

# Space Charge Tune Shift

JAI lectures - Hilary Term 2021

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## Goal of this course

- ▶ Brief introduction to a complex topic.
  - ▶ Understand the concept of Space Charge.
  - ▶ Distinguish different types of effects.
  - ▶ Understand the limitations imposed in the accelerator operation and performance.
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  - ▶ For more detailed derivations and more realistic cases please go to references.

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# References

## Specialized courses

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<https://indico.cern.ch/event/774280/contributions/3217261/>.
2. K.Schindl "Space charge", CAS lectures <https://cds.cern.ch/record/941316?>.
3. M.Migliorati, "Space Charge Effects and Instabilities",  
<https://indico.cern.ch/event/779575/contributions/3244564/>.

## Books

1. I. Hofmann "Space Charge Physics for Particle Accelerators", Springer 2017.
2. H. Wiedemann "Particle Accelerator Physics", Springer 2015.

# Introduction

The beam is a distribution of charged particles i.e. they create an EM fields that affect their dynamics.

We can distinguish three main contributions:

- ▶ Direct fields.
- ▶ Image fields.
- ▶ Wakefields (we will cover that in Instabilities lectures).

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# Space Charge Forces

Electric field generated by a point-like charge  $q$ :

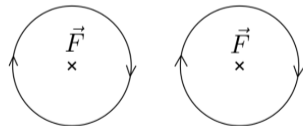
$$F_{\text{elec}} = \frac{e^2}{4\pi\epsilon r^2} \quad (1)$$



Repulsive!

Since the particle is moving with some speed  $v$ , this is equivalent to a current carrying wire with  $I = qv$ .

$$F_{\text{wire}} = \frac{\mu_0 I}{4\pi r^2} = \frac{v^2}{c^2} F_{\text{elec}} \quad (2)$$



Attractive!



## Space Charge Forces

The overall force is repulsive:

$$F_{\text{total}} = (1 - v^2/c^2)F_{\text{elec}} \quad (3)$$

we see that for  $v \rightarrow c$  the force  $F_{\text{total}}$  vanishes.

What does this mean?

## Space Charge Forces

Two main regimes exist to describe the effects of Coulomb interactions in a system with many particles.

Which regime are we? Debye length

$$\lambda_D = \sqrt{\frac{\epsilon_0 \gamma^2 k_B T}{q^2 n}} \quad (4)$$

### Collisional regime

Dominated by particle-on-particle collisions and described by single particle dynamics.

$$\lambda_D \gg a$$

### Space Charge regime

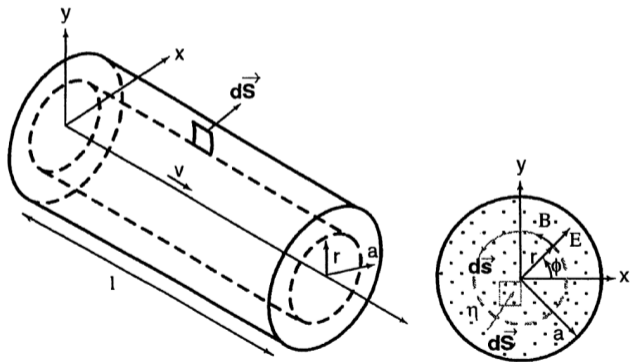
Dominated by the self fields of the particle distribution and it is described by collective effects.

$$\lambda_D \ll a$$

## Space Charge Forces

Simple model: beam as a continuous cylinder of charge  $q$ , length  $l$  and radius  $a$ .

$$\rho(r) = qn(r) = \frac{I_{\text{beam}}}{\pi a^2 v} \quad (5)$$



# Space Charge Forces

## Electric field

$$\nabla \cdot \vec{E} = \frac{\eta}{\epsilon_0} \quad (6)$$

Gauss' law:

$$\int_V \nabla \cdot \vec{E} dV = \int_S \vec{E} d\vec{S} \quad (7)$$

cylinder of radius  $r$  and length  $l$ :

$$\pi r^2 l \frac{\eta}{\epsilon_0} = E_R 2\pi r l \quad (8)$$

$$E_r = \frac{l}{2\pi\epsilon_0\beta c} \frac{r}{a^2} \quad (9)$$

## Magnetic field

$$\nabla \times \vec{B} = \mu_0 \vec{J} \quad (10)$$

Stoke's law:

$$\oint \vec{B} d\vec{S} = \int_S \nabla \times \vec{B} d\vec{S} \quad (11)$$

$$B_\phi 2\pi r = \mu_0 \pi r^2 \beta c \eta \quad (12)$$

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## Space Charge Forces

The force acting on a test particle is given by the Lorentz equation:

$$F_r = q(E_r - v_s B_\phi) \quad (14)$$

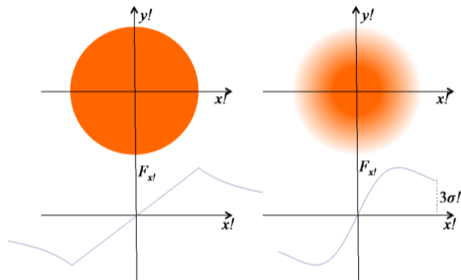
where:

$$F_r = \frac{el}{2\pi\epsilon_0\beta c\gamma^2} \frac{r}{a^2} \quad (15)$$

and in transverse coordinates:

$$F_x = \frac{el}{2\pi\epsilon_0\beta c\gamma^2 a^2} x, \quad F_y = \frac{el}{2\pi\epsilon_0\beta c\gamma^2 a^2} y \quad (16)$$

## Space Charge Forces: circular vs. Gaussian beam



**Figure:** Space charge force for a homogeneous circular beam (left) and a Gaussian-shaped beam (right).

## Self field tune shift

SC produces an extra defocusing. Let's include it in the Hill's equation:

$$x'' + (K(s) + K_{SC}(s))x = 0 \quad (17)$$

$$x'' + \left( K(s) - \frac{2r_0 I}{ea^2 \beta^3 \gamma^3 c} \right) x = 0 \quad (18)$$

Tune shift due to an error in focusing strength  $\Delta K$ :

$$\Delta Q_{x,y} = \frac{1}{4\pi} \int \Delta K(s) \beta_{x,y}(s) ds \quad (19)$$

In our case  $\Delta K = K_{SC}$ :

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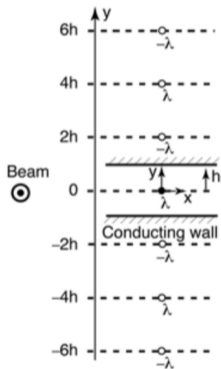
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## Image Effects

A second effect is coming from image currents due to conducting walls.



Electric field produced by a charge  $\lambda$  at a distance  $2n \cdot d$ :

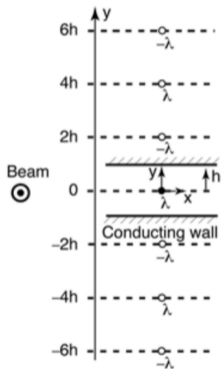
$$E_y = \frac{\lambda}{2\pi\epsilon_0} \frac{1}{d} \quad (23)$$

$$E_{2h} = \frac{\lambda}{2\pi\epsilon_0} \left( \frac{1}{2h-y} - \frac{1}{2h+y} \right) \quad (24)$$

$$E_{4h} = \frac{\lambda}{2\pi\epsilon_0} \left( \frac{1}{4h-y} - \frac{1}{4h+y} \right) \quad (25)$$

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## Image Effects

Let's do some algebra:

$$E_{inh} = \quad (26)$$

$$= (-1)^n \frac{\lambda}{2\pi\epsilon_0} \left( \frac{1}{2nh+y} - \frac{1}{2nh-y} \right) = \quad (27)$$

$$= (-1)^n \frac{\lambda}{4\pi\epsilon_0} \frac{y}{n^2 h^2} \quad (28)$$

$$E_{iy} = \sum_{n=1}^{\infty} E_{iny} = \frac{\lambda}{4\pi\epsilon_0 h^2} \sum_{n=1}^{\infty} \frac{(-1)^n}{n^2} y = \quad (29)$$

$$= \frac{\lambda}{4\pi\epsilon_0 h^2} \frac{\pi^2}{12} y \quad (30)$$

We obtain the corresponding fields and forces:

$$E_{ix} = -\frac{\lambda}{4\pi\epsilon_0 h^2} \frac{\pi^2}{12} x \quad (31)$$

$$F_{iy} = \frac{e\lambda}{\pi\epsilon_0 h^2} \frac{\pi^2}{48} y \quad (32)$$

$$F_{iy} = -\frac{e\lambda}{\pi\epsilon_0 h^2} \frac{\pi^2}{48} x \quad (33)$$

## Incoherent Tune Shift

The total contribution to the incoherent tune shift can be summarized:

$$\Delta Q_x = -\frac{2r_0 I_b R \langle \beta_x \rangle}{qc \beta^3 \gamma} \left( \frac{1}{2 \langle a^2 \rangle \gamma^2} - \frac{\pi^2}{48^2} \right) \quad (34)$$

$$\Delta Q_y = -\frac{2r_0 I_b R \langle \beta_y \rangle}{qc \beta^3 \gamma} \left( \frac{1}{2 \langle a^2 \rangle \gamma^2} + \frac{\pi^2}{48^2} \right) \quad (35)$$

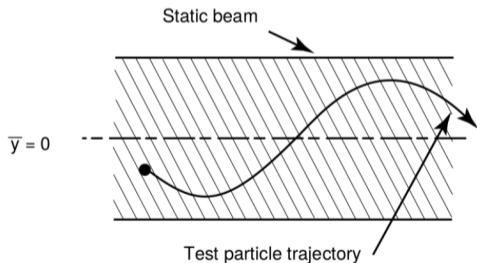
- ▶ Direct field.
- ▶ Image field.



# Coherent vs Incoherent effects

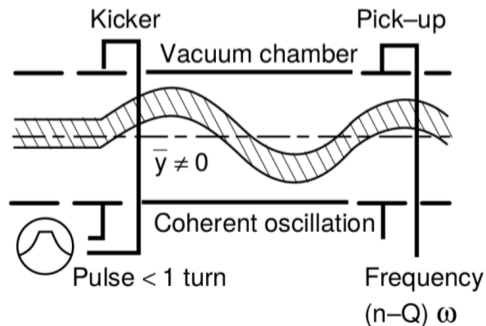
## Incoherent

Each particle is independent (has its own betatron oscillation, phase and tune). Impossible to observe any betatron motion. The beam "does not move".



## Coherent

The kick gets a fast deflection that affects the full distribution and starts to perform betatron oscillations as a whole. The source of space charge is now moving.



## Coherent Tune Shift

Taking  $\rho$  the beam pipe radius and  $\bar{x}$  the center of mass position.  
Image charge at  $b = \rho^2/\bar{x}$ .

$$E_{ix} = \frac{\lambda}{2\pi\epsilon_0} \frac{1}{b - \bar{x}} \approx \frac{\lambda}{2\pi\epsilon_0} \frac{1}{b} = \frac{\lambda}{2\pi\epsilon_0} \frac{\bar{x}}{\rho^2} \quad (36)$$

$$F_{ix} = \frac{e\lambda}{2\pi\epsilon_0} \frac{\bar{x}}{\rho^2} \quad (37)$$

$$\Delta Q_{x,y} = -\frac{r_0 R \langle \beta_{x,y} \rangle I}{ec\beta^3 \gamma \rho^2} = -\frac{r_0 \langle \beta_{x,y} \rangle N}{2\pi\beta^2 \gamma \rho^2} \quad (38)$$

- ▶ The force is linear in  $\bar{x}$ .
- ▶  $1/\gamma$  dependence.
- ▶ The coherent tune shift is never positive.
- ▶ Perfectly conducting beampipe assumed. Realistic effects are delicate.

## Laslett coefficients

A more realistic scenario is when we consider elliptic, unbunched uniformly distributed beams travelling at a speed  $\beta c$  through an elliptic vacuum chamber. For these geometries the tune shift can be expressed in terms of the "laslett coefficients".

Incoherent:  $\epsilon_{0,1,2}$ , Coherent:  $\xi_{1,2}$

$$\Delta Q_{y,inc.} = \frac{-Nr_0 \langle \beta_y \rangle}{\pi \beta^2 \gamma} \left( \frac{\epsilon_0^y}{b^2 \gamma^2} + \frac{\epsilon_1^y}{h^2} + \beta^2 \frac{\epsilon_2^y}{g^2} \right) \quad (39)$$

$$\Delta Q_{y,coh.} = \frac{-Nr_0 \langle \beta_y \rangle}{\pi \beta^2 \gamma} \left( \frac{\xi_1^y}{h^2} + \beta^2 \frac{\xi_2^y}{g^2} \right) \quad (40)$$

Laslett coefficients	Circular ( $a = b, w = h$ )	Elliptical (e.g. $w = 2h$ )	Parallel plates ( $h/w = 0$ )
$\epsilon_0^x$	1/2	$\frac{b^2}{a(a+b)}$	
$\epsilon_0^y$	1/2	$\frac{b}{a+b}$	
$\epsilon_1^x$	0	-0.172	-0.206
$\epsilon_1^y$	0	0.172	0.206
$\xi_1^x$	1/2	0.083	0
$\xi_1^y$	1/2	0.55	$0.617(\pi^2/16)$
$\epsilon_2^x$	$-0.411(-\pi^2/24)$	-0.411	-0.411
$\epsilon_2^y$	$0.411(\pi^2/24)$	0.411	0.411
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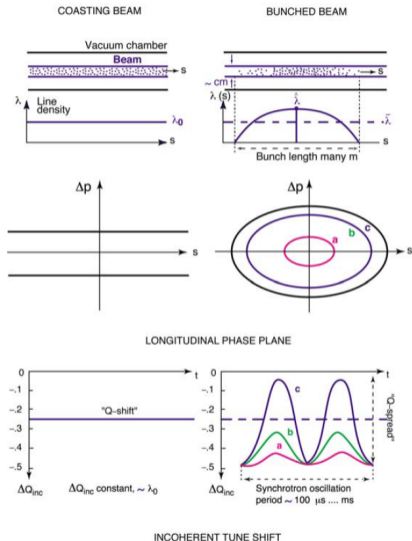
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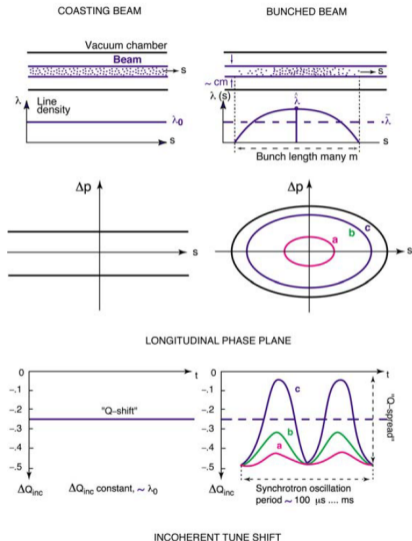
# Unbunched vs. bunched beams

- ▶ So far we just considered unbunched homogeneous beams. Create a constant tune shift. Easy to solve.
- ▶ When bunched beams are considered, the space charge effects are more notorious.
- ▶ In bunched beams, each "slice" of the beam feels a different space charge.
- ▶ Synchrotron oscillations modulate the space charge force felt by a single particle.
- ▶ This generates a tune spread.



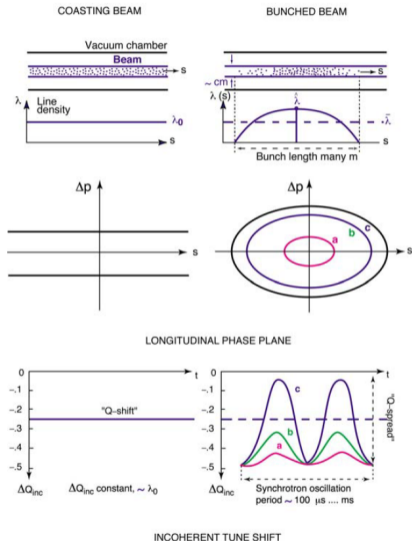
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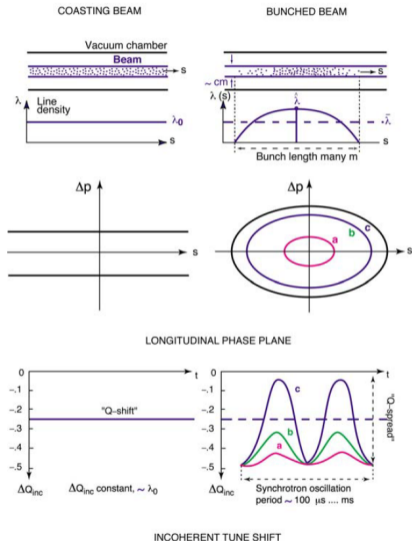
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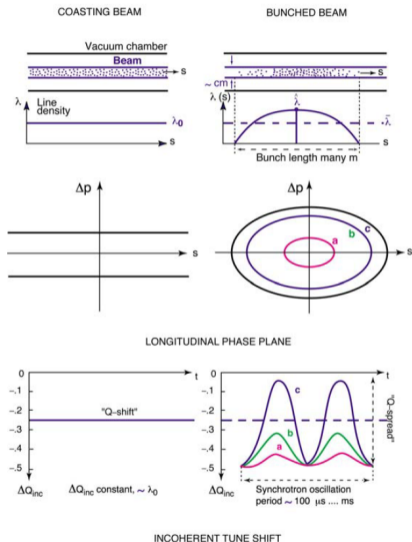
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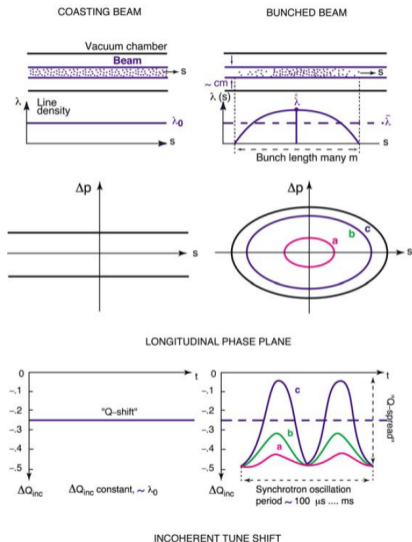
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## Space Charge Limit

Space Charge may limit the operation if the tune shift is too large and important resonances are crossed.

$$\Delta Q \sim \frac{N}{\beta^2 \gamma^2} \quad (41)$$

What can we do?

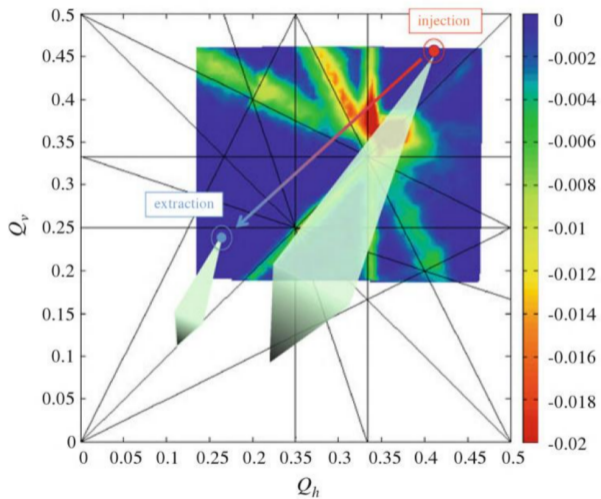
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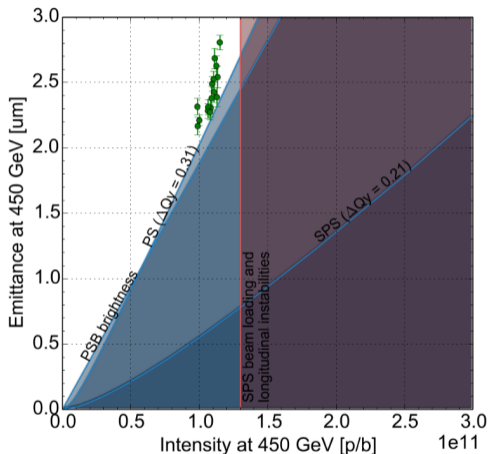
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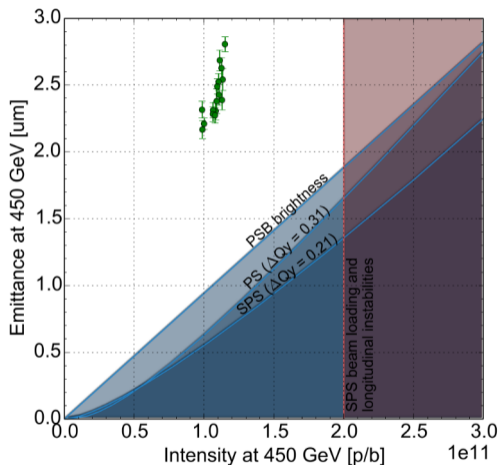


# Space Charge Limit: How to mitigate it



- ▶ At CERN accelerator chain, current injector configuration limits the bunch intensity.
- ▶ To overcome this limitation, a major upgrade of the injectors is required to achieve HL-LHC desired performance.
- ▶ In the example, we can see that the PSB, the PS and the SPS need to be upgraded.

# Space Charge Limit: How to mitigate it



## Linac

- ▶ Linac4 ( $H^-$ ) replaces Linac2 ( $H^+$ ).

## PSB

- ▶ Energy upgrade.
- ▶ Injection: 160 MeV (50 MeV).
- ▶ Extraction: 2 GeV (1.4 GeV).

## PS

- ▶ Replace 43 dipoles.

## SPS

- ▶ Cabling and Acceleration system.

# Summary

- ▶ Space Charge limits the performance of particle accelerators.
- ▶ Particular impact on low-energy hadron machines.
- ▶ We mainly focused on the induced tune shift and tune spread.
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Thank you!