



# Form Factor Effects in Higgs Couplings

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### **BSM Momentum Dependence**

Heavy BSM physics alters the SM Higgs couplings.

**i.e.** CHMs with 
$$\xi = \frac{v^2}{f^2}$$
, 2HDM with  $\tan \tilde{\alpha} \simeq \mathcal{O}\left(\frac{v^2}{M_{\pm}^2}\right)$ ,...  
In general an  $\mathcal{O}\left(\frac{v^2}{\Lambda^2}\right)$  effect.

Momentum effects assumed to decouple.

Might not happen for high off-shell momenta.

How to explore momentum dependence in a general manner?

What are the consequences of momentum dependent Higgs couplings?

We expect an enhancement of order  $p^2/\Lambda^2$  over the  $v^2/\Lambda^2$  coupling modification in <u>off-shell channels</u>.

## **Higgs Form Factors**



• 2-point functions: Spectral decomposition.

$$\Pi(p^2) = \int_0^\infty dm^2 \frac{\rho(m^2)}{p^2 - m^2}$$

• 3-point functions: ?

$$\Gamma(p_1^{\mu}, p_2^{\nu}) = \Gamma(p_1^2, p_2^2, p_1 \cdot p_2)$$

## Need to assume Higgs interactions with BSM states.

Dynamical Higgs couplings

### Higgs Form Factors: Examples





### **Higgs Form Factors: <u>Examples</u>**

$$\mathcal{L}_{\rm FF} = \sum_{i} \Gamma_i(p_1, p_2, p_h) \mathcal{L}_i^{\rm higgs}$$

i.e. Top-Yukawa form factor:

 $\Gamma_Y(p_1, p_2, p_h)\overline{q}_L \tilde{H} t_R$ 

**Expect modifications to SM mass-Yukawa relation** 

$$\underbrace{M(p_t^2)\overline{q}_L t_R}_{\substack{\swarrow \\ (\mathsf{Pole} \; \mathsf{Man})}} \underbrace{M(p_t^2)\overline{q}_L t_R}_{\substack{\swarrow \\ (\mathsf{Pole} \; \mathsf{Man})}} \underbrace{\kappa_\xi \frac{y_t^{SM}}{\sqrt{2}}}_{\substack{\swarrow \\ (\mathsf{Pole} \; \mathsf{Man})}} \underbrace{f(p_1^2, p_2^2, p_1 \cdot p_2)\overline{q}_L h t_R}_{\substack{\swarrow \\ (\mathsf{Pole} \; \mathsf{Man})}}$$

$$\underbrace{\mathsf{MCHM}_{\mathbf{5}}}_{\mathbf{f}_{t\bar{t}h}}(p_{1},p_{2}) = \frac{M(p_{1},p_{2})\left(1-2\xi\right)/\sqrt{2}}{\left(\Pi_{0}^{L}(p_{1})+\Pi_{1}^{L}(p_{1})\frac{1}{2}\left\langle S_{h}^{2}\right\rangle\right)\left(\Pi_{0}^{R}(p_{2})+\Pi_{1}^{R}(p_{2})\left\langle C_{h}^{2}\right\rangle\right)} \\
M(p_{1},p_{2}) = f^{2}y_{L}y_{R}\left(\frac{m_{4}}{p_{1}^{2}-m_{4}^{2}}-\frac{m_{1}}{p_{2}^{2}-m_{1}^{2}}\right) \qquad \kappa_{\xi}^{\mathbf{5}} = \frac{1-2\xi}{\sqrt{1-\xi}}$$

#### Absolute value of the form factors





• Implementation: Include form factors in UFO file. MADGRAPH @ parton level.

### **Fermionic resonances:** MCHM<sub>5</sub> Form factor



### **Modification of kinematic distribution shapes**



### **Vectorial resonances:** MCHM<sub>5</sub> Form factor



**Narrow resonances:** m<sub>o</sub>>2.5 TeV Eur. Phys. J. C (2014) 74:3181

Weaker bounds for broad resonances: m<sub>o</sub>>1 TeV Fully composite 3<sup>rd</sup> generation of quarks, g<sub>0</sub>>3 ArXiv:1901.01674

### **Local EFT expansion**

$$\Gamma(p_1, p_2, p_h) = \Gamma_0 + c_1 \frac{p_1^2}{\Lambda^2} + c_2 \frac{p_2^2}{\Lambda^2} + c_{12} \frac{p_1 \cdot p_2}{\Lambda^2} + \dots$$



### Conclusions

- Form factor couplings in CHMs promote a **competition between** <u>misalignment suppression</u> and <u>momentum enhancement.</u>
- Signal enhancement valid for generic BSM states.
- Modification of kinematic distribution shapes  $\rightarrow$  HL-LHC strategy
- We verify the <u>off-shell enhancement</u> of order  $p^2/\Lambda^2$  over  $v^2/\Lambda^2$  coupling modifications.
- Presence of momentum effects in local and non-local formulations.

#### Form Factor Effects in Higgs Couplings (Soon!)

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# **Thank You!**

### References

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# Backup

$$U[h] = \exp\left(\frac{i\sqrt{2}}{f}h^{\hat{a}}T^{\hat{a}}\right) \qquad q_{L}^{5} = \begin{pmatrix} -ib_{l} \\ -b_{l} \\ -it_{l} \\ t_{l} \\ 0 \end{pmatrix} \qquad t_{r}^{5} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \\ t_{r} \end{pmatrix} \qquad b_{r}^{5} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \\ b_{r} \end{pmatrix}.$$

Partial Compositeness:

$$\mathcal{L}_{int}^{F} = f \left[ y_{L1}(\overline{q_{L}^{5}}U[\pi])_{5}\psi_{1} + y_{L4}(\overline{q_{L}^{5}}U[\pi])_{j}\psi_{4,j} \right] + h.c.$$
  
+  $f \left[ y_{R1}(\overline{t_{r}^{5}}U[\pi])_{5}\psi_{1} + y_{R4}(\overline{t_{r}^{5}}U[\pi])_{j}\psi_{4,j} \right] + h.c.$ 



Vector Meson dominance/ Hidden Local Symmetry:

$$\begin{aligned} \mathcal{L}_{CS}^{V} &= -\frac{1}{4} \rho_{\mu\nu}^{a} \rho^{a,\mu\nu} + \frac{m_{\rho}^{2}}{2} \rho_{\mu}^{a} \rho^{a,\mu} + g_{\rho} \rho_{\mu}^{a} J^{a,\mu} + \frac{g_{\rho}^{2}}{2} \rho_{\mu}^{a} \rho^{a,\mu} h^{2}, \\ \mathcal{L}_{ES}^{V} &= -\frac{1}{4} W_{\mu\nu}^{a_{L}} W^{a_{L},\mu\nu} + g_{0} W_{\mu}^{a_{L}} J^{a_{L},\mu} + \frac{g_{0}^{2}}{2} W_{\mu}^{a_{L}} W^{a_{L},\mu} h^{2} \\ &- \frac{1}{4} B_{\mu\nu} B^{\mu\nu} + g_{0}' B_{\mu} J^{3_{R},\mu} + \frac{g_{0}'^{2}}{2} B_{\mu} B^{\mu} h^{2}, \\ \mathcal{L}_{int}^{V} &= \frac{1}{2} \frac{g_{0}}{g_{\rho}} W_{\mu\nu}^{a_{L}} \rho^{a_{L},\mu\nu} + \frac{1}{2} \frac{g_{0}}{g_{\rho}} B_{\mu\nu} \rho^{3_{R},\mu\nu}. \end{aligned}$$



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#### **Form Factor Lagrangian**

$$\begin{aligned} \mathcal{L}_{eff}^{F} = &\overline{q_{l}} \not p \left( \Pi_{0}^{L}(p) + \Pi_{1}^{L}(p_{1}, p_{2})\Sigma_{i}\Sigma^{i} \right) q_{l} + \overline{t_{r}} \not p \left( \Pi_{0}^{R}(p) + \Pi_{1}^{R}(p_{1}, p_{2})\Sigma_{i}\Sigma^{i} \right) t_{r} + \\ &+ \overline{q_{l}} \left( M_{1}(p_{1}, p_{2})\Gamma^{i}\Sigma_{i} \right) t_{r} + h.c. \end{aligned}$$
$$\begin{aligned} \mathcal{L}_{eff}^{V} = &\frac{1}{2} \mathcal{P}^{\mu\nu} \left( \Pi_{0}(p) \operatorname{Tr}(A_{\mu}A_{\nu}) + \Pi_{1}(p_{1}, p_{2})\Sigma^{T}A_{\mu}A_{\nu}\Sigma \right) \end{aligned}$$

$$\begin{split} \Pi_0^L(p) &= 1 + \Pi_4^L(p) = 1 + \frac{f^2 |y_L|^2}{p^2 - m_4^2}, \qquad \Pi_0^R(p) = 1 + \Pi_1^L(p) = 1 + \frac{f^2 |y_L|^2}{p^2 - m_1^2} \\ \Pi_1^L(p_1, p_2) &= \Pi_1^L(p_1) - \Pi_4^L(p_2) = f^2 |y_L|^2 \left(\frac{1}{p_1^2 - m_1^2} - \frac{1}{p_2^2 - m_4^2}\right), \qquad L \leftrightarrow R \\ M(p_1, p_2) &= M_4(p_1) - M_1(p_2) = f^2 y_L y_R \left(\frac{m_4}{p_1^2 - m_4^2} - \frac{m_1}{p_2^2 - m_1^2}\right), \end{split}$$

### **Heavy scalars**

$$V = M_a^2 H_a^{\dagger} H_a + M_b^2 H_b^{\dagger} H_b + \mu^2 (e^{i\theta} H_a^{\dagger} H_b + h.c.)$$
$$\frac{\lambda}{2} (H_a^{\dagger} H_a + H_b^{\dagger} H_b)^2 + \lambda' (H_a^{\dagger} H_b H_b^{\dagger} H_a).$$

Form Factor

$$\Gamma_{hVV}(p^2) = g_V^2 v \left( 1 - \frac{\mu^4}{4M_b^2} \frac{1}{p^2 - M_b^2} \right)$$

### **Broad Resonances**

$$G^{(2)}(p^2) = \frac{iZ}{p^2 - M_X^2(p^2) + i\sqrt{s}\Gamma(s)}.$$

