

Precision Calculation of Dark Radiation from Spinning Primordial Black Holes and Early Matter Dominated Eras

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PHENO

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Outline

- Formation and evaporation of spinning primordial black holes (PBHs) in an early matter dominated era (EMDE)
- Precision calculation of dark radiation (graviton) emitted by spinning PBHs
- Constraints on PBHs from BBN, CMB, and the projected sensitivity of CMB Stage 4

PBHs: motivated by many different scenarios

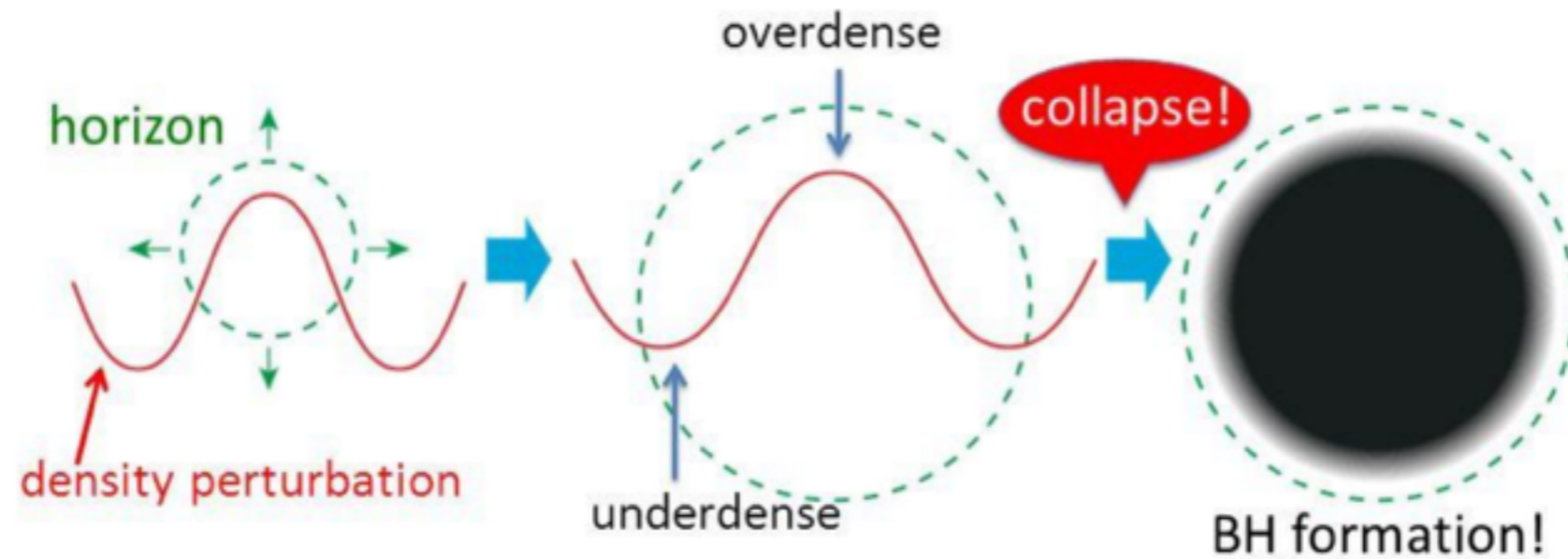


image credit: <https://slideplayer.com/slide/7773485/>

BHs can be fully characterized by their mass, spin, and charge

mass:

$$M_i \sim \frac{4}{3}\pi H(t_i)^{-3} \rho(t_i)$$

mass distribution:

- 1) monochromatic, e.g., if the PBH production occurs at a precise time.
- 2) extended, e.g., if the power spectrum of primordial inhomogeneities embeds a wide peak around some spatial scale

spin (more involved):

Radiation-dominated: small angular momentum

Matter-dominated: large angular momentum

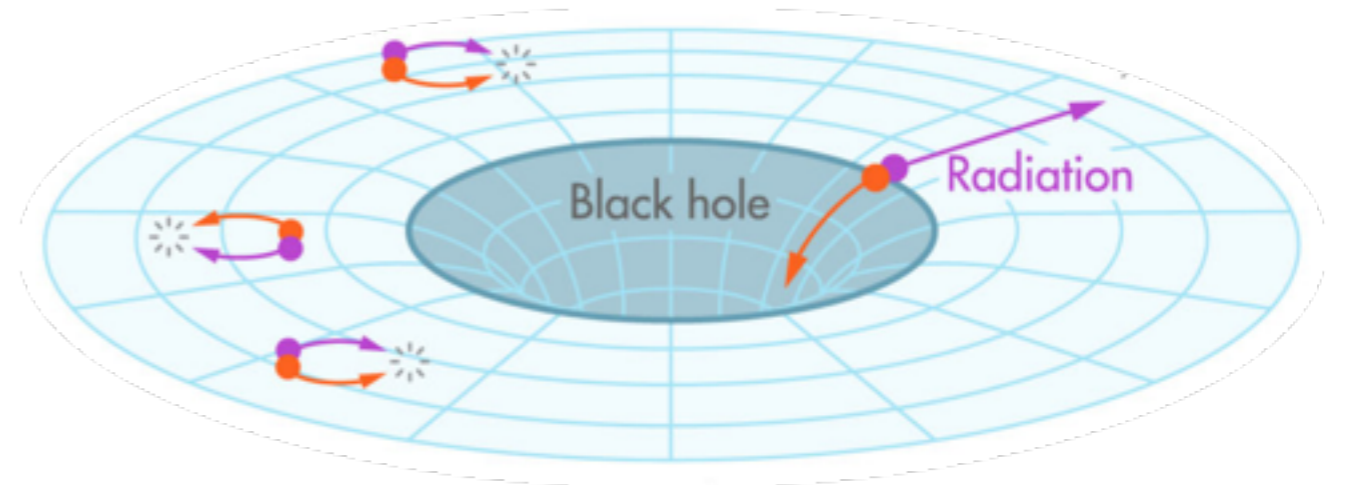
can also gain spin either through early accretion or hierarchical merger

Hawking evaporation of Schwarzschild BHs:

$$T_{\text{BH}} = \frac{M_{\text{Pl}}^2}{8\pi M_{\text{BH}}}$$

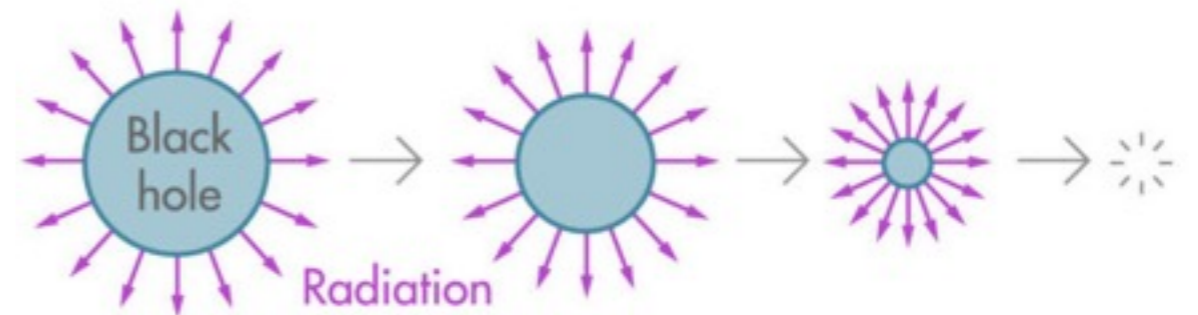
image credit: Quantamagazine

$$\frac{d^2 u_i(E, t)}{dt dE} = \frac{g_i}{8\pi^2} \frac{E^3}{e^{E/T_{\text{BH}}} \pm 1}$$



$$M(t) = M_i \left(1 - \frac{(t - t_i)}{\tau}\right)^{1/3}, \quad \tau = \frac{10240\pi}{g_*(T_{\text{BH}})} \frac{M_i^3}{M_{\text{Pl}}^4}$$

$$N_i = \frac{120 \zeta(3)}{\pi^3} \frac{g_i}{g_*(T_{\text{BH}})} \frac{M_{\text{BH}}^2}{M_{\text{Pl}}^2}, \quad T_{\text{BH}} > m_i$$



$$N_i = \frac{15 \zeta(3)}{8\pi^5} \frac{g_i}{g_*(T_{\text{BH}})} \frac{M_{\text{Pl}}^2}{m_i^2}, \quad T_{\text{BH}} < m_i$$

Hawking evaporation of Kerr BHs:

$$T_K = \frac{1}{2\pi} \left(\frac{r_+ - M_{\text{BH}}}{r_+^2 + a^{*2} M_{\text{BH}}^2} \right)$$

$$a^* \equiv L/M_{\text{BH}}^2$$

$$r_+ \equiv M_{\text{BH}}(1 + \sqrt{1 - a^{*2}})$$

$$\frac{d^2 N_i}{dt dE} = \frac{1}{2\pi} \frac{\Gamma_{s_i}^{l,m}}{e^{E'/T_K} - (-1)^{2s_i}}$$

$$E' \equiv E - m\Omega = E - ma^*/2r_+$$

$$\frac{dM_{\text{BH}}}{dt} = - \frac{f(M_{\text{BH}}, a^*)}{M_{\text{BH}}^2},$$

$$\frac{da^*}{dt} = \frac{a^* [2f(M_{\text{BH}}, a^*) - g(M_{\text{BH}}, a^*)]}{M_{\text{BH}}^3}$$

$$f(M_{\text{BH}}, a^*) \equiv M_{\text{BH}}^2 \int_0^{+\infty} \sum_i \sum_{\text{dof}} \frac{E}{2\pi} \frac{\Gamma_{s_i}^{l,m}(E, M_{\text{BH}}, a^*)}{e^{E'/T_K} - (-1)^{2s_i}} dE,$$

$$g(M_{\text{BH}}, a^*) \equiv \frac{M_{\text{BH}}}{a^*} \int_0^{+\infty} \sum_i \sum_{\text{dof}} \frac{m}{2\pi} \frac{\Gamma_{s_i}^{l,m}(E, M_{\text{BH}}, a^*)}{e^{E'/T_K} - (-1)^{2s_i}} dE$$

BlackHawk

[Arbey, Auffinger1905.04268](#)

Kerr BHs have enhanced emission of particles with higher spin.

Gravitons: impact the effective number of relativistic degrees of freedom ΔN_{eff}

$$\Delta N_{\text{eff}}$$

assuming PBHs are abundant enough to initiate an EMDE and reheat the universe via Hawking evaporation:

$$\Delta N_{\text{eff}} = \frac{\rho_{\text{DR}}(t_{\text{EQ}})}{\rho_{\text{R}}(t_{\text{EQ}})} \left[N_{\nu} + \frac{8}{7} \left(\frac{11}{4} \right)^{4/3} \right]$$

$$\frac{\rho_{\text{DR}}(t_{\text{EQ}})}{\rho_{\text{R}}(t_{\text{EQ}})} = \frac{\rho_{\text{DR}}(t_{\text{RH}})}{\rho_{\text{R}}(t_{\text{RH}})} \left(\frac{g_*(T_{\text{RH}})}{g_*(T_{\text{EQ}})} \right) \left(\frac{g_{*,S}(T_{\text{EQ}})}{g_{*,S}(T_{\text{RH}})} \right)^{4/3} \quad N_{\nu} = 3.046$$

Hooper, Krnjaic, McDermott 1905.01301

current experimental limits and projected sensitivities of future experiments:
CMB(two are taken from the Planck)

CMB¹ (TT+low E, conservative)

CMB² (TT,TE,EE+low E, more stringent)

BBN: AlterBBN

the sensitivity of the future CMB Stage 4 (CMB-S4) experiment

precision calculation:

- extended spin distribution

$$\int_0^1 \frac{dn}{da^*} da^* = 1$$

$$\rho_{\text{DR/SM}}(t_{\text{RH}}) = \int_0^1 da^* \frac{dn}{da^*} \int_0^{t_{\text{RH}}} dt \int_0^{+\infty} dE E \frac{d^2 N_{\text{DR/SM}}}{dt dE}(M, a^*)$$

- reheating temperature

- 1) RH time = time at which last PBH (with lowest spin evaporates)
- 2) RH time = the average PBH lifetime, weighted by the spin distribution

- degrees of freedom

$$g_*(T), g_{*,S}(T)$$

SuperIso Relic

benchmark spin distributions:

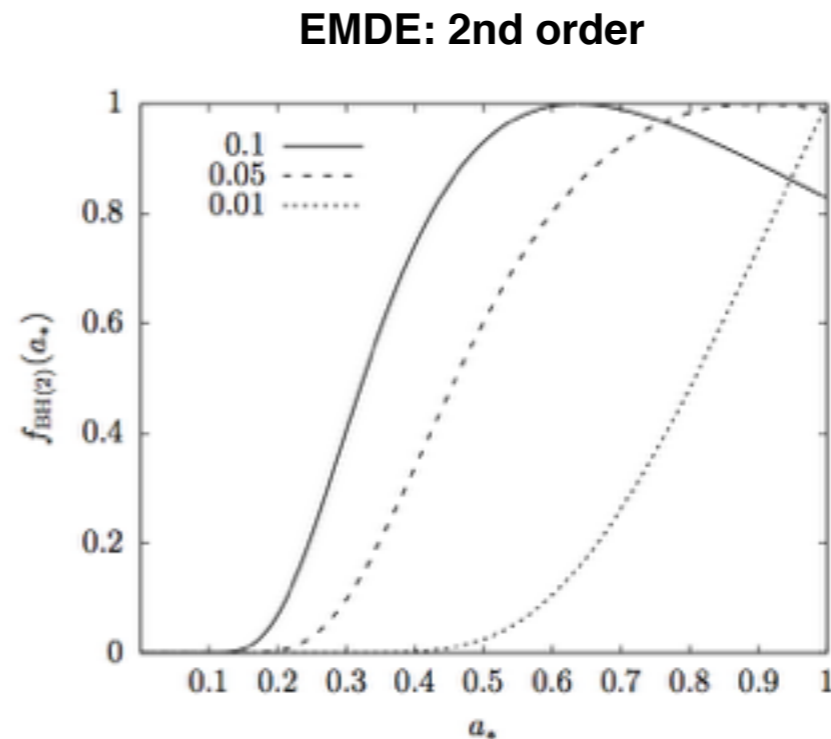
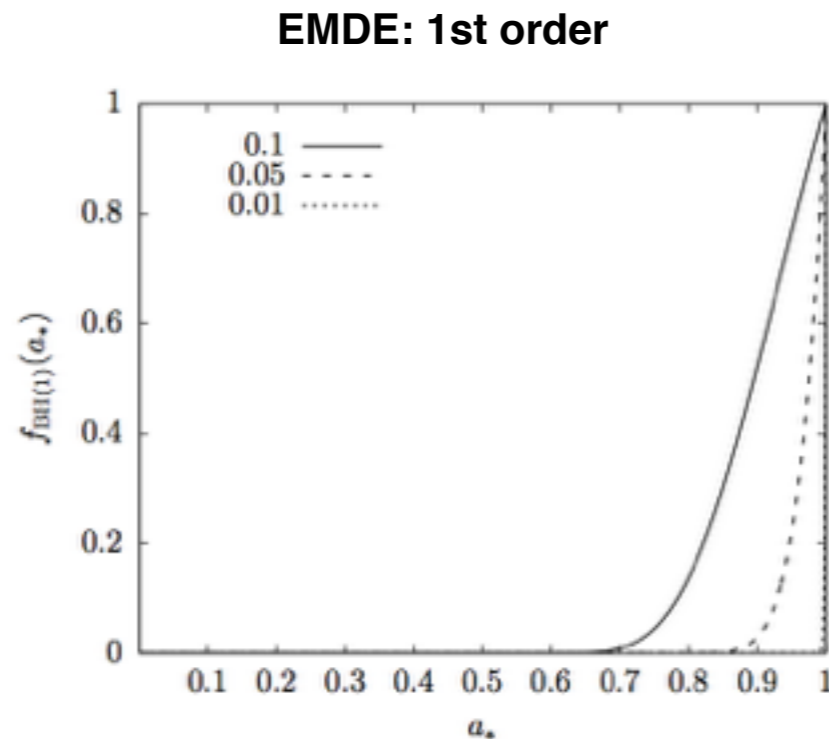
1) Early matter domination era (EMDE):

spin source:

- the first-order contribution: from deviation of the boundary of the volume from a sphere.
- the second-order contribution: density fluctuations in the comoving region.

The first-order effect usually dominates

(when the initial deviation of the boundary of collapsing region from a sphere is large)



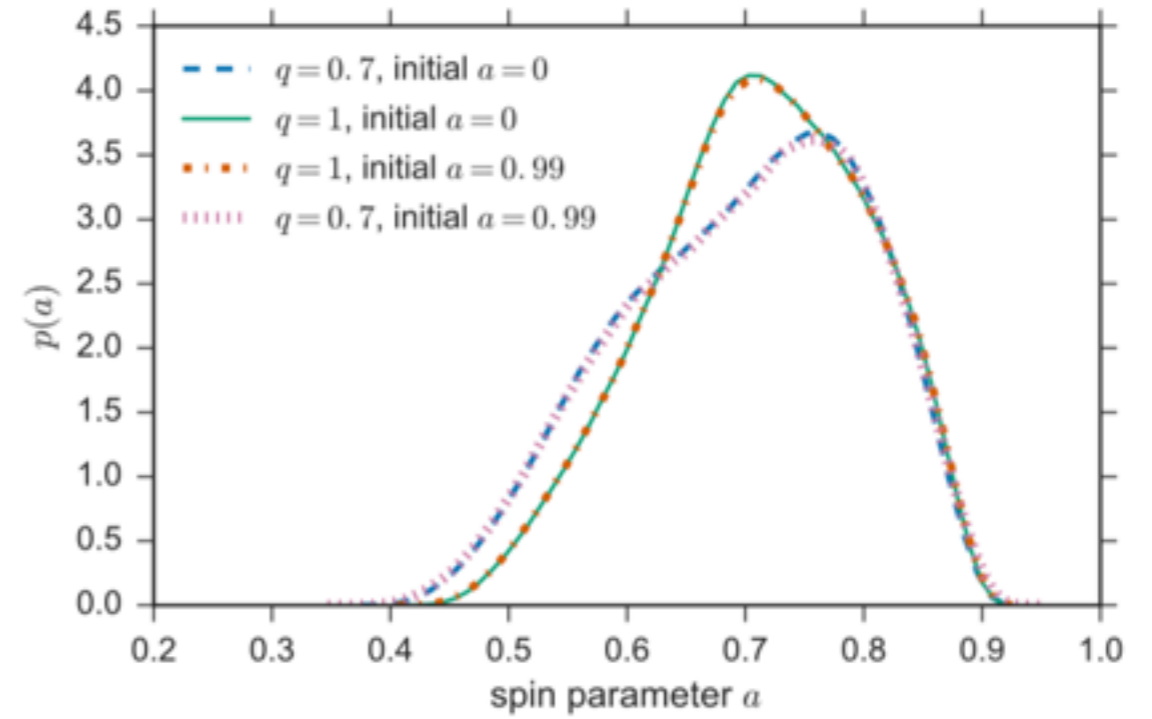
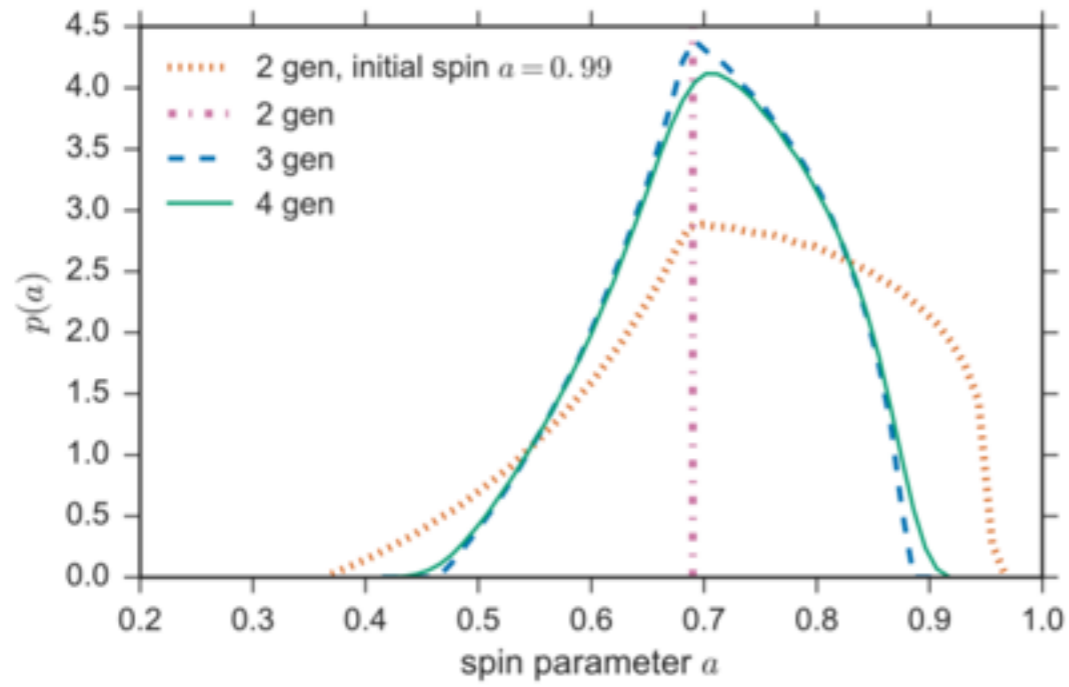
Harada, Yoo, Kohri, Nakao 1707.03595

In both cases, spin distribution is a function of the mean variance of the density perturbations at horizon entry

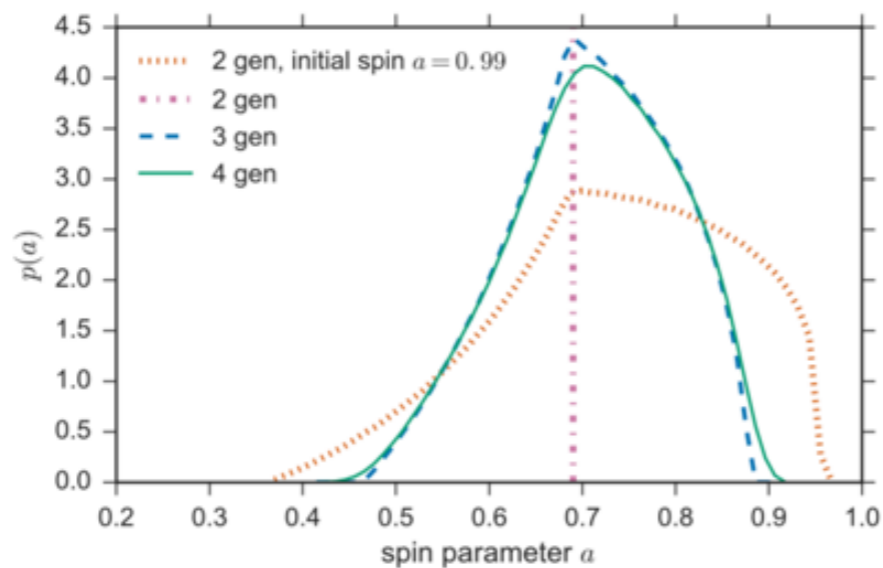
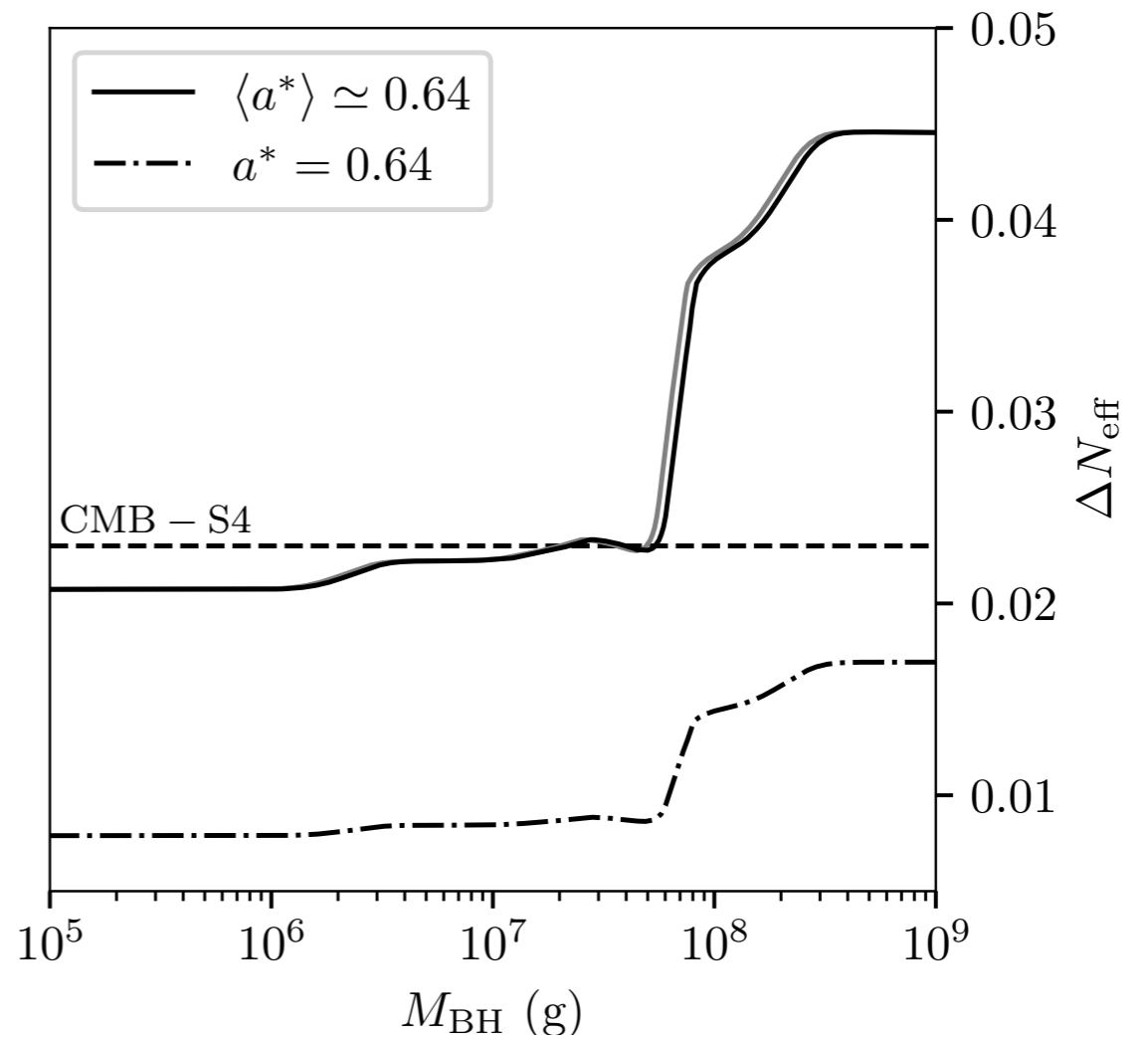
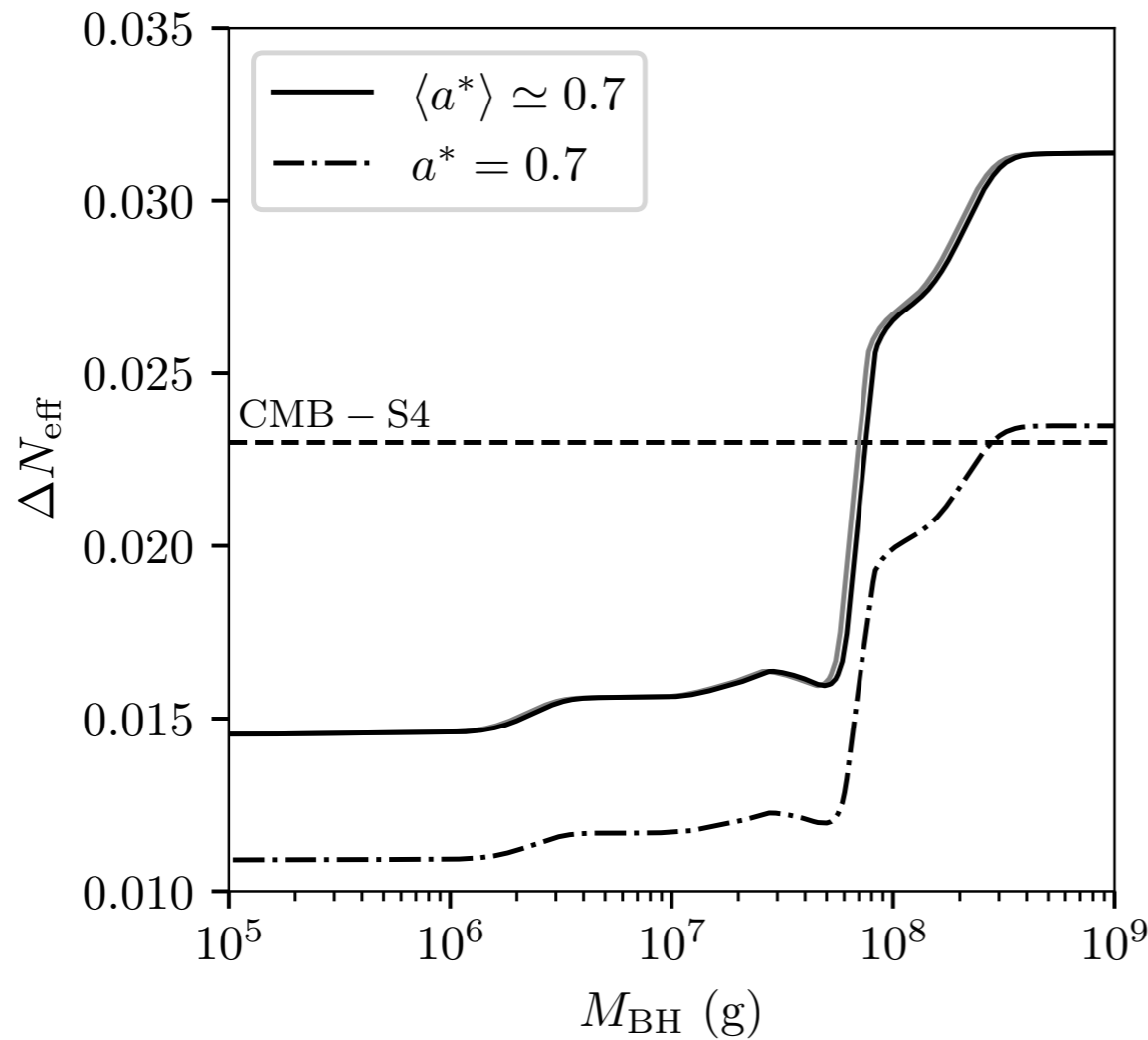
$$\sigma_H = \langle \delta_s(t_H)^2 \rangle^{1/2}$$

2) distribution from inspirals:

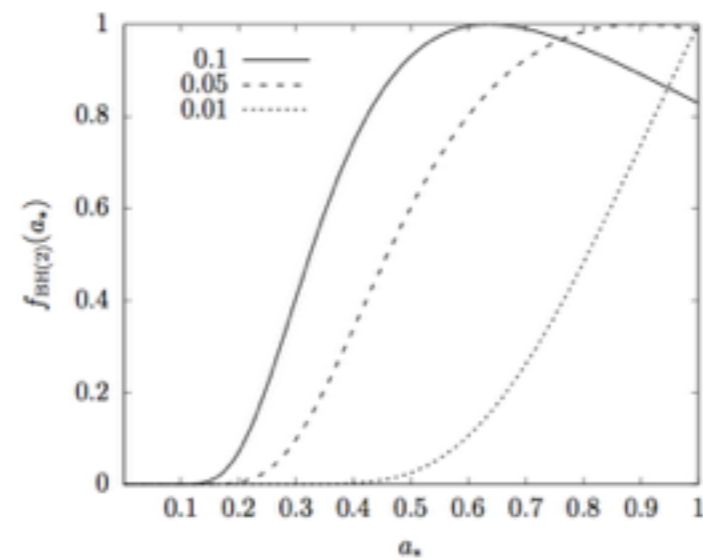
Fishbach, Holz, Farr 1703.06869



PBH spins compares to the monochromatic approximation



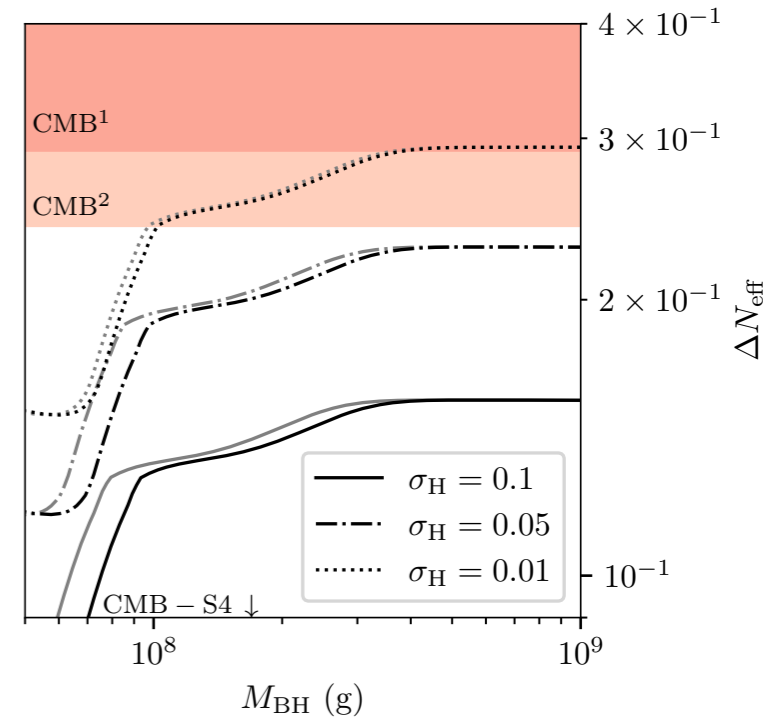
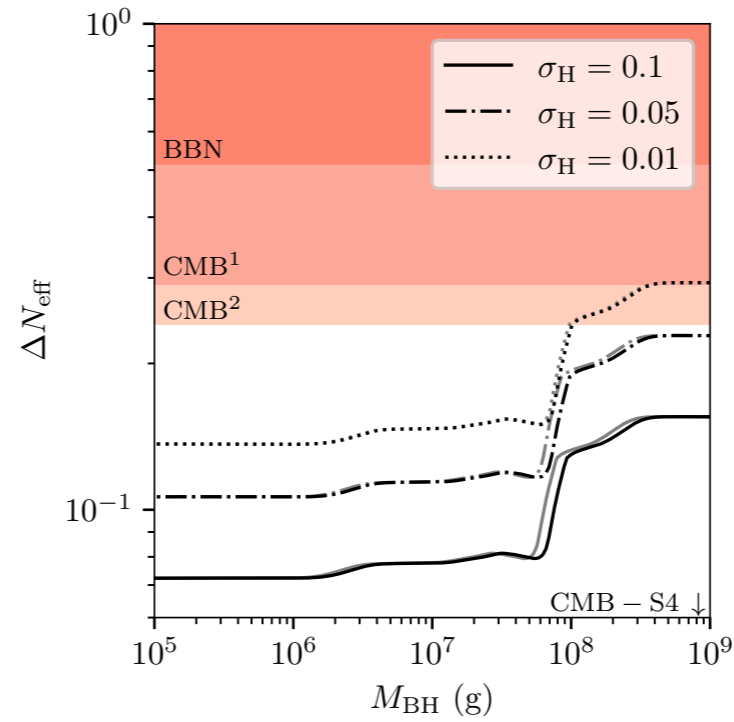
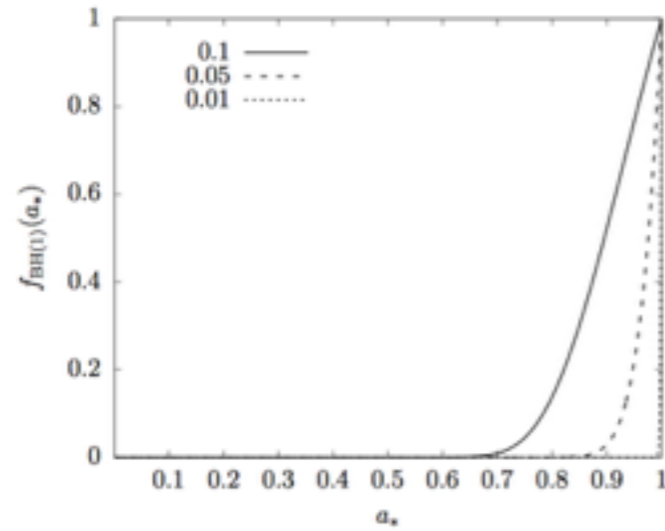
mergers



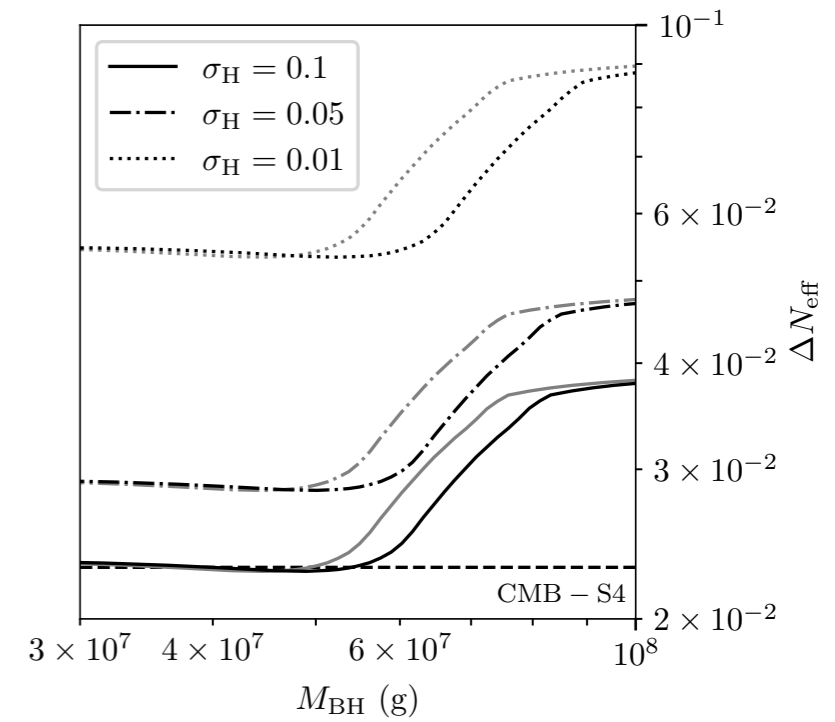
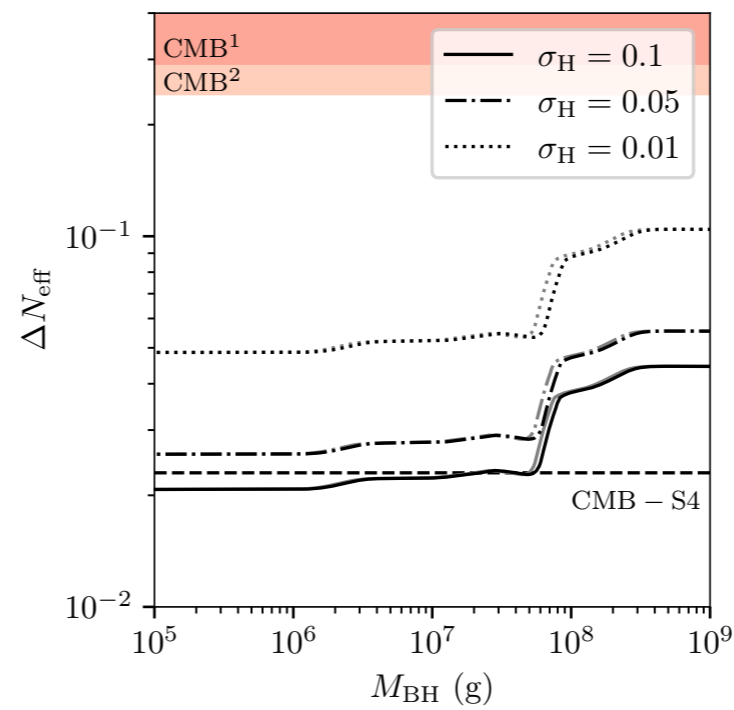
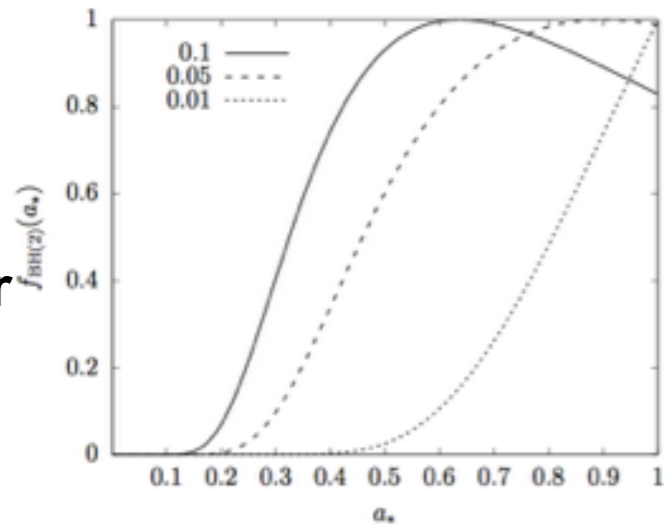
EMDE: 2nd order

Early Matter domination:

EMDE:
1st order

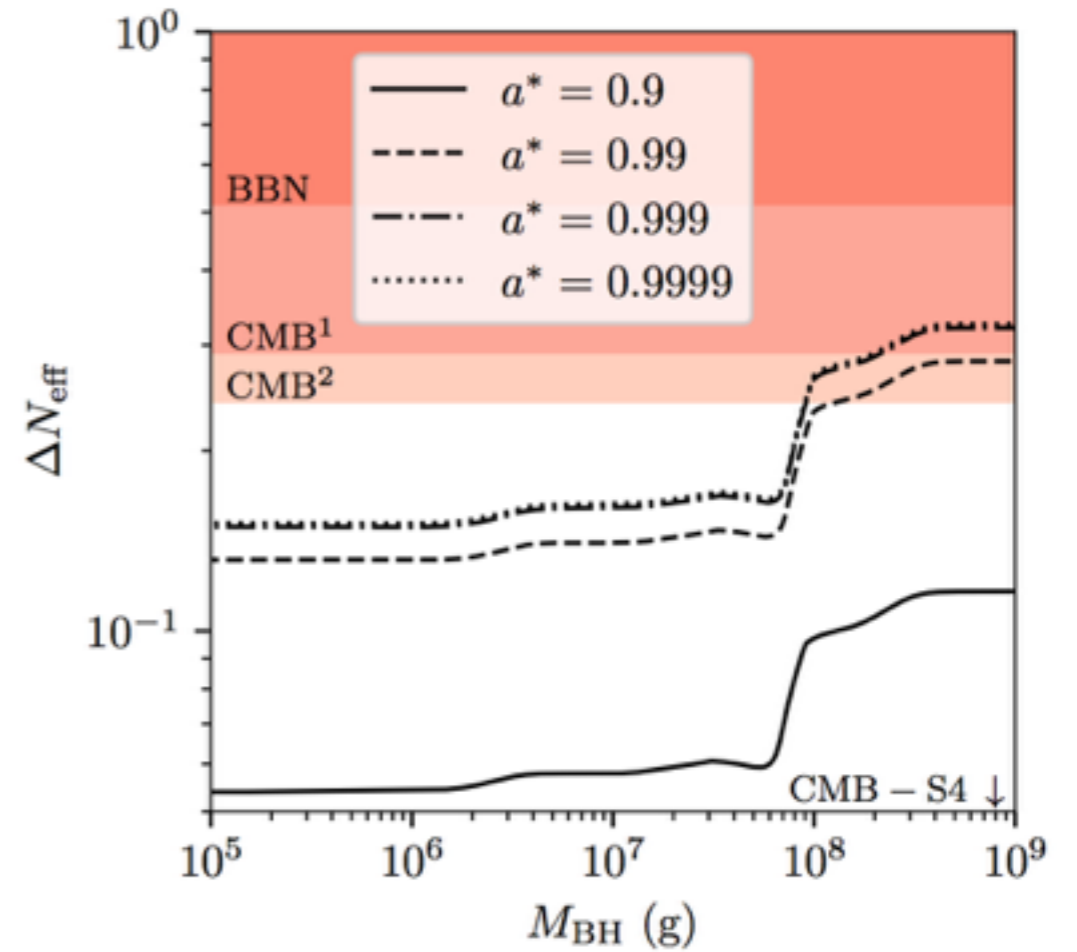


EMDE:
2nd order



near-extremal PBH spins (monochromatic spin distribution):

- Target PBH masses already excluded.
- CMB-S4 will probe all near-extremal monochromatic spin distributions.



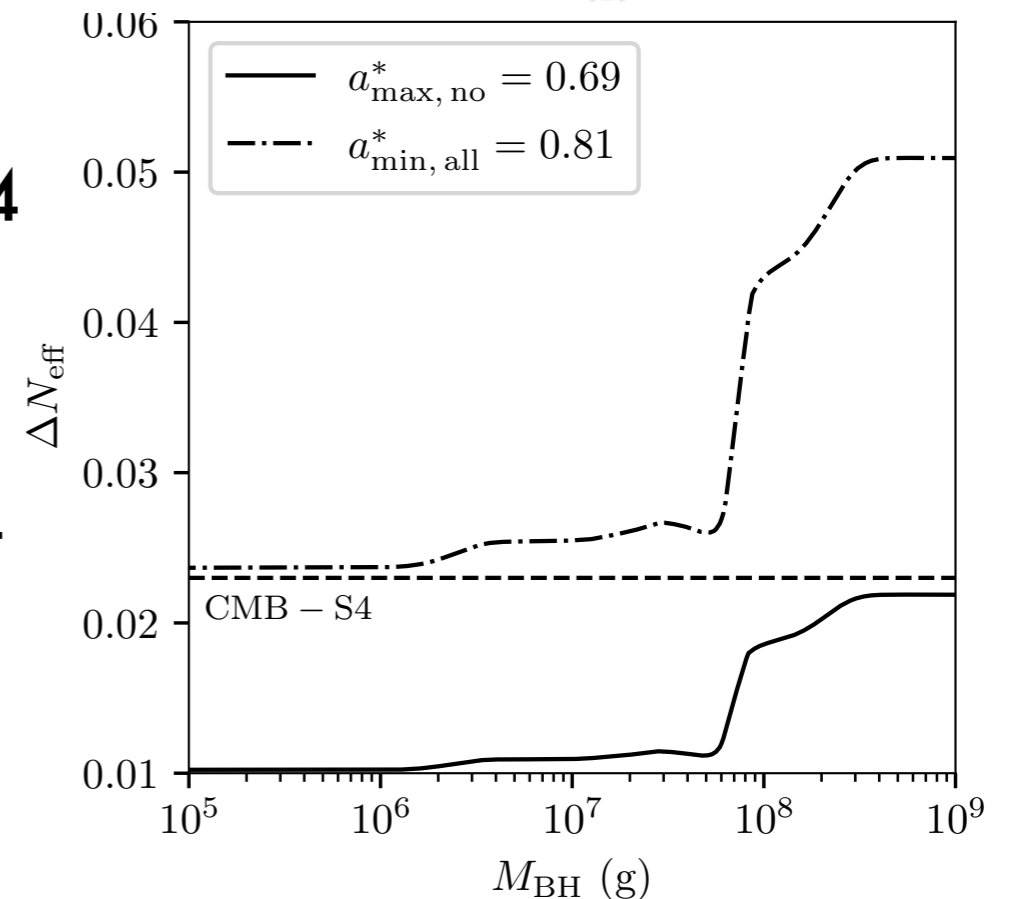
Low and high cut-off for CMB-S4:

- The smallest monochromatic spin for which CMB-S4 will be sensitive to the entire mass range:

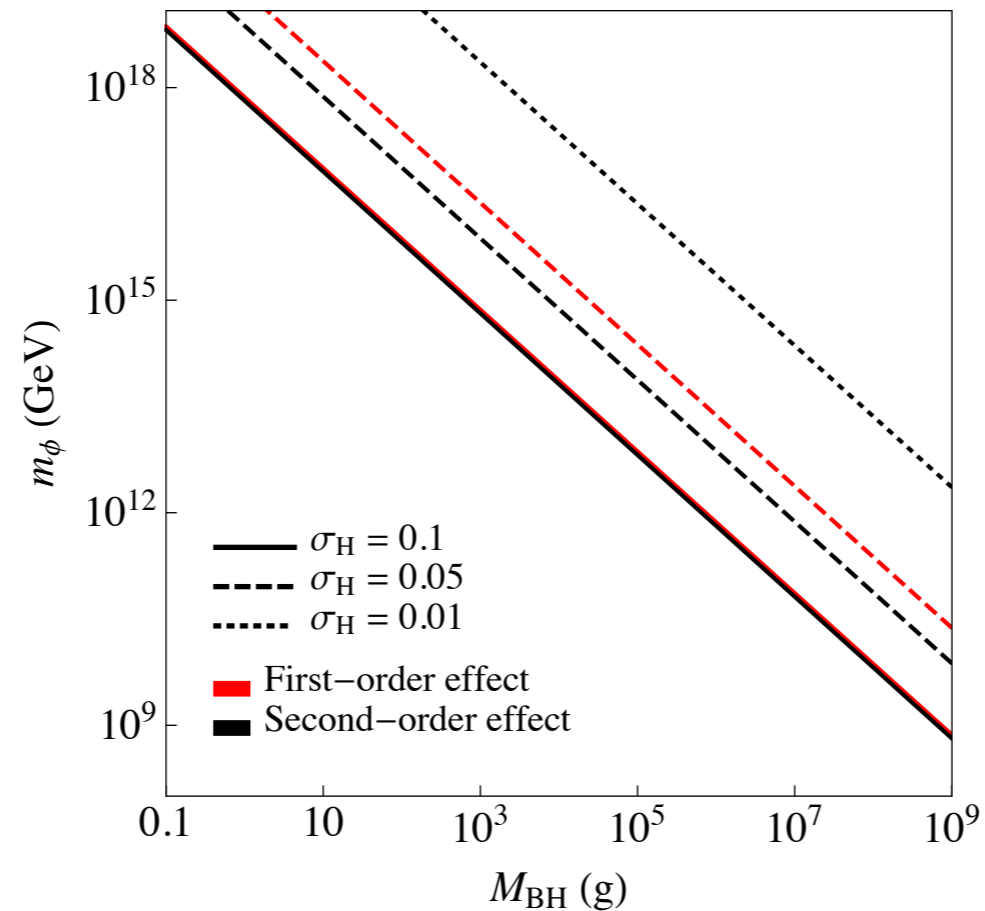
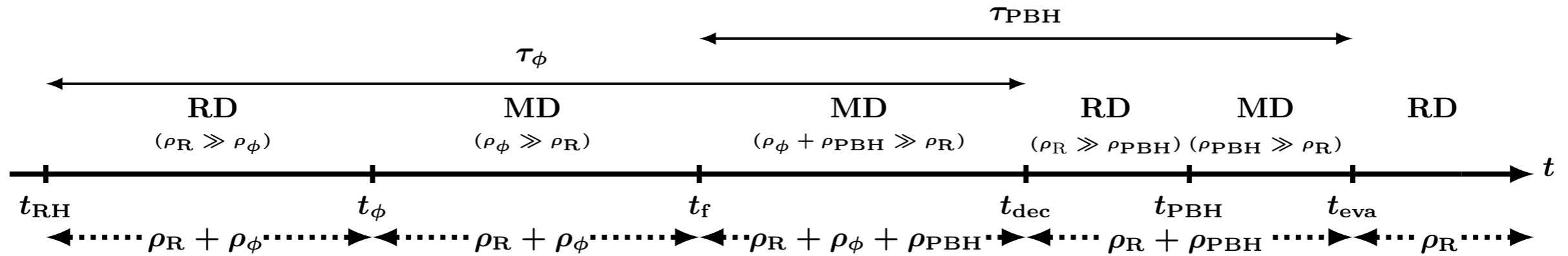
$$a_{\text{min, all}}^* = 0.81$$

- The largest monochromatic spin for which CMB-S4 will not be sensitive to any part of the mass range:

$$a_{\text{max, no}}^* = 0.69$$



Early matter domination era leading to a PBH-dominated era: a gravitationally coupled modulus field



Conclusions:

- The first precision study of dark radiation emission by Kerr PBHs
- Precision study includes: extended spin distributions, careful prescription for reheating time, precise evaluation of degrees of freedom.
- precision results for PBHs formed during an EMDE.
- Provided PBHs dominate the energy density of the Universe before evaporation, CMB-S4 will be sensitive to nearly all EMDE distributions and monochromatic spins greater than 0.69.