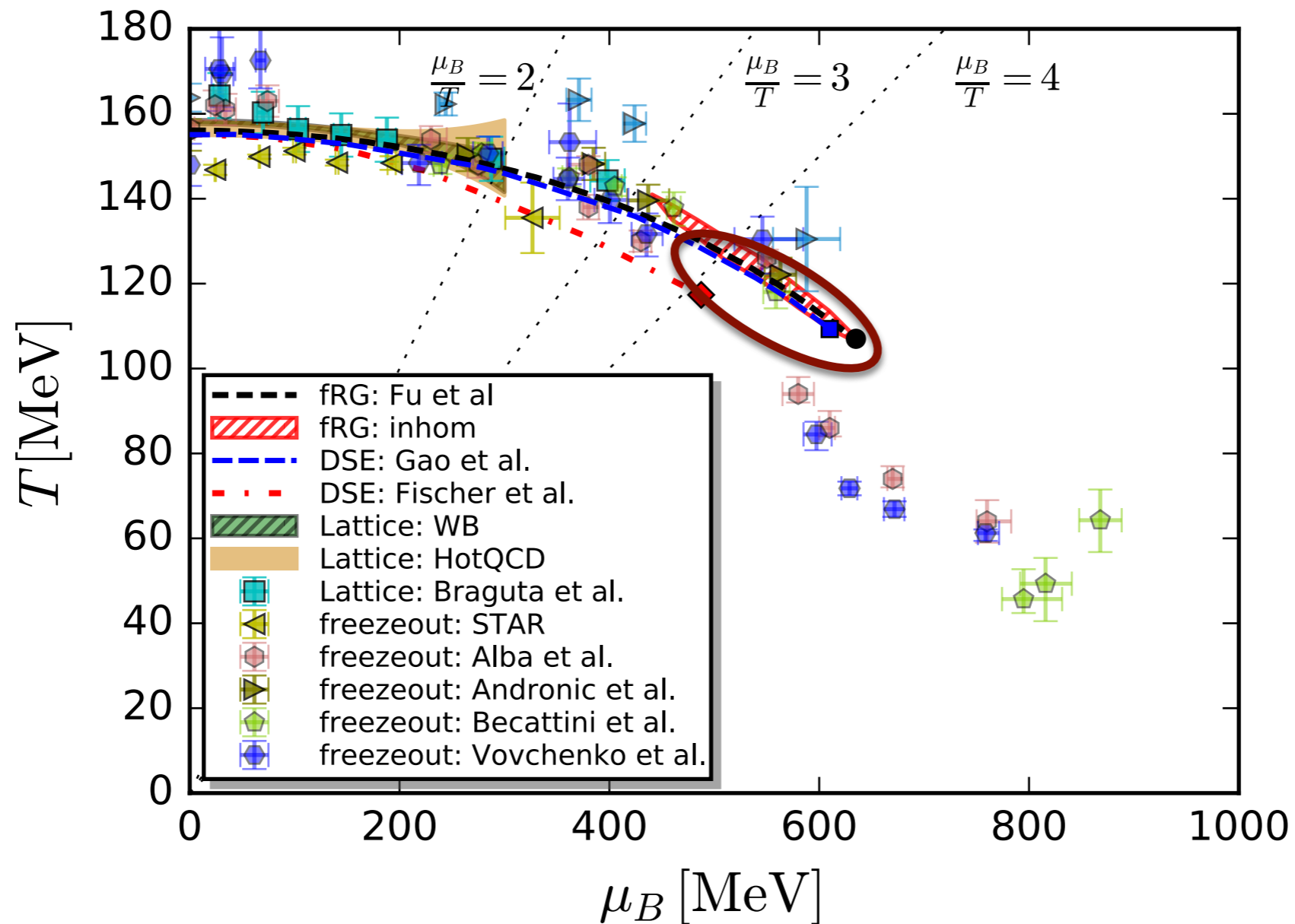


QCD phase structure



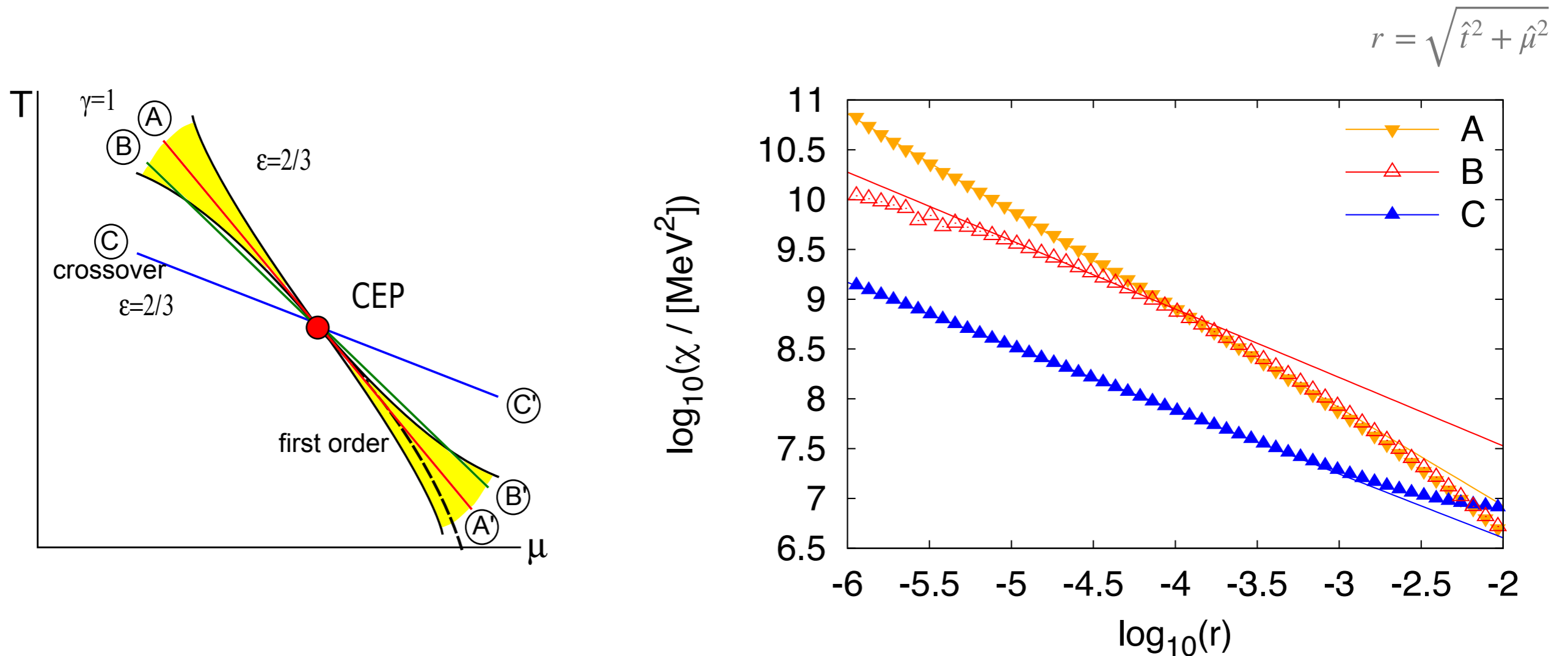
$$(135, 450) \text{ MeV} \lesssim (T_{\text{CEP}}, \mu_{B, \text{CEP}}) \lesssim (100, 650) \text{ MeV}$$

Suggests: no criticality for $\frac{\mu_B}{T_c} \lesssim 4$

fRG: [Fu, JMP, Rennecke, PRD 101, \(2020\) 054032](#)

DSE: [Fischer, PPNP 105 \(2019\) 1, Gao, JMP, arXiv:2010.137005](#)

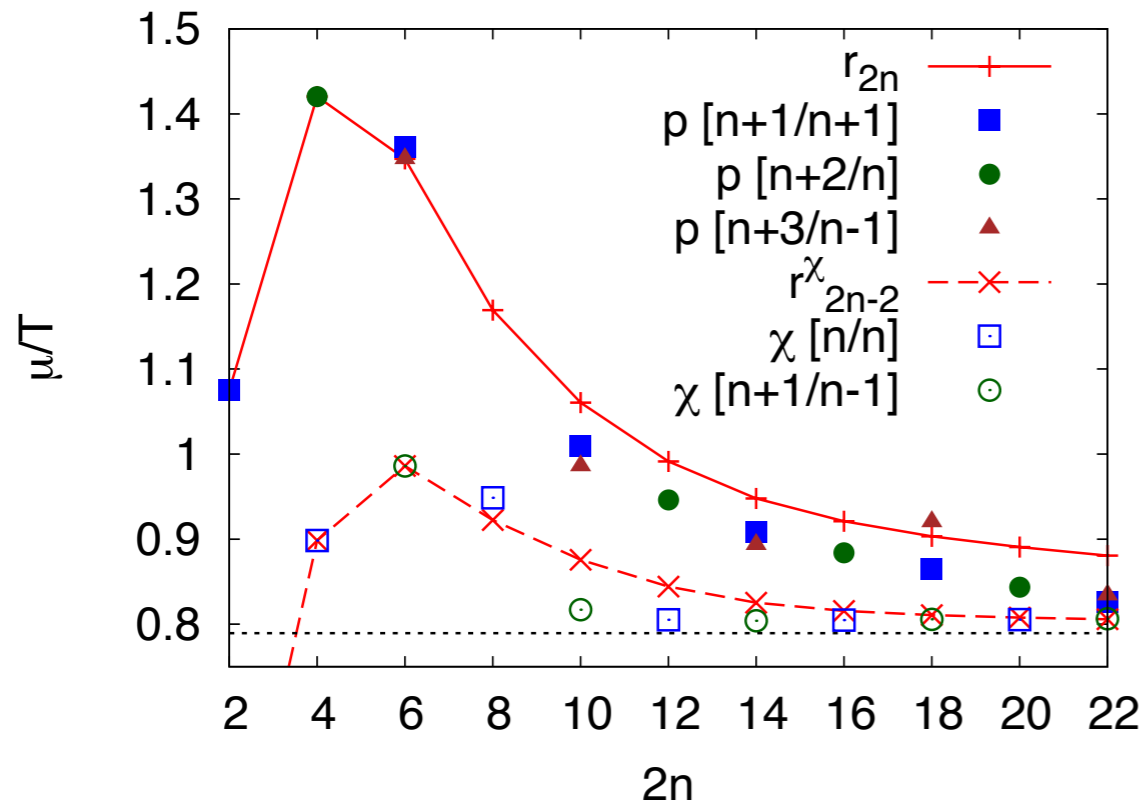
SIZE OF THE CRITICAL REGION



- mean-field PQM: critical region $\lesssim 10 \text{ MeV}$ (tangential to transition line (A)) [Schaefer, Wagner (2012)]
- including quantum fluctuations: critical region shrinks to $< 1 \text{ MeV}$ in T and μ_B direction

[Schaefer, Wambach (2006)]
[Chen, Wen, Fu (2021)]

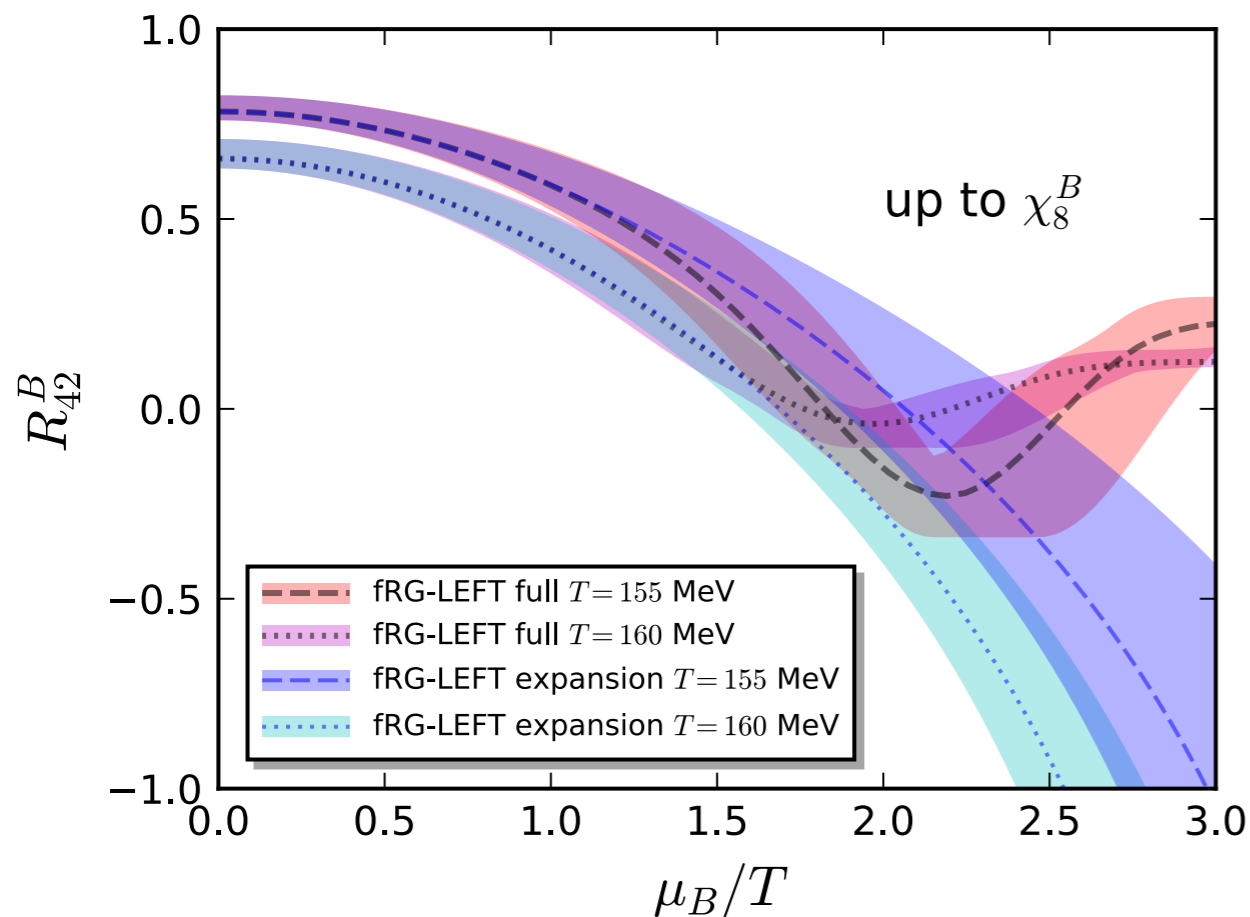
VALIDITY OF THE TAYLOR EXPANSION



- estimate of location of 2nd order phase transition based on Taylor expansion about $\mu = 0$ for different expansion orders n

very high orders necessary for good estimates

[Karsch, Schaefer, Wagner, Wambach (2011)]



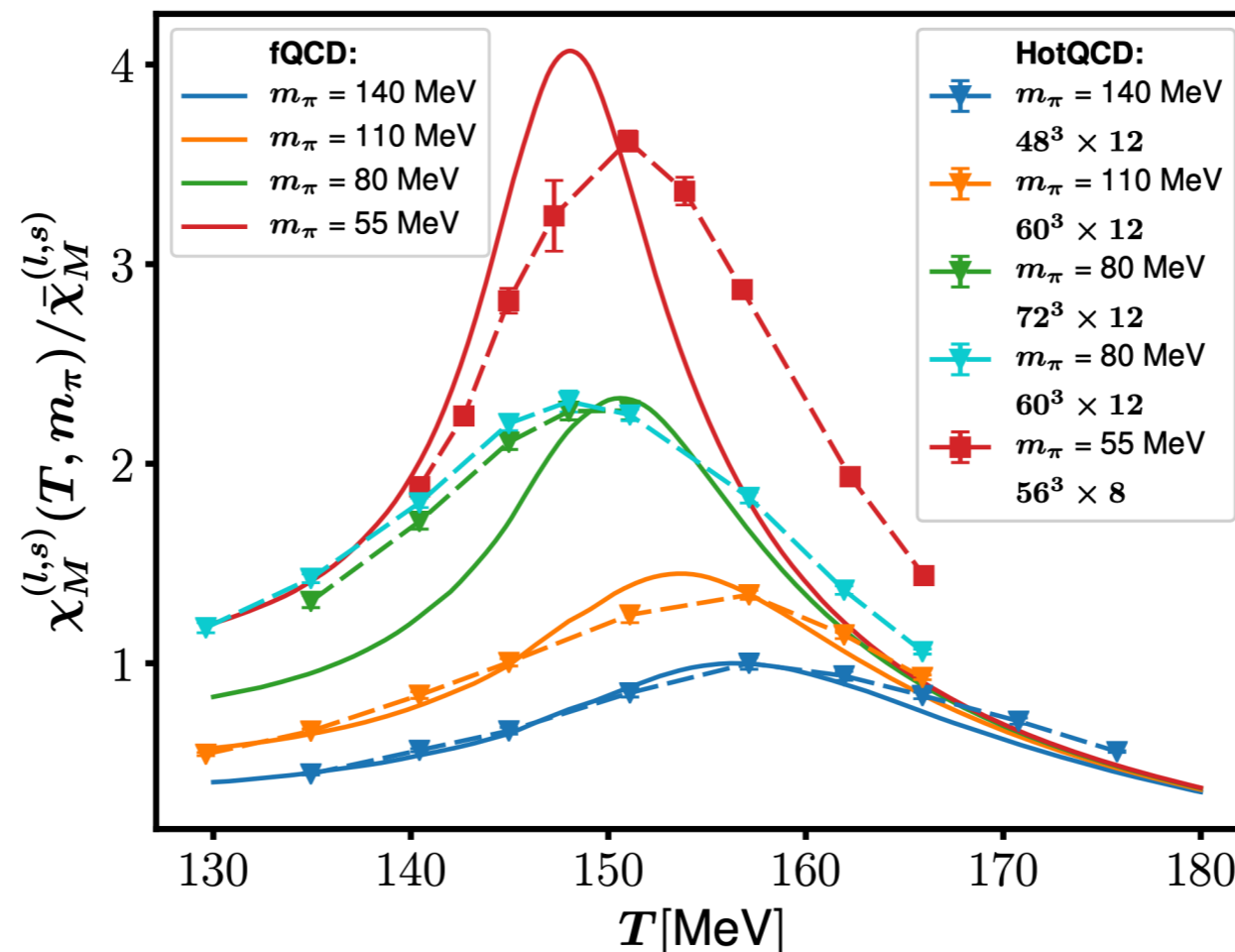
- Taylor expansion about $\mu_B = 0$ vs direct computation

Taylor expansion fails to capture qualitative features at finite μ_B

[Fu, Luo, Pawłowski, Rennecke, Wen (2021)]

Scaling analysis from functional QCD

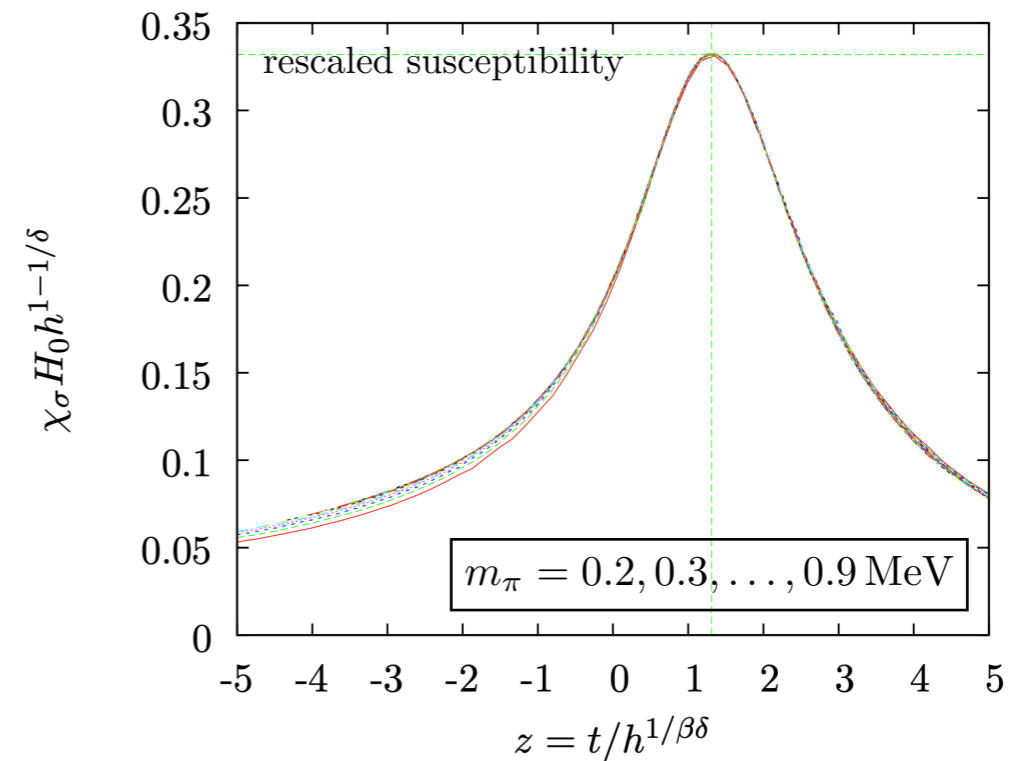
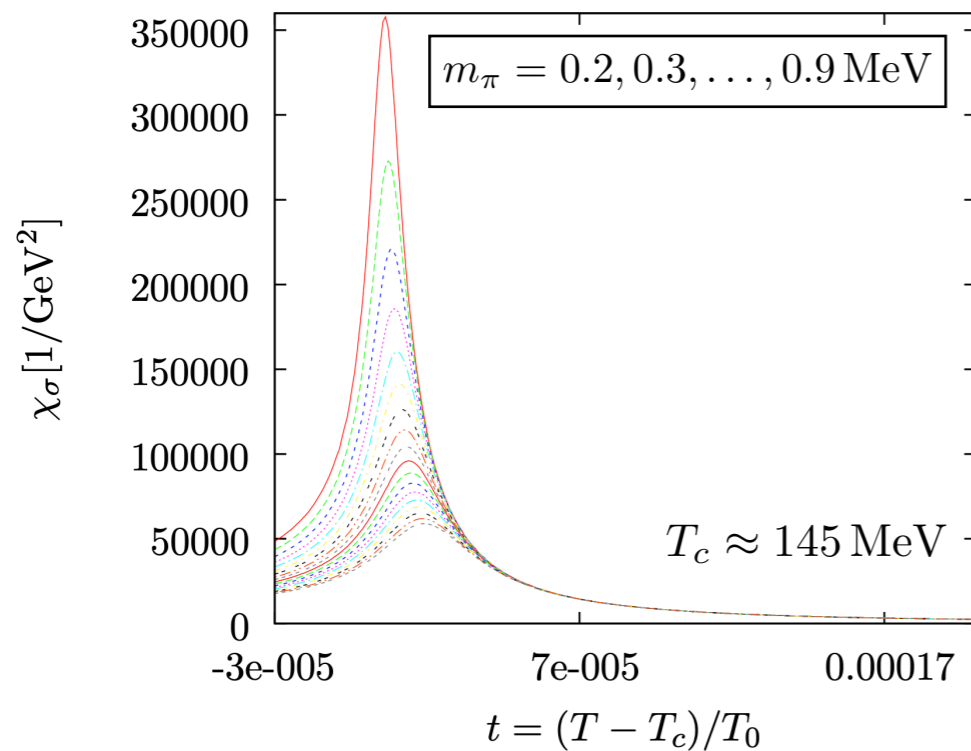
Braun, Fu, Pawłowski, Rennecke, Rosenblüh, Yin, PRD 102 (2020) 056010



- Very good agreement with lattice QCD results from the hotQCD collaboration around physical pion masses and above
- Susceptibilities do not show indications for scaling for $m_\pi \gtrsim 30$ MeV

Size of the scaling region: quark-meson model

Braun, Klein, Piasecki, EPJ C 71 (2011) 1576

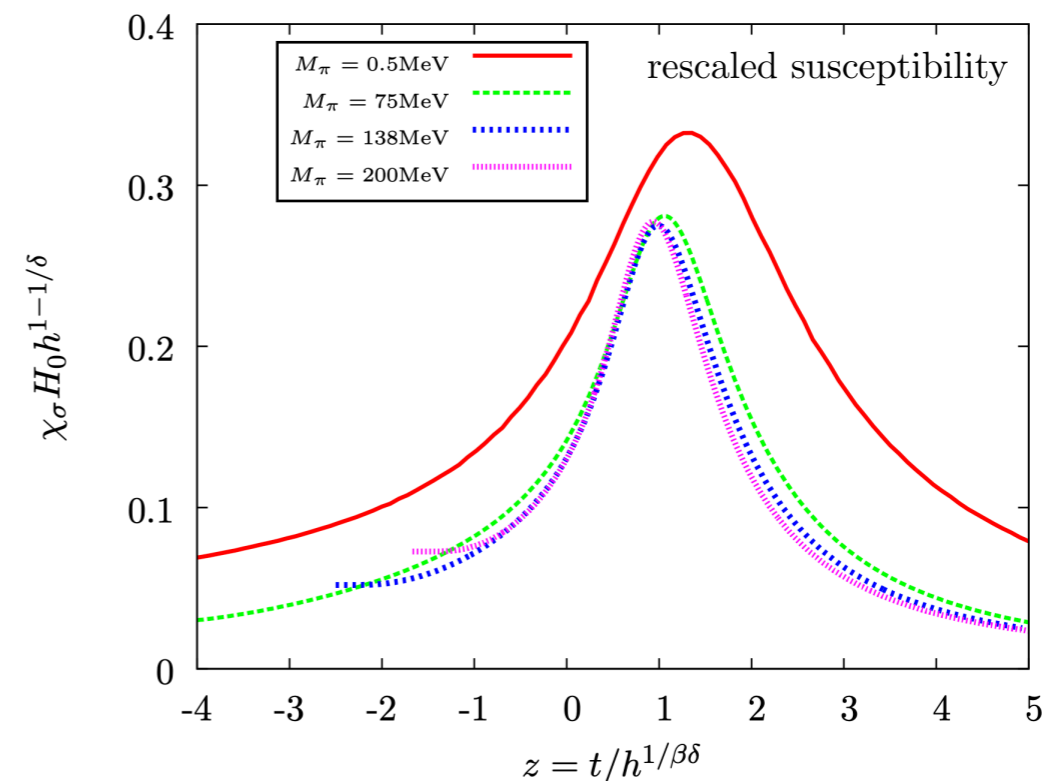


- Size of the true scaling region:

$$m_\pi \lesssim 1 \text{ MeV}$$

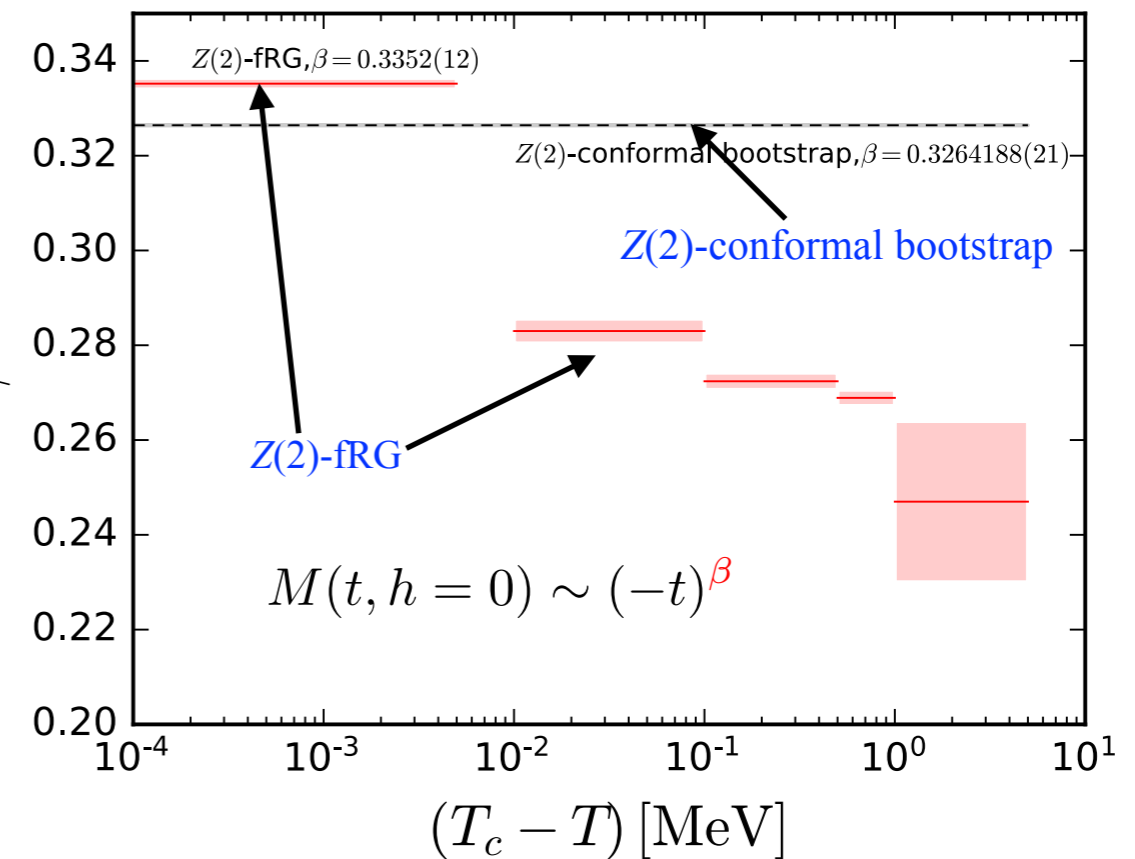
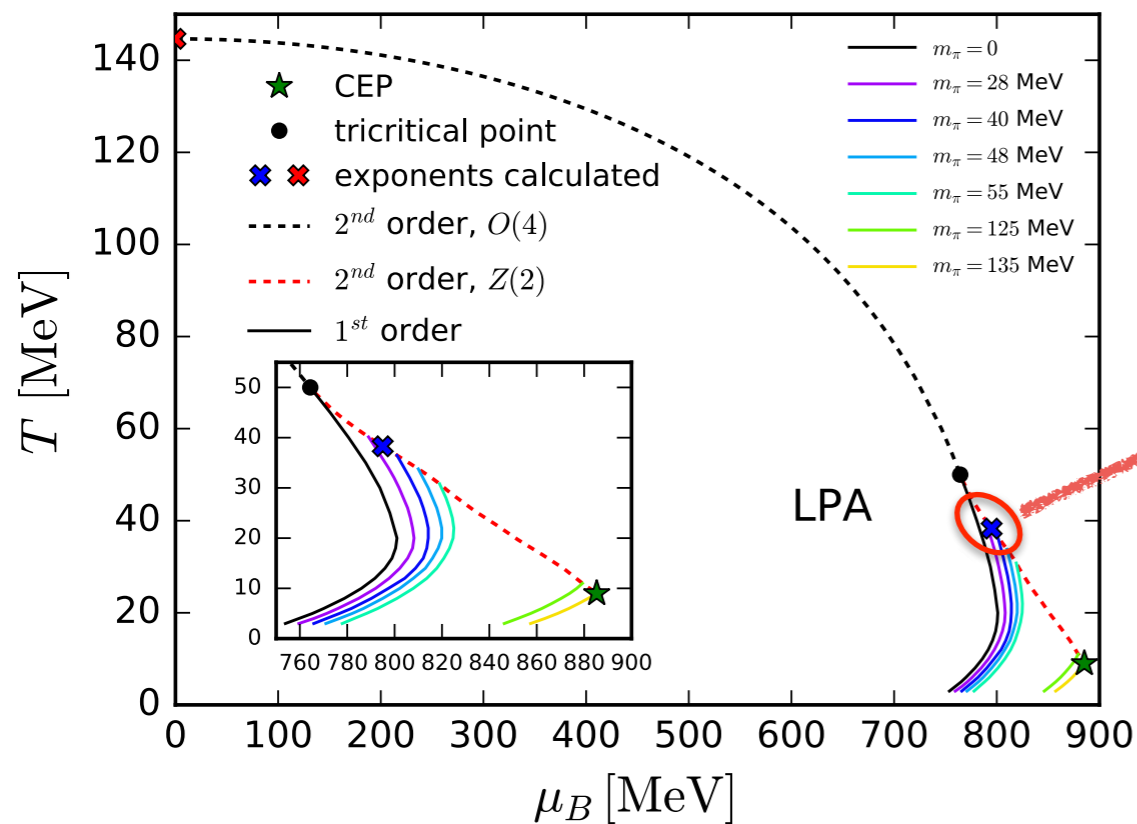
- “Seeming” scaling region:

$$75 \text{ MeV} \lesssim m_\pi \lesssim 200 \text{ MeV}$$



Critical region near the CEP with Z(2) symmetry

Chen, Wen, Fu, arXiv: 2101.08484



- Critical exponent β is close to the CEP
- Critical region (in temperature direction) smaller than **1 MeV!**

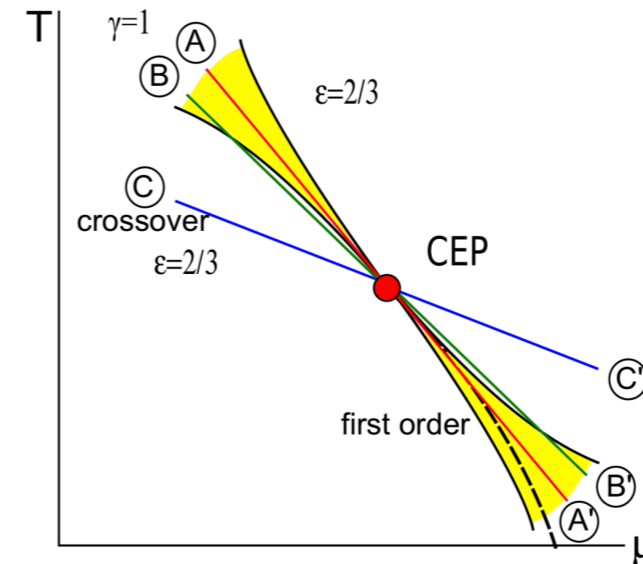
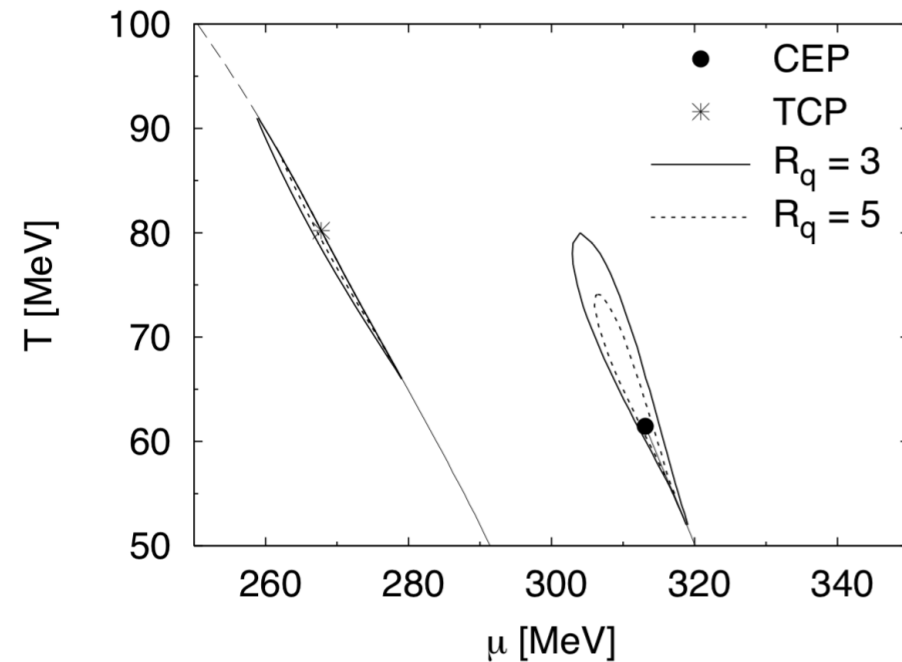
See also: [Schaefer, Wambach, PRD 75 \(2007\) 085015](#)

[Schaefer, Wagner, PRD 85 \(2012\) 034027](#)

Critical region from RG flows

[B.-J. Schaefer, J. Wambach, Phys. Rev. D 75 (2007)]

[B.-J. Schaefer, M. Wagner, Phys. Rev. D (2012)]



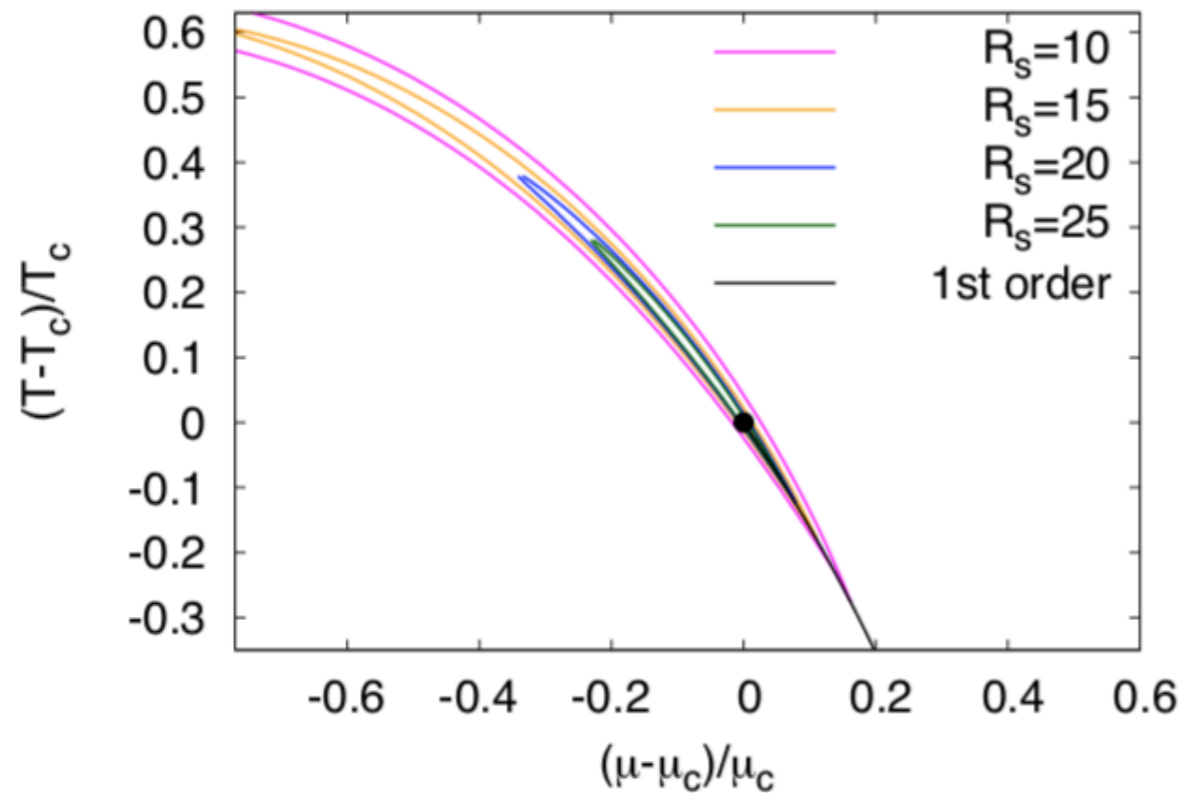
- Contour regions for quark number susceptibility
(chiral limit: TCP physical pion mass: CEP)

- Critical region around CEP
 → criticality depends on the path towards CEP
 → crossover between different universality classes

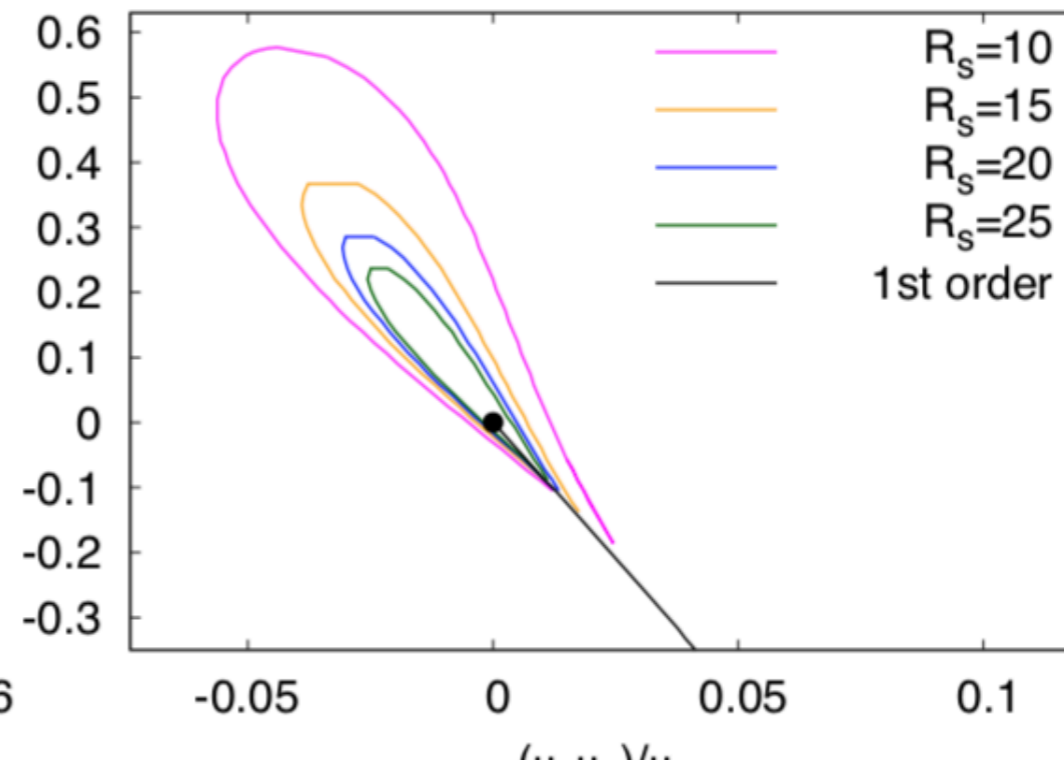
Size of the critical region: quark-meson model

[B.-J. Schaefer, J. Wambach, Phys. Rev. D 75 (2007)]

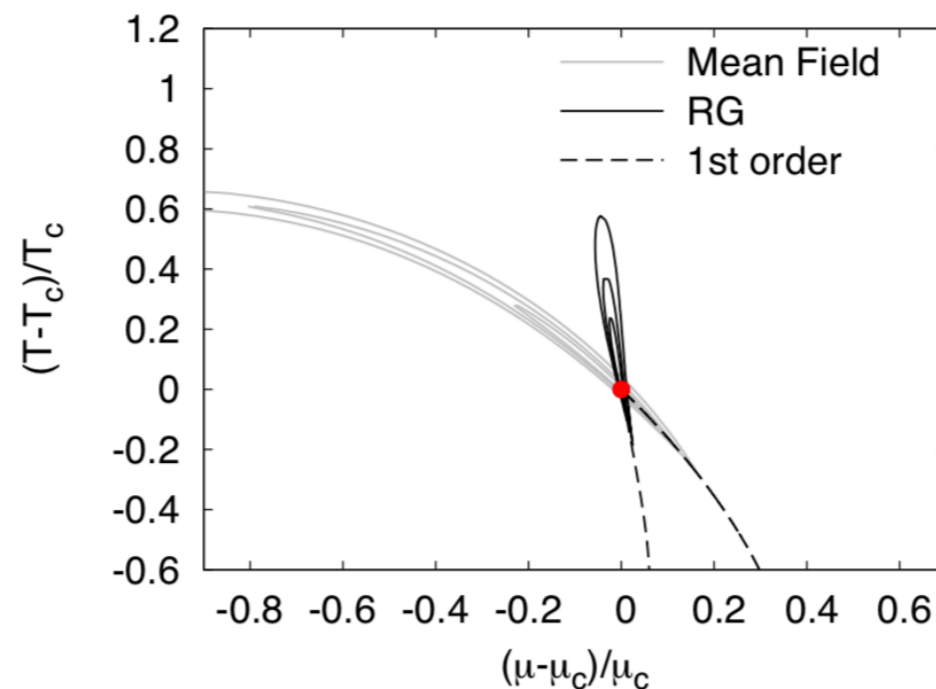
Mean-field



FRG

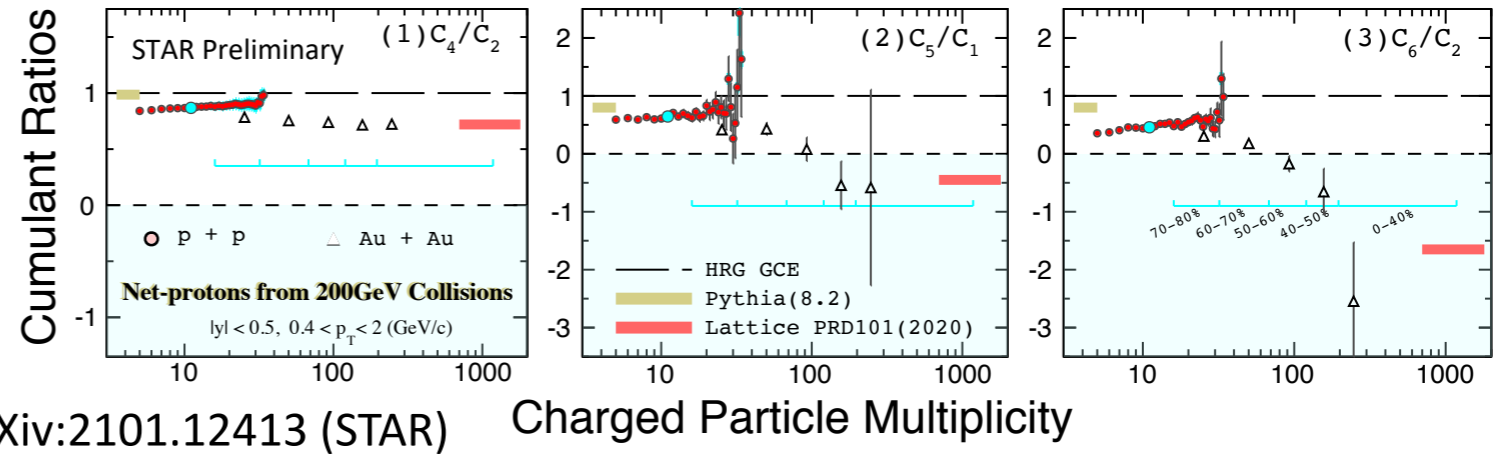


- Size much more compressed with FRG

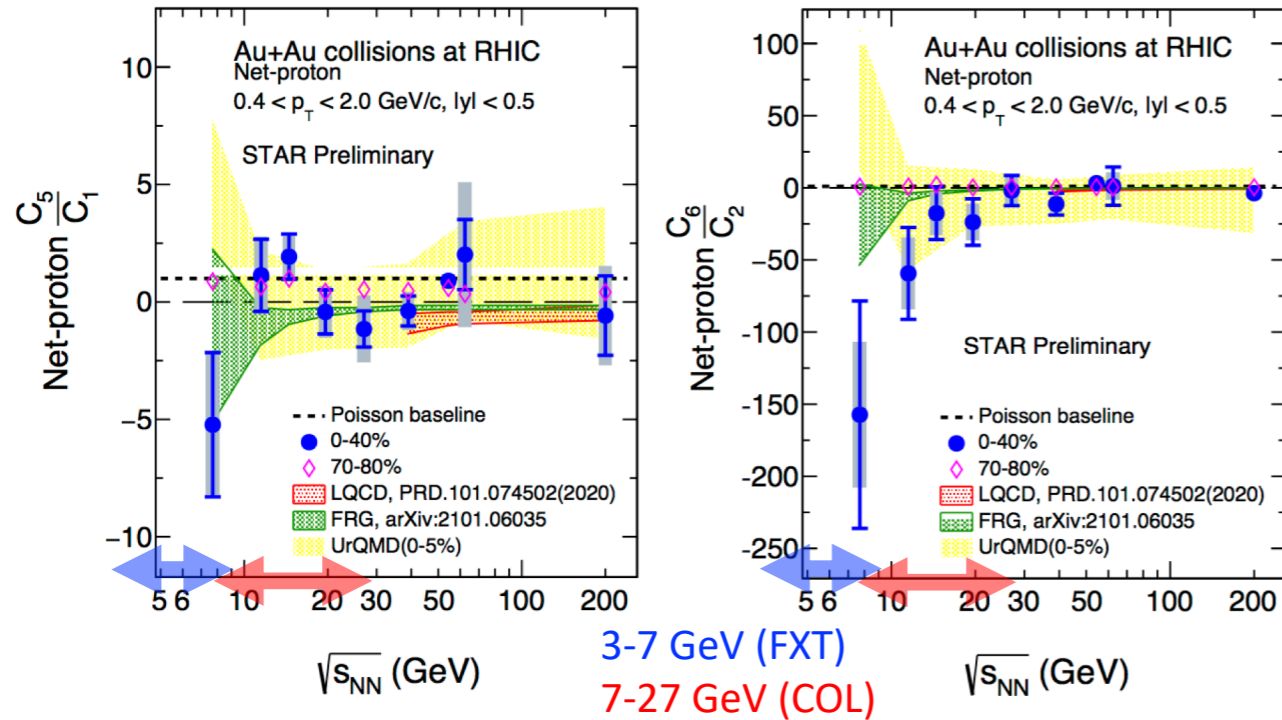


The 5th, 6th order fluctuation of net-proton including pp collisions

New data of net-p cumulants at 200 GeV pp collisions will be shown by Risa N.

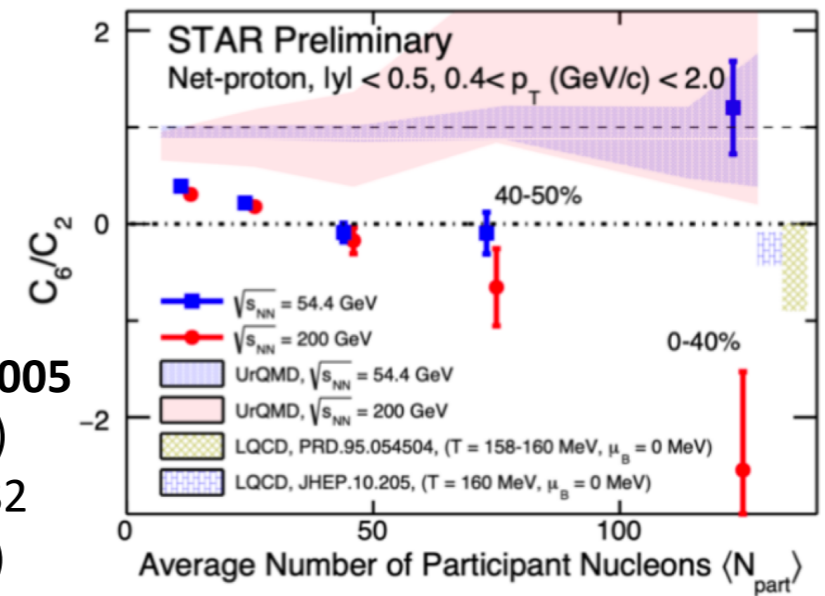


New data of net-proton C_5, C_6 at BES-I will be shown by Ashish P.

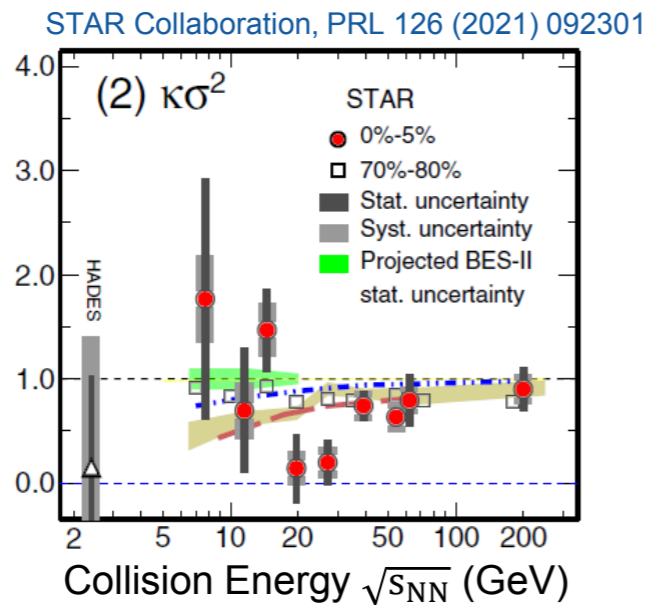


negative c_6 could be taken as an indication of cross-over transition at small μ_B

NPA 1005 (2021) 121882 (STAR)



Critical Fluctuations



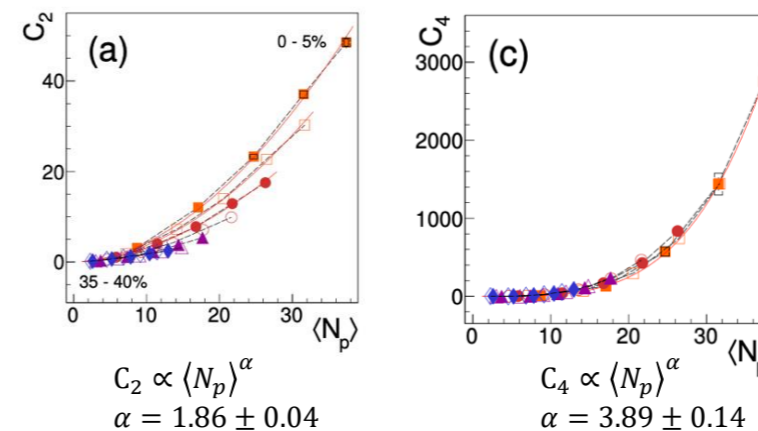
Ling, Stephanov, PRC 93, 034915 (2016)

Cumulants k_n hold information on multi-particle correlators C_n

Bzdak, Koch, Strodthoff, PRC 95, 054906 (2017)

Investigate C_n vs. $\langle N_p \rangle$ to isolate relevant physics, $C_n \propto \langle N_p \rangle^\alpha$

HADES Collaboration, PRC 102 (2020) 2, 024914
<https://www.hepdata.net/record/ins1781493>



$\alpha \approx n \rightarrow$ signature of multi-particle correlations ($\Delta y_{\text{corr}} > 1$)

Direct link to EoS

$$\frac{1}{VT^3} k_n = \frac{\partial^n \hat{p}}{\partial \hat{\mu}^n}$$

$\hat{p} = \frac{p}{T^4}$ reduced pressure
 $\hat{\mu} = \frac{\mu}{T}$ reduced chemical potential

cf. B. Friman *et al.*, EPJC 71 (2011) 1694
 M. Stephanov, PRL107 (2011) 052301

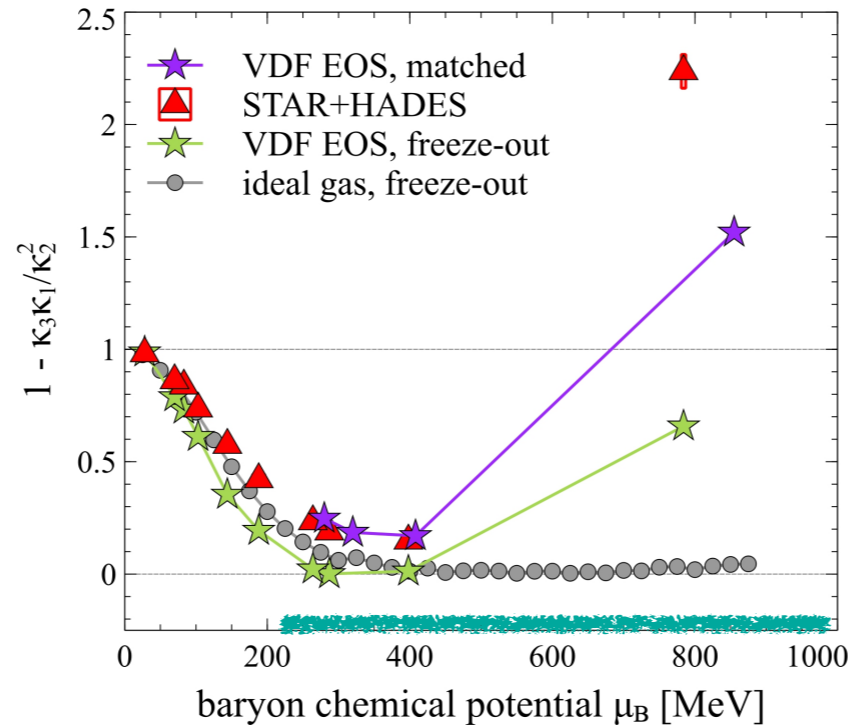
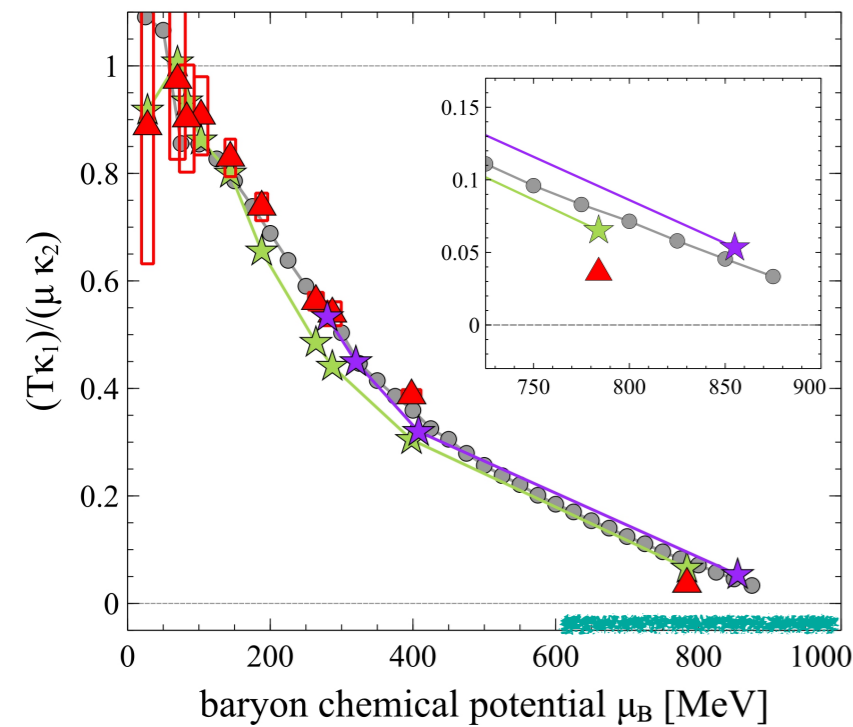
Further progress with data:

- Applying identity method A. Rustamov, M.I. Gorenstein, PRC 86 (2012) 044906
- Analyzing new Ag+Ag data

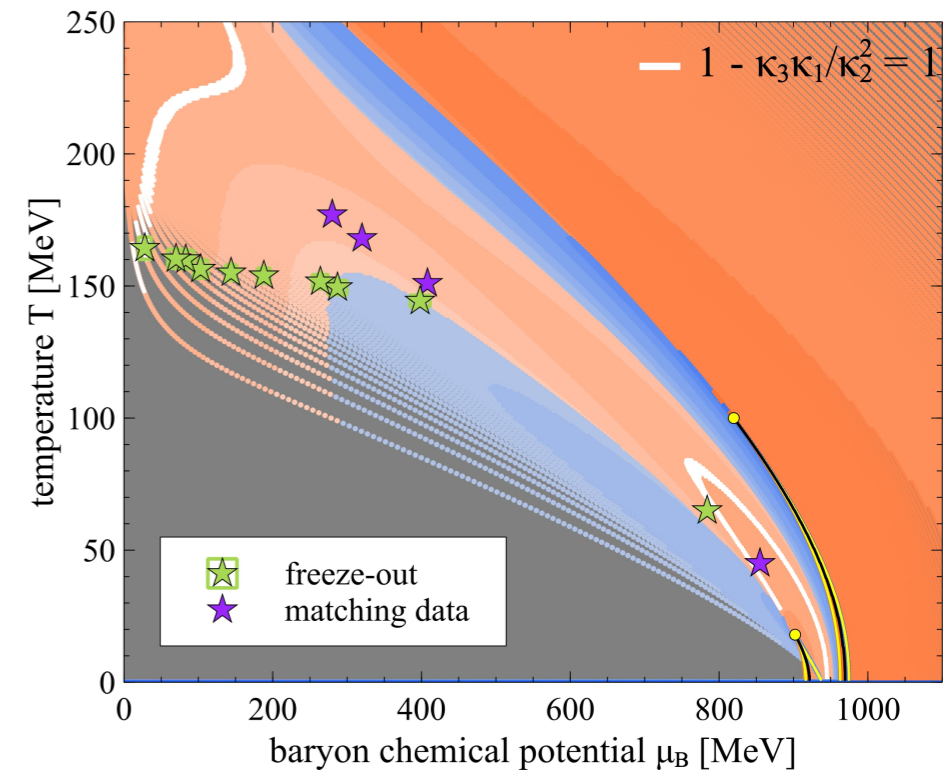


Questions

STAR + HADES data



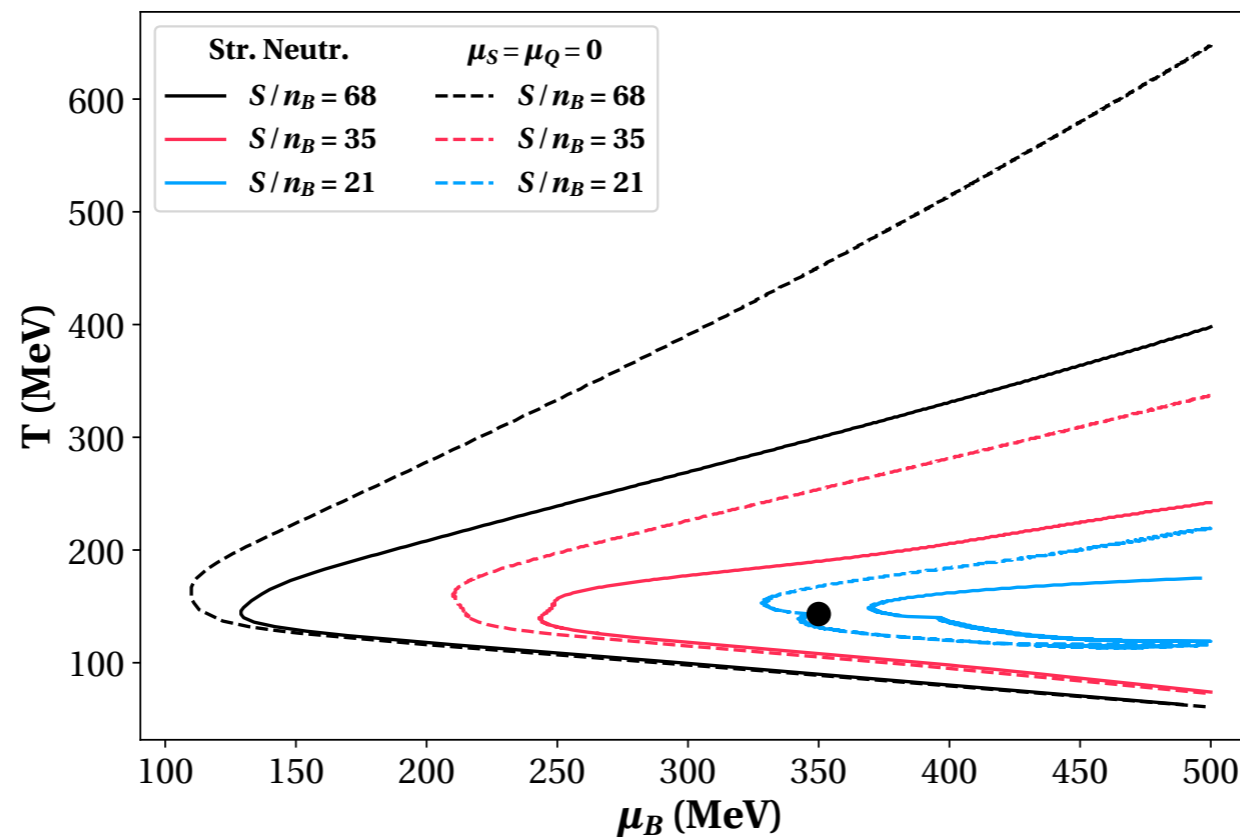
VDF model



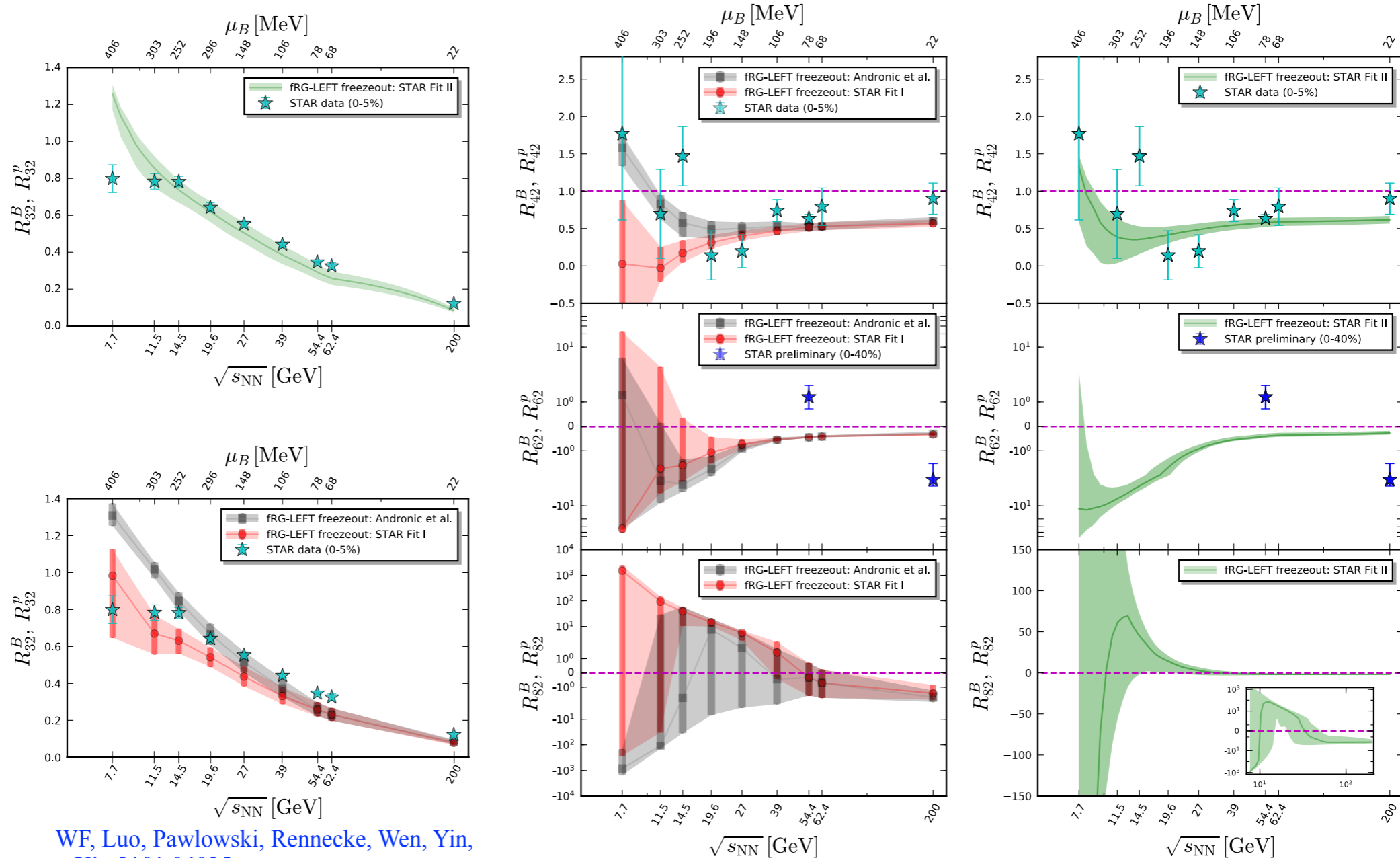
- what happens in the STAR BES fixed target region?
- is behavior of the cumulants at low energies dominated by hadronic effects and the nuclear liquid-gas phase transition?
- can we study c_T^2 below $\sqrt{s} = 2.4$ GeV?

Isentropic Trajectories

- ▶ Isentropes show the path of the HIC system through the phase diagram in the absence of dissipation
 - ▶ Different path when conserved charge conditions applied



Fluctuations on the freeze-out curve



WF, Luo, Pawlowski, Rennecke, Wen, Yin, arXiv:2101.06035

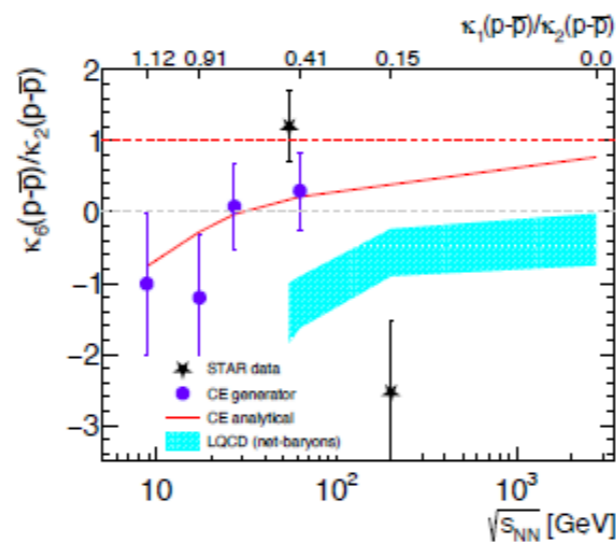
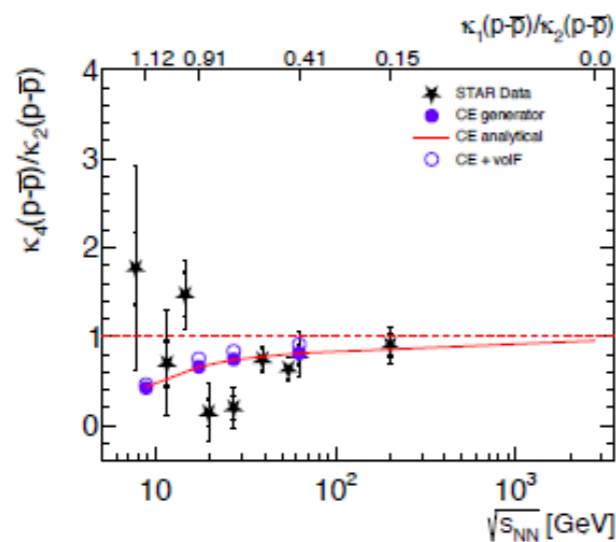
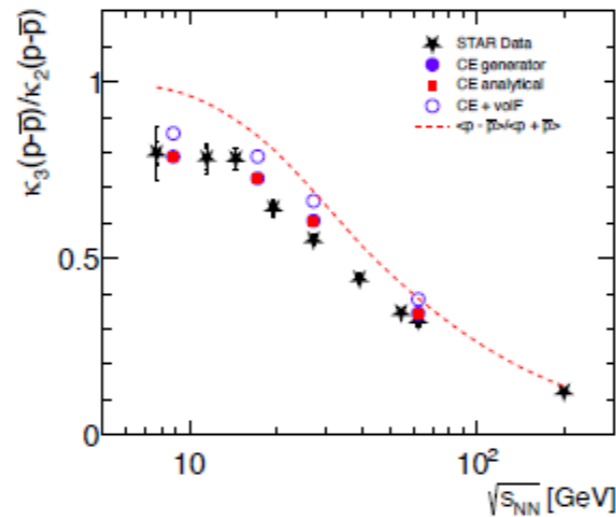
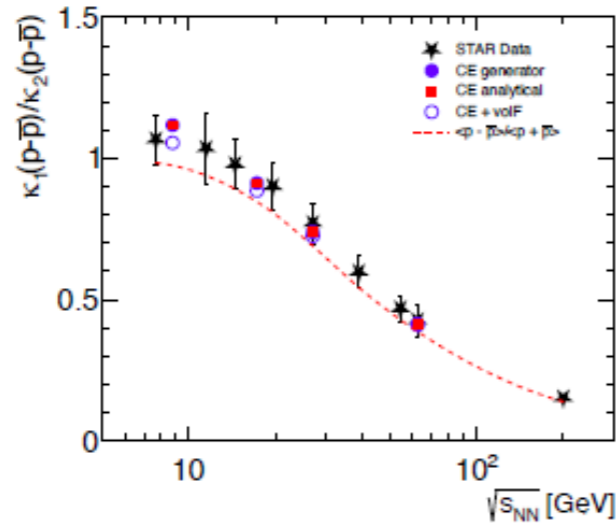
J. Adam *et al.* (STAR), *PRL* 126 (2021), 092301
M. Abdallah *et al.* (STAR), arXiv:2101.12413

T. Nonaka (STAR), *NPA* 1005 (2021) 121882; A. Pandav (STAR), arXiv: *NPA* 1005 (2021) 121936

Ratios of cumulants: STAR data versus model

Stars: STAR data

Adam et al. arXiv 2001.02852v2



- Broken lines calculated from

Skellam distribution

$$\frac{\kappa_1}{\kappa_2} = \frac{\kappa_3}{\kappa_2} = \frac{\langle p - \bar{p} \rangle}{\langle p + \bar{p} \rangle}$$

- Open circle include MC volume fluctuations

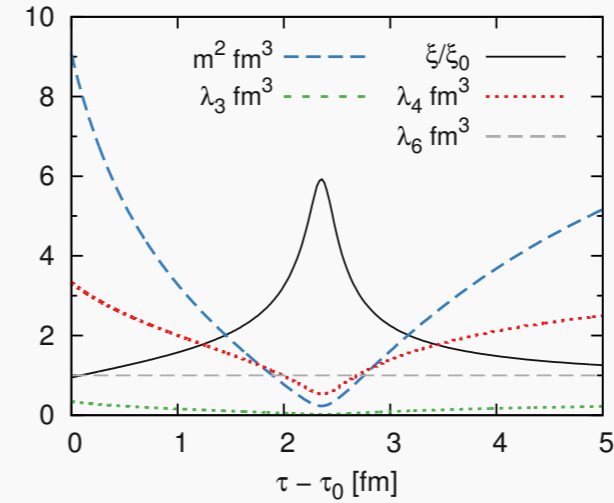
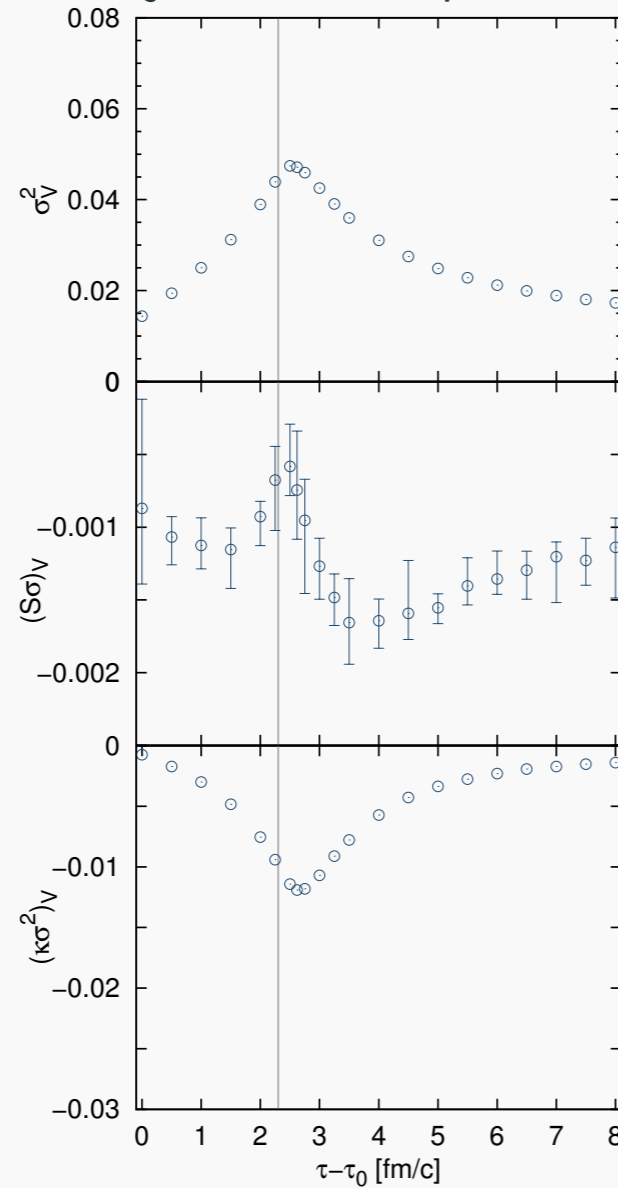
P. Braun-Munzinger, A. Rustamov and J. Stachel, Nucl. Phys. A960, 114 (2017).

V. Skokov, B. Friman and K. Redlich, Phys. Rev. C88, 034911 (2013)

- Cumulants up to $n < 4$ order follow the STAR data
- Kurtosis data exhibit interesting deviations, however *not necessarily* of statistical significant

Dynamics: time evolution of critical fluctuations

For a Bjorken-like temperature evolution:



- shift of extrema for variance/kurtosis (retardation effects) to later times corresponding to $T(\tau) < T_c$
- |extremal values| in dyn simulations < equilibrium values (nonequilibrium effects):

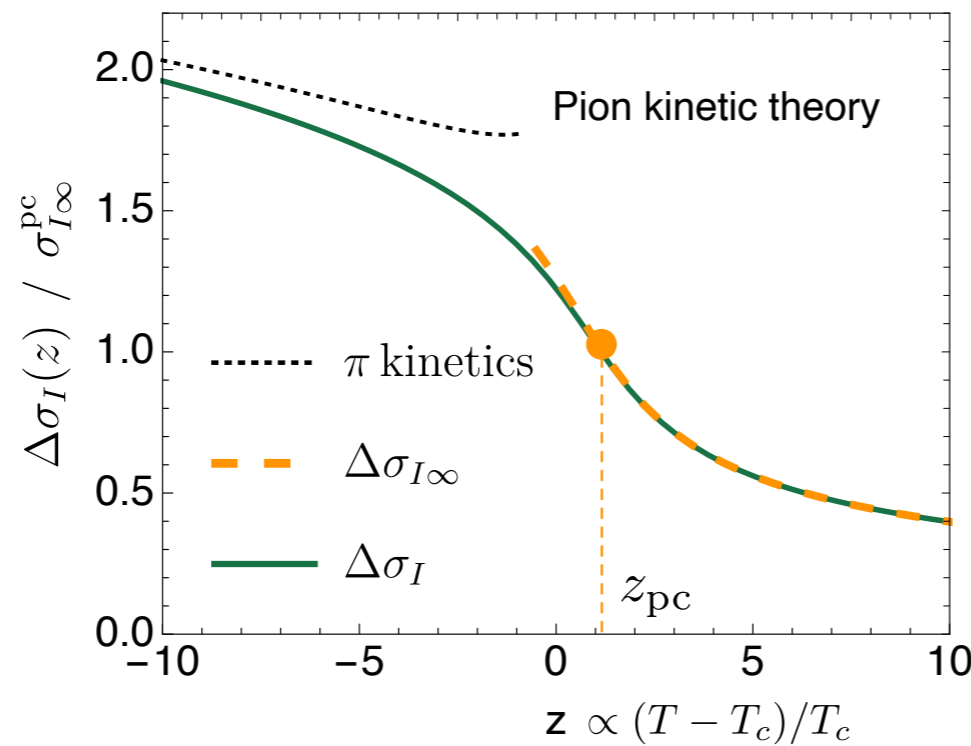
$$(\sigma_V^2)_{\text{dyn}}^{\text{max}} \approx 0.75 (\sigma_V^2)_{\text{eq}}^{\text{max}}$$

$$((\kappa\sigma^2)_V)_{\text{dyn}}^{\text{min}} \approx 0.5 (\kappa\sigma_V^2)_{\text{eq}}^{\text{min}}$$

- expected behavior with varying D and c_s^2 (expansion rate)

The conductivity through T_c

$$\sigma_I = \sigma_{I\text{reg}} + \underbrace{\Delta\sigma_I}_{\text{critical part}}$$

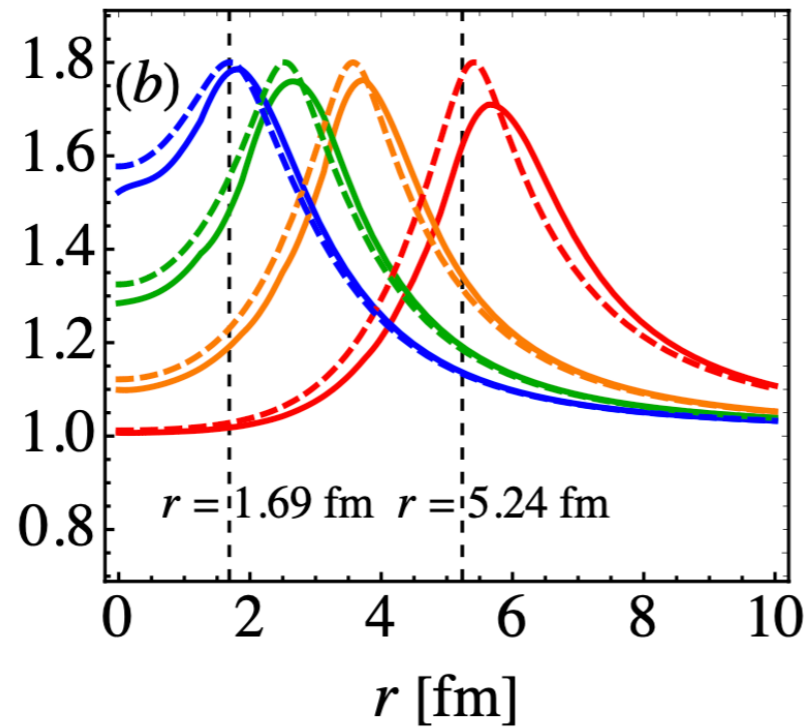


$$\sigma_{I\infty}^{\text{pc}} \equiv \left. \frac{T}{16\pi D_m m} \right|_{z=z_{\text{pc}}}$$

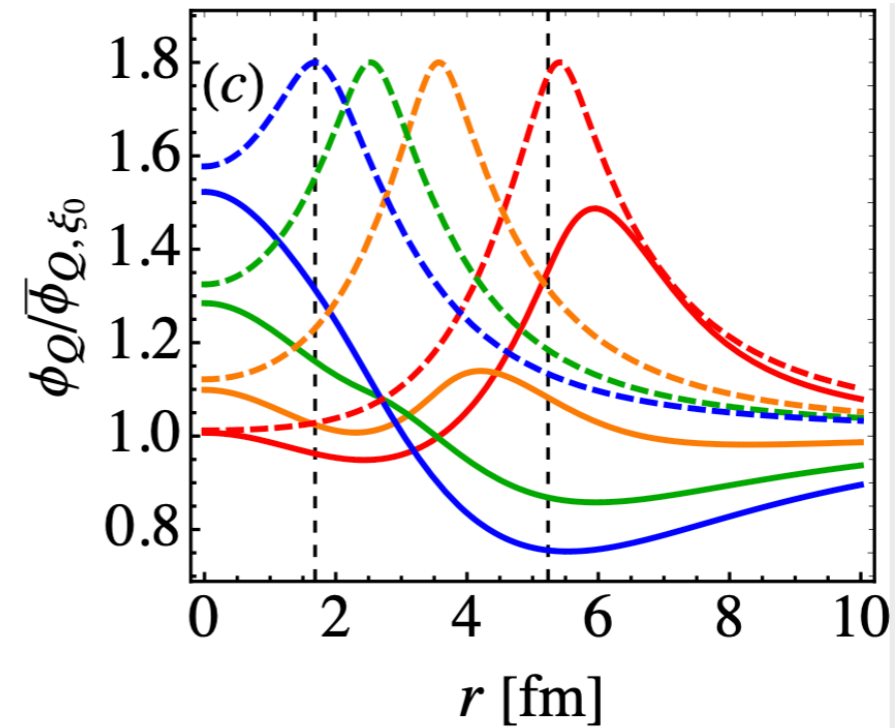
Estimate of the absolute magnitude for $\sigma_{I\infty}^{\text{pc}}$ with $T_{\text{pc}} \simeq 155 \text{ MeV}$

$$\Delta D_I = \left(\frac{\Delta\sigma_I}{\chi} \right) = \frac{0.50}{2\pi T} \times \left(\frac{1.3}{m_{\text{pc}}/T} \right) \left(\frac{0.4}{\chi_Q/T^2} \right) \left(\frac{3.0}{2\pi T D_m} \right)$$

Critical slowing down and advection

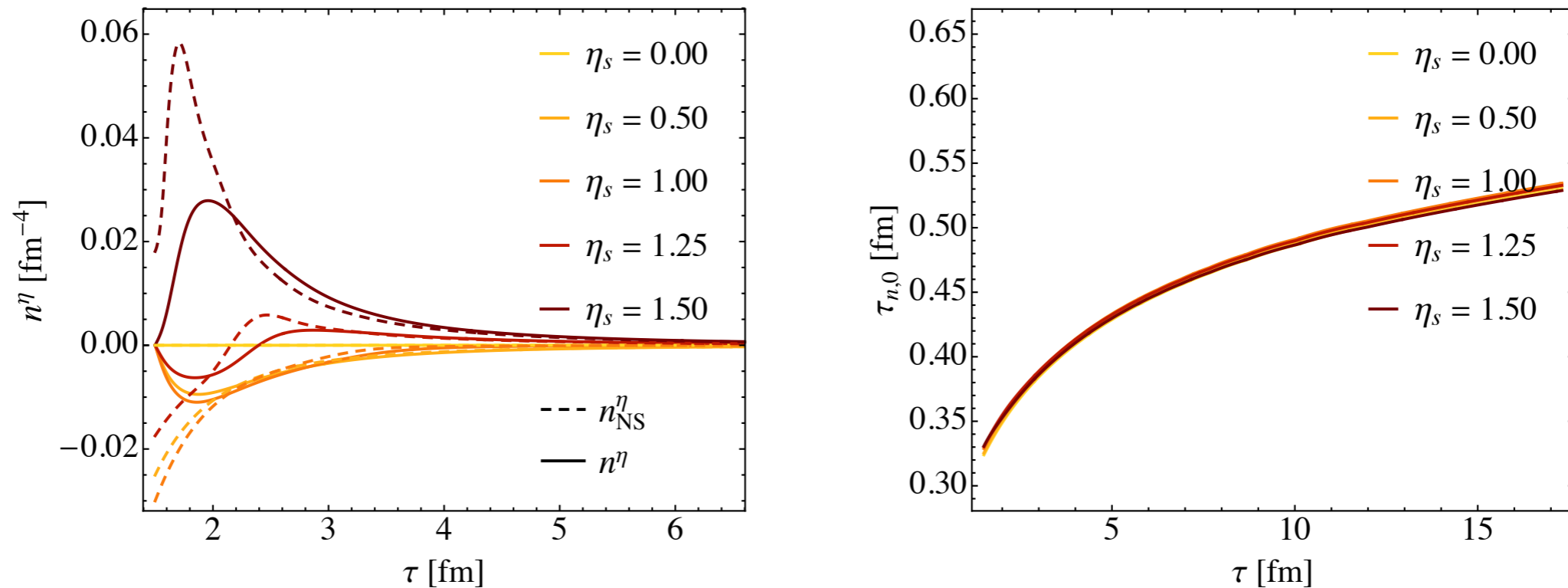


Turning off radial flow



Turning on radial flow

Time evolution of baryon diffusion current



- With the fireball cooling down, the driving force of diffusion current ($\kappa_n \nabla(\mu/T)$) decreases:
 - Two reasons: (a) gradient $\nabla(\mu/T)$ gets smoothed; (b) κ_n decreases.
- Response to the driving force also gets slower, because of the growing relaxation time;
- Critical slowing-down ($\tau_n \simeq 6 \tau_{n,0}$) would help n^η to stay at (almost) zero, even if $\kappa_n \nabla(\mu/T)$ got affected by the critical point.

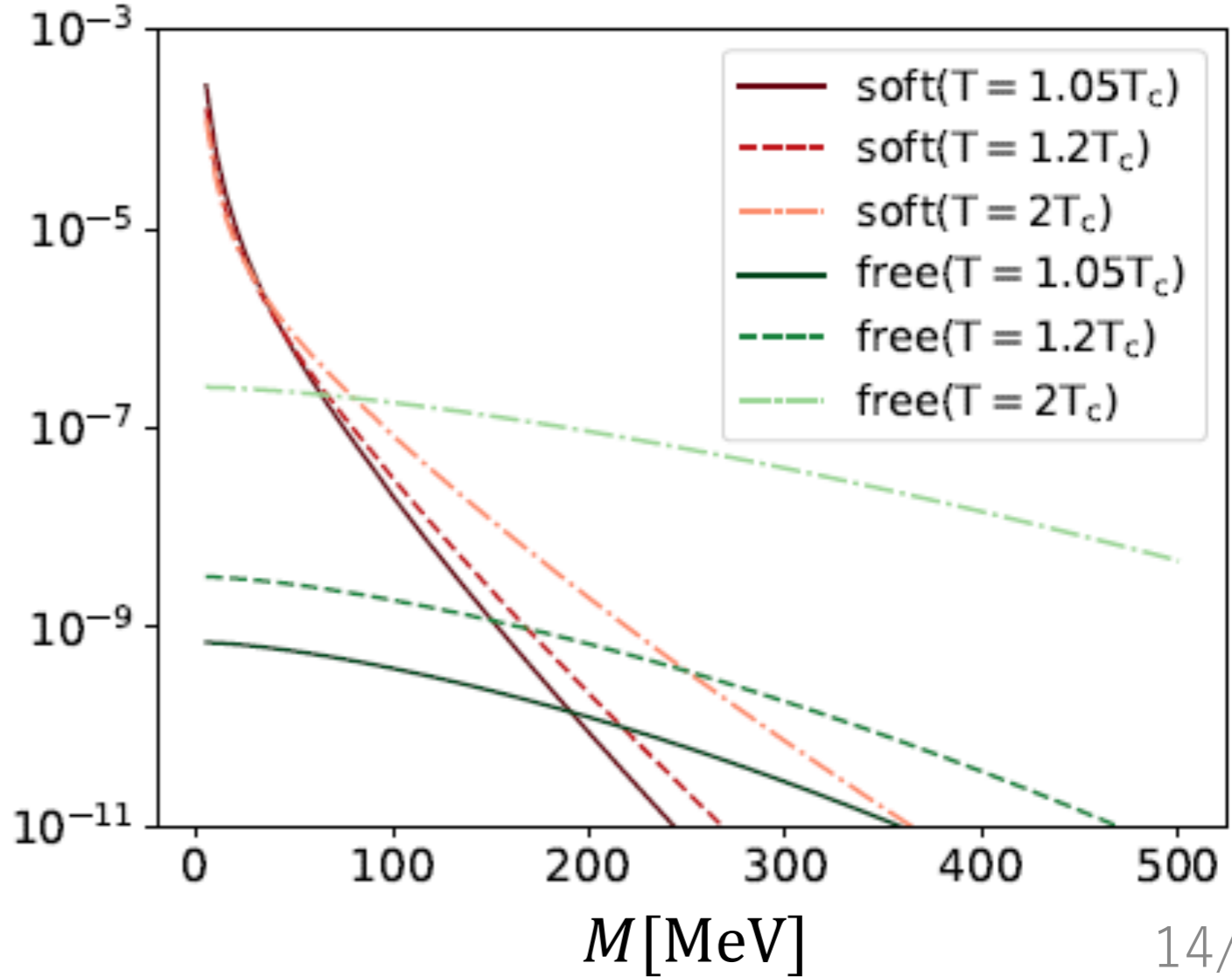
■ Invariant mass spectrum $\frac{d\Gamma}{dM^2} = \frac{\alpha}{6\pi^3 M^2} \int dk \frac{k^2}{\omega} \frac{\text{Im}\Pi_{\mu}^{\mu R}(\mathbf{k}, \omega)}{e^{\beta\omega} - 1}$

Comparable with experiments

Enhancement in low-invariant mass region

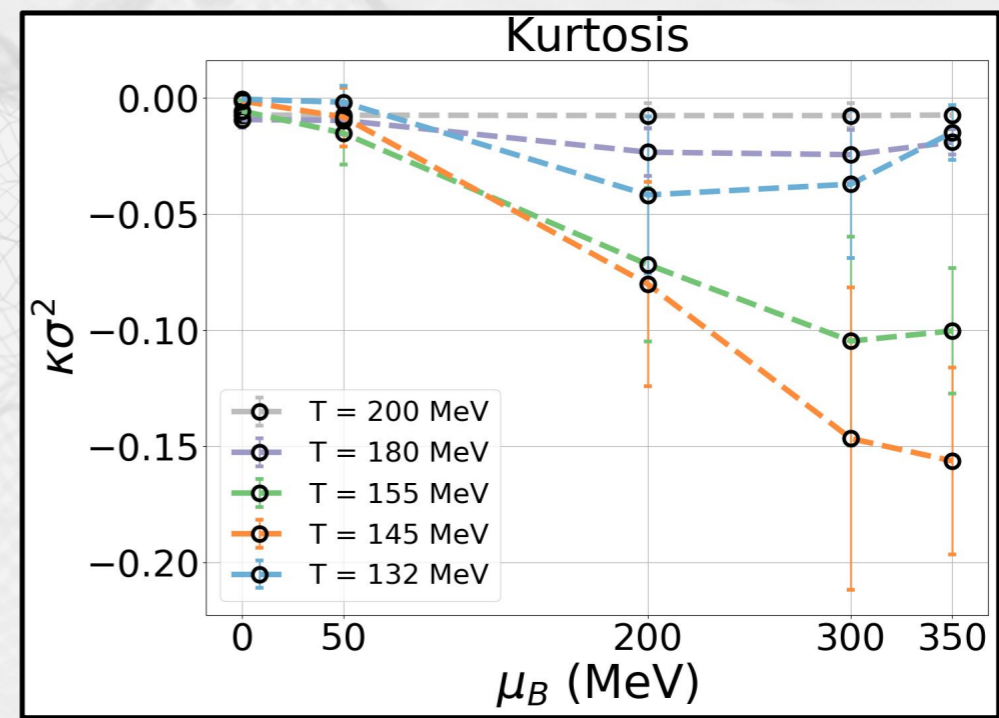
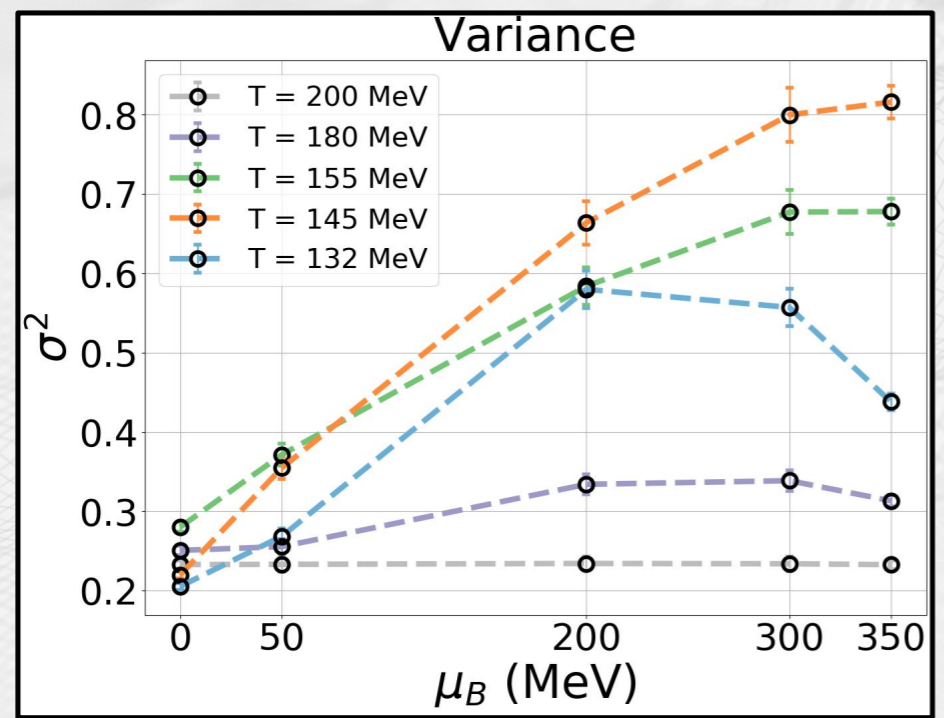
Red : AL, MT terms
Green : free quark

$\frac{d\Gamma}{dM^2} [\text{GeV}^{-2} \text{fm}^{-4}]$



Including non-linear coupling : Variance and Kurtosis

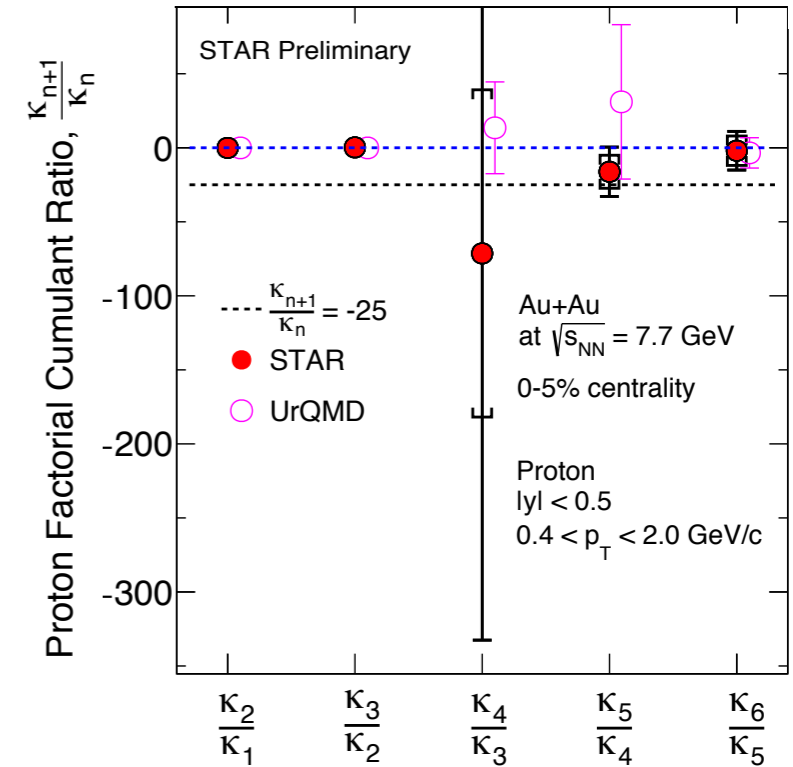
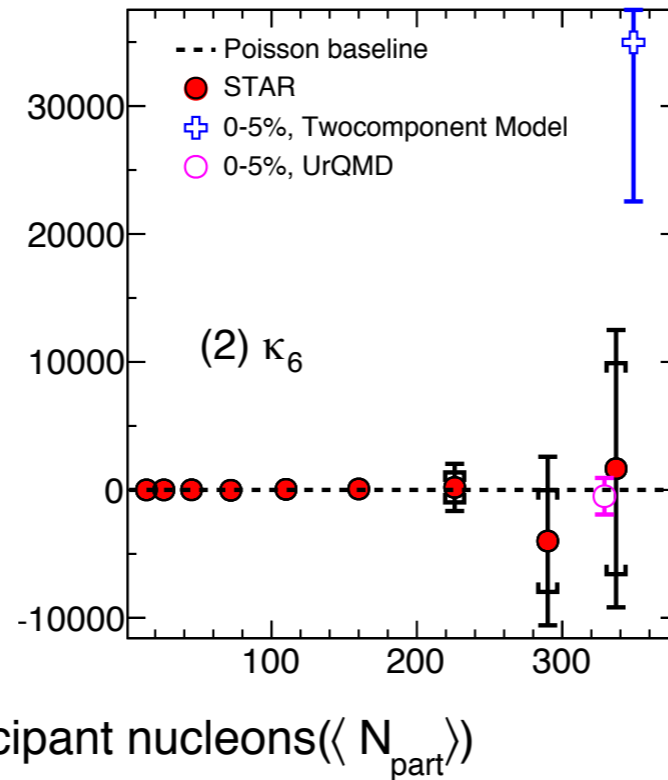
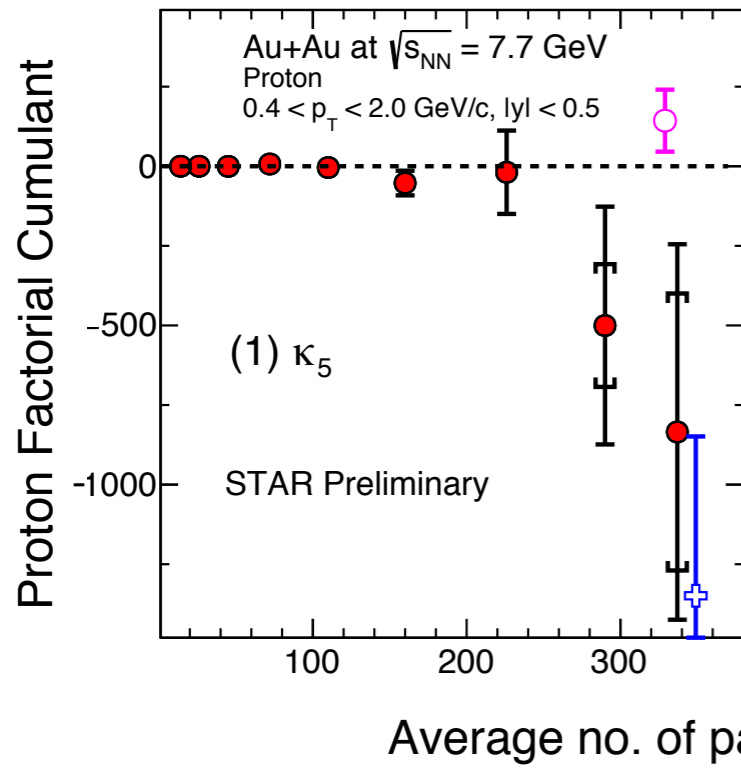
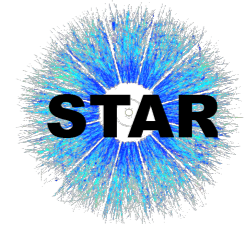
$$F(T, \delta n_B) = T \int \left[\frac{m_{\text{eff}}^2(T)}{2 n_c^2} \delta n_B^2 + \frac{K}{2 n_c^2} (\nabla \delta n_B)^2 + \frac{\lambda_{3,\text{eff}}(T)}{2 n_c^3} \delta n_B^3 + \frac{\lambda_{4,\text{eff}}(T)}{2 n_c^4} \delta n_B^4 + \frac{\lambda_{6,\text{eff}}}{2 n_c^6} \delta n_B^6 \right] d\vec{x}$$



The growth of cumulants near the critical point can be seen at freeze-out

Very sensitive to the freeze-out temperature

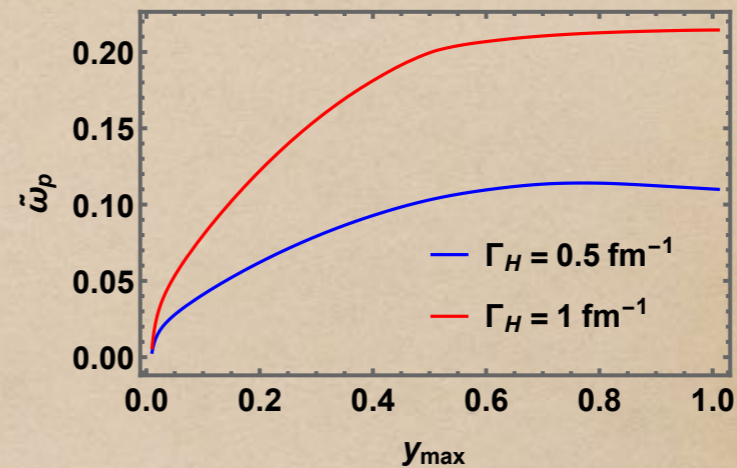
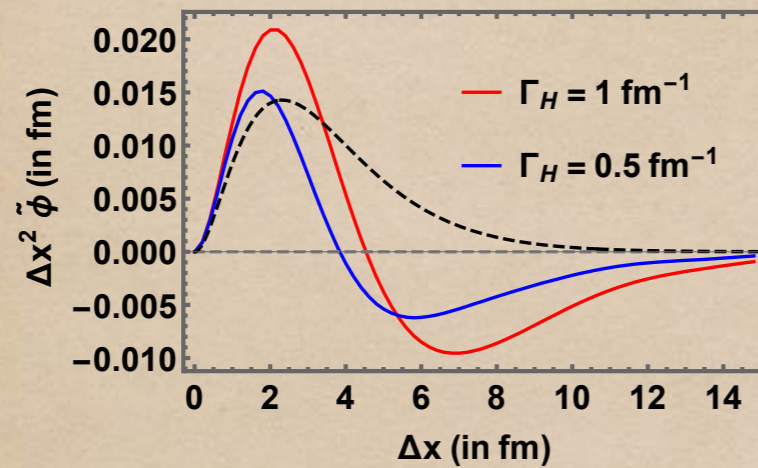
Proton Factorial Cumulant κ_5 and κ_6 at 7.7 GeV



PHYSICAL REVIEW C100, 051902(R) (2019)

κ_5 (0-5%) consistent with two component model expectation within uncertainties while κ_6 (0-5%) remains 1.8σ away. The ratios κ_5/κ_4 and κ_6/κ_5 (0-5%) consistent with zero.

Interplay of effects in dynamical evolution of QGP

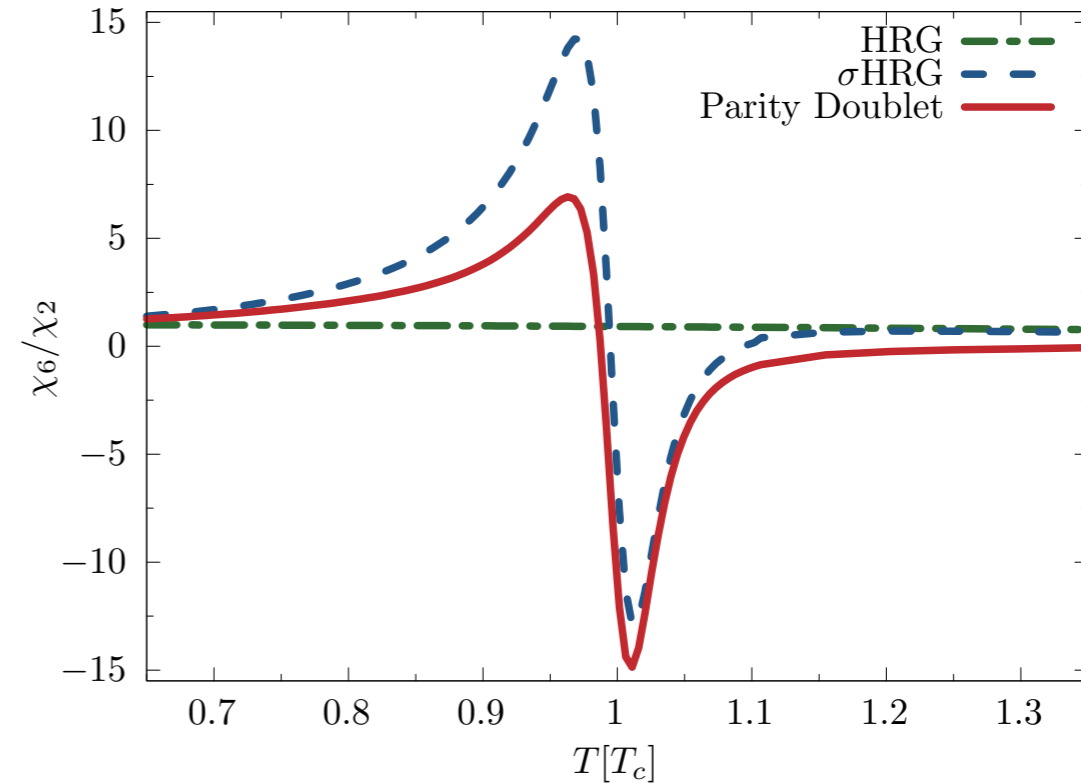
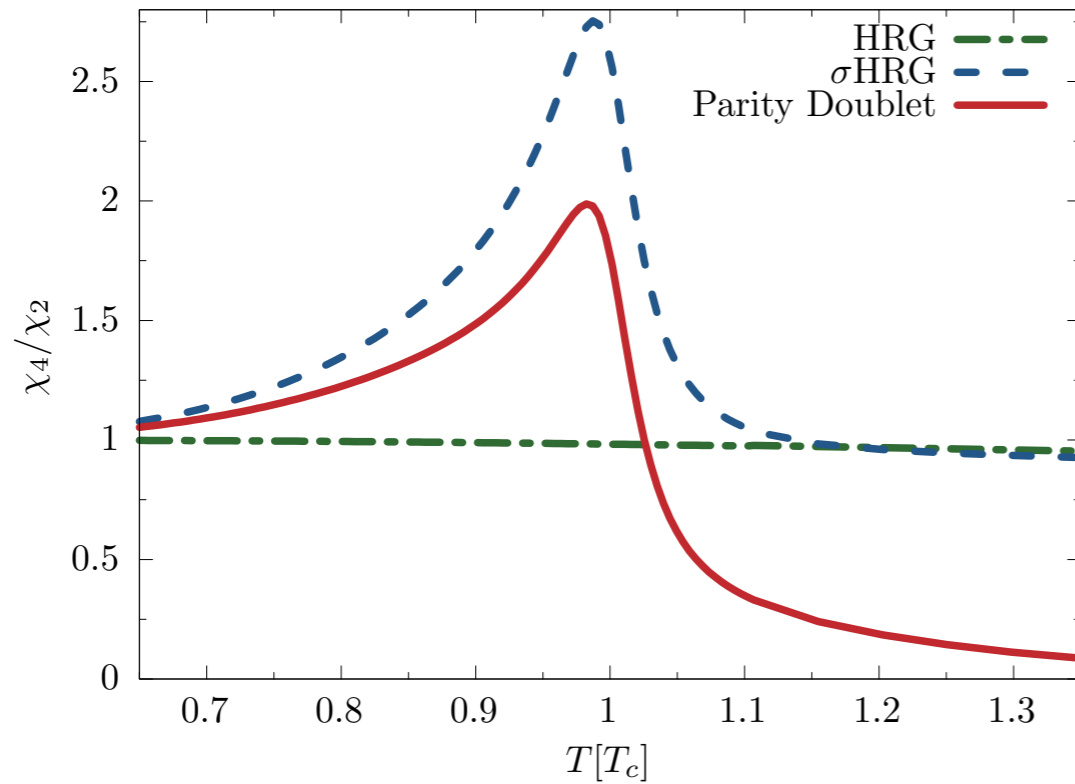


$$\Gamma(x) = \frac{\Gamma_H \xi_0^3}{\xi^3} K(x), K(x) \sim x^2 \text{ for } x \ll 1$$

$$\tilde{\omega}_A(y_{\text{max}}) = \left(\frac{\langle \delta N_A^2 \rangle_{\sigma, \text{eq}}}{\langle N_A \rangle} \right)^{-1} \frac{\langle \delta N_A^2(y_{\text{max}}) \rangle_{\sigma}}{\langle N_A(y_{\text{max}}) \rangle}$$

Ratios of higher-order cumulants: kurtosis and χ_6/χ_2

interactions \rightarrow strong deviations from the HRG baseline

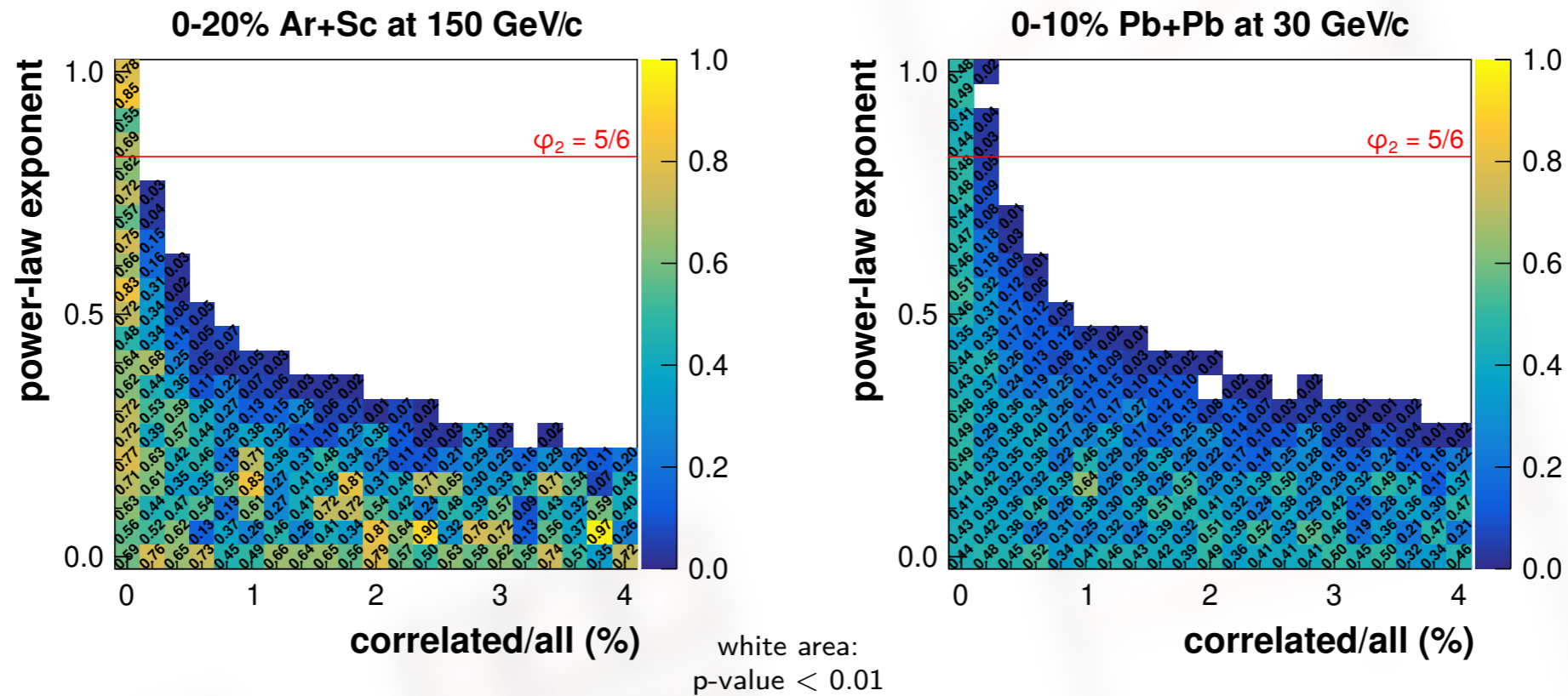


- structure dictated by the chiral symmetry
- no chiral-critical behavior encoded in β

- χ_4/χ_2 and χ_6/χ_2 suppressed by repulsion, but qualitative structure the same

Exclusion plot

Comparison with simple power-law model



exclusion plots for parameters of simple power-law model

The intermittency index φ_2 for a system freezing out at the QCD critical endpoint is expected to be $\varphi_2 = 5/6$ assuming that the latter belongs to the 3-D Ising universality class.