

Higgs and Flavour

Admir Greljo

Outline

• Part I: Flavor physics of the Higgs Boson (2/3)

• Part II: Implications of flavor anomalies for the Higgs Boson (1/3)

Part I

ullet Consider $\mathscr{L}_{\mathrm{SM}}$ sans Yukawa

$$G_{\text{global}}^{\text{SM}}(Y^{u,d,e} = 0) = SU(3)^3 \times SU(3)^2 \times U(1)^5$$

Three identical copies of five gauge representations: q, U, D, l, E

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Yukawa sector

$$-\mathcal{L}_{Yuk} = \bar{q}Y^{u}\tilde{H}U + \bar{q}Y^{d}HD + \bar{l}Y^{e}HE$$

Fermion masses and mixings

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The breaking spurions

$$Y^{u} = (3,\bar{3},1,1,1)$$
 $Y^{d} = (3,1,\bar{3},1,1)$ $Y^{e} = (1,1,1,3,\bar{3})$

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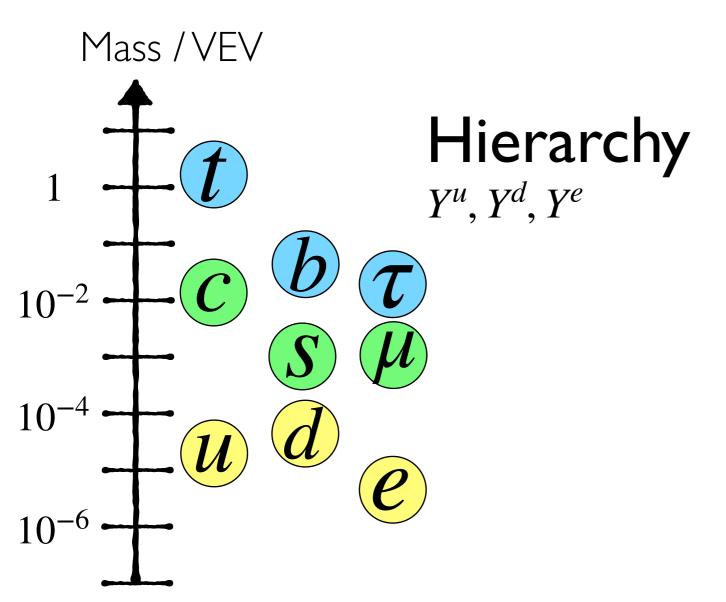
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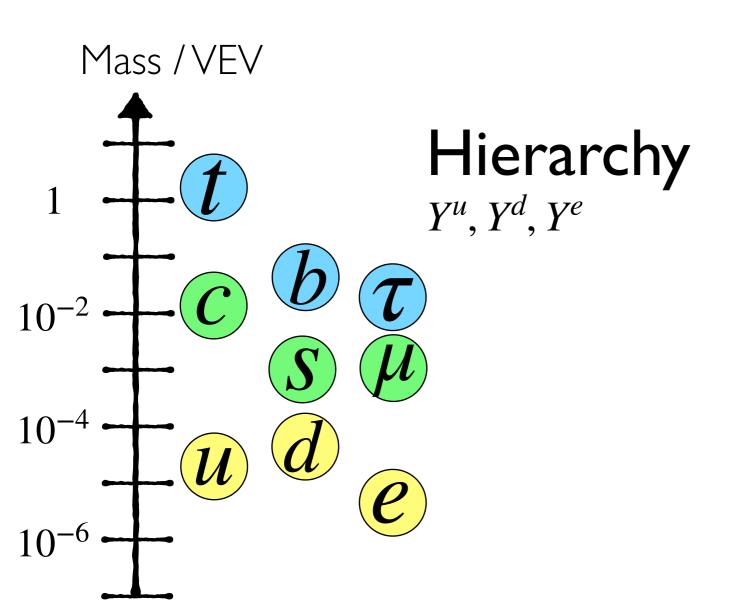
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The birth of flavor physics

$$G_{\text{global}}^{\text{SM}}(Y^{u,d,e} \neq 0) = U(1)_B \times U(1)_e \times U(1)_\mu \times U(1)_\tau$$

[Success: No proton decay, no cLFV]

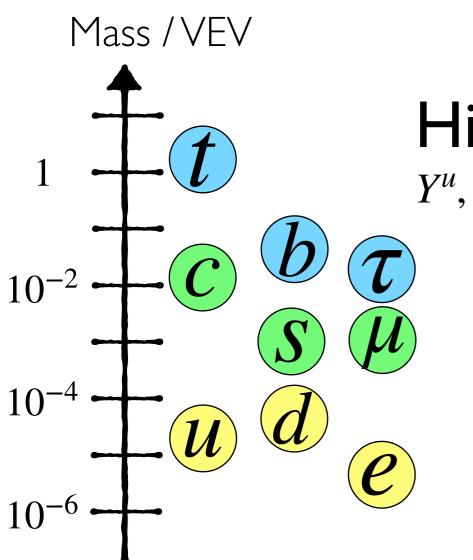




The CKM mixing

$$V_{CKM} \sim \begin{bmatrix} 1 & 0.2 & 0.2^3 \\ 0.2 & 1 & 0.2^2 \\ 0.2^3 & 0.2^2 & 1 \end{bmatrix}$$

Alignment Yu & Yd



Hierarchy

 Y^u, Y^d, Y^e

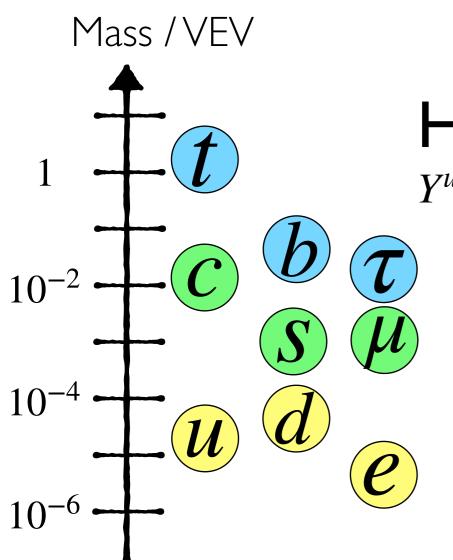
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$$\det[Y^dY^{d\dagger}, Y^uY^{u\dagger}] \approx \mathcal{O}(10^{-22})$$



Hierarchy Y^u, Y^d, Y^e

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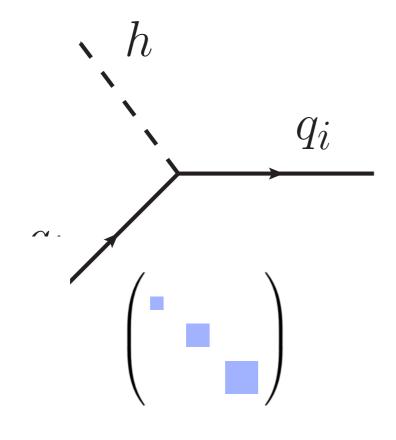
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In the SM

$$H_0 \to v + h$$

$$\mathcal{L}_{\text{Yuk}} = -\frac{h}{v} \left(m_e \,\overline{e_L} \, e_R + m_\mu \,\overline{\mu_L} \,\mu_R + m_\tau \,\overline{\tau_L} \,\tau_R \right. \\ \left. + m_u \,\overline{u_L} \,u_R + m_c \,\overline{c_L} \,c_R + m_t \,\overline{t_L} \,t_R + m_d \,\overline{d_L} \,d_R + m_s \,\overline{s_L} \,s_R + m_b \,\overline{b_L} \,b_R + \text{h.c.} \right)$$



Diagonal

Non-universal

Proportional to the fermion masses

Real in the mass basis

Beyond the SM

New sources of flavour and (or) EWS breaking would change these predictions!

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• 2HDM example

Add another Higgs doublet H_i where i = 1,2

$$-\mathcal{L}_{Yuk} = \bar{f} Y_i^f H_i F$$

$$M^f = Y_1^f v_1 + Y_2^f v_2$$

$$h = h_1 \cos \alpha + h_2 \sin \alpha$$

In general, the Higgs boson can have couplings that are neither proportional to the mass matrix nor diagonal, nor CP conserving.

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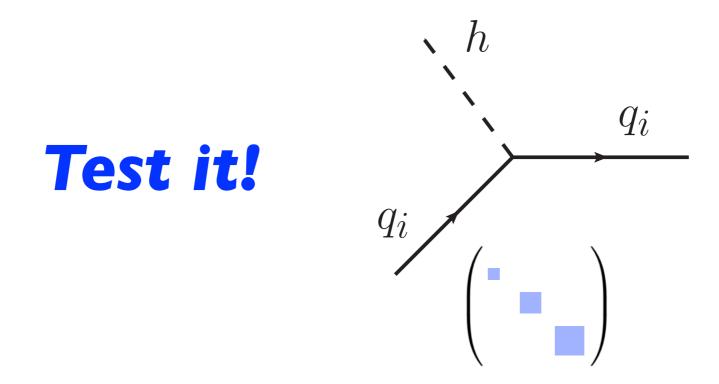
SM EFT example

Add a dim-6 SM EFT correction

$$-\mathcal{L}_{Yuk} = \bar{f} Y_1^f HF + \frac{1}{\Lambda^2} \bar{f} Y_2^f HF H^{\dagger} H$$

$$M^f \propto Y_1^f + Y_2^f \frac{v^2}{\Lambda^2} \qquad h: Y_1^f + 3 Y_2^f \frac{v^2}{\Lambda^2}$$

In general, the Higgs boson can have couplings that are neither proportional to the mass matrix nor diagonal, nor CP conserving.



- Diagonal couplings?
- Off-diagonal couplings?
- CP violation?

Diagonal couplings

$$\kappa_t = 1.43 \pm 0.23,$$

$$\kappa_s < 65$$
,

$$\kappa_{\tau} = 0.88 \pm 0.13,$$

$$\kappa_b = 0.60 \pm 0.18,$$

$$\kappa_d < 1.4 \cdot 10^3,$$

$$\kappa_{\mu} = 0.2^{+1.2}_{-0.2},$$

$$\kappa_c \lesssim 6.2$$
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$$\kappa_u < 3.0 \cdot 10^3,$$

$$\kappa_e \lesssim 630.$$

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1610.07922, Section IV.6.2.c, LHC Higgs Cross Section Working Group

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Charm Yukawa

- Exclusive Higgs decays to mesons:
 1407.6695, 1406.1722, 1505.03870
- Vh associated production:
 1503.00290,1505.06689,1505.06689
- Higgs differential distributions: 1606.09253, 1606.09621

HL-LHC sensitivity
$$\mathcal{O}(y_c)$$

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HL-LHC sensitivity $\mathcal{O}(y_c)$

Muon Yukawa

 1.2 ± 0.6 , ATLAS 2007.07830.

 1.2 ± 0.4 , CMS CMS-PAS-HIG-19-006.

The observation at the end of Run 3?

Off-diagonal couplings

Quarks

- Neutral meson mixing provide stringent constraints

Off-diagonal couplings

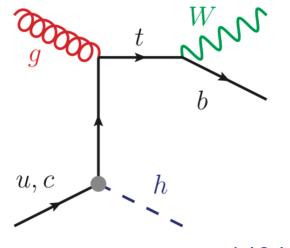
Quarks

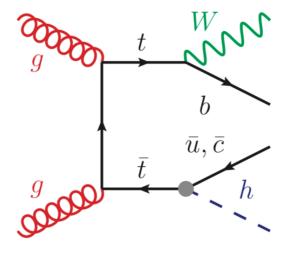
- Neutral meson mixing provide stringent constraints

$$K - \bar{K}$$
 $Br(h \rightarrow s\bar{d} + d\bar{s}) < 4.2 \times 10^{-7}$
 $D - \bar{D}$ $Br(h \rightarrow c\bar{u} + u\bar{c}) < 3.7 \times 10^{-6}$
 $B - \bar{B}$ $Br(h \rightarrow b\bar{d} + d\bar{b}) < 1.7 \times 10^{-5}$
 $B_s - \bar{B}_s$ $Br(h \rightarrow b\bar{s} + s\bar{b}) < 1.3 \times 10^{-3}$

1610.07922, Section IV.6.2.c, LHC Higgs Cross Section Working Group

- Top decays and tH production





$$Br(t \to ch) < 0.11 \%$$
ATLAS, 1812.11568

$$Br(t \to ch) < 0.47 \%$$

1404.1278

Off-diagonal couplings

Leptons

 $\mu \to e \gamma$ implies stringent constraints on $h \to \mu e$

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Leptons

 $\mu \to e \gamma$ implies stringent constraints on $h \to \mu e$

ullet For $h o au \mu$ and h o au e the best constraints are from Higgs decays

$$Br(h \to \tau \mu) < 0.25 \%$$

$$Br(h \to \tau e) < 0.61 \%$$

$$Br(h \to \tau \mu) < 0.28 \%$$

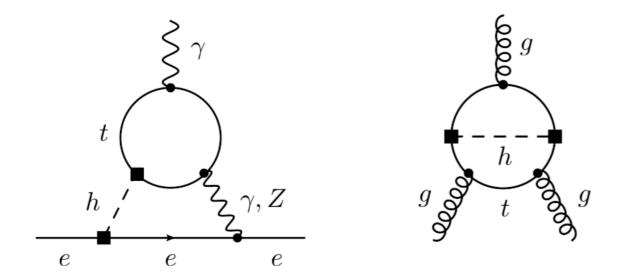
$$Br(h \to \tau e) < 0.47 \%$$

CMS 1712.07173

ATLAS 1907.06131

CP violation

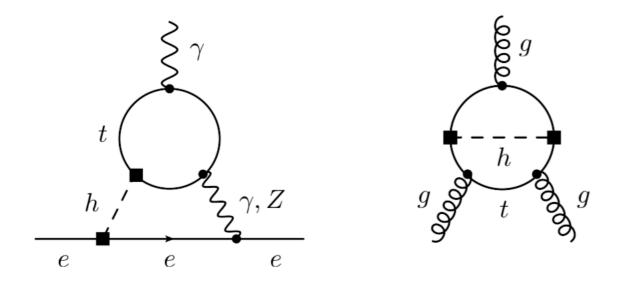
• EDMs



1310.1385, 1503.04830, 1510.00725

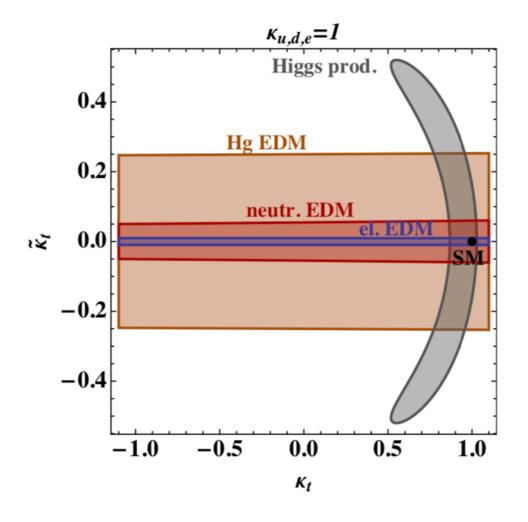
CP violation

• EDMs



1310.1385, 1503.04830, 1510.00725

• EDMs versus LHC interplay

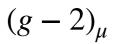


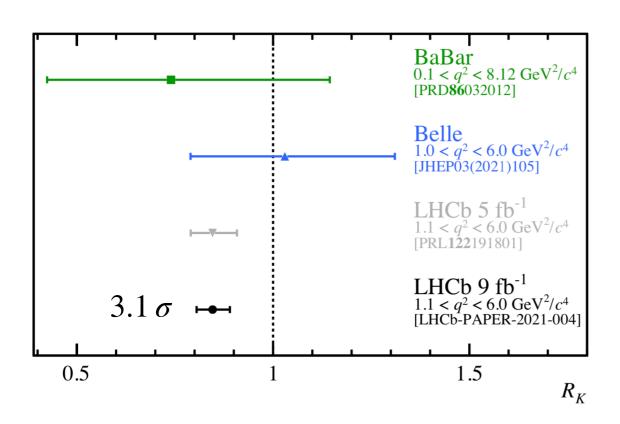
1310.1385

Part II

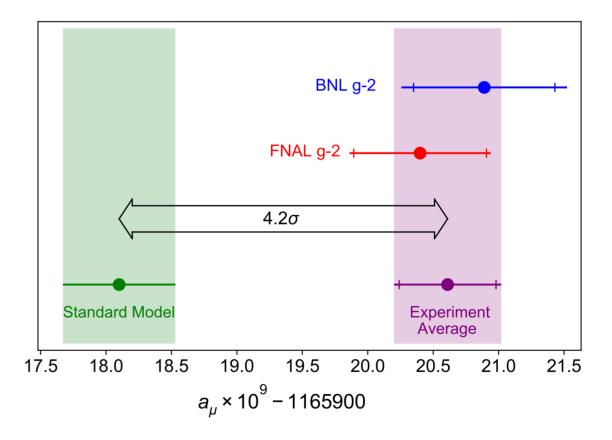
Hot topic in flavour physics: Muon Anomalies

$$\frac{b \to s\mu\mu}{b \to see}$$





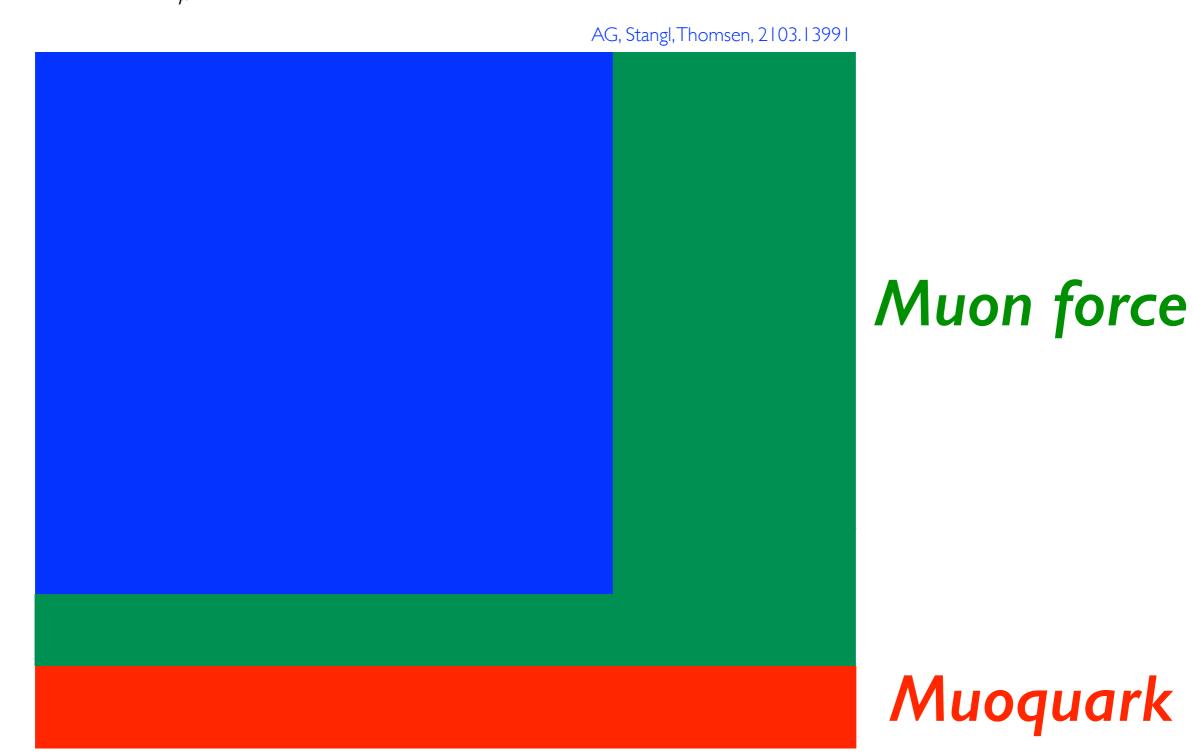
LHCb, CERN, 2103.11769



The Muon g-2, Fermilab, 2104.03281

• $SM \times U(1)_{B-3L_{\mu}}$ gauge symmetry

SM



• $SM \times U(1)_{B-3L_{\mu}}$ gauge symmetry

SM

AG, Stangl, Thomsen, 2103.13991

	AG, Stangl, I homsen, 2103.1399				
	SU(3) _e	SU(2)L	U(1) _Y		
QL	3	2	1/6		
L	1	2	-1/2		
UR	3	ı	2/3		
dR	3	l	-1/3		
V _R	1	l	0		
$\mathcal{C}_{\mathcal{R}}$	l	1	-1		
H	1	2	1/2		

Muon force

Muoquark

• $SM \times U(1)_{B-3L_{\mu}}$ gauge symmetry

SM

AC Stand Thomson 2103 13991

	AG, Stangl, Thomsen, 2103.13991				
	SU(3) _e	SU(2)L	U(1) _Y	U(1) B-3 Lm	
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_ dr	3	l	-1/3	1/3	
V_R	1	ı	0	₹0,-3,0g	
e _R	l	l	-1	₹0,-3,0g	
+	1	2	1/2	0	
豆	1	l	0	3	

Muon force

* Minimal type-I seesaw for the neutrino masses

Muoquark

• $SM \times U(1)_{B-3L_{\mu}}$ gauge symmetry

SM

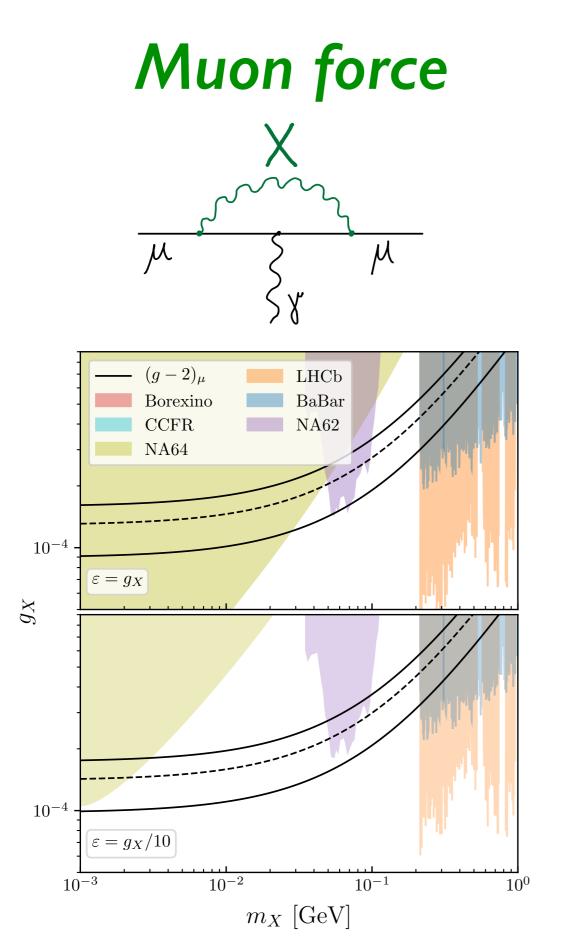
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	SU(3)c	SU(2)L	U(1) _Y	U(1) B-3 Lm
Q_{L}	3	2	1/6	1/3
L	1	2	-1/2	₹0,-3,03
UR	3		2/3	1/3
clr	3		-1/3	1/3
V_R	1		0	₹0,-3,0g
$\mathcal{C}_{\mathcal{R}}$	l		-1	₹0,-3,03
+	1	2	1/2	0
更			0	3
S_3	3	3	1/3	8/3

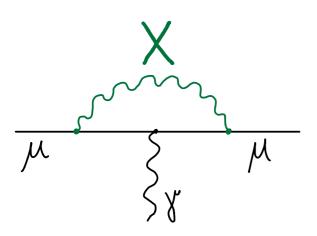
Muon force

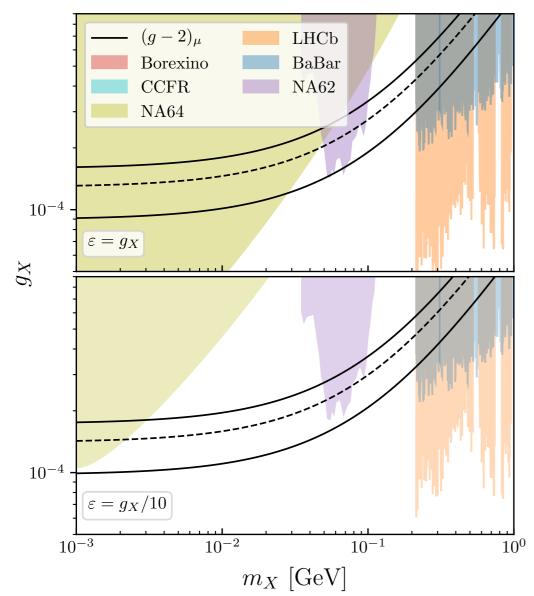
Muoquark

$$\mathcal{L} \supset Q_L L_L^{(2)} S_3$$

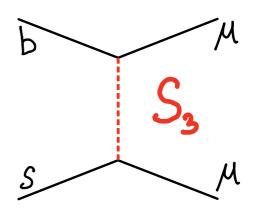


Muon force





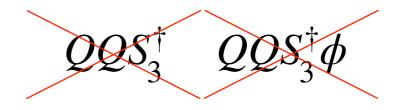
Muoquark



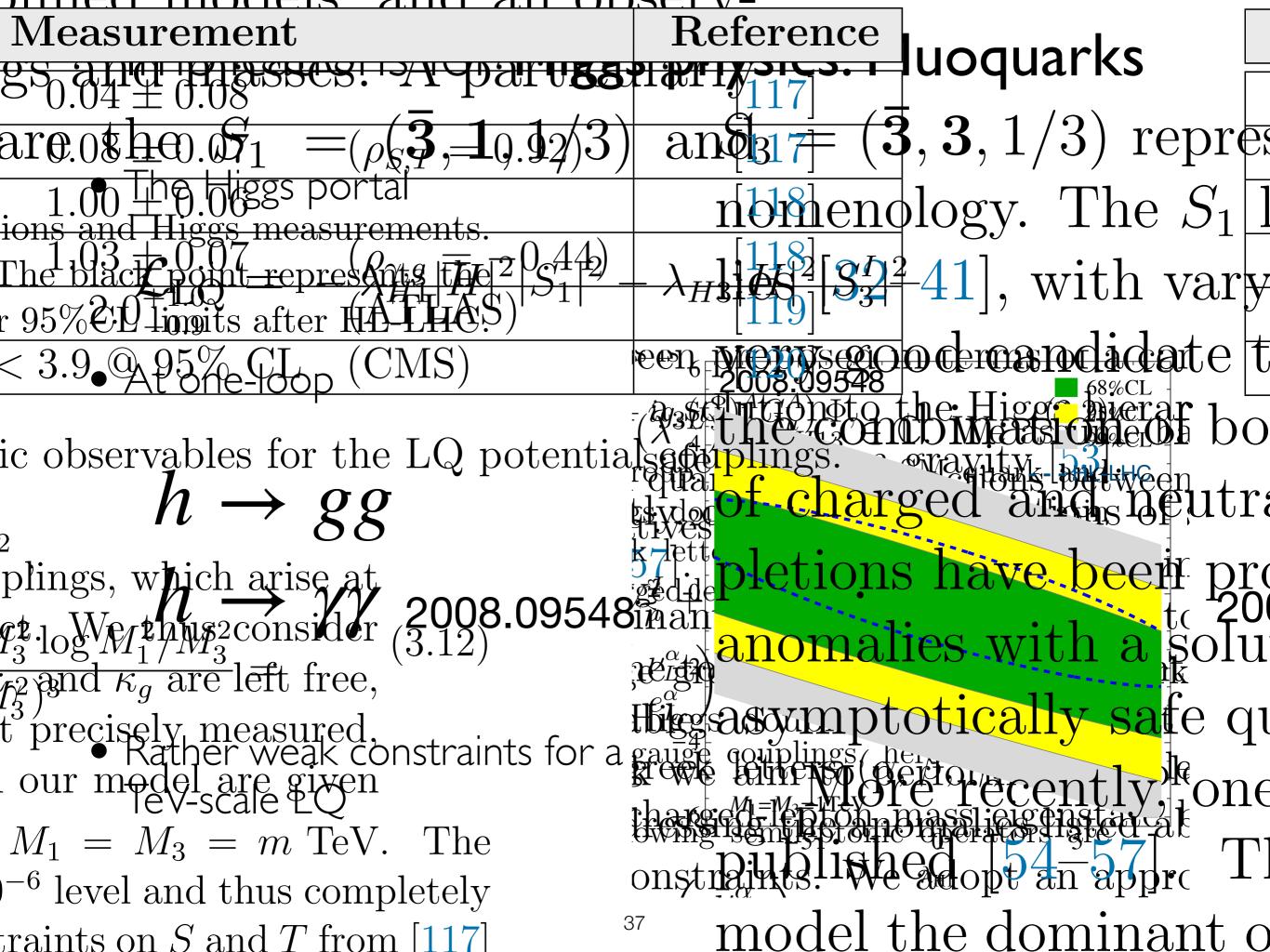
- What $U(1)_{X_{\mu}}$ does to a leptoquark?
 - Interacts only with muons

$$\mathcal{L} \supset Q_L L_L^{(2)} S_3$$

No proton decay up to dim-6



Juoquarks gs and masses: A parusse $\frac{1}{3}$ (3, 3, 1/3) represent $are.0$heo.6 = (\rho(3.10.12)$ the higgs portal <u>nbinenology</u>. The S_1 $H^{\frac{3}{2}}[S_3^2] + 41]$, with vary r 95% EV Timpits after HALLAC very good candidate t < 3.9ic observables for the LQ potential couplings. of charged and neutra $h \rightarrow gg$ pletions have been pro plings, which arise at 2008.09548 1×1008 1×10 anomalies with a solu $\frac{1}{2}$ and $\frac{1}{\kappa_g}$ are left free, asymptotically safe qu t³ precisely measured, our model are given More recently, one $M_1 = M_3 = m \text{ TeV.}$ The published [54–57]. T 0^{-6} level and thus completely model the dominant of traints on S and T from [117]



Implications for Higgs physics: Muon force

$$V_{H\Phi} = -\mu_H^2 |H|^2 - \mu_\Phi^2 |\Phi|^2 + \frac{1}{2}\lambda_H |H|^4 + \frac{1}{4}\lambda_\Phi |\Phi|^4 + \lambda_{\Phi H} |\Phi|^2 |H|^2$$

• From $(g-2)_{\mu}$ we have $g_X \sim 10^{-4}$ and $m_X \in [10,200]\,{
m MeV}$.

$$v_{\Phi} = \sqrt{2}m_X/|q_{\Phi}|g_X \sim 60 \,\mathrm{GeV}/|q_{\Phi}|$$

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• Mixing between real scalars h and ϕ .

$$g_X: X \to \nu_{\mu} \bar{\nu}_{\mu}$$
 $\lambda_{\Phi}: \phi \to XX$
 $h \to inv$

• This scenario has a chance to leave observable imprints in the overall Higgs couplings or in the invisible Higgs decays.

Conclusions

- Flavor physics of the Higgs Boson is a rich subject
 - Diagonal couplings
 - Off-diagonal couplings
 - CP violation
- Flavor anomalies: Model specific predictions for Higgs physics