



# Radio-frequency Dark Photon Dark Matter across the Sun

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2010.15836 [PRL 126 (2021) 181102]  
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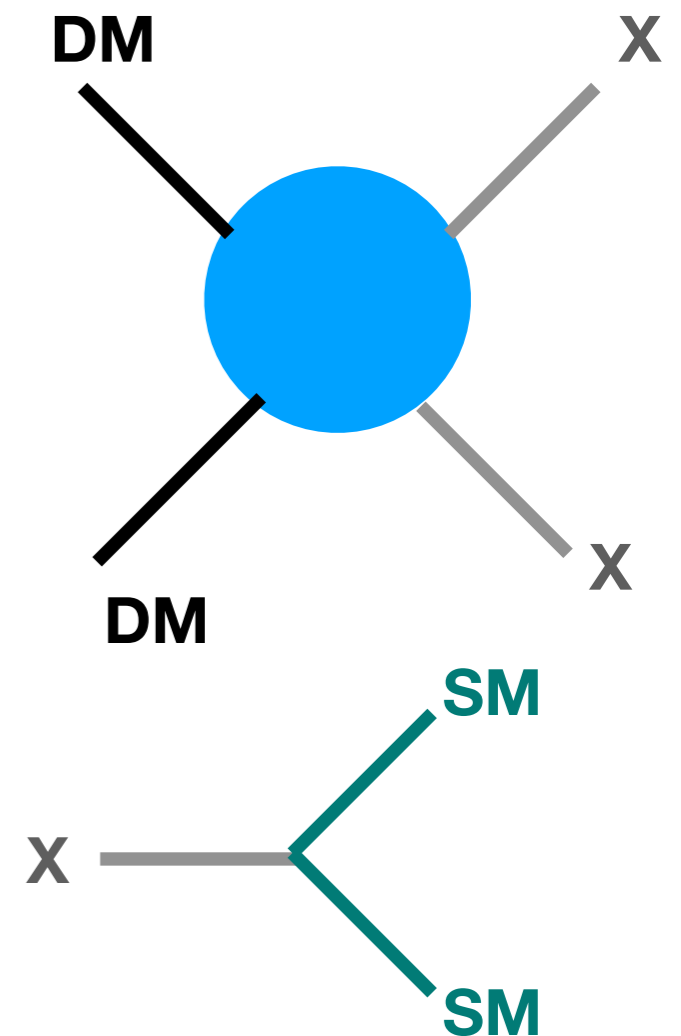
Phenomenology 2021 Symposium @ University of Pittsburgh  
05/27/2021

# The outline

- From dark matter to the dark sector
- The dark photon DM and the Sun
- Summary

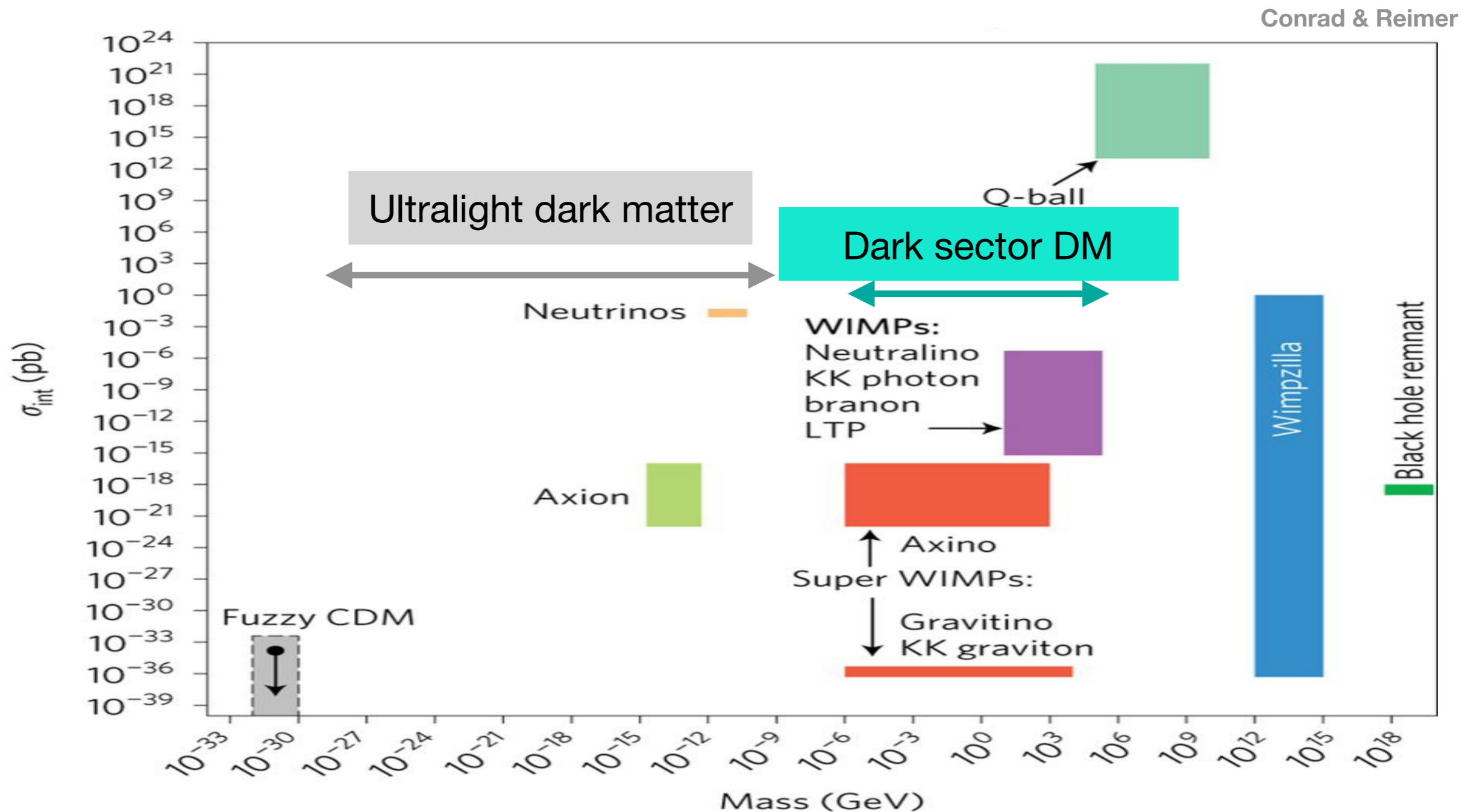
# The physics motivation of Dark Sector (X)

- 1. Existence of dark matter
  - do not interact with strong, weak, or electromagnetic forces
  - A zoo of similar particles in the dark sector as in the visible sector
- 2. The null detection of dark matter
  - Secluded annihilation:  $DM + DM \rightarrow X + X$
  - X is light and weakly coupled to visible sector



# The physics motivation of Dark Sector (X)

- 3. The experiment status
  - Technically difficult to increase E
  - Easier to accumulate higher luminosity



# The examples of dark sector models

- Coupling through gauge singlet operators of SM

- Kinetic mixing portal- Dark Photon  $B_{\mu\nu}F'^{\mu\nu}$

- Higher dimensional operators- Axion  $\frac{a}{\Lambda}\tilde{F}F, \frac{a}{\Lambda}\tilde{G}G$

- Neutrino portal  $LH$

- Higgs portal  $H^\dagger H$

# The outline

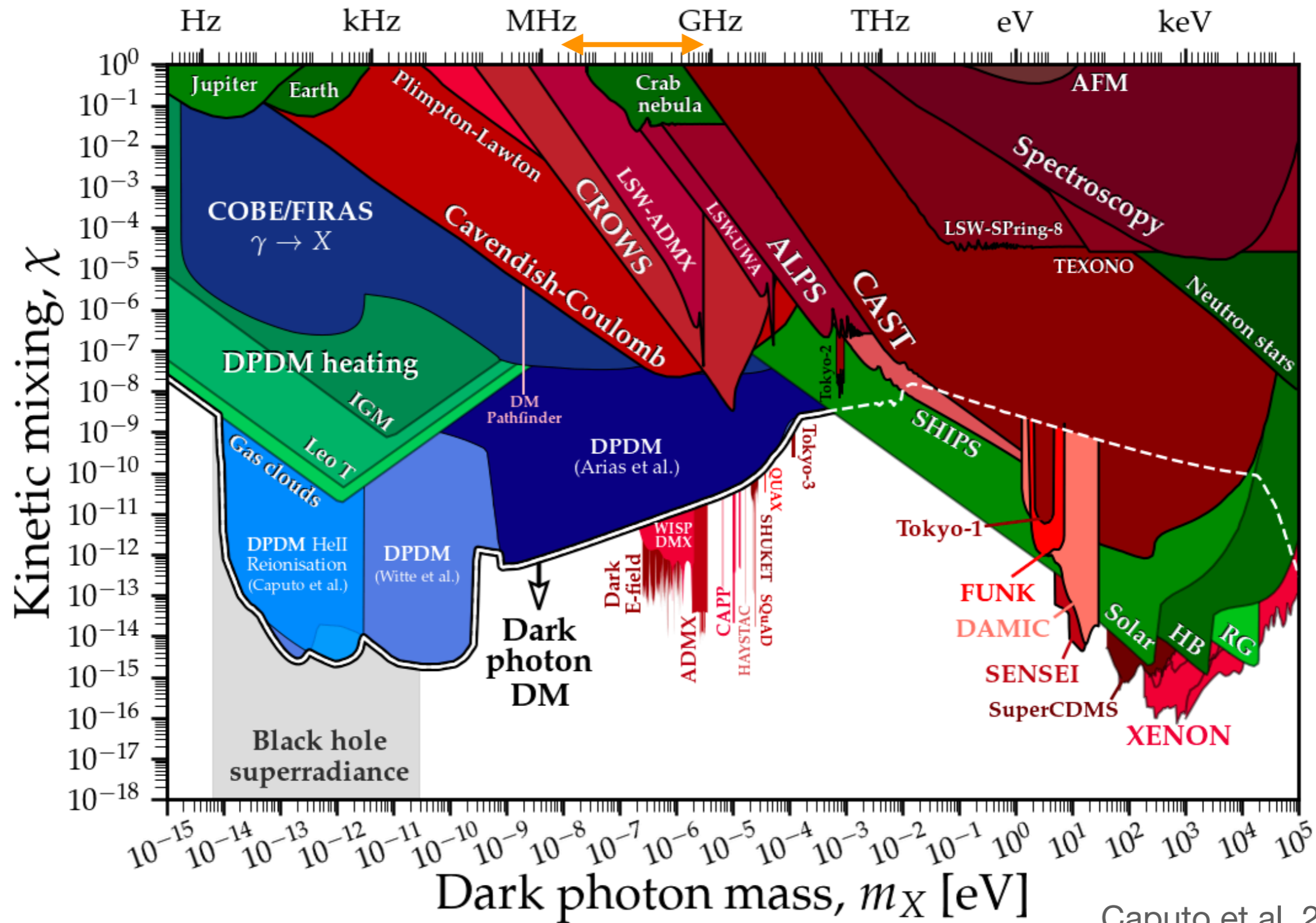
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# Motivation for the dark photon

- A simple extension for NP from marginal operator portal
- An ultralight Dark Matter candidate
- A unique dark force carrier example, similar interaction as photon

$$\mathcal{L} = -\frac{1}{4}F'_{\mu\nu}F'^{\mu\nu} - \frac{1}{2}m_{A'}^2 A'_\mu A'^\mu - \frac{1}{2}\epsilon F_{\mu\nu}F'^{\mu\nu}$$
$$\supset e\epsilon A'_\mu J_{\text{EM}}^\mu - \frac{1}{2}m_{A'}^2 A'_\mu A'^\mu$$

# The searches for dark photon



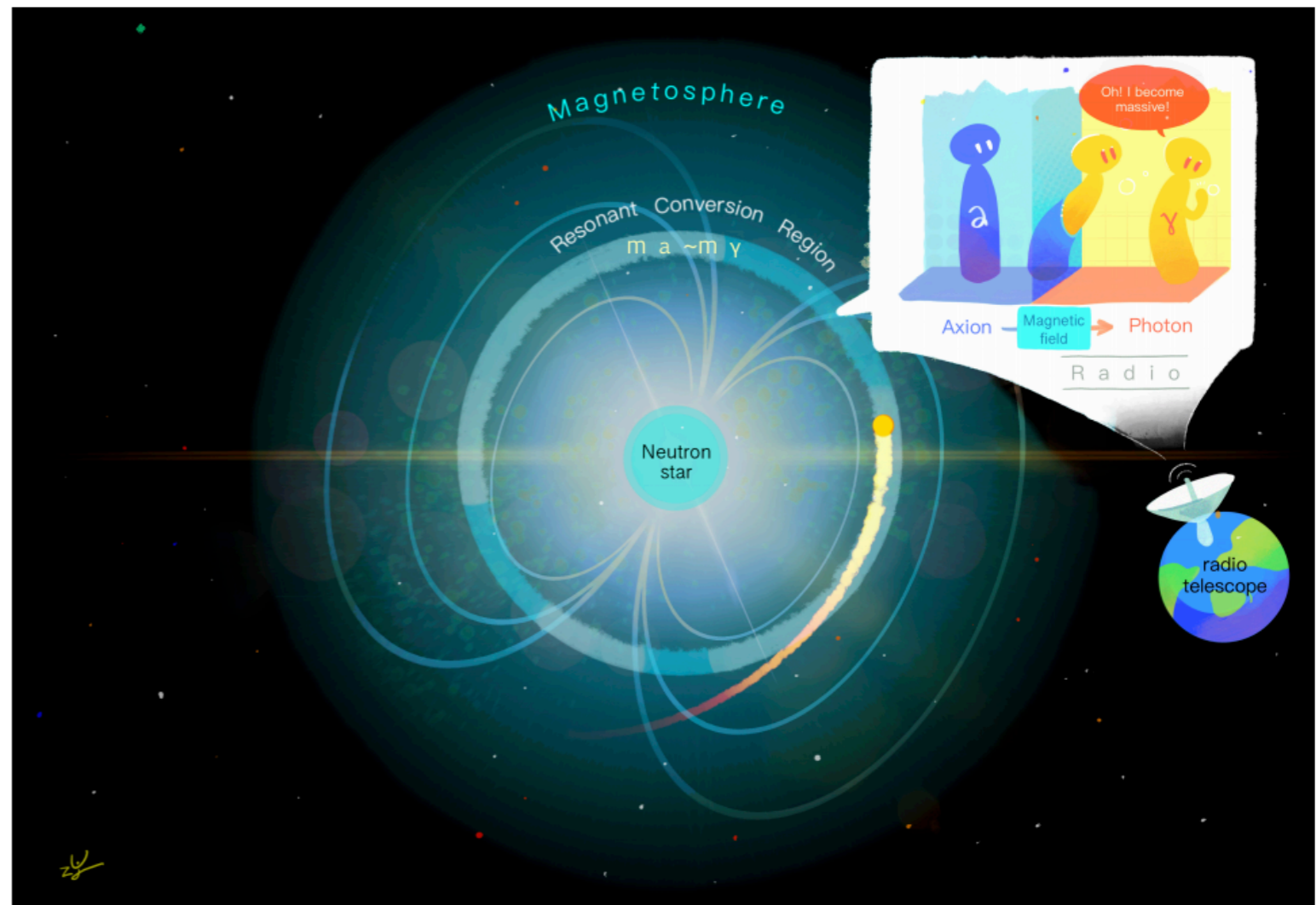
Caputo et al. 2105.04565

- Similar story to axion, but no B field needed.



# Particle physics meets astrophysics

- Recent progress in axion searches with astrophysics telescopes
  - Axion DM conversion > radio telescope (magnetized astrophysical objects, e.g. neutron star, white dwarf) (see [1803.08230](#), [1804.03145](#), [1811.01020](#), [2004.00011](#))
  - Axion DM stimulated decay > radio telescope (photon rich environment) (see [1811.08436](#))
  - Axion conversion in magnetic WD > X-ray telescope (B-field) (see [1903.05088](#))



Courtesy of Fapeng Huang

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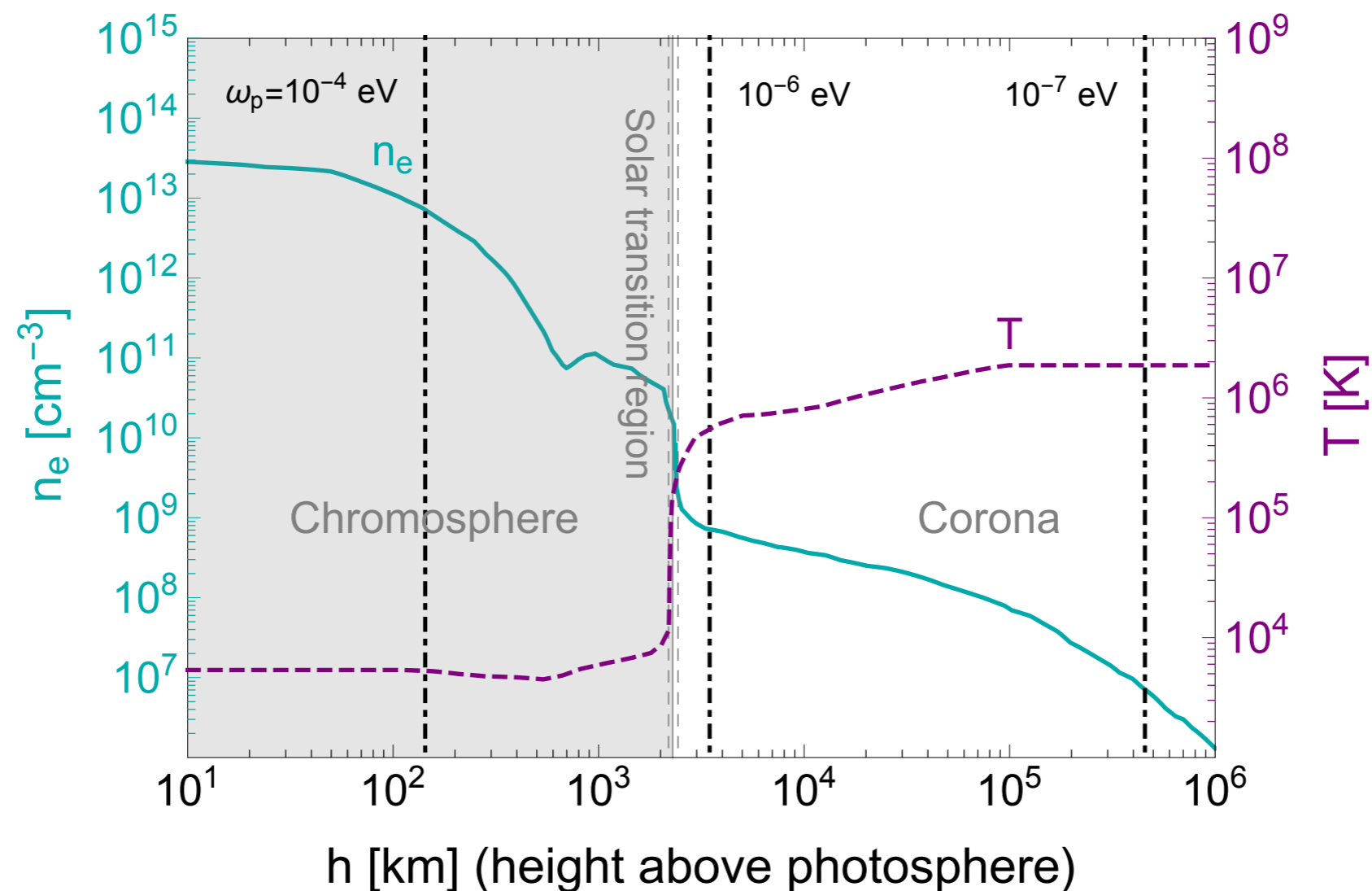
Courtesy of Haipeng An

**Our difference: probing A' dark matter**  
**No need of extreme B field**  
**What about the closest star, the Sun?**

# The dark photon dark matter conversion at Sun

- The plasma frequency

$$\omega_p = \left( \frac{4\pi\alpha n_e}{m_e} \right)^{1/2} = \left( \frac{n_e}{7.3 \times 10^8 \text{ cm}^{-3}} \right)^{1/2} \mu\text{eV}$$



# The conversion calculation using QFT

- Resonant conversion probability  $A' \rightarrow \gamma$  ( $1 \rightarrow 1$ )



$$P_{A' \rightarrow \gamma}(v_r) = \frac{1}{3} \int \frac{dt}{2\omega} \frac{d^3p}{(2\pi)^3 2\omega} (2\pi)^4 \delta^4(p_{A'}^\mu - p_\gamma^\mu) \sum_{\text{pol}} |\mathcal{M}|^2$$

$$= \frac{2}{3} \times \pi \epsilon^2 m_{A'} v_r^{-1} \left| \frac{\partial \ln \omega_p^2(r)}{\partial r} \right|_{\omega_p(r)=m_{A'}}^{-1}$$

$$\mathcal{M} = -\epsilon m_{A'}^2 \left( \xi_\gamma^*(p) \cdot \xi_{A'}(p) \right)$$

$$\frac{1}{3} \sum_{\text{pol}} |\mathcal{M}|^2 = \frac{2}{3} \epsilon^2 m_{A'}^4$$

$$\int dt \delta(E_{A'} - E_\gamma) = 2\omega^{-1} \left( \frac{\partial \ln \omega_p^2}{\partial t} \right)^{-1}$$

- Due to the forced 4-momentum conservation, it applies to resonant conversion only.

# The conversion calculation using wave method

- Eliminating kinetic mixing term by redefinition

$$\left[ \omega^2 - k^2 - \begin{pmatrix} \omega_p^2 & -\epsilon m_{A'}^2 \\ -\epsilon m_{A'}^2 & m_{A'}^2 \end{pmatrix} \right] \begin{pmatrix} A(r, t) \\ A'(r, t) \end{pmatrix} = 0$$

- Deplete the time dependence

$$\omega^2 - k^2 = m_A^2, \quad A(r, t) = e^{i(\omega t - rk)} \tilde{A}(r) \quad A'(r, t) = e^{i(\omega t - rk)} \tilde{A}'(r)$$

- Substitute  $\omega \rightarrow -i \frac{\partial}{\partial t}$  and  $k \rightarrow i \frac{\partial}{\partial r}$

- Use **WKB approximation**  $|\partial_r^2 \tilde{A}(r)| \ll |k \partial_r \tilde{A}(r)|$

- Obtain **linearized wave equation**

$$[-i \partial_r + H_0 + H_I] \begin{pmatrix} \tilde{A}(r) \\ \tilde{A}'(r) \end{pmatrix} = 0,$$

$$H_0 = \begin{pmatrix} \frac{m_{A'}^2 - \omega_p^2}{2k} & 0 \\ 0 & 0 \end{pmatrix}, \quad H_I = \begin{pmatrix} 0 & -\frac{\epsilon m_{A'}^2}{2k} \\ -\frac{\epsilon m_{A'}^2}{2k} & 0 \end{pmatrix}$$

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- 1st order solution for conversion probability

$$P_{A' \rightarrow \gamma} = \left| \int_0^\infty dr \frac{-\epsilon m_{A'}^2}{2k} e^{-i \int_0^r d\tilde{r} \frac{m_{A'}^2 - \omega_p^2(\tilde{r})}{2k}} \right|^2$$

**Worked for both resonant and Non-resonant conversion**

- Further simplification using Saddle point approximation

$$\int_{-\infty}^{\infty} dr e^{-f(r)} \approx e^{-f(r_0)} \sqrt{\frac{2\pi}{f''(r_0)}}$$



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- Further simplification using Saddle point approximation

$$\int_{-\infty}^{\infty} dr e^{-f(r)} \approx e^{-f(r_0)} \sqrt{\frac{2\pi}{f''(r_0)}} \quad \leftarrow \begin{matrix} f'(r_0) = 0 \\ f(r) \approx f(r_0) + \frac{1}{2}(r - r_0)^2 f''(r_0) \end{matrix}$$

$$f(r) = i \int_0^r d\tilde{r} \frac{m_{A'}^2 - \omega_p^2(\tilde{r})}{2k}$$

**QFT method = Wave method !**

# The dark photon dark matter conversion at Sun

- The radiation power per solid angle at conversion radius  $r_c$

$$\begin{aligned} \frac{d\mathcal{P}}{d\Omega} &\approx 2 \times \frac{1}{4\pi} \rho_{\text{DM}} v_0 \int_0^b dz 2\pi z P_{A' \rightarrow \gamma}(v_r) \\ &= P_{A' \rightarrow \gamma}(v_0) \rho_{\text{DM}} v(r_c) r_c^2 \end{aligned}$$

- $z$  is impact parameter for incoming  $A'$
  - $b$  is the max impact parameter which can reach  $r_c$
  - $v_0 \sim 220\text{km/s}$  is the DM local velocity dispersion
- The spectral power flux density per solid angle

$$S_{\text{sig}} = \frac{1}{1\text{AU}^2} \frac{1}{\mathcal{B}} \frac{d\mathcal{P}}{d\Omega}$$

$$\begin{aligned} \mathcal{B} &= \max(B_{\text{sig}}, B_{\text{res}}) \\ B_{\text{sig}} &\approx \frac{m_{A'} v_0^2}{2\pi} \sim 130 \text{ Hz} \times \frac{m_{A'}}{\mu\text{eV}} \end{aligned}$$



# The photon propagation

- Photon out-going direction

$$n(\omega) = (1 - \omega_p^2/\omega^2)^{1/2}$$

$$n_{\text{res}} \sim 10^{-3} - 10^{-2}$$

$$\sin \theta_{\text{out}} = \frac{n_{\text{res}}}{n_{\text{out}}} \times \sin \theta_{\text{res}} \lesssim 10^{-3} - 10^{-2} .$$

- Absorption from inverse bremsstrahlung process

$$\Gamma_{\text{inv}} \approx \frac{8\pi n_e n_N \alpha^3}{3\omega^3 m_e^2} \left( \frac{2\pi m_e}{T} \right)^{1/2} \log \left( \frac{2T^2}{\omega_p^2} \right) (1 - e^{-\omega/T})$$

- The Compton scattering

$$\Gamma_{\text{Com}} = \frac{8\pi\alpha^2}{3m_e^2} n_e$$

- Survival probability

$$P_s \equiv e^{-\int \Gamma_{\text{att}} dt} \simeq \exp \left( - \int_{r_c}^{r_{\text{max}}} \Gamma_{\text{att}} dr / v_r \right)$$

# Sensitivity of Radio Telescopes

- The system equivalent flux density

$$\text{SEFD} = 2k_B \frac{T_{\text{sys}} + T_{\odot}^{\text{nos}}}{A_{\text{eff}}}$$

- The minimum detectable flux density

$$S_{\text{min}} = \frac{\text{SEFD}}{\eta_s \sqrt{n_{\text{pol}} \mathcal{B} t_{\text{obs}}}}$$



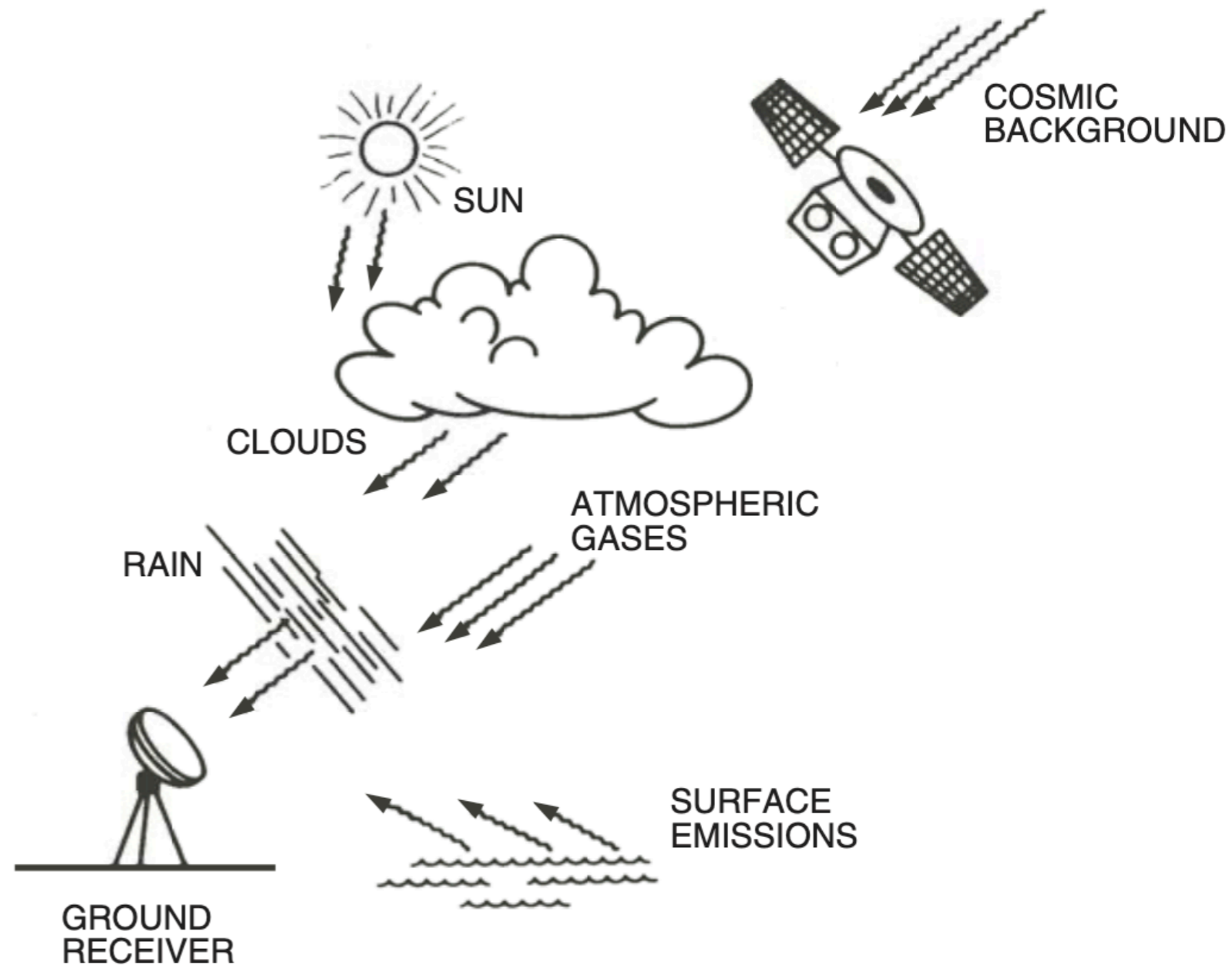
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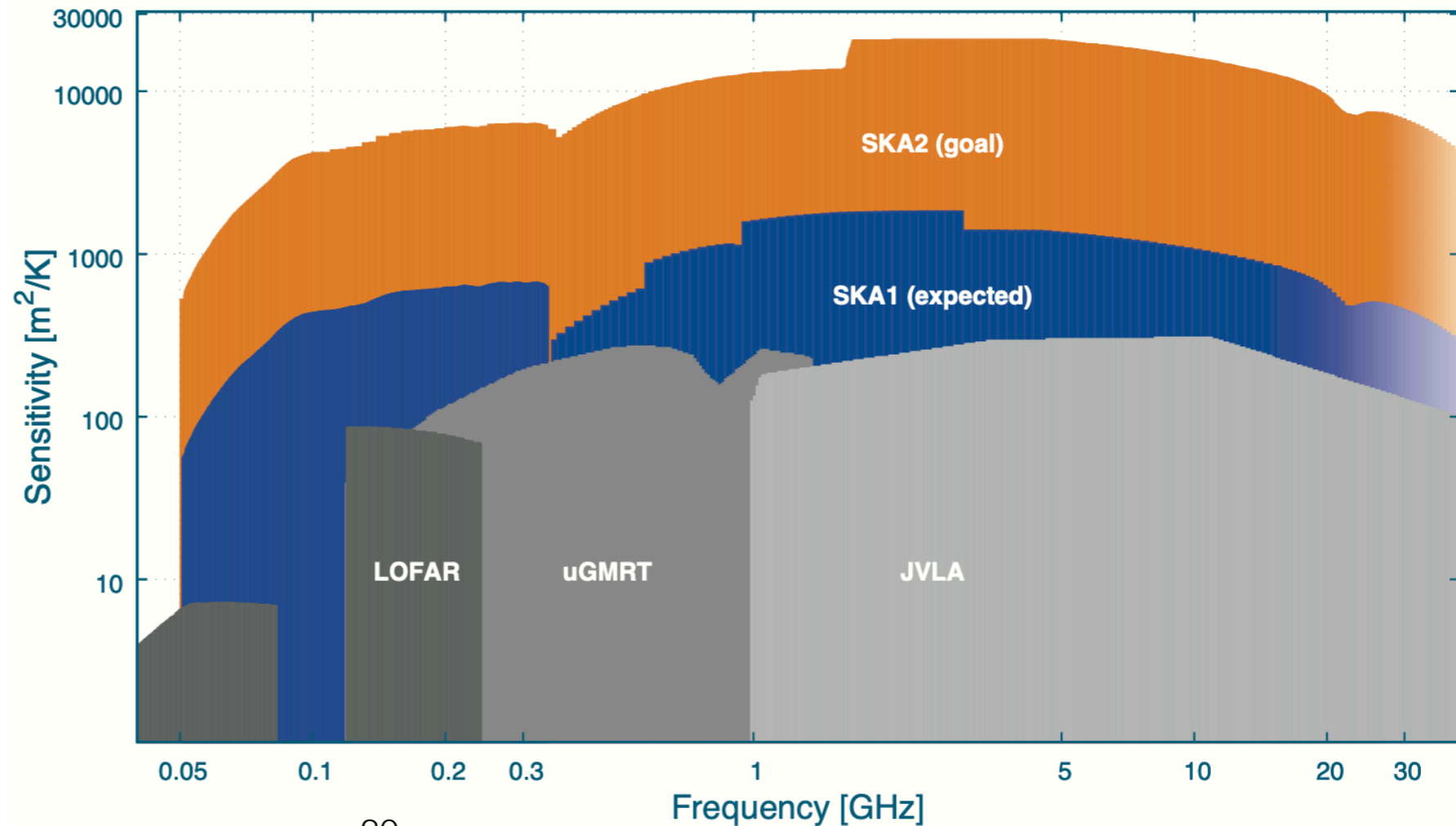
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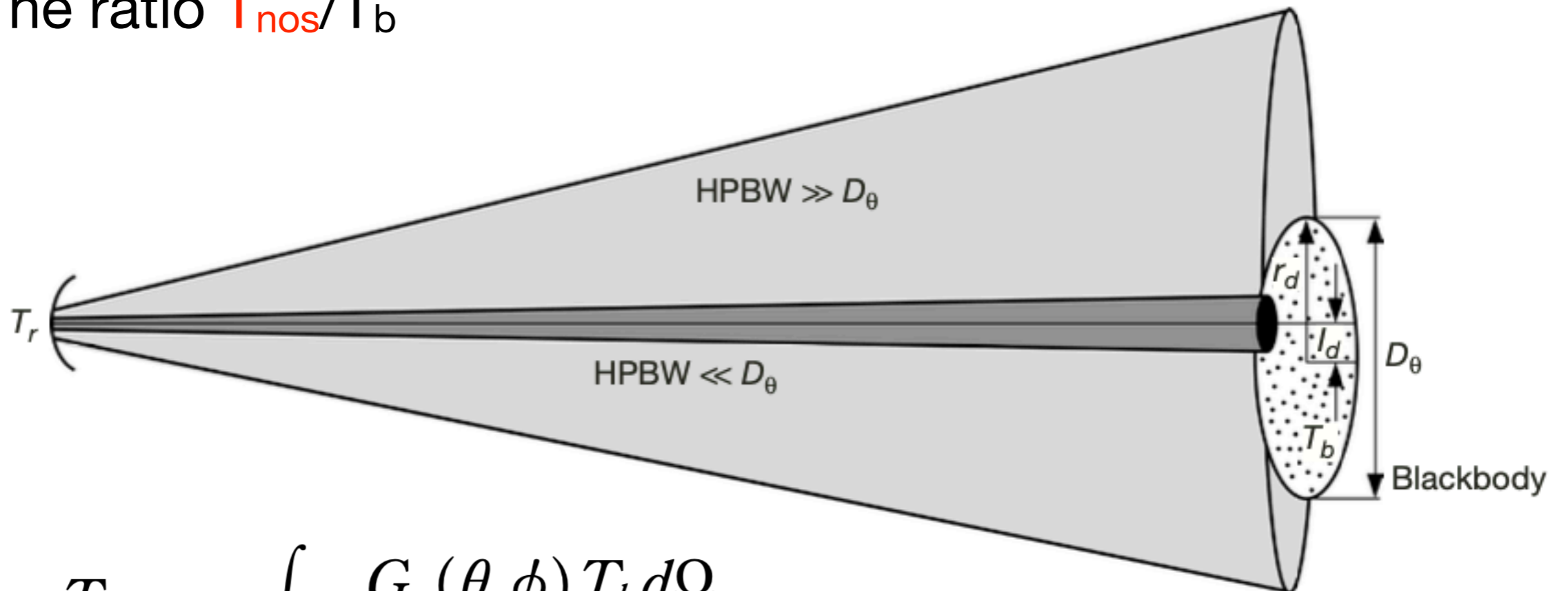
Name	$f$ [MHz]	$B_{\text{res}}$ [kHz]	$\langle T_{\text{sys}} \rangle$ [K]	$\langle A_{\text{eff}} \rangle$ [m <sup>2</sup> ]
SKA1-Low	(50, 350)	1	680	$2.2 \times 10^5$
SKA1-Mid B1	(350, 1050)	3.9	28	$2.7 \times 10^4$
SKA1-Mid B2	(950, 1760)	3.9	20	$3.5 \times 10^4$
LOFAR	(10, 80)	195	28,110	1,830
LOFAR	(120, 240)	195	1,770	1,530

Radio Interferometer Sensitivity Comparison



# The noise from the Sun

- The ratio  $T_{\text{nos}}/T_b$

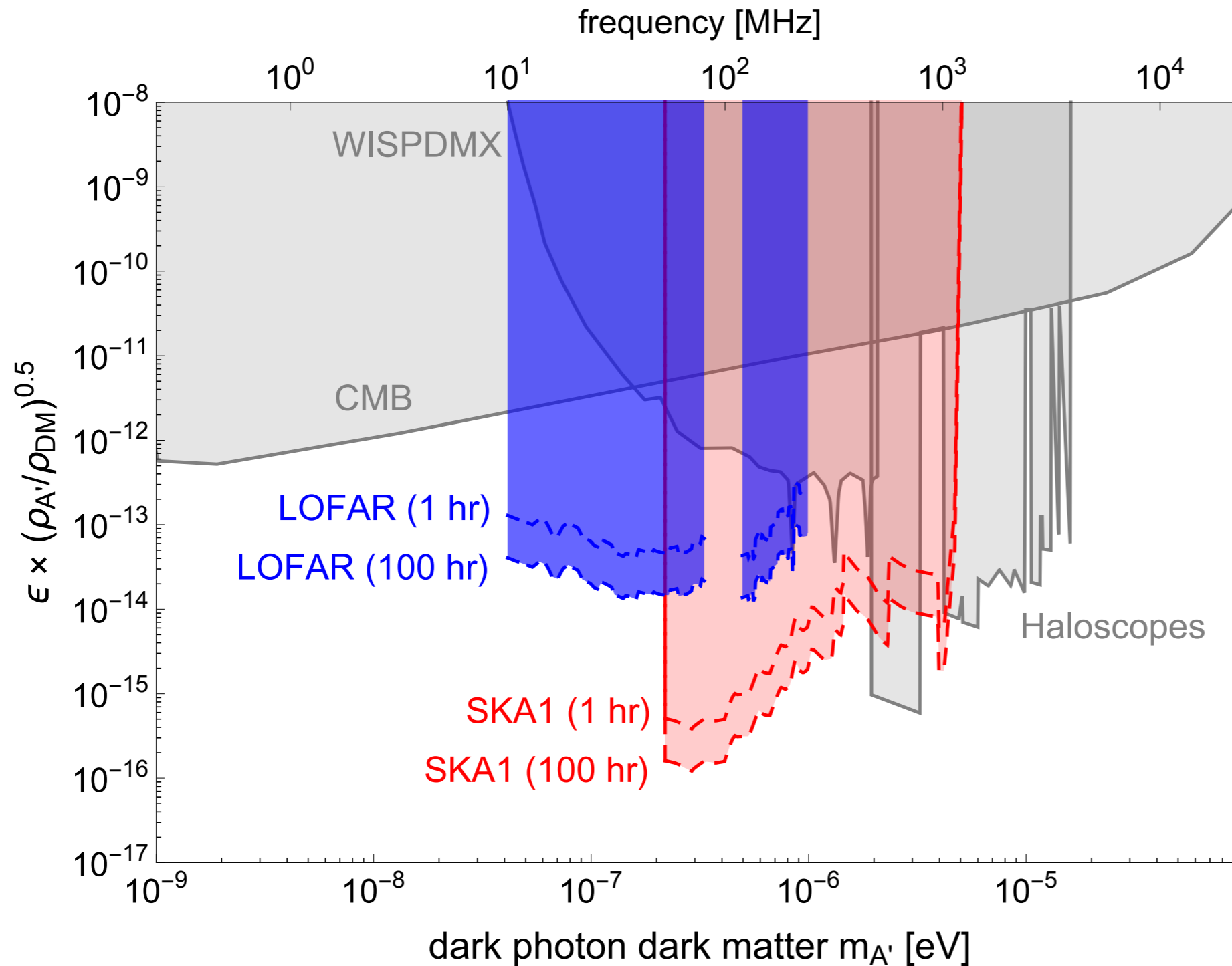


$$\frac{T_{\text{incr}}}{T_b} = \frac{\int_{\text{disk}} G_r(\theta, \phi) T_b d\Omega}{\int_{4\pi} G_r(\theta, \phi) T_b d\Omega}$$

**HPBW = Half Power Beam Width**

- For narrow HPBW (SKA1),  $T_{\text{nos}}/T_b = 1$
- For wide HPBW (LOFAR),  $T_{\text{nos}}/T_b \ll 1$

# The physics reach for dark photon dark matter



$$S_{\text{sig}} \times P_s = S_{\text{min}}$$

- 10 MHz lower end from LOFAR, 1 GHz higher end due to opacity

# Summary

- The inverse Primakoff effect should be **included** for solar axion detection
- Stellar cooling tension is mildly alleviated but still a problem, and the **uncertainties** in astrophysics calculation should be included
- We proposed search for radio frequency  $A'$  dark matter from 10 — 1000 MHz
- $A'$  DM **resonantly converts** into radio photon
- Only conversion from **solar corona** can propagate out of the Sun
- **1 hour** solar observation from LOFAR and SKA1 can provide strong sensitivity
- **Future** experiments like Arecibo, JVLA, FAST, TianLai etc can further explore the scenario if having solar program

**Thank you!**