

Neutrino masses from a pseudo-Dirac bino and its detection prospects

Julia Gehrlein

Brookhaven National Laboratory



Phenomenology symposium 2021

25. May 2021

Neutrino masses

- ▶ observation of neutrino oscillations \Rightarrow non-vanishing neutrino masses
- ▶ explain smallness of neutrino mass & testable model
 \rightarrow variants of type I seesaw like **inverse seesaw**
- ▶ ISS: right-handed neutrinos form pseudo-Dirac pairs, lepton number is (approximately) conserved
 \rightarrow lepton number violating effects (light neutrino masses) are naturally suppressed



Neutrino masses

- ▶ explain smallness of neutrino mass & testable model
→ variants of type I seesaw like **inverse seesaw**
- ▶ ISS: right-handed neutrinos form pseudo-Dirac pairs, lepton number is (approximately) conserved
→ lepton number violating effects (light neutrino masses) are naturally suppressed
- ▶ **challenge**: explanation of the texture needed in the neutrino mass matrix, suppression of lepton number violating effects
possible solutions: identify lepton number breaking parameter with scalar vev in gauged $B - L$ or $U(1)_D$ symmetry [De Romeri, Fernandez-Martinez, JG, Machado, Niro '17, Ballett, Hostert, Pascoli '19, JG, Pierre '19]

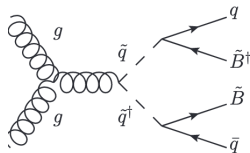


- ▶ augment $U(1)_{R-L}$ -symmetric SUSY model with Dirac gauginos to $U(1)_{R-L}$ symmetry [Coloma, Ipek '16]
- ▶ small gaugino Majorana masses are generated by $U(1)_{R-L}$ -violating effects (gravitino mass) \rightarrow suppression of light neutrino masses
- ▶ bino and its Dirac partner \rightarrow automatically form a pseudo-Dirac pair, play the role of RH neutrinos in ISS, the **bi ν o**
- ▶ **differences** to traditional ISS: no right handed neutrinos are needed, particles needed for neutrino mass generation can be produced in decays of colored particles

Based on " *$B\nu o$ phenomenology at the LHC*"
in collaboration with Seyda Ipek and Patrick J. Fox
JHEP **03** (2019), 073, arXiv:1901.09284
and
"*Long-lived $b\nu o$ at the LHC,*"
in collaboration with Seyda Ipek
JHEP **05** (2021), 020, arXiv:2103.0125.



- ▶ leading $B\nu$ production channel



- ▶ $B\nu$ decay channels: $BR(\tilde{B} \rightarrow Wl/Z\nu/h\nu) \approx 1/3$

- ▶ decay rates:

- ▶ light $B\nu$

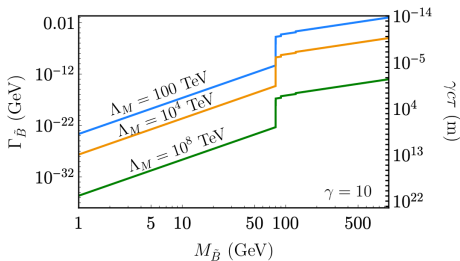
$$(m_{\tilde{B}} < m_Z, m_W, m_h):$$

$$\Gamma^{light} \simeq \kappa \frac{G_F^2 M_{\tilde{B}}^7}{192\pi^3 \Lambda_M^2}$$

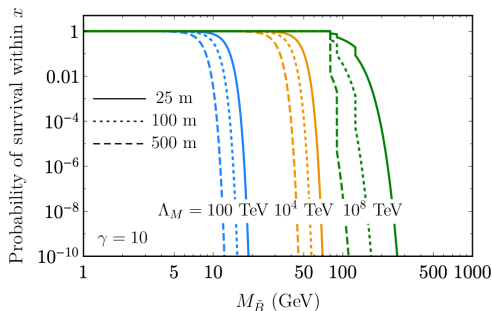
- ▶ heavy $B\nu$

$$(m_{\tilde{B}} > m_Z, m_W, m_h):$$

$$\Gamma^{light} \simeq \frac{M_{\tilde{B}}^3}{\Lambda_M^2}$$

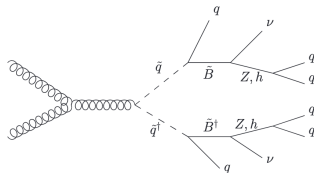


► probability of bi ν survival



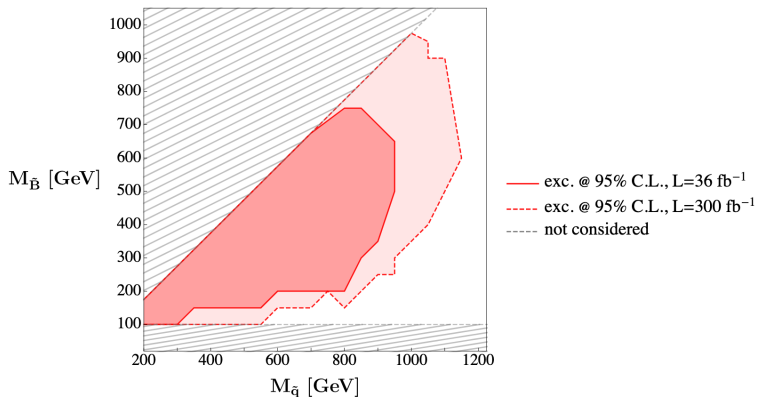
► lengths corresponding to ATLAS, MATHUSLA, FASER

- ▶ constraints on low messenger scale $\Lambda_M \sim 100$ TeV \Rightarrow bi ν o decays inside ATLAS
- ▶ largest BR $\sim 20\%$ to 6j+MET \neq standard search for particles involved in neutrino mass mechanism (e.g. heavy neutrinos)

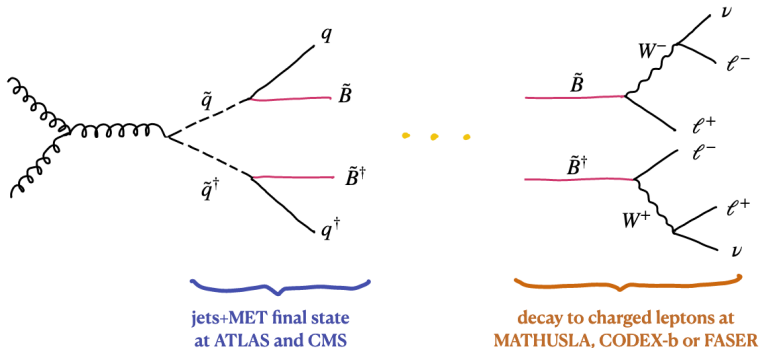


- ▶ recast 6jet+MET search from ATLAS [ATLAS-CONF-2017-022]

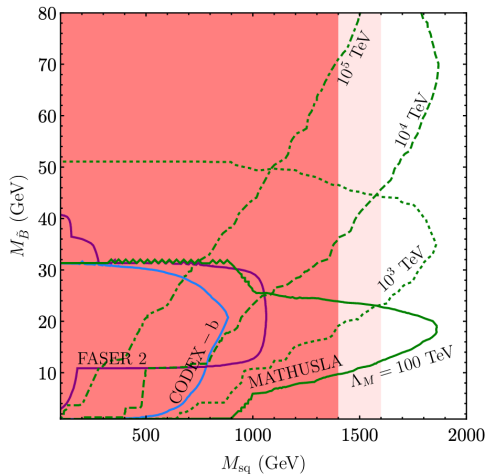
- ▶ recast 6jet+MET search from ATLAS [ATLAS-CONF-2017-022] using Madanalysis code from [Asadi, Buckley, DiFranzo, Monteux, Shih '17]



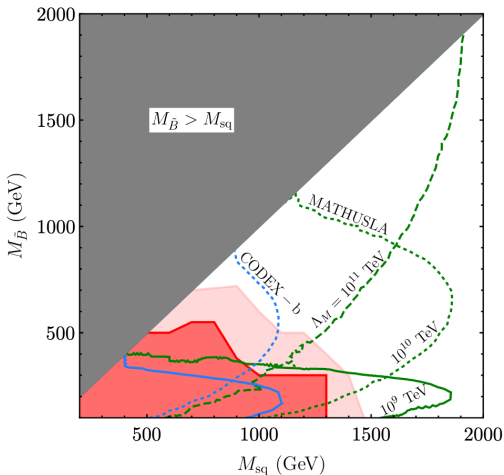
higher messenger scale \Rightarrow bi ν decays inside LLP experiments



► light bi ν



► heavy bi ν o



- ▶ constraints from lepton flavour violation
- ▶ $\text{BR}(\mu \rightarrow e\gamma)$ constrains $v^2/(2\Lambda_M^2)f_e f_\mu^* \rightarrow \Lambda_M > 30 \text{ TeV}$

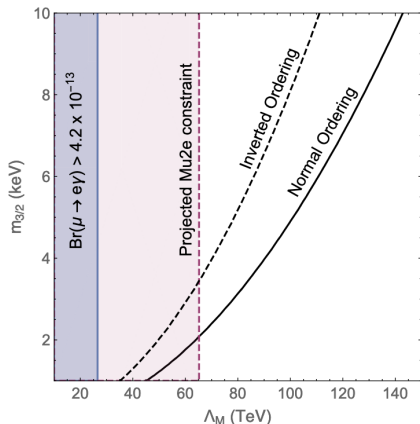
- ▶ light neutrino masses

$$m_1 = 0$$

$$m_2 = \frac{m_{3/2} v^2}{\Lambda_M^2} (1 - \rho)$$

$$m_3 = \frac{m_{3/2} v^2}{\Lambda_M^2} (1 + \rho),$$

$$\rho \simeq 0.7$$



- ▶ presented SUSY model which naturally gives rise to ISS neutrino mass matrix pattern
- ▶ particles related to neutrino mass mechanism testable at LHC and dedicated LLP experiments
- ▶ complementarity between collider searches and LFV searches to probe parameter space

Thank you for your attention!



[All artwork by Sandbox Studio, Chicago]

- ▶ in $U(1)_R$ -symmetric supersymmetric models: superpartners have +1 R -charges, SM fields are not charged under $U(1)_R \rightarrow$ Gaugino Majorana masses forbidden
- ▶ Dirac masses for gauginos: adjoints with opposite R -charges are introduced
- ▶ extend $U(1)_R$ to $U(1)_{R-L}$ to allow for neutrino mass terms

- ▶ SUSY broken in hidden sector that communicates with the visible sector at messenger scale $\Lambda_M \rightarrow$ **Dirac mass** for the bino \tilde{B} , Dirac partner singlino S

$$\int d^2\theta c \frac{W'_\alpha}{\Lambda_M} W_{\tilde{B}}^\alpha \Phi_S \rightarrow \frac{cD}{\Lambda_M} \tilde{B}S = M_{\tilde{B}} \tilde{B}S$$

- ▶ $U(1)_{R-L}$ broken by supergravity \rightarrow Majorana masses for bino and singlino via anomaly mediation generated (proportional to gravitino mass $m_{3/2}$)

$$m_{\tilde{B}} \sim m_S \sim \frac{1}{16\pi^2} m_{3/2} \ll M_{\tilde{B}}$$

- ▶ singlino and bino form pseudo Dirac fermion $\Psi^T = (\tilde{B}, S^\dagger) \rightarrow$ **bi ν**
- ▶ interactions of bino with lepton doublet

$$\frac{f_i}{\Lambda_M^2} \int d^2\theta W'_\alpha W_{\tilde{B}}^\alpha H_u L_i \rightarrow \frac{f'_i M_{\tilde{B}}}{\Lambda_M} h_u \tilde{B} l_i$$

- ▶ not enough freedom to give mass to at least two SM neutrinos \rightarrow singlino interaction like $\Phi_S H_u L_i$
- ▶ not charged under $U(1)_{R-L}$, needs to be induced a $d = 5$ Kähler potential term with $\phi = 1 + \theta^2 m_{3/2}$

$$\int d^2\theta d^2\bar{\theta} \phi^\dagger \frac{d_i \Phi_S H_u L_i}{\Lambda_M} \rightarrow \frac{d_i m_{3/2}}{\Lambda_M} h_u S l_i$$

- ▶ neutrino mass matrix

$$M_\nu = \begin{pmatrix} 0_{3 \times 3} & Y\nu & G\nu \\ Y^T\nu & m_{\tilde{B}} & M_{\tilde{B}} \\ G^T\nu & M_{\tilde{B}} & m_S \end{pmatrix}$$

$Y_i = f_i M_{\tilde{B}} / \Lambda_M$, $G_i = d_i m_{3/2} / \Lambda_M$, for ISS: $G \ll Y$

- ▶ $m_1 = 0$, $m_2 = \frac{m_{3/2} v^2}{\Lambda_M^2} (1 - \rho)$, $m_3 = \frac{m_{3/2} v^2}{\Lambda_M^2} (1 + \rho)$
- ▶ parametrically

$$m_\nu \simeq (2 - 20) \times 10^{-2} \text{ eV} \left(\frac{m_{3/2}}{10 \text{ keV}} \right) \left(\frac{100 \text{ TeV}}{\Lambda_M} \right)^2$$