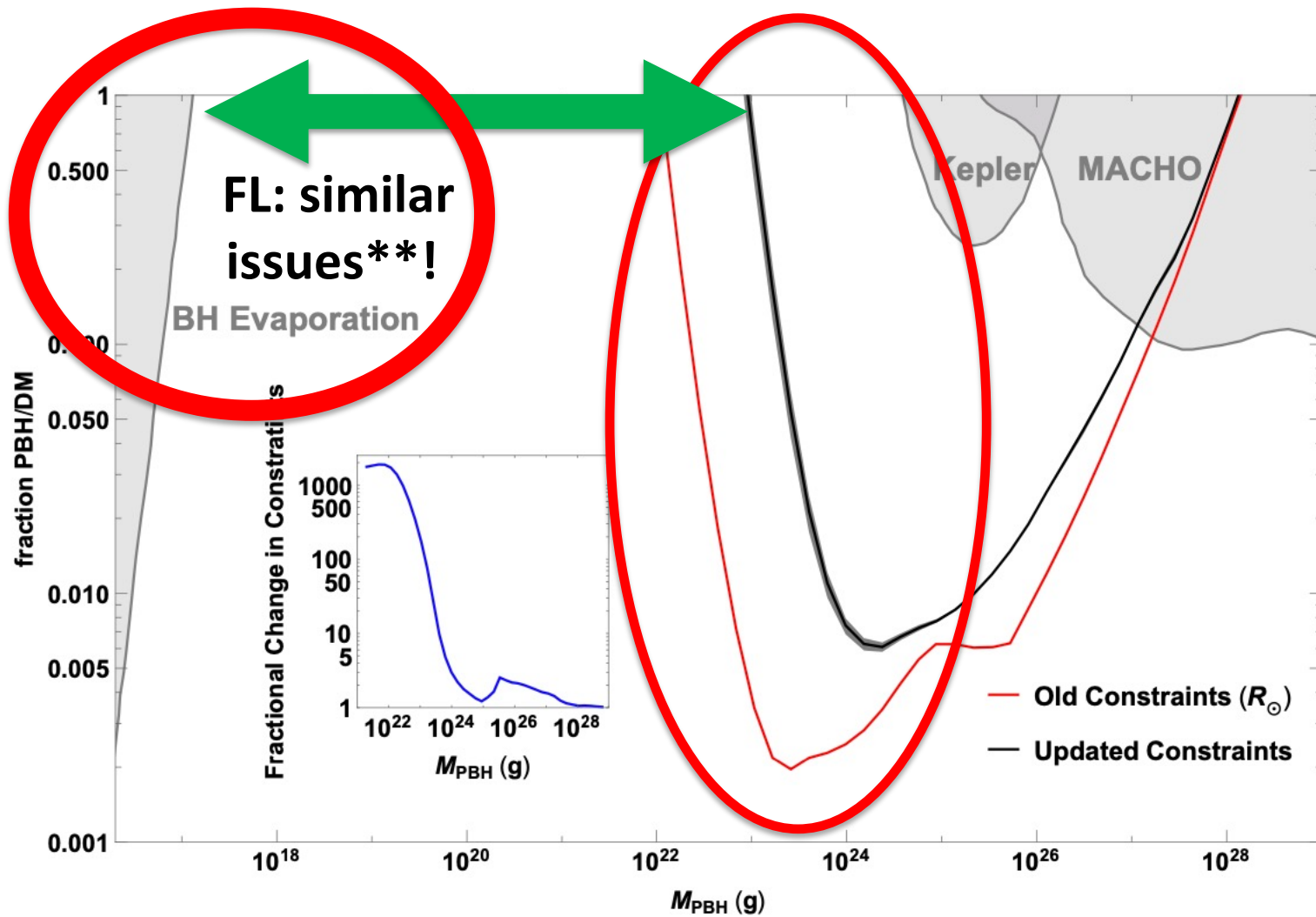


The bigger the star, the more important  
finite-**source-size** effects!



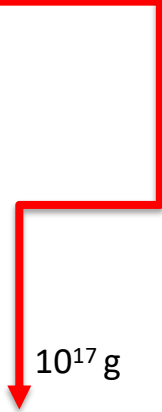
# The bigger the star, the more important finite-source-size effects!





How do we **go after** them? Capture and perturbation around **PSR**?

**“Pyramid-Mass”  
( $10^{16}$  g)  
“Evanescent” Black Holes**



$10^{-3}$  g

$10^7$  g

$10^{17}$  g

$10^{27}$  g

$10^{37}$  g

$10^{30}$  eV

$10^{40}$  eV

$10^{50}$  eV

$10^{60}$  eV

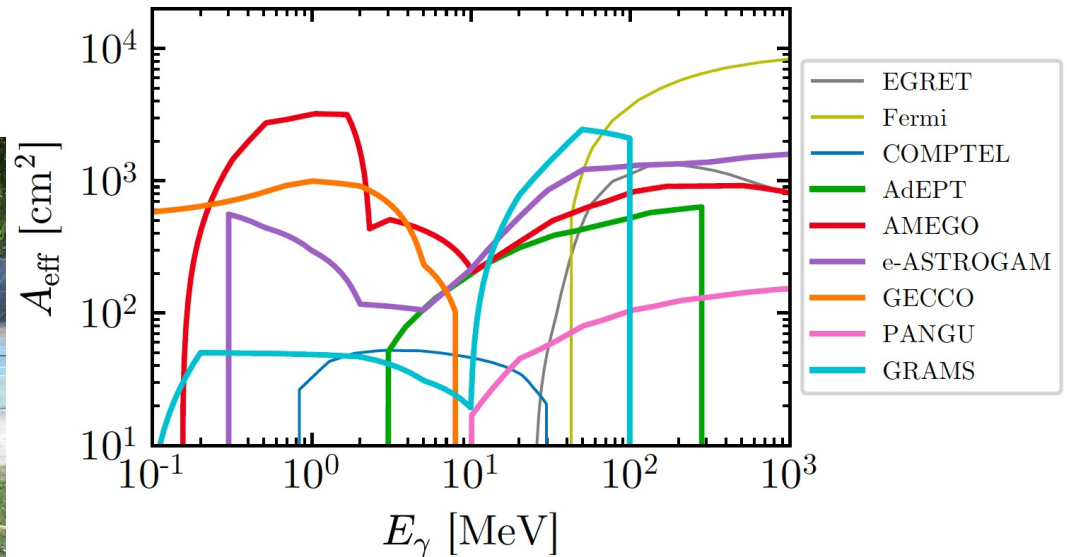
$10^{70}$  eV

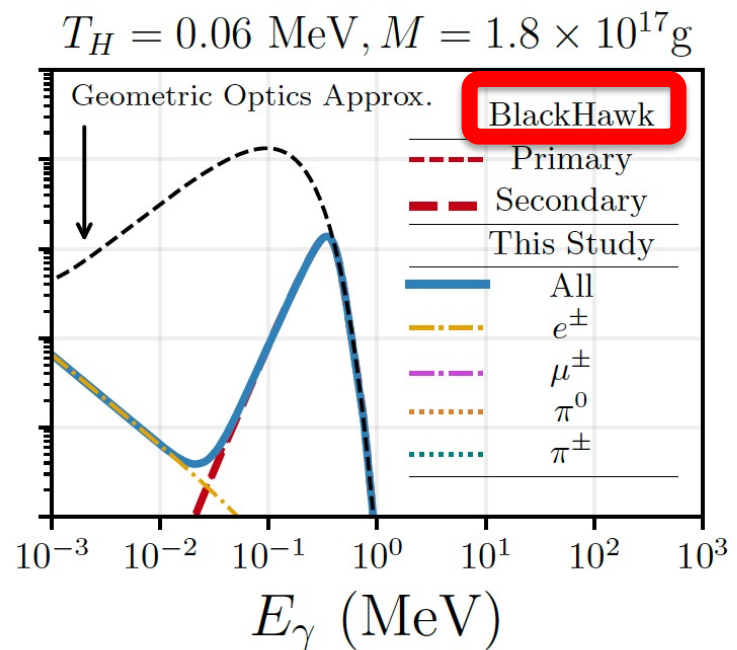
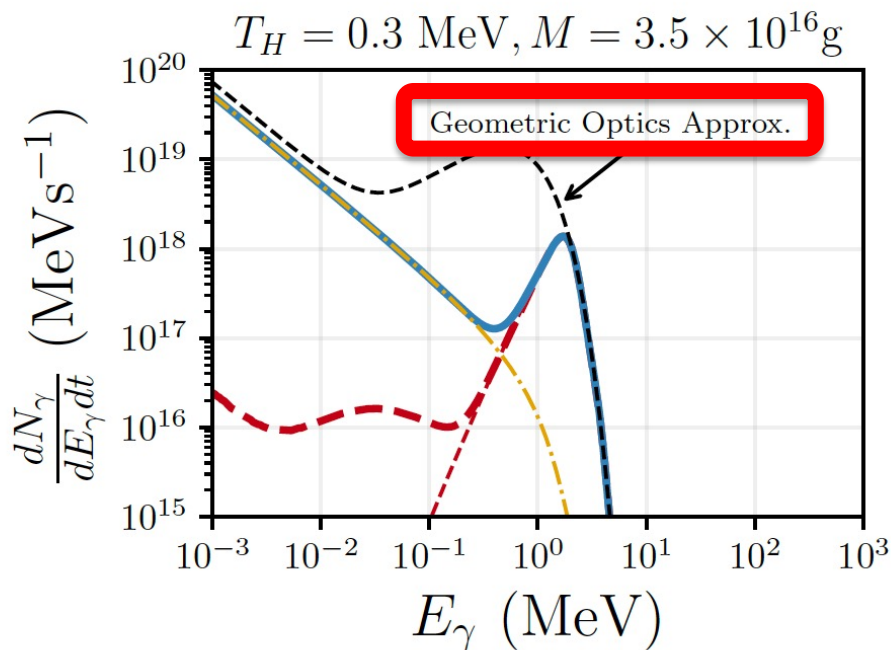
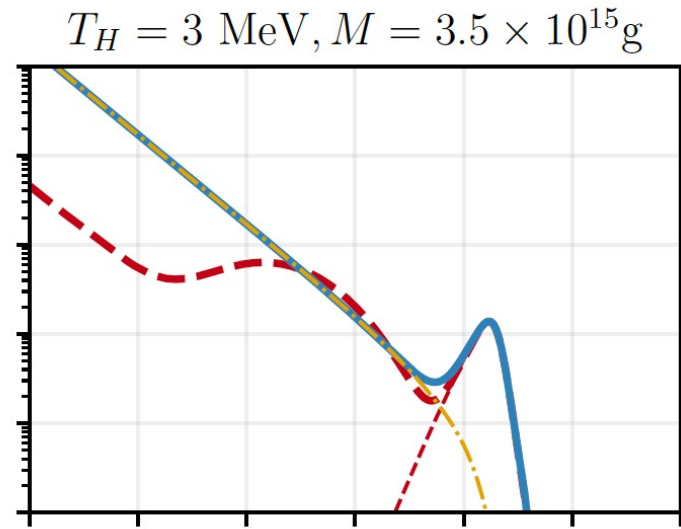
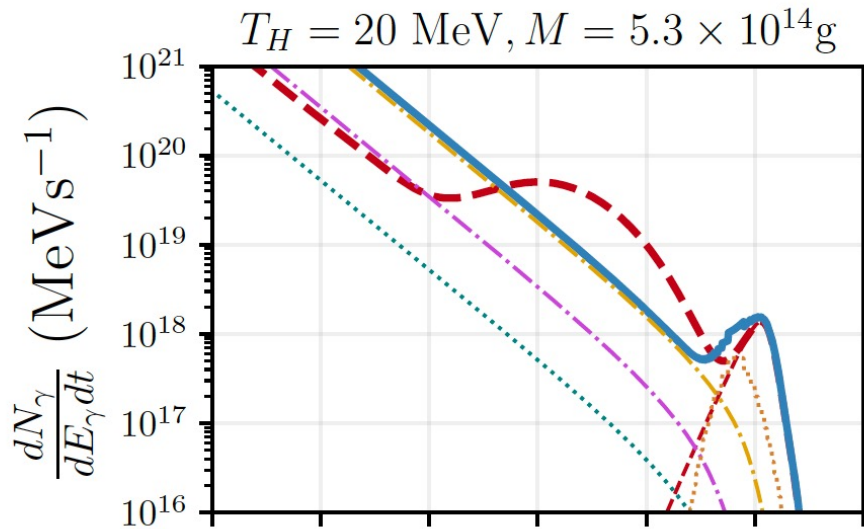


# Lightest PBH that can be dark matter...

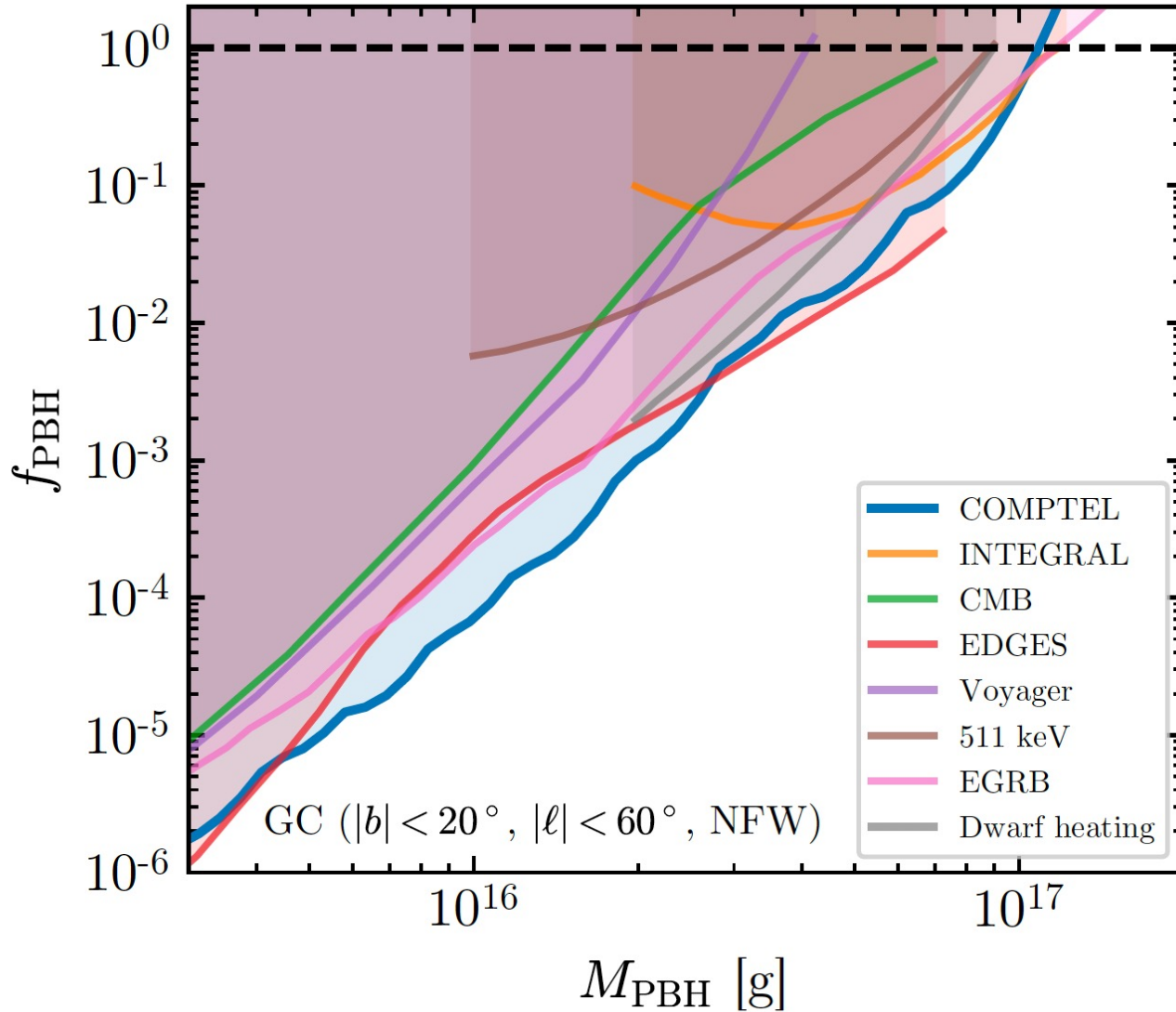
$$\tau(M) \simeq 200 \tau_U \left( \frac{M}{10^{15} \text{ g}} \right)^3 \simeq 200 \tau_U \left( \frac{10 \text{ MeV}}{T_H} \right)^3$$

- are  $\sim$  asteroid/comet/**PYRAMID** mass
- can't be much hotter than **10 MeV**

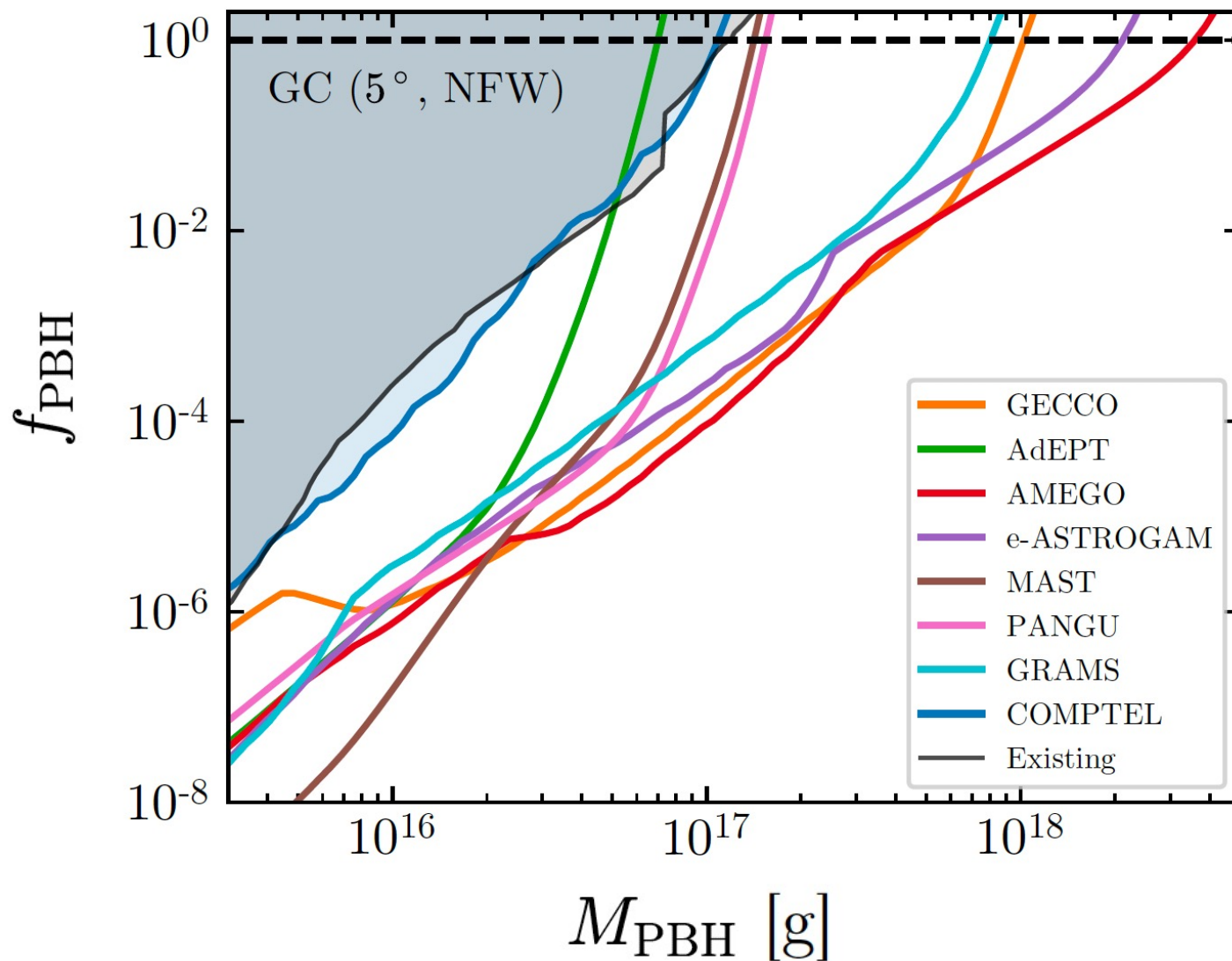




# Our new **COMPTEL** constraints are among **strongest/robust**

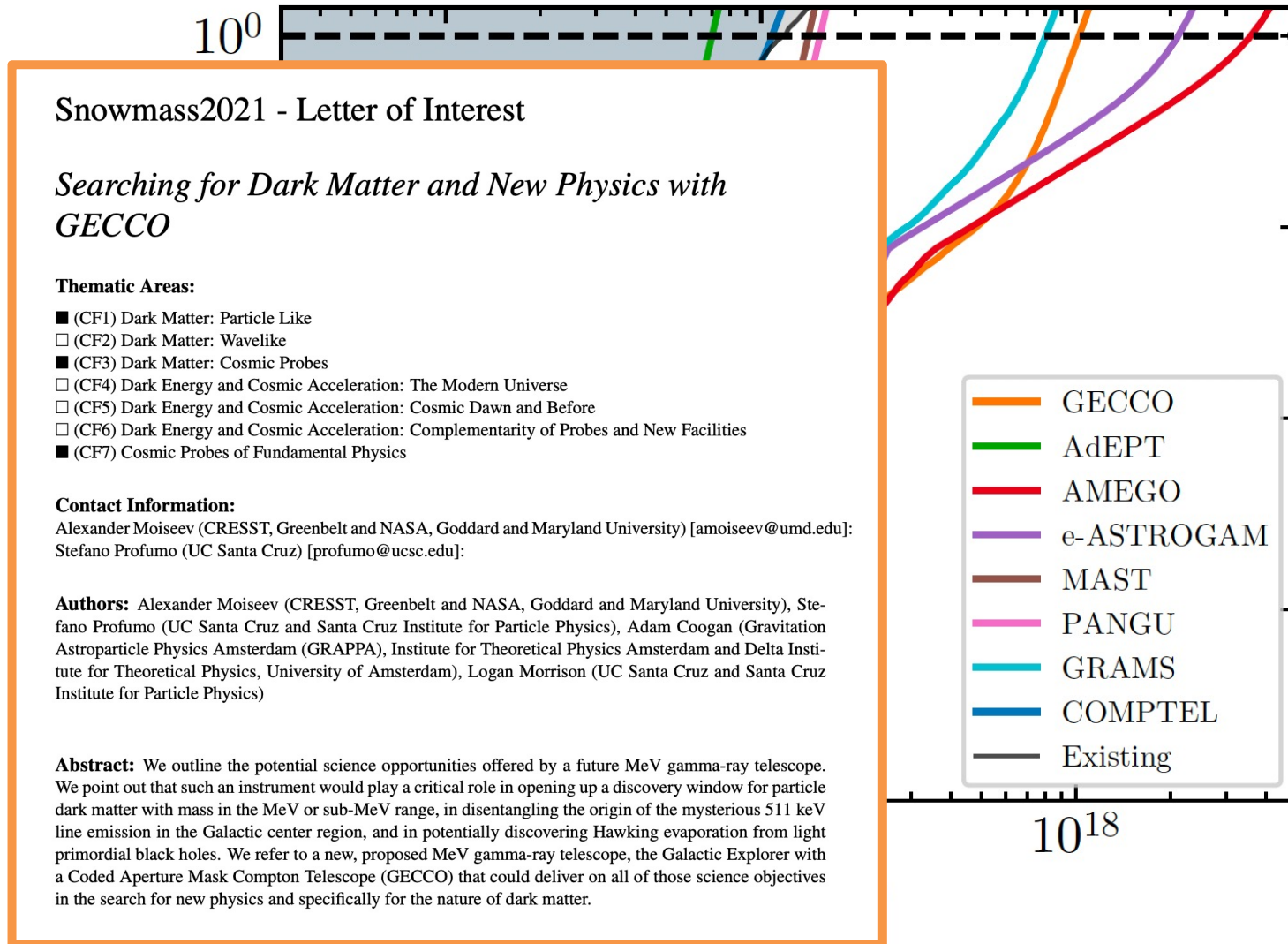


# New MeV Telescopes could discover Hawking evaporation!

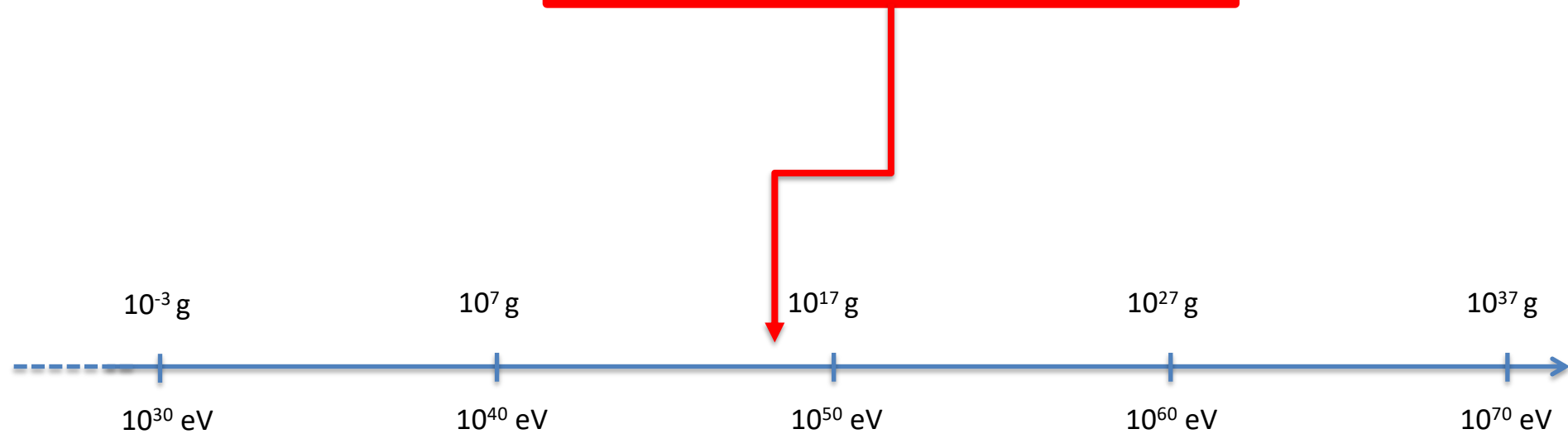




# New MeV Telescopes could discover Hawking evaporation!



**“Pyramid-Mass”  
( $10^{16}$  g)  
“Evanescent” Black Holes**

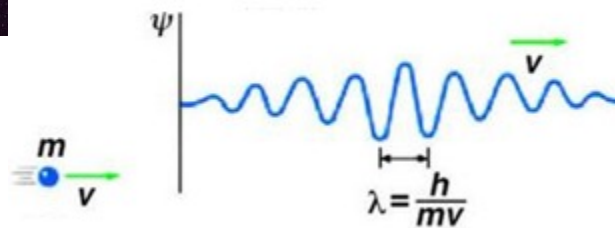
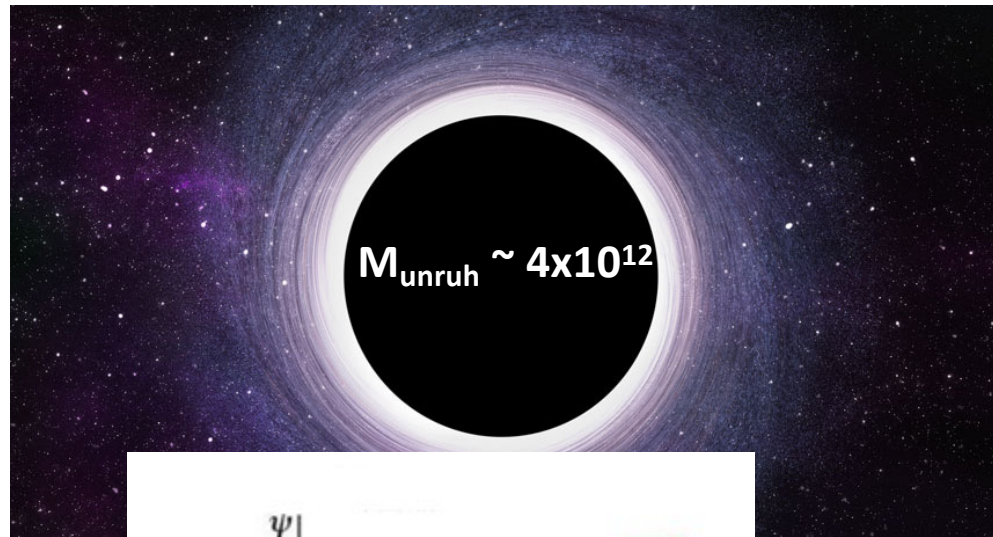


- ✓ **Best constraints: COMPTEL**
- ✓ **Future MeV telescopes**
- ✓ **NS quantum death!**

Hot off the press!!

**Neutron Star** Quantum Death by **Small** black holes

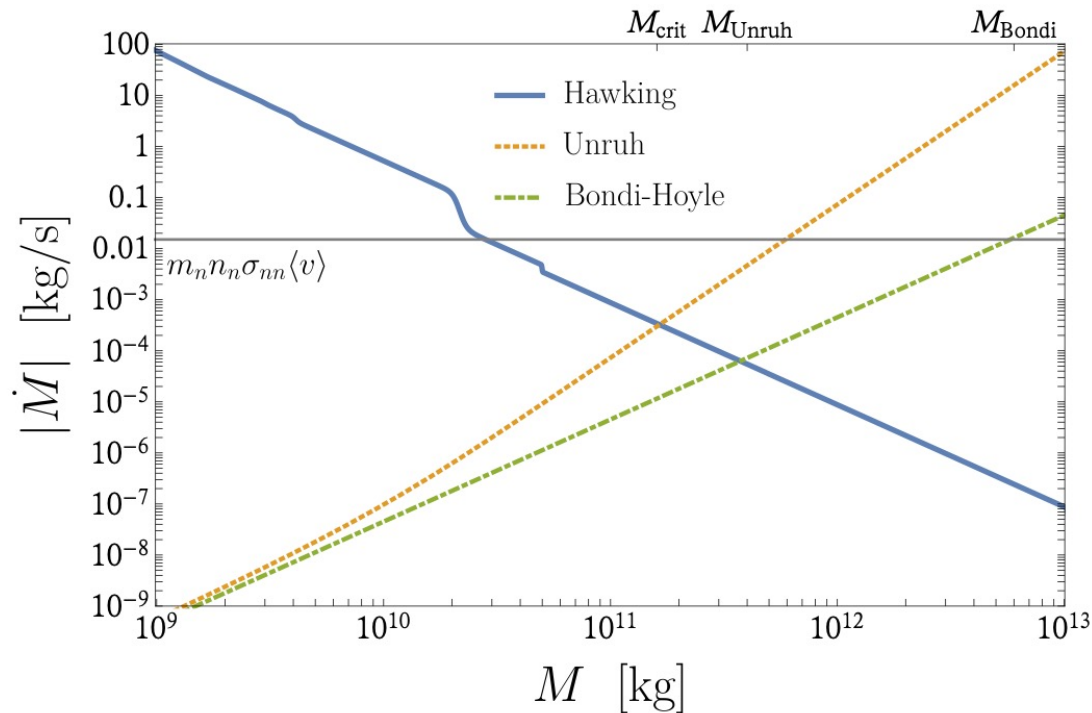
**Bondi** spherical fluid **accretion** breaks down if the accreting **black hole** has **size**  $\sim$  **neutron** de Broglie **wavelength**!



# Hot off the press!!

## Neutron Star Quantum Death by Small black holes

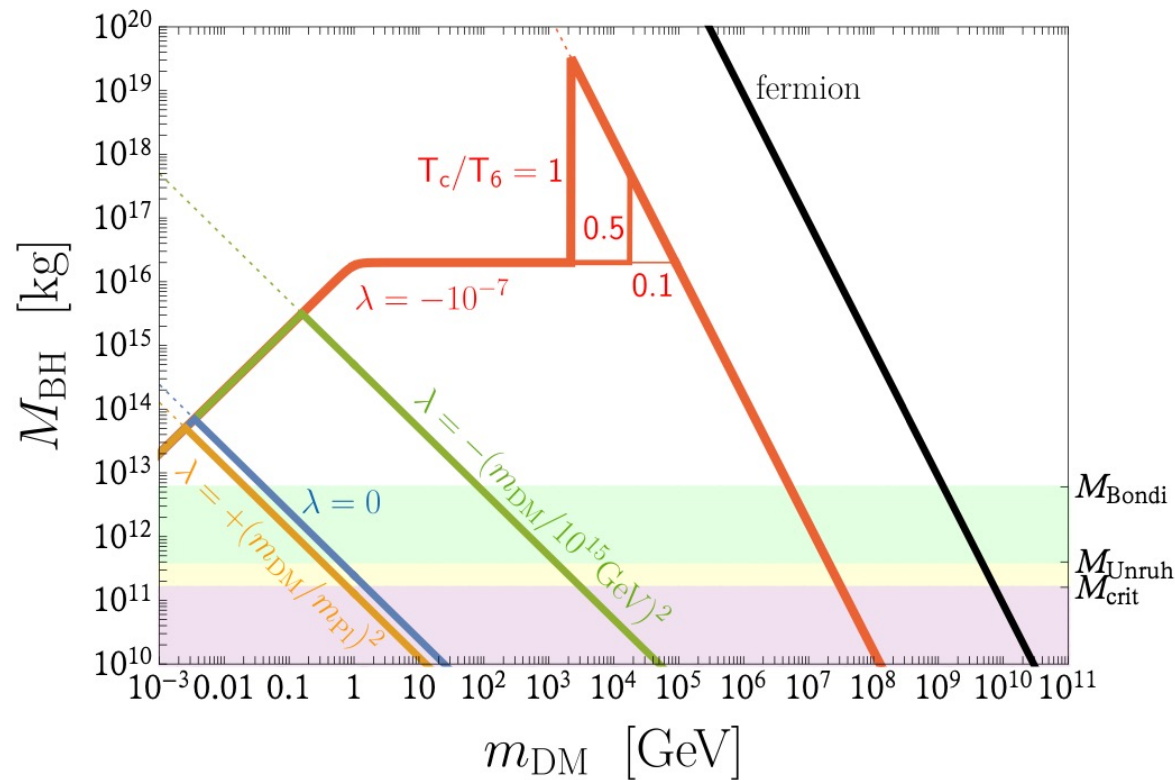
**Bondi** spherical fluid **accretion** breaks down if the accreting **black hole** has **size**  $\sim$  **neutron de Broglie wavelength!**



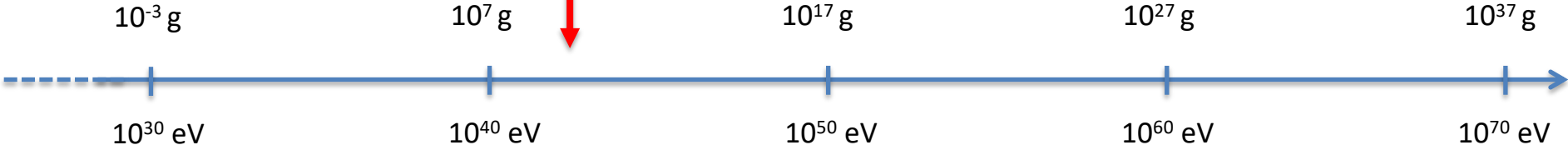
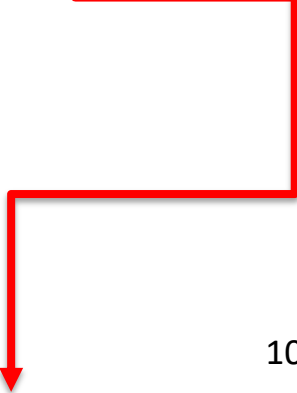
# Hot off the press!!

## Neutron Star Quantum Death by **Small** black holes

The **initial size** of the **black hole** in a NS depends on the dark matter **spin/mass/interaction** properties



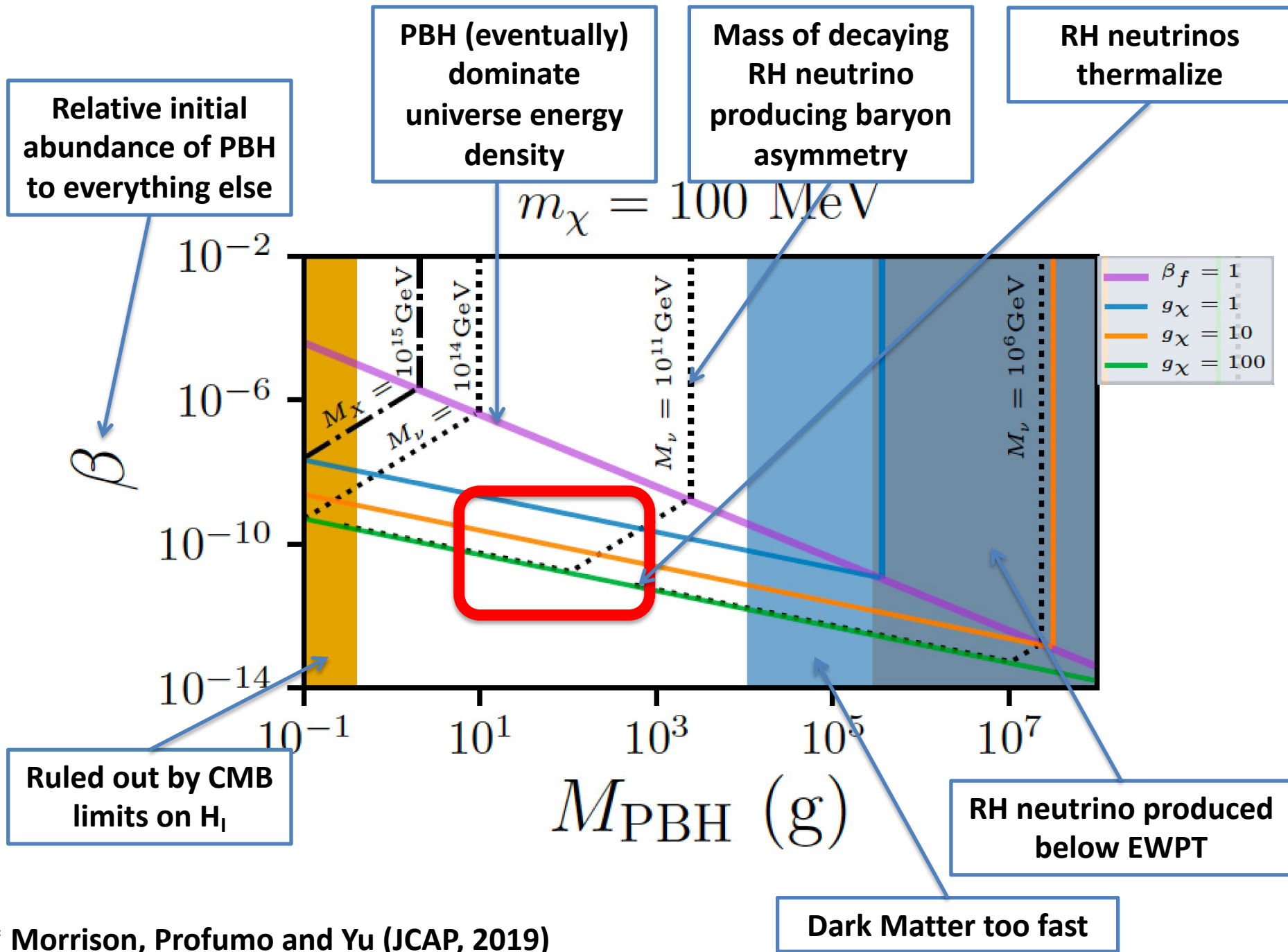
**Ton-size  
"Space-cow"  
Black Holes**



...even if PBH are **NOT** the dark matter, they can **PRODUCE** the dark matter via **Hawking evaporation!**

Mass (g)	$T_H$ (GeV)	$\tau$ (s)	$T_{\text{evap}} = T(\tau)$ (GeV)
$5M_P \simeq 10^{-4}$	$1.7 \times 10^{17}$	$10^{-41}$	$2 \times 10^{17}$
1	$1.7 \times 10^{13}$	$4 \times 10^{-29}$	$2 \times 10^{11}$
$10^3$	$1.7 \times 10^{10}$	$4 \times 10^{-20}$	$6 \times 10^6$
$10^6$	$1.7 \times 10^7$	$4 \times 10^{-11}$	200
$10^9$	$1.7 \times 10^4$	0.04	0.006
$10^{12}$	17	$4 \times 10^7 \sim 1 \text{ yr}$	$\sim 1 \text{ keV}$

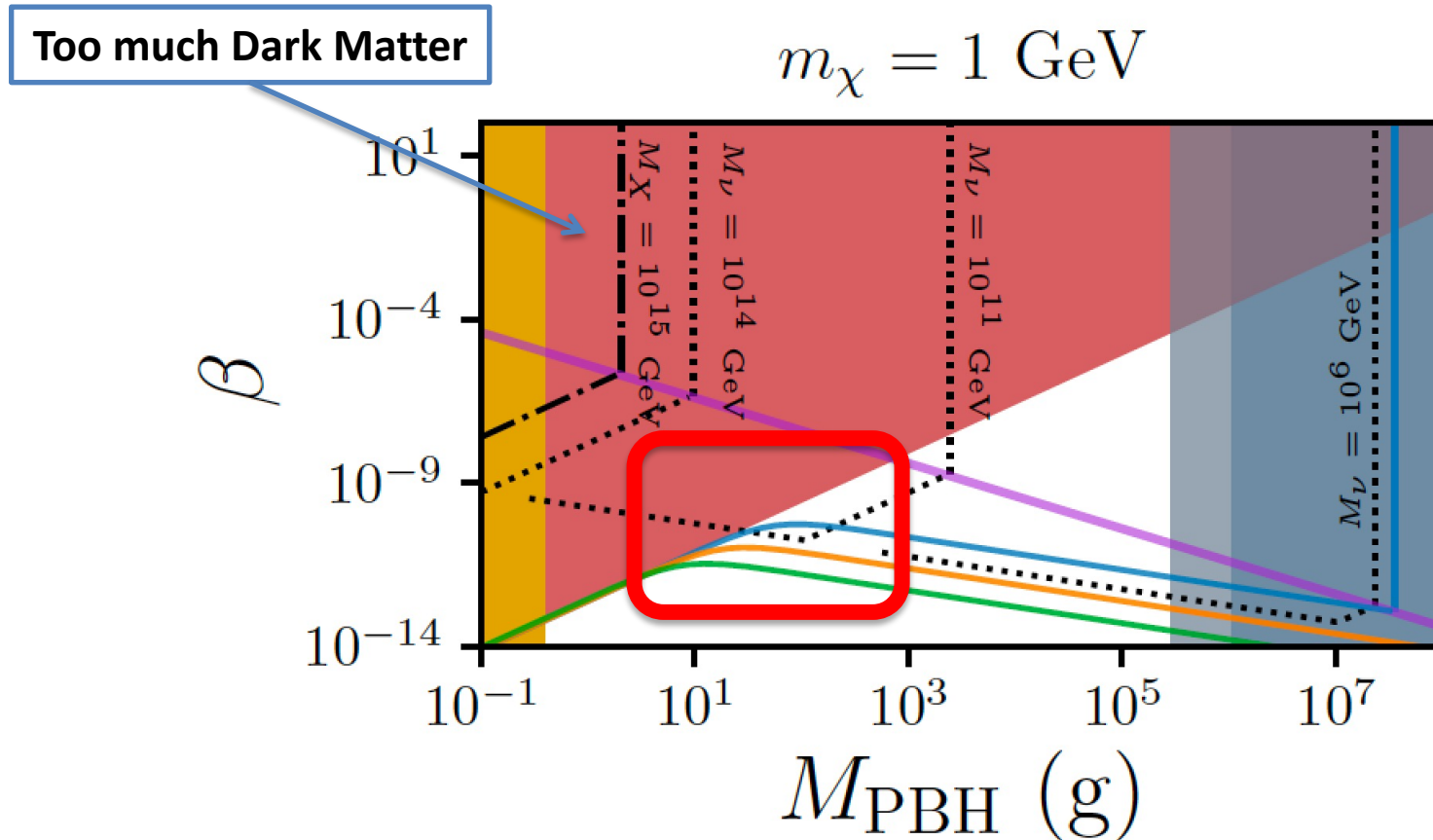
ruled out by **BBN** (more on that later!)



\* Morrison, Profumo and Yu (JCAP, 2019)



Dark Matter can be a **mix of Planck-scale relics** from PBH evaporation, and stuff the PBH **evaporated into!**



As BH approach the **Planck scale**, they can acquire a significant **relic electric charge**

(under simple **assumptions**)  $P(Q) \sim \exp(-4\pi\alpha(Q/e)^2)$   
the relic charge is  
approximately **Gaussian**\*  $(8\pi\alpha)^{-1/2} \approx 2.34$

If evaporation **stops** around the Planck scale  
(because of **extremality**, or because of **quantum gravity**)  
we are left with a population of **charged, Planck-scale relics!**

\* Page, 1977

\*\* Lehmann, Johnson, Profumo and Schwemberger, 1906.06348 (JCAP10(2019)046)

# Grain-of-Salt Black Holes

$10^{-3}$  g

$10^7$  g

$10^{17}$  g

$10^{27}$  g

$10^{37}$  g

$10^{30}$  eV

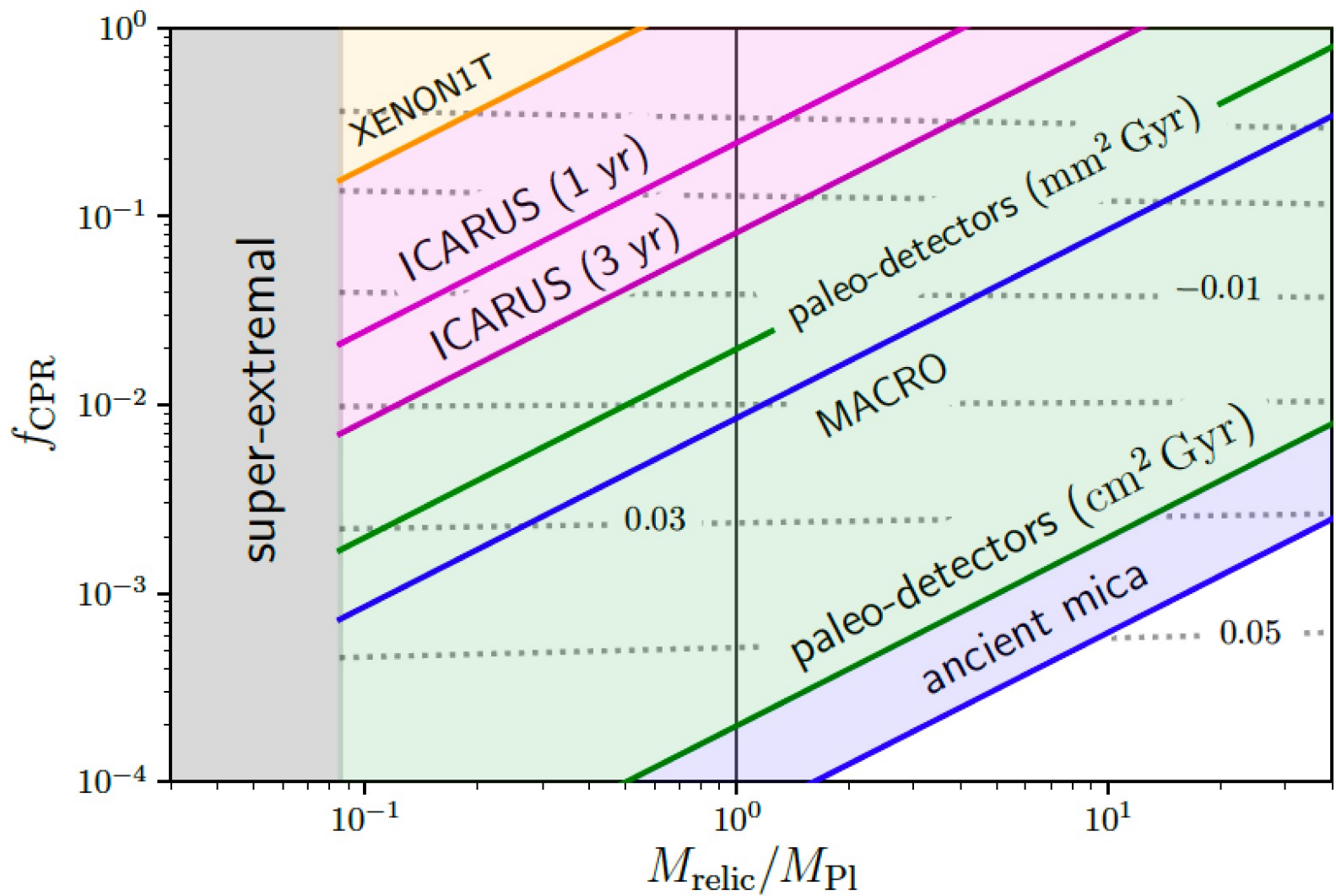
$10^{40}$  eV

$10^{50}$  eV

$10^{60}$  eV

$10^{70}$  eV





\* Lehmann, Johnson, Profumo and Schwemberger, 1906.06348 (JCAP10(2019)046)

# Hawking-Radiation Recoil of Microscopic Black Holes

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<sup>1</sup>Department of Theoretical Physics, Palacký University Olomouc, Czech Republic

## Black hole remnants are not too fast to be dark matter

Benjamin V. Lehmann<sup>1,2,\*</sup> and Stefano Profumo<sup>1,2,†</sup>

<sup>1</sup>Department of Physics, University of California Santa Cruz, 1156 High St., Santa Cruz, CA 95064, USA

<sup>2</sup>Santa Cruz Institute for Particle Physics, 1156 High St., Santa Cruz, CA 95064, USA

### Abstract

The Hawking radiation would cause black holes evaporate rapidly from many astrophysical contexts. It has been argued that the expansion of space would alter this behavior of a Planck-size black hole. The remnant left behind is a Planck-mass black hole with a section on the order of  $10^{-70} r_s^2$ .

Such black hole remnants have been identified as possible dark matter candidates. Here we argue that the final stage of the evaporation has a recoil effect which would give the microscopic black hole velocity on the order of  $10^{-1}c$  which is in disagreement with the cold dark matter cosmological model.

We comment on recent claims that recoil in the final stages of Hawking evaporation gives black hole remnants large velocities, rendering them inviable as a dark matter candidate. We point out that due to cosmic expansion, such large velocities at the final stages of evaporation are not in tension with the cold dark matter paradigm so long as they are attained at sufficiently early times. In particular, the predicted recoil velocities are robustly compatible with observations if the remnants form before the epoch of big bang nucleosynthesis, a requirement which is already imposed by the physics of nucleosynthesis itself.

of the striking difference compared to the ordinary black hole theory is that the Hawking temperature [\[9\]](#) defined to be proportional to the surface gravity at the horizon does not grow indefinitely but instead drops to zero at small but positive mass, resulting in a microscopic black hole remnant.

Black holes remnants have been considered as

**...true only if evaporation stops very late (much later than BBN), which cannot happen!**

17v1 [gr-qc] 11 Feb 2021

**“Stellar-Mass”  
( $10^{35}$  g)  
Black Holes**

$10^{-3}$  g

$10^7$  g

$10^{17}$  g

$10^{27}$  g

$10^{37}$  g

$10^{30}$  eV

$10^{40}$  eV

$10^{50}$  eV

$10^{60}$  eV

$10^{70}$  eV

- ✓ **Spins look a lot like PBH!**
- ✓ **...or maybe they are low because of superradiance?**
- ✓ **Sub-Chandrasekhar goldilocks!!**

**“Asteroid-Mass”  
( $10^{22}$  g)  
Black Holes**

$10^{-3}$  g

$10^7$  g

$10^{17}$  g

$10^{27}$  g

$10^{37}$  g

$10^{30}$  eV

$10^{40}$  eV

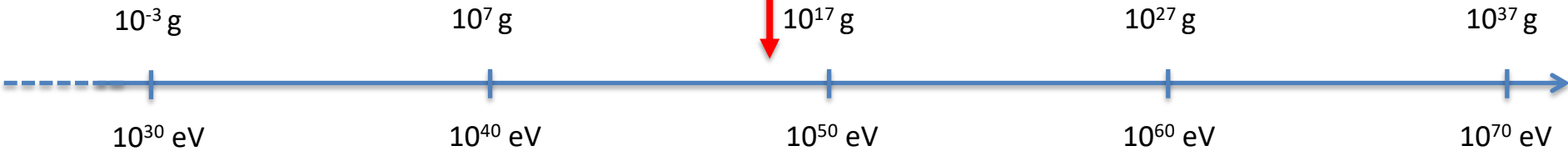
$10^{50}$  eV

$10^{60}$  eV

$10^{70}$  eV

✓ **Microensing a lot trickier  
than previously thought!**

**“Pyramid-Mass”  
( $10^{16}$  g)  
“Evanescent” Black Holes**



- ✓ **Best constraints: COMPTEL**
- ✓ **Future MeV telescopes**
- ✓ **NS quantum death!**



**Ton-size  
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$10^{-3}$  g

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$10^{50}$  eV

$10^{60}$  eV

$10^{70}$  eV

✓ **Decays can produce DM,  
BAU, Planck relics**

## Grain-of-Salt Black Holes

$10^{-3}$  g

$10^7$  g

$10^{17}$  g

$10^{27}$  g

$10^{37}$  g

$10^{30}$  eV

$10^{40}$  eV

$10^{50}$  eV

$10^{60}$  eV

$10^{70}$  eV

- ✓ Likely (partly) charged
- ✓ Detectable!
- ✓ Not too fast!

In the era of **gravitational wave** astronomy,  
the physics of **macroscopic** DM candidates  
offers many **opportunities** for the ingenuity  
of **theorists** and the craft of **observers**

$10^{-3}$  g

$10^7$  g

$10^{17}$  g

$10^{27}$  g

$10^{37}$  g

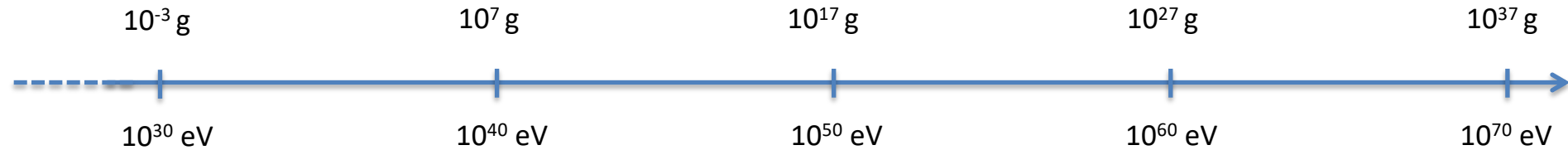
$10^{30}$  eV

$10^{40}$  eV

$10^{50}$  eV

$10^{60}$  eV

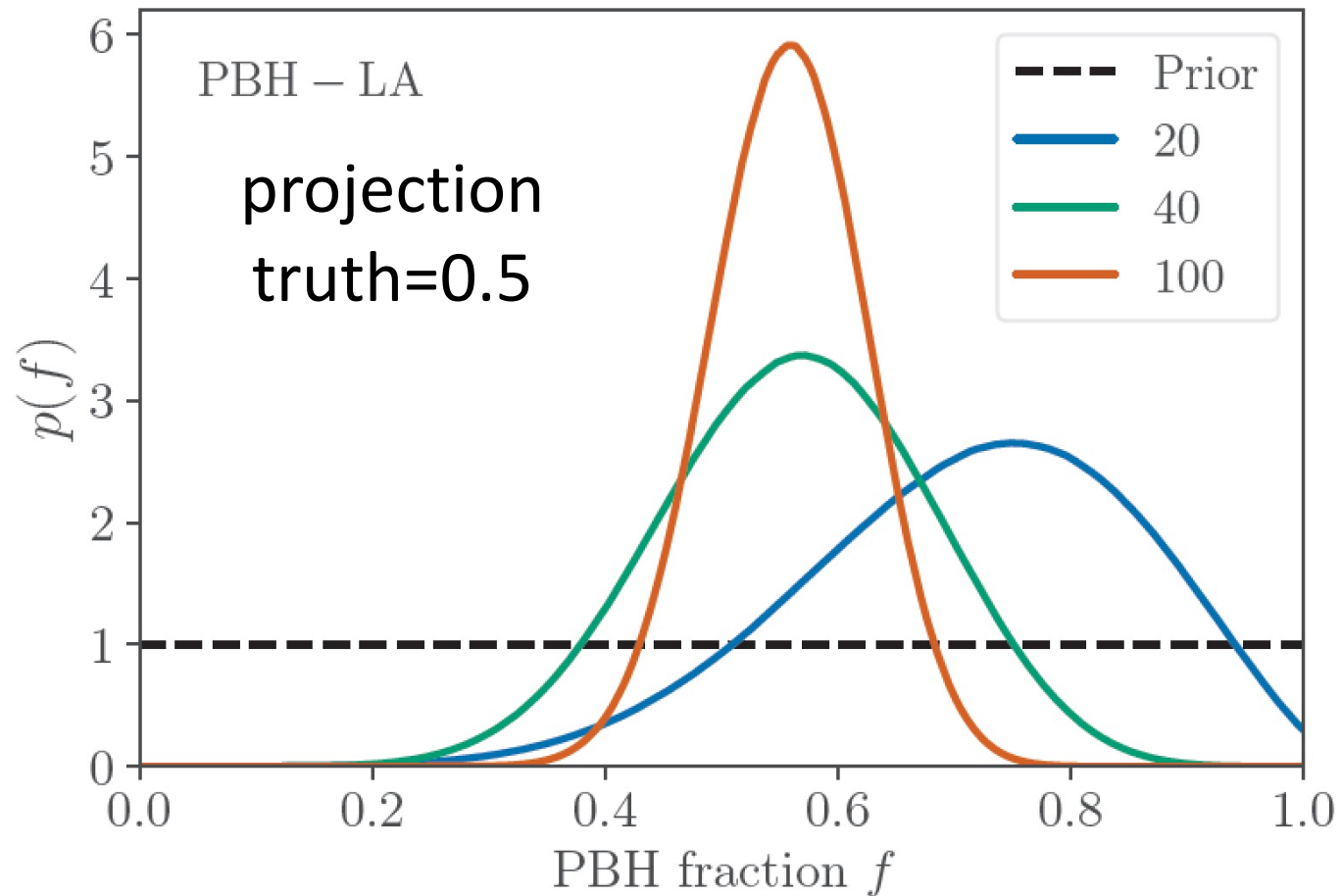
$10^{70}$  eV

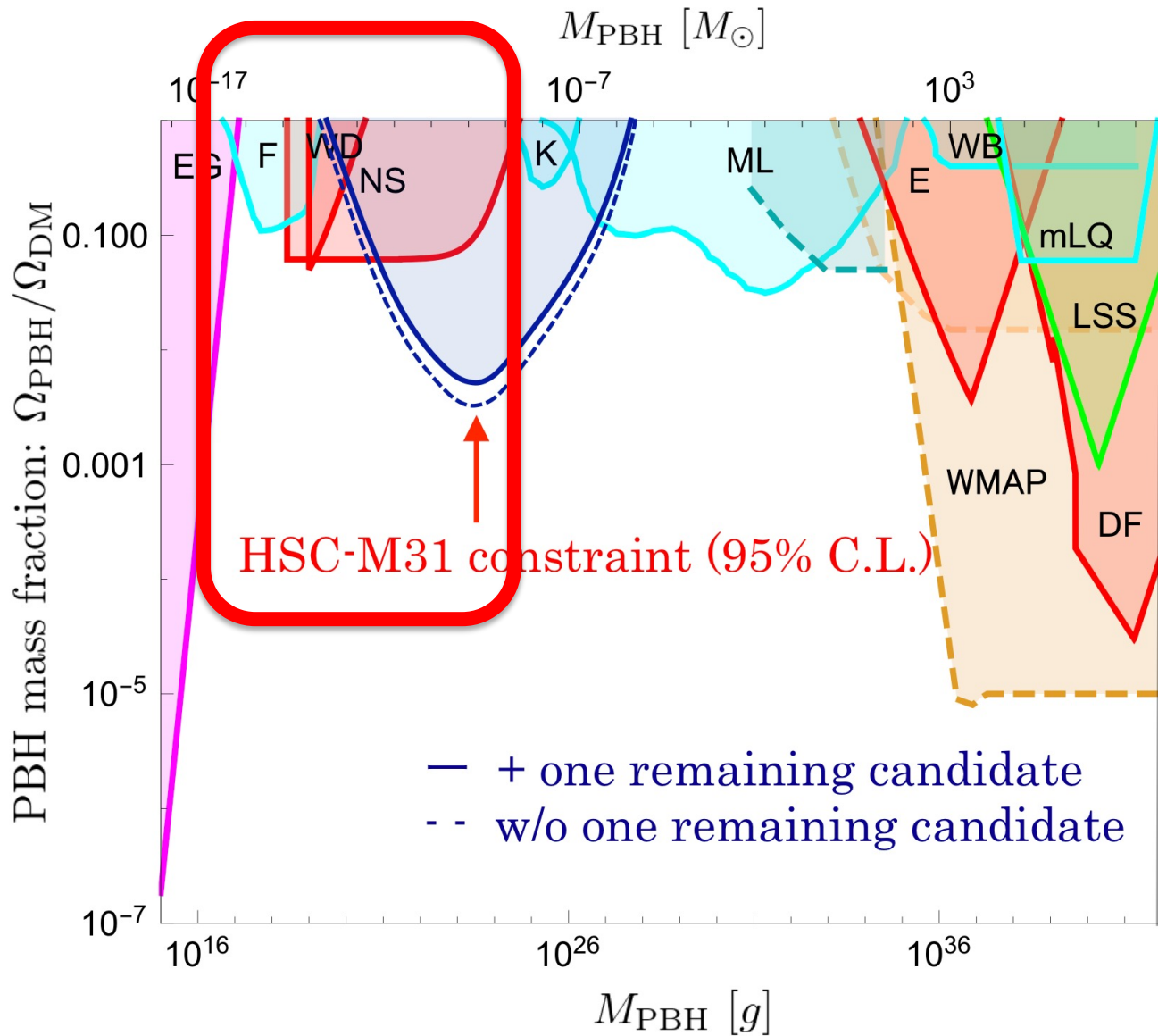




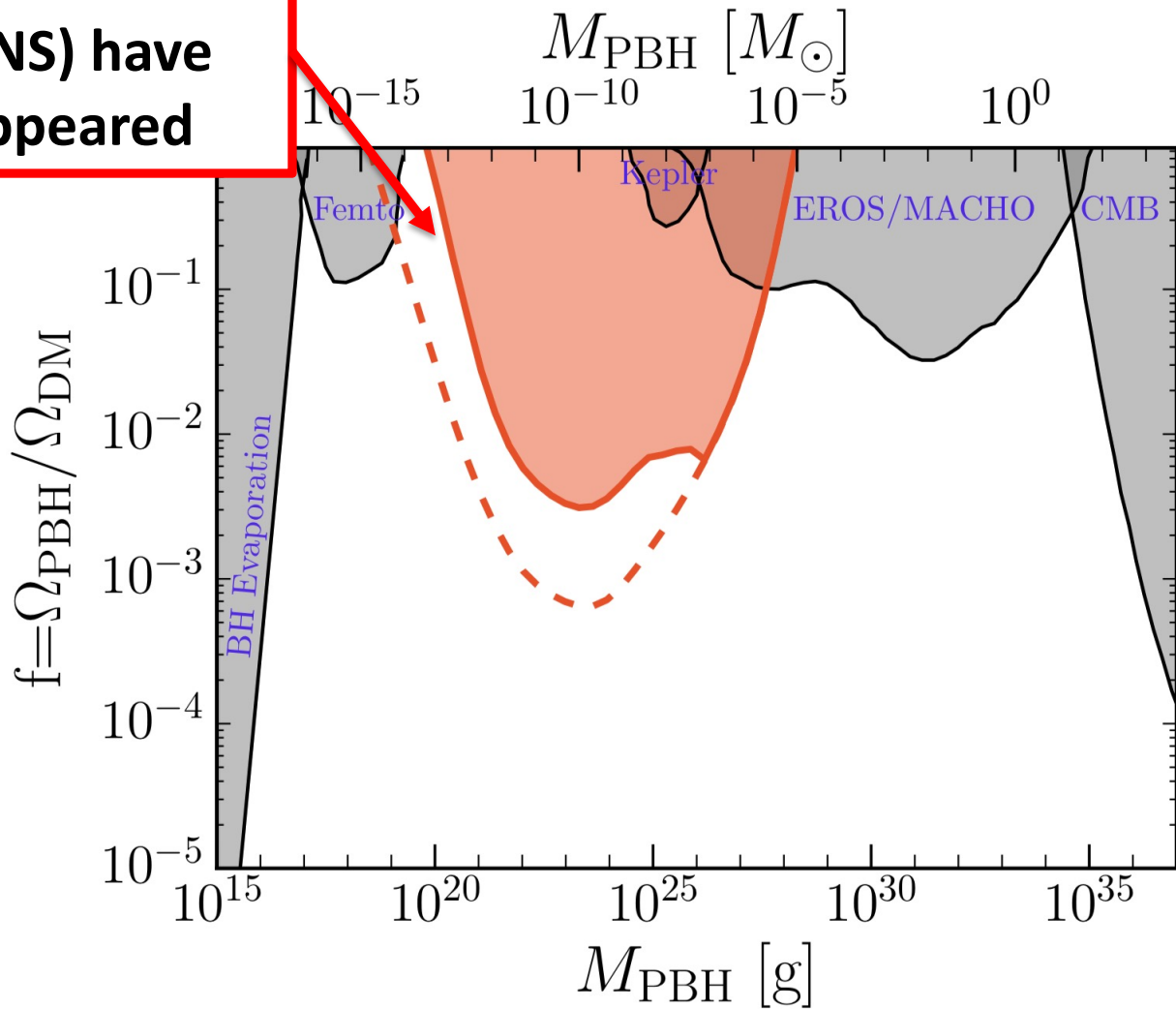
# What about mixed models?

# What about mixed models?





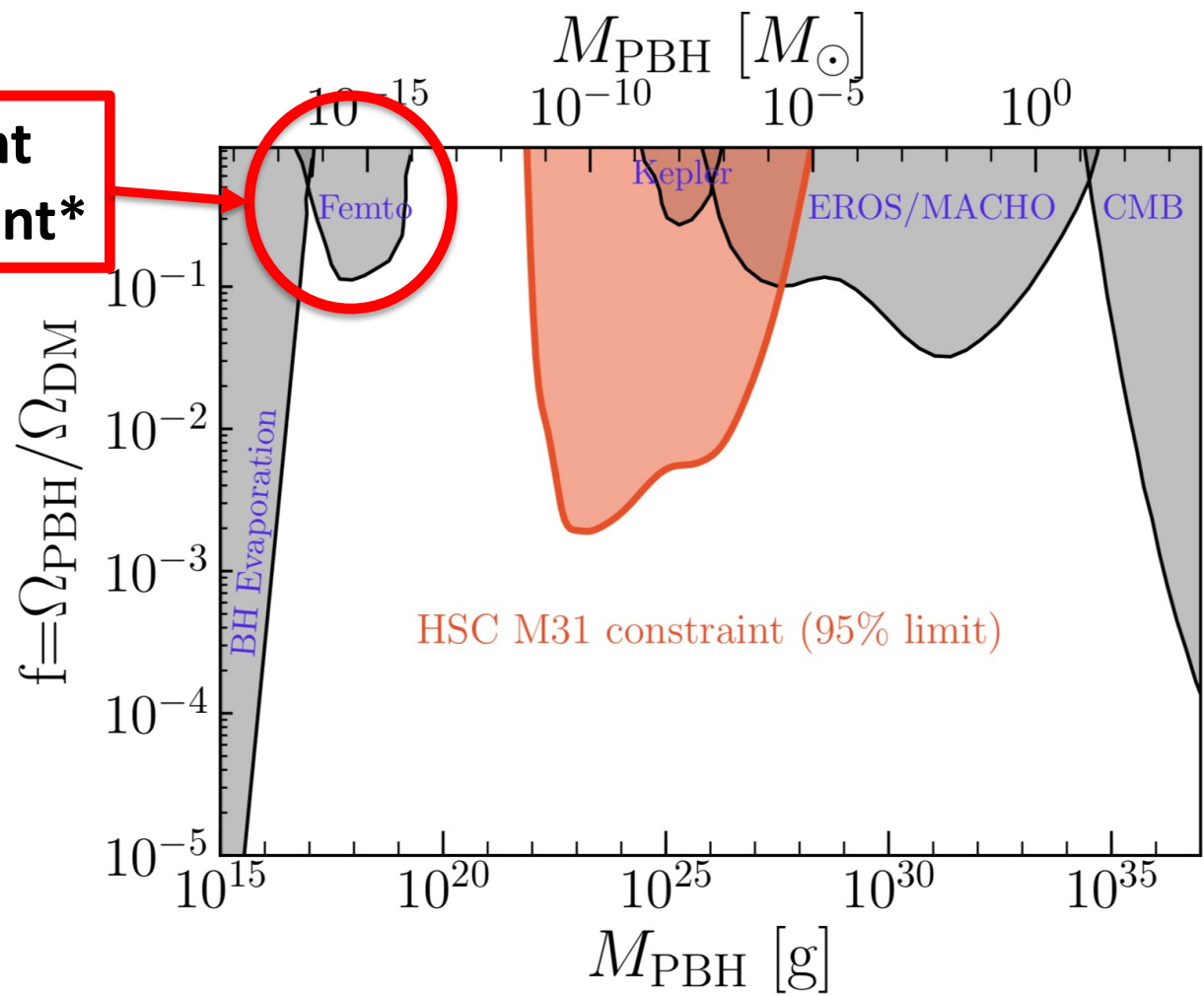
**wacky constraints  
(WD, NS) have  
disappeared**





\* Katz et al, 1807.11495

This constraint also non-existent\*



SUBARU HSC microlensing, **VERSION 3: finite source AND wave effects**

...but assuming all stars have  **$R = R_{sun}$**  !