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Flow fluctuations in heavy-ion collisions measured with ALICE



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THE VELUX FOUNDATIONS

Anisotropic flow

- Partial overlap ⇒ Spatial anisotropy ⇒
 ⇒ Different pressure gradients ⇒
 ⇒ Particles "flow"
- Fourier series decomposition of azimuthal distribution of emitted particles:

$$\frac{\mathrm{d}N}{\mathrm{d}\varphi} \propto 1 + 2\sum_{n=1}^{\infty} v_n \cos n(\varphi - \Psi_n)$$

- Ψ_n flow symmetry plane (defined by *xz*)
- v_n flow coefficients
 - Nowadays typically calculated using *m*-particle correlations, $v_n\{m\}$
- Together they make flow vector $\overrightarrow{V}_n = v_n e^{in\Psi_n}$







Measuring flow vector

- Two ways how to measure flow vector
 - Multi-particle correlations
 - Event-shape engineering
- What information can measuring of flow vector provide?
 - Initial-state conditions (initial geometry)
 - QGP properties, e.g. transport
 coefficients (shear viscosity η/s, bulk
 viscosity ξ/s) as a function of
 temperature



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- Event-by-event v_2 distribution
- Event-by-event fluctuations of the flow vector
- Correlations between different flow vectors



ALICE detector

- General-purpose heavy-ion experiment at the Large Hadron Collider
- Pb-Pb collisions at 5.02 TeV (taken in 2015 and 2018)





Flow of charged hadrons

- Ordering $v_2 > v_3 > v_4$
- Elliptic flow (v₂) is dominant in non-central heavy-ion collision where the overlap of colliding nuclei has an almond-like shape
- v_3 originates from fluctuations
- v₄ originates from the initial geometry and fluctuations
- Hydrodynamic description in a very good agreement with data





Multi-particle cumulants

- Initial geometry fluctuations result in fluctuations in the final state
 - $\langle v_2^k \rangle \neq \langle v_2 \rangle^k$
 - Probability density function of v_n can be studied using multi-particle cumulants (= genuine multi-particle correlations)
- Different orders can provide different information
 - $v_2\{2, |\Delta \eta| > 1\} > v_2\{4, 6, 8\}$, flow fluctuations
 - v_2 {4,6,8} $\approx v_2$ {4,6,8, $|\Delta \eta| > 0$ }, less sensitive to non-flow
 - Non-flow = correlations not associated with the common symmetry plane, originating e.g. from jets or resonance decays







Probability density function of v_2

- Bessel-Gaussian *p.d.f.* if $v_2{4} = v_2{6} = v_2{8}$
 - Agrees above approx. 3.5 GeV/c
- Possible to calculate higher-order moments of the p.d.f.
 - Skewness γ_1 negative at low $p_{\rm T}$
 - Kurtosis γ_2 positive at low $p_{\rm T}$
 - Consistent with Bessel-Gaussian *p.d.f.* above approx. 3.5 GeV/*c*
- Probability density function of v₂ has non-trivial transverse momentum dependence





P.d.f. with event-shape engineering

- Event-shape engineering (ESE) [1] allows a selection of initial geometry, e.g. specific eccentricities (larger/smaller)
- q_2 selection by constructing reduced flow vector, $q_n = \frac{|Q_n|}{\sqrt{M}}$ $|Q_n| = \sqrt{Q_{n,x}^2 + Q_{n,y}^2}$
- Selecting q₂ ⇒ selecting
 eccentricity ⇒ different v₂
 distribution (e.g. more skewed
 probability density function)





[1] Schukraft, Timmins, Voloshin PLB 719, 4–5, 394-398 (2013)

Flow vector fluctuations

 $\langle v_n(p_T^a) \cdot v_n(p_T^b) \rangle \neq \langle v_n(p_T^a) \rangle \cdot \langle v_n(p_T^b) \rangle$





Angle — final symmetry plane

How to disentangle contributions to fluctuations from flow magnitude and flow angle?

Hydrodynamic models show its $p_{\rm T}$ dependence





Flow magnitude and angle fluctuations

 $F(\Psi_2^a, \Psi_2)$

0.9

0.8

Magnitude: $\frac{\langle \langle \cos n(\varphi_1^a + \varphi_2 - \varphi_3^a - \varphi_4) \rangle \rangle}{\langle \langle \cos n(\varphi_1^a - \varphi_3^a) \rangle \rangle \langle \langle \cos n(\varphi_2 - \varphi_4) \rangle \rangle} = \frac{\langle v_n^2(p_T^a) v_n^2 \rangle}{\langle v_n^2(p_T^a) \rangle \langle v_n^2 \rangle}$

Angle: $F(\Psi_n^a, \Psi_n) = \frac{\langle \langle \cos[n(\varphi_1^a + \varphi_2^a - \varphi_3 - \varphi_4)] \rangle \rangle}{\langle \langle \cos[n(\varphi_1^a + \varphi_2 - \varphi_3^a - \varphi_4)] \rangle \rangle} \approx \langle \cos 2n[\Psi_n(p_T^a) - \Psi_n] \rangle$

- Both flow magnitude and angle fluctuations are $p_{\rm T}$ dependent (deviation from $\langle v_n^4 \rangle / \langle v_n^2 \rangle^2$ and unity, respectively)
- Deviation most significant in the most central collisions ($\sim 5\sigma$ for magnitude and $> 5\sigma$ for angle fluctuations) \Rightarrow discovery







Correlations between harmonics with ESE

- Event-shape engineering (ESE) selects specific initial eccentricity
- q_2 selection by constructing reduced flow vector
- Negative correlations between v_2 and v_3





Linear and non-linear flow modes

- Positive correlations between v₂ and v₄
- Extracting linear and non-linear responses by fitting v_4 with $v_4 = \sqrt{c_0^2 + (c_1 v_2^2)^2}$ (in each centrality class)
- Non-linear part (e₂ contribution) shows strong centrality dependence while linear part (e₄ contribution) stays almost constant with increasing centrality







Summary

- Flow coefficients $v_n\{m\}$ of different harmonics and different orders of multiparticle correlations were measured in heavy-ion collisions with ALICE
- Probability density function of v_2 was measured using multi-particle cumulants and has non-trivial transverse momentum dependence
- New observables allow to differ between the contributions from flow magnitude and flow angle fluctuations to the flow vector fluctuations
- Using event-shape engineering, it is possible to select subsets of events with different shapes of *p.d.f.*, obtain important information on correlations between different flow harmonics, and also extract linear and non-linear flow modes

Thank you for you attention!



BACK UP

Multi-particle correlations

- Two particles with φ_1, η_1 and φ_2, η_2
- Two-particle correlation: $\langle \langle 2 \rangle \rangle = \langle \langle \cos n(\varphi_1 \varphi_2) \rangle \rangle$
- Four-particle **correlation**: $\langle \langle 4 \rangle \rangle = \langle \langle \cos n(\varphi_1 + \varphi_2 - \varphi_3 - \varphi_4) \rangle \rangle$

- Two-particle **cumulant**: $c_n\{2\} = \langle \langle 2 \rangle \rangle$
- Four-particle **cumulant**: $c_n\{4\} = \langle \langle 4 \rangle \rangle 2 \langle \langle 2 \rangle \rangle^2$
 - $\langle \langle 4 \rangle \rangle$ contains (lower order) correlations between $\varphi_1 \varphi_3$, $\varphi_2 \varphi_4$ and $\varphi_1 \varphi_4$, $\varphi_2 \varphi_3$ that have to be subtracted



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ALICE Pb-Pb $0.2 < p_T < 3 \text{ GeV/c}$ $h_I < 0.8$ $2.76 \quad 5.02 \text{ TeV}$ $\circ \quad \circ c_2 \text{(m)}$ $0 \quad 10 \quad 20 \quad 30 \quad 40 \quad 50 \quad 60 \quad 70$ Centrality (%) ALICE, JHEP 1807 (2018) 103

ALICE, JHEP 1807 (2018) 103

3₂{2}



Multi-particle cumulants

•
$$c_n\{2\} = \langle \langle 2 \rangle \rangle_{n,-n}$$

•
$$c_n\{4\} = \langle \langle 4 \rangle \rangle_{n,n,-n,-n} - 2 \langle \langle 2 \rangle \rangle_{n,-n}^2$$

•
$$c_n\{6\} = \langle \langle 6 \rangle \rangle_{n,n,n,-n,-n,-n} - 9 \langle \langle 4 \rangle \rangle_{n,n,-n,-n} \langle \langle 2 \rangle \rangle_{n,-n} + 12 \langle \langle 2 \rangle \rangle_{n,-n}^3$$

•
$$c_n\{8\} = \langle \langle 8 \rangle \rangle_{n,n,n,n,n,-n,-n,-n} - 16 \langle \langle 6 \rangle \rangle_{n,n,n,-n,-n,-n} \langle \langle 2 \rangle \rangle_{n,-n} - 18 \langle \langle 4 \rangle \rangle_{n,n,-n,-n}^2 + 144 \langle \langle 4 \rangle \rangle_{n,n,-n,-n} \langle \langle 2 \rangle \rangle_{n,-n}^2 - 144 \langle \langle 2 \rangle \rangle_{n,-n}^4$$



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Multi-particle cumulants

- Flow coefficients:
 - $v_n\{2\} = \sqrt{c_n\{2\}}$
 - $v_n\{4\} = \sqrt[4]{-c_n\{4\}}$

•
$$v_n\{6\} = \sqrt[6]{\frac{1}{4}c_n\{6\}}$$

•
$$v_n\{8\} = \sqrt[8]{-\frac{1}{33}c_n\{8\}}$$

- Difference (at the same energy) caused by non-flow and fluctuations
- Non-flow: flow-like correlations caused by jets and decays of resonances



Linear and non-linear flow modes

- No correlations between linear and non-linear components
- Linear part of v₄ stays unchanged with changing the event shape
- Significant increase of non-linear part of v₄ when modifying the event shape



