

Higgs in the Swampland

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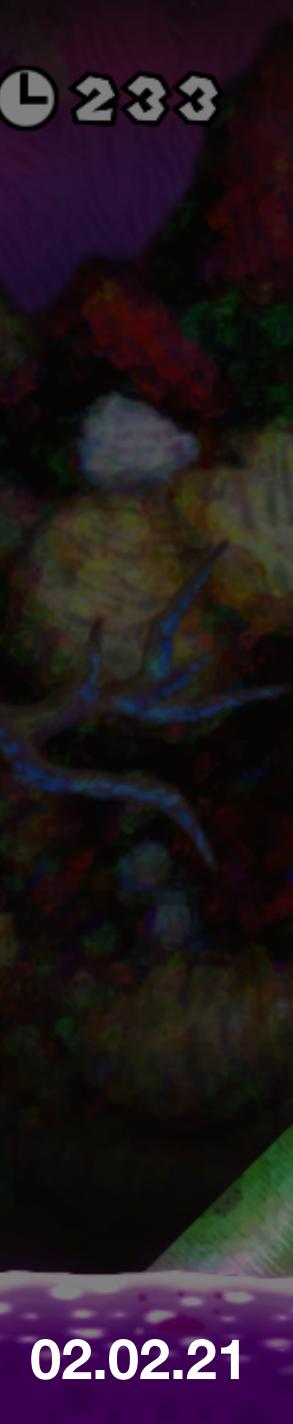
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CAU BSM 2021

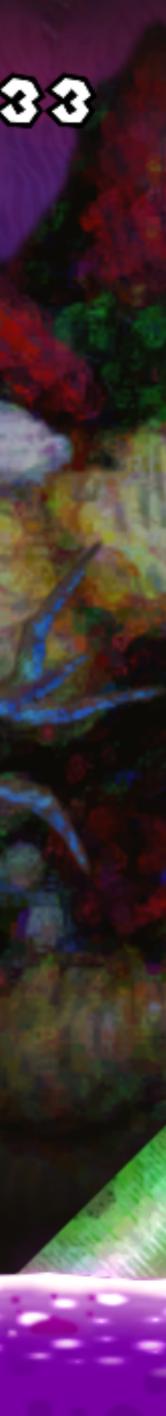
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The Naturalness Strategy

Param	UV sensitivity	Natural if	NP	Scale	Natural?
"me"	$e^2\Lambda$	Λ ≲ 5 MeV	Positron	511 keV	
m _{π±} ² - m _{π0} ²	$\frac{3\alpha}{4\pi}\Lambda^2$	Λ ≲ 850 MeV	Rho	770 MeV	
MKL-MKS	$\frac{s_c^2 f_K^2 m_{K_L^0}}{24\pi^2 v^4} \Lambda^2$	Λ ≲ 2 GeV	Charm	1.2 GeV	
m _H 2	$-\frac{6y_t^2}{16\pi^2}\Lambda^2 + \dots$	Λ ≲ 500 GeV	?	?	?

From the "naturalness strategy" to new physics

At this level, we expect

New physics below the TeV scale...

...coupling to the Higgs

Strong motivation for BSM Higgs physics! But maybe too broad to be useful...

Implementation is up to us

- The Standard Model coupled to gravity is a 1. generic EFT.
- 2. The solutions to the hierarchy problem involve symmetries, low cutoffs, or anthropics.
- Symmetries imply new particles charged 3. under the SM.

We've refined this strategy using some rules of thumb, for example...

In turn, this tells us what kind of NP to expect: SUSY, CHM,...



 $\tilde{H}\tilde{H}, \tilde{H} \rightarrow h\tilde{G}/Z\tilde{G}$

Stable g R-hadron

2018-024

018-020

Direct $\tilde{\chi}_1^+ \tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^+$

 $\tilde{\chi}_1^{\pm} \tilde{\chi}_1^{\mp} / \tilde{\chi}_1^0$, $\tilde{\chi}_1^{\pm} \rightarrow Z \ell \rightarrow \ell \ell \ell$

$$\begin{split} \tilde{\chi}_{1}^{\pm} \tilde{\chi}_{1}^{\mp} / \tilde{\chi}_{2}^{0} &\to WW/Z\ell\ell\ell\ell\ell\nu\nu\\ \tilde{g}\tilde{g}, \, \tilde{g} \to qq\tilde{\chi}_{1}^{0}, \, \tilde{\chi}_{1}^{0} \to qqq \end{split}$$

Metastable \tilde{g} R-hadron, $\tilde{g} \rightarrow qq \tilde{\chi}_1^0$

 $\mathsf{LFV}\ pp {\rightarrow} \tilde{\nu}_{\tau} + X, \tilde{\nu}_{\tau} {\rightarrow} e\mu/e\tau/\mu\tau$

						$\int \mathcal{L} dt = (\mathcal{L} - 1)$	58, 0	∖ =0, ∖
	Model	ℓ,γ	Jets†	E_{T}^{miss}	∫£ dt[fb			Referen
Extra dimensions	ADD $G_{KK} + g/q$ ADD non-resonant $\gamma\gamma$ ADD QBH ADD BH high $\sum p_T$ ADD BH multijet RS1 $G_{KK} \rightarrow \gamma\gamma$ Bulk RS $G_{KK} \rightarrow WW/ZZ$ Bulk RS $G_{KK} \rightarrow WW \rightarrow troc$ Bulk RS $g_{KK} \rightarrow tt$ 2UED / RPP	1 <i>e</i> ,µ ≥	$1 - 4j$ $-$ $2j$ $\geq 2j$ $\geq 3j$ $-$ 1 $2j / 1 J$ $\geq 1 b, \geq 1J/2$ $\geq 2 b, \geq 3j$		36.1 36.7 37.0 3.2 3.6 36.7 36.1 139 36.1 36.1	Min 8.9 TeV n = 6 Min 8.2 TeV n = 6, / Min 9.55 TeV n = 6, / Grad mass 2.3 TeV k/Mp/ Grad mass 2.0 TeV k/Mp/ grad mass 3.8 TeV I/mp/	= 1.0 = 1.0	1711.033 1707.041 1703.091 1606.022 1512.025 1707.041 1808.023 2004.146 1804.106 1804.106
Gauge bosons	$\begin{array}{c} \mathrm{SSM}\; Z' \to \ell\ell \\ \mathrm{SSM}\; Z' \to \tau\tau \\ \mathrm{Leptophobic}\; Z' \to bb \\ \mathrm{Leptophobic}\; Z' \to tt \\ \mathrm{SSM}\; W' \to \ell\nu \\ \mathrm{SSM}\; W' \to \tau\nu \\ \mathrm{HVT}\; W' \to WZ \to \ell\nu qq qm m \\ \mathrm{HVT}\; V' \to WH \to qq qq m m \\ \mathrm{HVT}\; V' \to WH / ZH model \\ \mathrm{HVT}\; V' \to WH model \\ \mathrm{LRSM}\; W_R \to tb \\ \mathrm{LRSM}\; W_R \to \mu N_R \end{array}$	$\begin{array}{c} 1 \ e, \mu \\ 1 \ \tau \\ \end{array}$ odel B 1 e, μ odel B 0 e, μ B multi-channel	≥ 1 b, ≥ 2 J	Yes Yes Yes –	139 36.1 36.1 139 36.1 139 36.1 139 36.1 139 36.1 80	Z' mass 5.1 TeV Z' mass 2.42 TeV Z' mass 2.1 TeV Z' mass 4.1 TeV V' mass 6.0 TeV W' mass 3.7 TeV W' mass 2.93 TeV W' mass 3.8 TeV V' mass 2.93 TeV W' mass 3.2 TeV Wr mass 3.2 TeV Wr mass 5.0 TeV	3	1903.062 1709.072 1805.092 2005.051 1906.056 1801.069 2004.146 1906.085 1712.065 CERN-EP-20 1807.104 1904.126
CI	Cl qqqq Cl ℓℓqq Cl ℓℓtt	 2 e, μ ≥1 e,μ	2 j ≥1 b, ≥1 j	_ _ Yes	37.0 139 36.1		8 TeV η _{LL} 35.8 TeV η _{LL}	1703.091 CERN-EP-20 1811.023
MQ	Axial-vector mediator (Dirac Colored scalar mediator (Dir $VV_{\chi\chi}$ EFT (Dirac DM) Scalar reson. $\phi \rightarrow t\chi$ (Dirac	racDM) 0 e, μ 0 e, μ	$\begin{array}{c} 1-4 \ j \\ 1-4 \ j \\ 1 \ J, \leq 1 \ j \\ 1 \ b, \ 0\mbox{-}1 \ J \end{array}$	Yes Yes Yes Yes	36.1 36.1 3.2 36.1	m _{med} 1.67 TeV g=1.0, r M. 700 GeV m(\chi)	5, g_{χ} =1.0, $m(\chi) = 1 \text{ GeV}$ $m(\chi) = 1 \text{ GeV}$ < 150 GeV 4, $\lambda = 0.2$, $m(\chi) = 10 \text{ GeV}$	1711.033 1711.033 1608.023 1812.097
۲Ø	Scalar LQ 1 st gen Scalar LQ 2 nd gen Scalar LQ 3 rd gen Scalar LQ 3 rd gen	1,2 e 1,2 μ 2 τ 0-1 e, μ	≥ 2 j ≥ 2 j 2 b 2 b	Yes Yes – Yes	36.1 36.1 36.1 36.1		ightarrow b au) = 1 ightarrow t au) = 0	1902.003 1902.003 1902.081 1902.081
Heavy quarks	$ \begin{array}{l} VLQ\;TT \rightarrow Ht/Zt/Wb + X \\ VLQ\;BB \rightarrow Wt/Zb + X \\ VLQ\;T_{5/3}T_{5/3}T_{5/3} \rightarrow Wt + \\ VLQ\;Y \rightarrow Wb + X \\ VLQ\;Y \rightarrow Wb + X \\ VLQ\;B \rightarrow Hb + X \\ VLQ\;QQ \rightarrow WqWq \end{array} $	multi-channel - X 2(SS)/≥3 e,µ	l ₂ ≥1 b, ≥1 j ≥ 1 b, ≥ 1j	Yes	36.1 36.1 36.1 36.1 79.8 20.3		doublet $\rightarrow Wt$)= 1, $c(T_{5/3}Wt)$ = 1 $\rightarrow Wb$)= 1, $c_R(Wb)$ = 1	1808.023 1808.023 1807.118 1812.073 ATLAS-CONF-2 1509.042
Excited fermions	Excited quark $q^* \rightarrow qg$ Excited quark $q^* \rightarrow q\gamma$ Excited quark $b^* \rightarrow bg$ Excited lepton ℓ^* Excited lepton ν^*	- 1 γ - 3 e,μ 3 e,μ,τ	2j 1j 1b,1j –	- - -	139 36.7 36.1 20.3 20.3			1910.084 1709.104 1805.092 1411.292 1411.292
Other	Type III Seesaw LRSM Majorana v Higgs triplet $H^{\pm\pm} \rightarrow \ell\ell$ Higgs triplet $H^{\pm\pm} \rightarrow \ell\tau$ Multi-charged particles Magnetic monopoles $\sqrt{s} = 8 \text{ TeV}$	1 e,μ 2μ 2,3,4 e,μ(SS 3 e,μ,τ - - - - - -	≥ 2 j 2 j) – – – – –	Yes 	79.8 36.1 36.1 20.3 36.1 34.4	H** mass 870 GeV DY prod H** mass 400 GeV DY prod multi-charged particle mass 1.22 TeV DY prod monopole mass 2.37 TeV DY prod	$)=$ 4.1 TeV, $g_L=g_R$	ATLAS-CONF-2 1809.111 1710.097 1411.293 1812.036 1905.101
	vs = o lev	partial data	full da	ata		10 ⁻¹ 1 ¹⁰ M	lass scale [TeV]	

*Only a selection of the available mass limits on new states or phenomena is shown. *†Small-radius (large-radius) jets are denoted by the letter j (J).*

ΒP	$ \begin{array}{l} \widetilde{ti}_{1}, \widetilde{t} \rightarrow \widetilde{k}_{1}^{0}, \widetilde{\chi}_{1}^{0} \rightarrow tbs \\ \widetilde{ti}_{1}, \widetilde{t} \rightarrow b\widetilde{k}_{1}^{+}, \widetilde{\chi}_{1}^{+} \rightarrow bbs \\ \widetilde{t}_{1}, \widetilde{t}_{1}, \rightarrow bs \\ \widetilde{t}_{1}, \widetilde{t}_{1}, \rightarrow q\ell \end{array} $	$\begin{array}{c} \mbox{Multiple} \\ \geq 4b \\ 2 \mbox{ jets } + 2 \ b \\ 2 \ e_* \mu & 2 \ b \\ 1 \ \mu & \mbox{DV} \end{array}$	36.1 139 36.7 36.1 136	
phér	a selection of the available r nomena is shown. Many of th lified models, c.f. refs. for the	ne limits are based on	r 1	0 ⁻¹

0 e,μ 4 e,μ

3 *e*,μ

 $e\mu,e au,\mu au$

139 139

36.

36.1

36.

36. 36.

 $\begin{array}{ccc} \geq 3 \ b & E_T^{\text{miss}} & \textbf{36.1} \\ \textbf{0 jets} & E_T^{\text{miss}} & \textbf{139} \end{array}$

Disapp. trk 1 jet Emiss 36

Multiple Multiple

4 e, μ 0 jets E_T^{miss}

4-5 large-*R* jets Multiple

Mass scale [TeV]

 $BR(\tilde{t}_1 \rightarrow be/b\mu) > 20\%$ $BR(\tilde{t}_1 \rightarrow q\mu) = 100\%, \cos\theta_t = 1$

 $m(\tilde{\chi}_1^0) =$

 $\mathsf{m}(\tilde{\ell}^0) = 0$ $\mathsf{m}(\tilde{\ell}) - \mathsf{m}(\tilde{\ell}^0) = 10 \text{ GeV}$

 $\mathsf{BR}(\tilde{\chi}^0_1 \to h\tilde{G}) =$ $\mathsf{BR}(\tilde{\chi}^0_1 \to Z\tilde{G}) =$

Pure higgsir

 $m(\tilde{\chi}_1^0)=100 \text{ GeV}$

 $m(\tilde{\chi}_1^0)=100 \text{ GeV}$

m(X̃[±]₁)=500 Ge\

 $\lambda'_{311}=0.11, \lambda_{132/133/233}=0.07$

m($\tilde{\chi}_1^0$)=200 GeV, bino-like

Pure Wino

Large λ_{112}'' $m(\tilde{\chi}_1^0)$ =200 GeV, bino-like

1710.05544 2003.11956

1911.06660

1908.08215 1911.12606

1806.04030 ATLAS-CONF-2020-040

719 091

ATL-PHYS-PUB-2017-019

1902.01636,1808.04095

1710.04901,1808.04095

ATLAS-CONF-2020-009

1607.08079

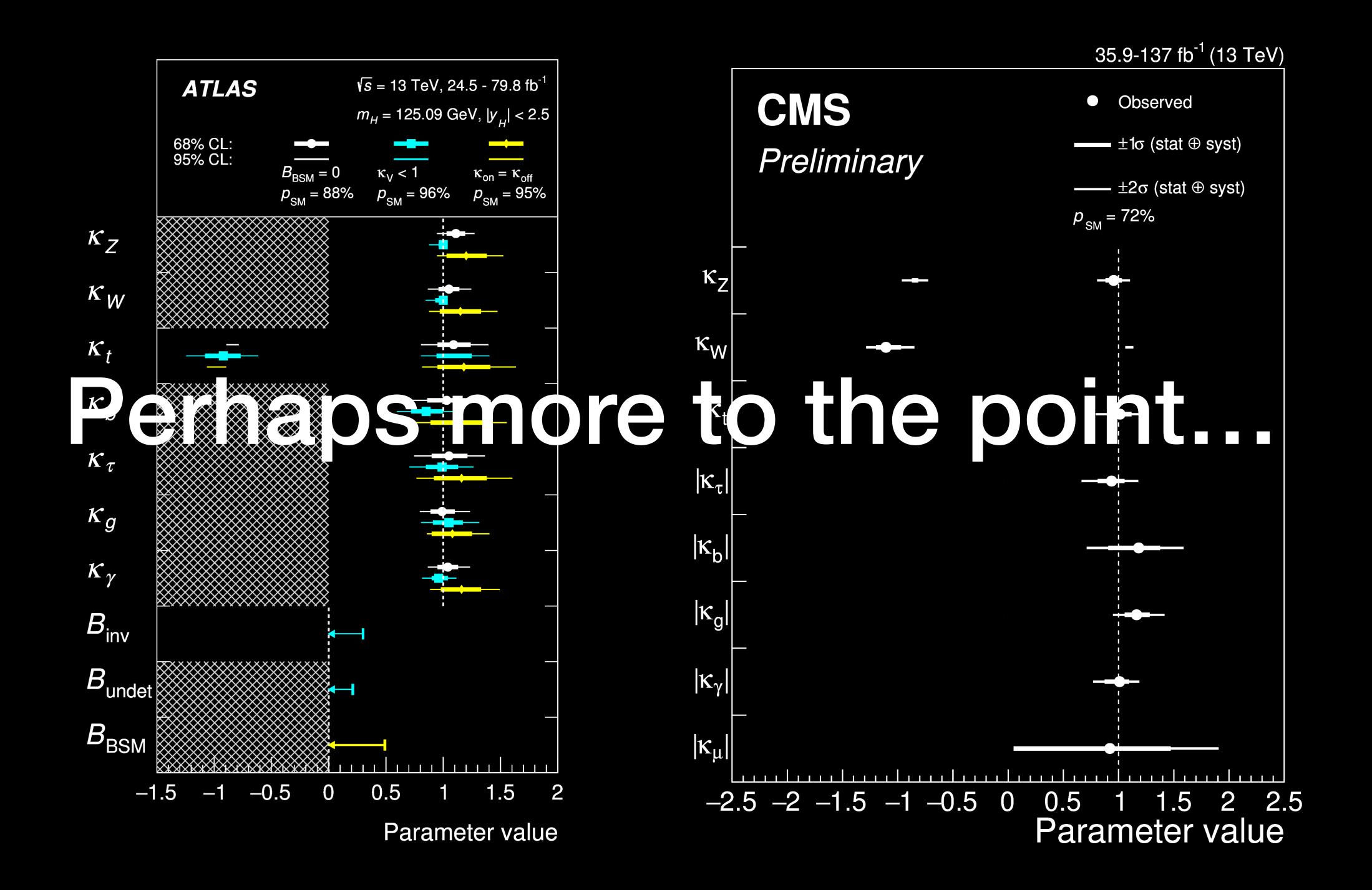
1804.03602

1804.03568 ATLAS-CONF-2018-003

ATLAS-CONF-2018-003

ATLAS-CONF-2020-016

1710.07171



On the wrong track?

What if some of the rules of thumb are wrong?

- 1
- 3.

The Standard Model coupled to gravity is a —> TODAY generic EFT.

2. The solutions to the hierarchy problem involve symmetries, low cutoffs, or anthropics.

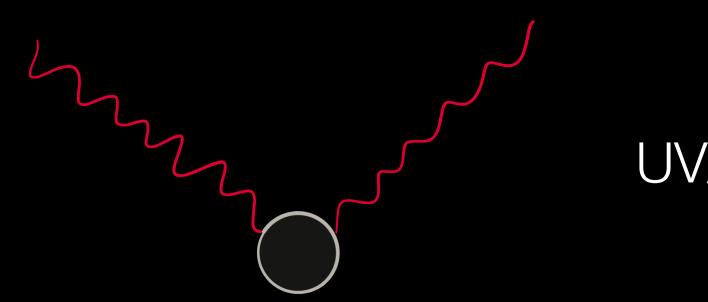
> Symmetries imply new particles charged under the SM.

Relaxion

Twin Higgs, discrete symmetries, "neutral naturalness", ...

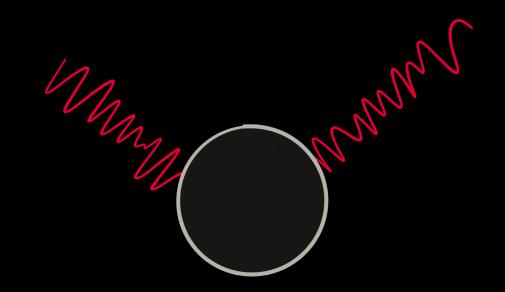
The End of EFT?

Parameter space of EFTs shaped by consistent coupling to gravity



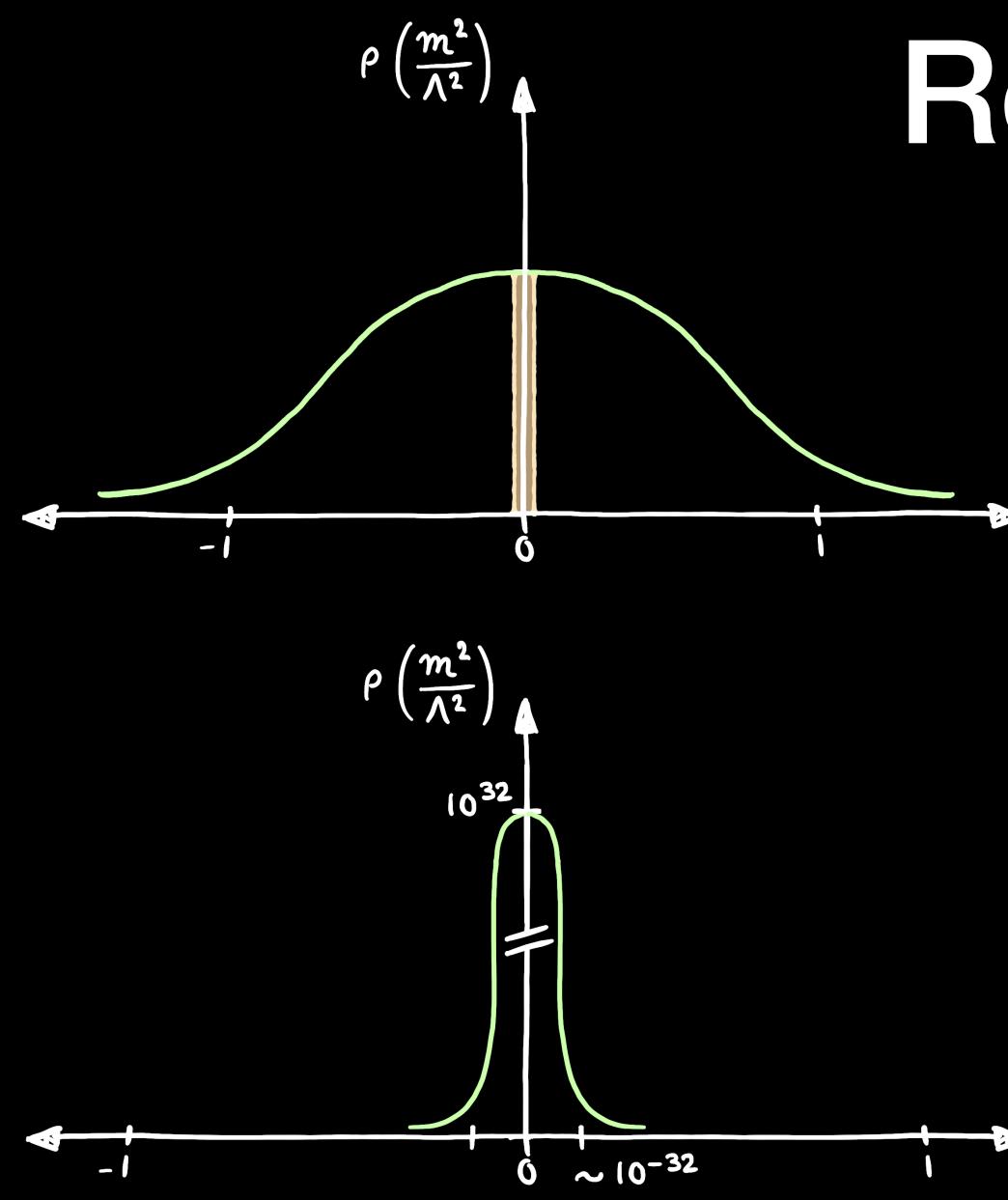
Various (conjectured) consequences:

- No global symmetries
- Charge quantization
- Completeness hypothesis



UV/IR mixing...

- No (meta)stable dS vacua
- Infinite states @ infinite distances
- "Gravity is the weakest force"



[Isabel Garcia Garcia, BSM Pandemic seminar 07/20]

M

Relevance to BSM?

Usual (EFT) logic of hierarchy problem: uncorrelated UV contributions give broad distribution of possible values of m_h up to cutoff; m_h well below cutoff "unlikely"

Usual (EFT) logic of hierarchy solution: lower the cutoff.

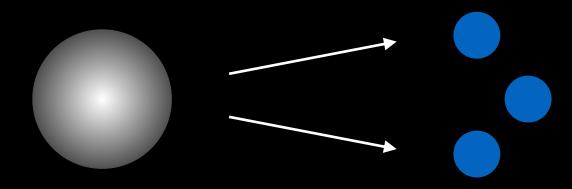
Alternately: consistency with gravity orchestrates correlations among UV parameters to satisfy bounds, changing the distribution.

The Weak Gravity Conjecture

(Electric) weak gravity conjecture: an charge q and mass m satisfying

abelian gauge theory must contain a state of

[Arkani-Hamed, Motl, Nicolis, Vafa '07]



Then BH satisfies

Extremal BH (Z=1) stable unless there exists a state with z > 1

 $\Rightarrow q > m/M_{Pl}$ to avoid stable black holes, remnants, in conflict w/ holography

$$g > \frac{m}{M_{\rm Pl}}$$

"Justification": consider BH of charge Q, mass M decaying to this particle

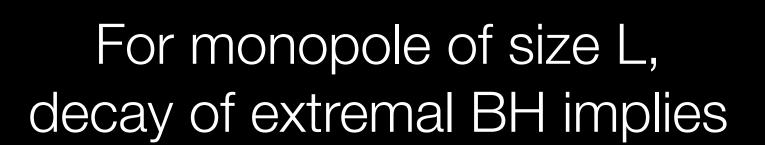
particles produced = Q/q

Energy conservation: mQ/q < M

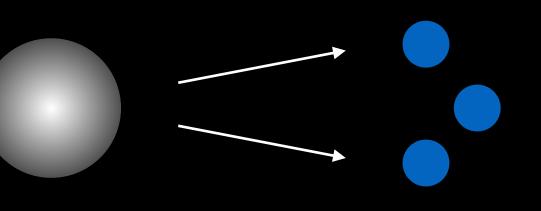
$Z = Q M_{Pl}/M < z = q M_{Pl}/m$

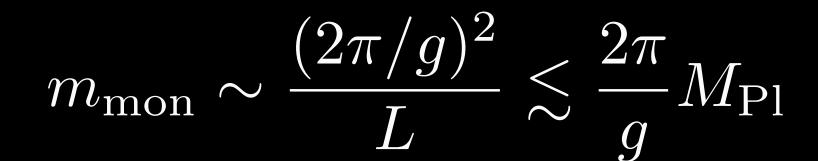
Magnetic Weak Gravity Conjecture

Analogous argument for BHs carrying magnetic charge:



Note: cutoff need not imply appearance of quantum gravity, only physics underlying monopole structure





 $\Rightarrow \Lambda \equiv L^{-1} \leq g M_{\rm Pl}$

A Family of Conjectures

Electric WGC:

[Arkani-Hamed, Motl, Nicolis, Vafa '07]

Magnetic WGC: [Arkani-Hamed, Motl, Nicolis, Vafa '07]

+Scalar WGC:

dS WGC: [Montero, Van Riet, Venken '19]

Axion WGC:

[Arkani-Hamed, Motl, Nicolis, Vafa '07]

New hierarchies from EFT + gravity.

- $m \leq (gq)M_{\rm Pl}$

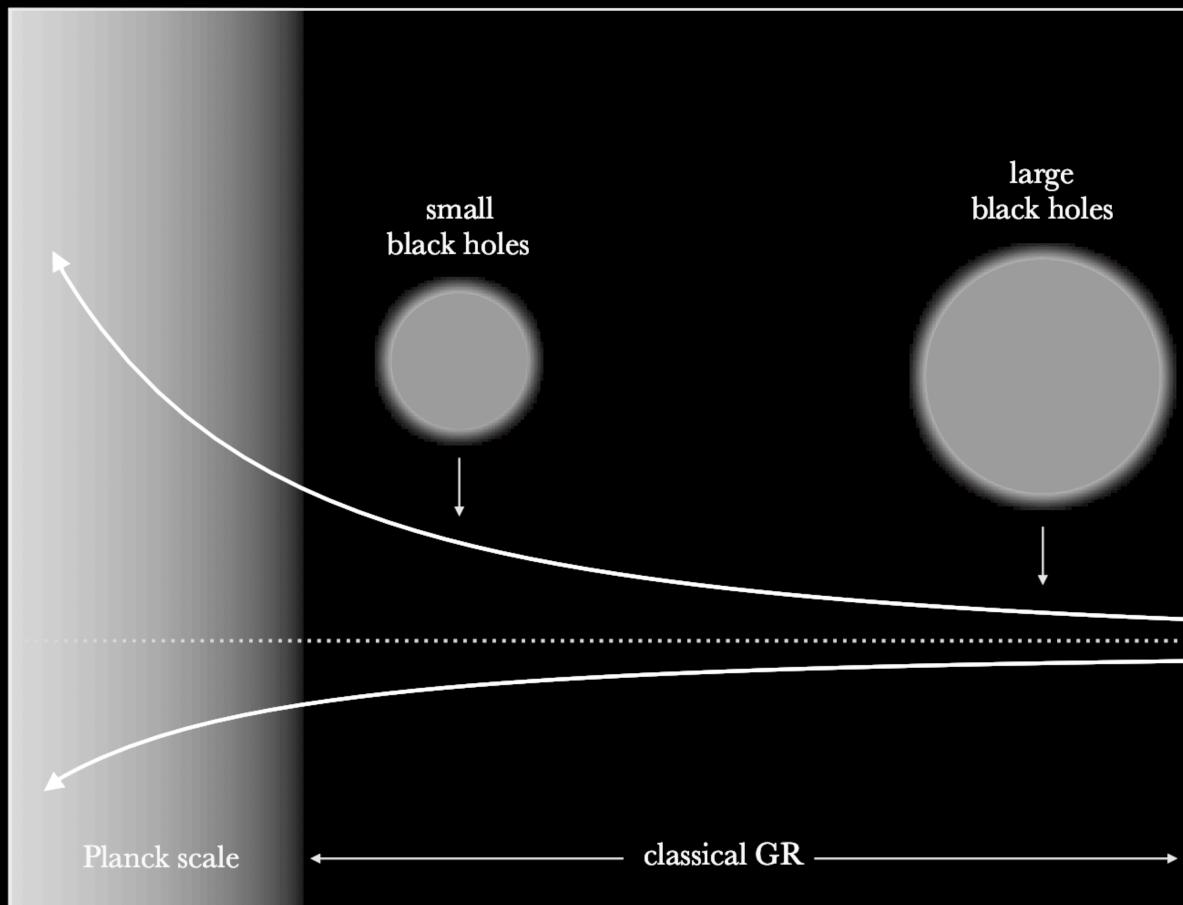
 $\Lambda \lesssim g M_{\rm Pl}$

- [Palti '17]

- $m \leq \sqrt{g^2 q^2 \mu^2 M_{\rm Pl}}$
- $m^2 \gtrsim gq M_{\rm Pl} H$
- $f \leq (1/S)M_{\rm Pl}$

rrelevance to BSM?

[Cheung, Liu, Remmen '18]



m

large black holes

Higher-dimensional operators deform extremality curve in direction that allows larger extremal black holes to decay into smaller extremal black holes, "self-satisfying" WGC.

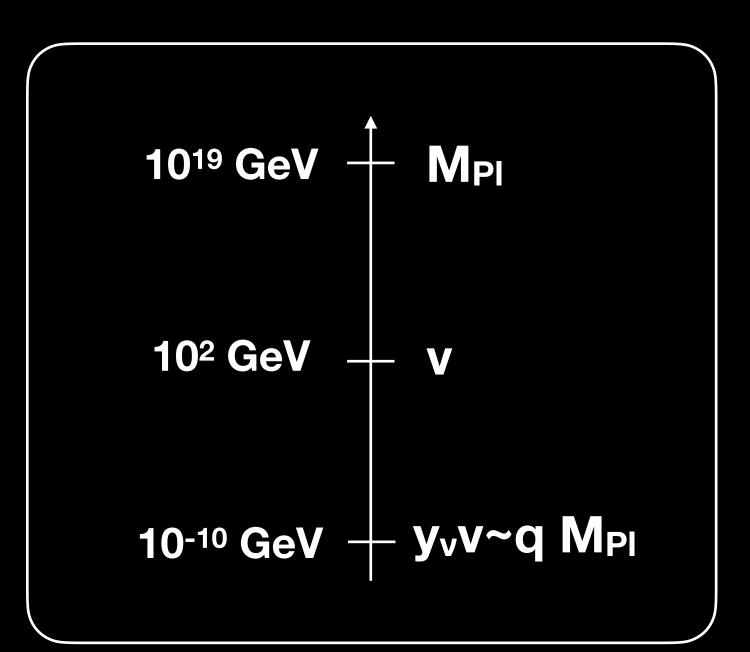
Could still expect arguments to hold for sub-Planckian states, in which case WGC still relevant to particle physics; status unclear.

[Cheung, Remmen '14]: If mass of WGC particle is UV sensitive, then for fixed UVinsensitive parameters, satisfying the WGC enforces fine-tuning. (Or: would orchestrate correlations among UV contributions)

Application to SM: charge SM fermions under weakly gauged (unbroken) $U(1)_{B-L}$ (bounds currently $q \leq 10^{-24}$). Cancel anomalies with RHN v_R

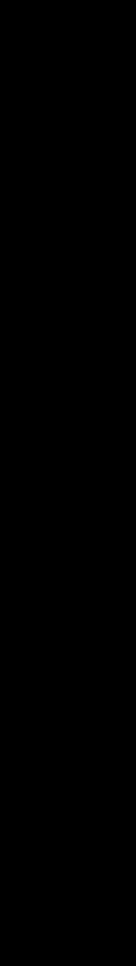
Neutrino mass from EWSB $y_{\nu}HL\nu_R \rightarrow m_{\nu} \sim y_{\nu}v$

For fixed y, q, satisfying WGC places an upper bound on v



If lightest neutrino is WGC particle, m_v ~ 0.1 eV, q≥10⁻²⁹

See also: [Ibañez, Martin-Lozano, Valenzuela '17,...]



- WGC could be satisfied by states outside EFT
- Satisfying WGC could compel the appearance of a new light state that enforces apparent UV correlations (e.g. relaxion)
- Apparently UV-sensitive parameters might control apparently UV-insensitive ones (e.g. emergent gauge fields)

Magnetic WGC implies cutoff of U(1) at $\Lambda \leq q M_{Pl}$

Things that could go wrong:

- Thing that certainly goes wrong:

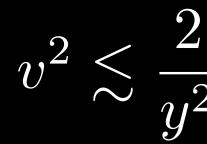
First order of business: can m, Λ be raised to the weak scale?

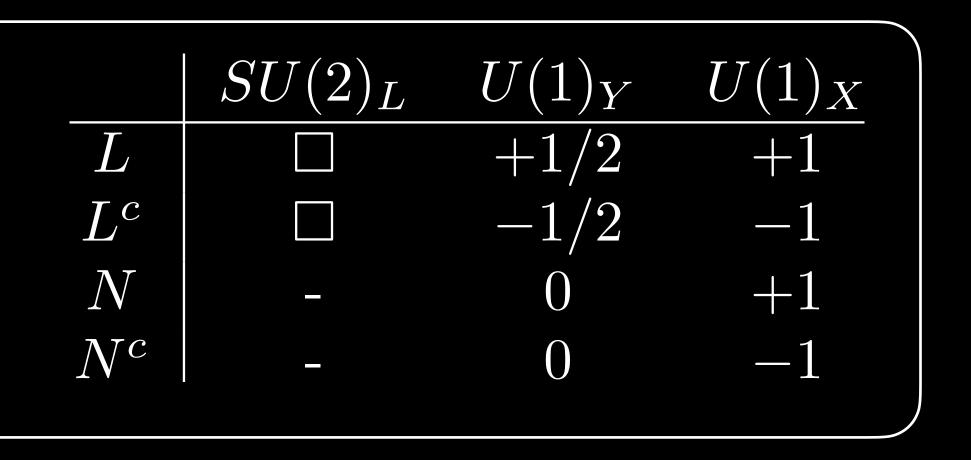
New $U(1)_X$ plus matter acquiring some mass from the Higgs. E.g...

[NC, Garcia Garcia, Koren '19]

Best option: $m_N < m_L$, lightest mass eigenstate χ_1 is WGC particle

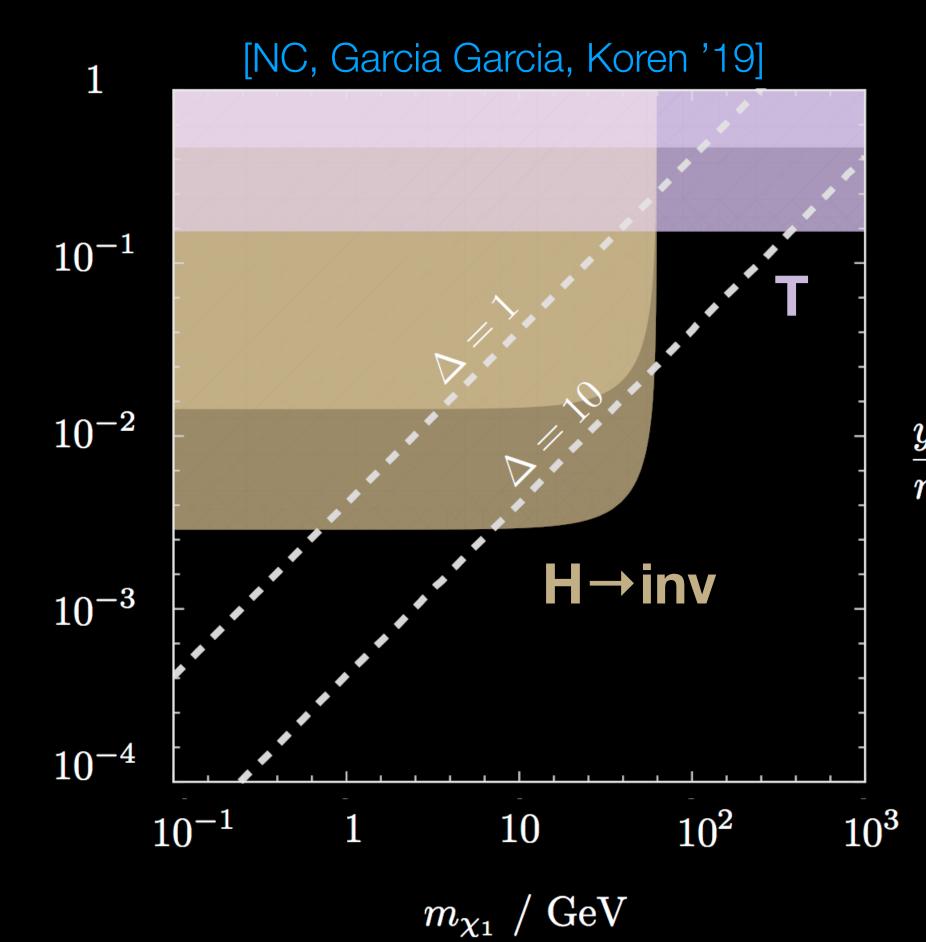
Then for fixed (technically natural) g, m_L, m_N, y,



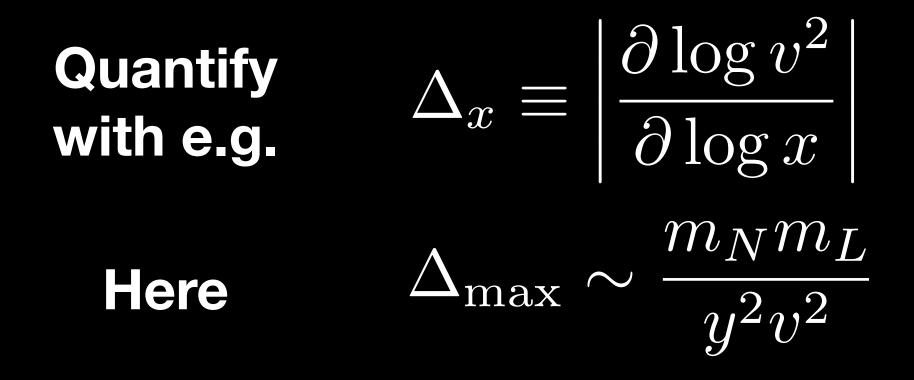


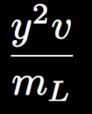
 $-\mathcal{L} \supset \left\{ m_L L L^c + m_N N N^c + y H^{\dagger} L N^c + \tilde{y} H L^c N \right\} + \text{h.c.}$

$$\frac{1}{2}\left(m_{\chi_1}^2 + m_{\chi_1}(m_L - m_N) - m_L m_N\right)$$



Still have a notion of sensitivity of the weak scale to parameters involved in the bound





Not surprising: WGC particle should get "most of" its mass from EWSB.

Surprisingly predictive: look for new singlet fermions coupled to the Higgs at/below the weak scale.

DM story interesting...

Lightest particle charged under $U(1)_X$ is stable \Rightarrow dark matter candidate

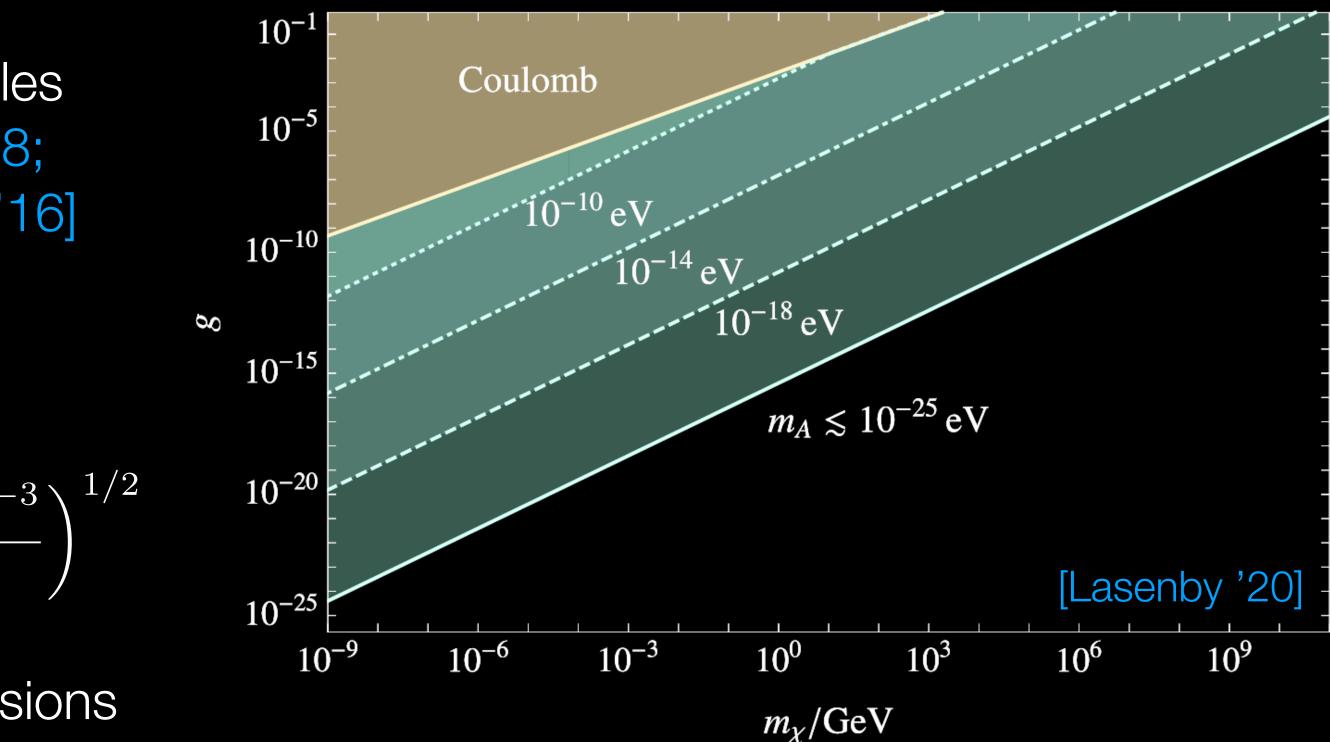
U(1)_X gives a very weak, long-range force, too weak to influence individual collisions but relevant on scale of galaxy clusters

Galaxy cluster collisions can trigger plasma instabilities, making DM collisional on large scales [Ackerman, Buckley, Carroll, Kamionkowski '08; Heikinheimo, Raidal et al '15; Spethmann et al '16]

> Timescale of plasma fluctuations set by plasma frequency,

$$\omega_p = \sqrt{\frac{g^2 \rho}{m^2}} \ge \frac{\sqrt{\rho}}{M_{\rm Pl}} \qquad \omega_p^{-1} \lesssim 10^{15} \,\mathrm{s} \times \left(\frac{0.04 \,\mathrm{GeV \, cm^2}}{\rho}\right)$$

C.f. $\tau \sim 1 \, \text{Gyr} \sim 10^{16} \, \text{s}$ for galaxy cluster collisions



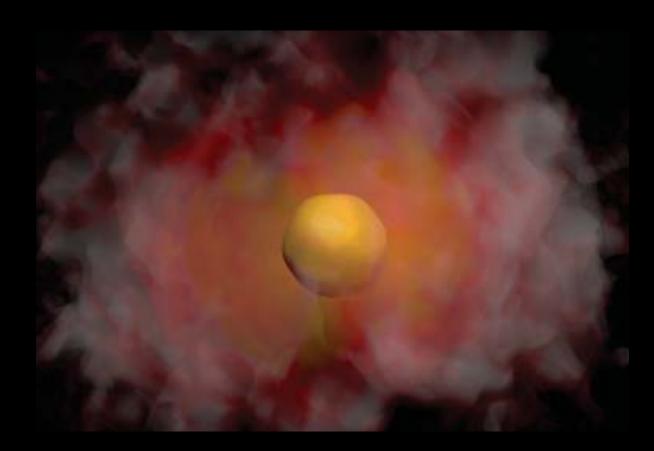
Second order of business: can the magnetic WGC scale be something less severe than the SM cutoff? Only confident that $\Lambda \sim$ scale associated w/ structure of magnetic monopoles

E.g. t' Hooft-Polyakov monopoles SU

" Λ " = $m_W = g_2 f = 2gf \leq 2gM_{\rm Pl}$

W's would trivialize bound from vanilla electric WGC, but not e.g. unit charge version (charge ± 2 under U(1)_x)

$$V(2)_X \mathop{
ightarrow}_{\langle {
m Adj}
angle} U(1)_X$$



Resolution of physics at $\Lambda \sim$ weak scale implies additional exotic physics coupling directly or indirectly to the Higgs.

The de Sitter Conjecture

de Sitter Conjecture: Low energy effective scalar potential in any consistent theory of quantum gravity must satisfy

[Obied, Ooguri, Spodyneiko, Vafa '18]

Arguments from string theory examples, entropy of dS + distance conjecture Challenging for inflation, suggests dark energy should be dynamical

$$\frac{|\nabla V|}{V} \simeq \frac{|\nabla V_{\text{quintessence}}|}{\Delta V_H} \sim \frac{10^{-120}}{10^{-65}} \frac{1}{M_{\text{Pl}}} \sim \frac{10^{-55}}{M_{\text{Pl}}}$$

Either quintessence field couples nontrivially to Higgs, or conjecture too strong...



[Denef, Hebecker, Wrase '18]: Badly violated by SM Higgs boson + dynamical DE

The Refined dS Conjecture

Refined dS Conjecture: Low

effecive scalar potential in any c theory of quantum gravity mus either the dS Conjecture

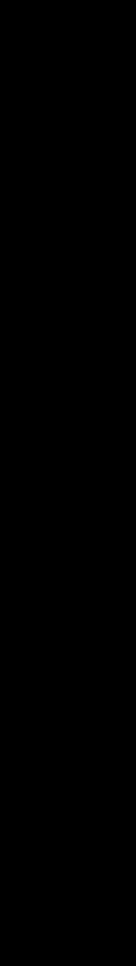
Consistent with SM Higgs,

New wrinkle: essentially forbids metastable vacua, since satisfying RdSC in metastable minimum implies quintessence evolving too rapidly in our vacuum, reaching deep AdS state ("big crunch") within ~fraction of a Hubble time

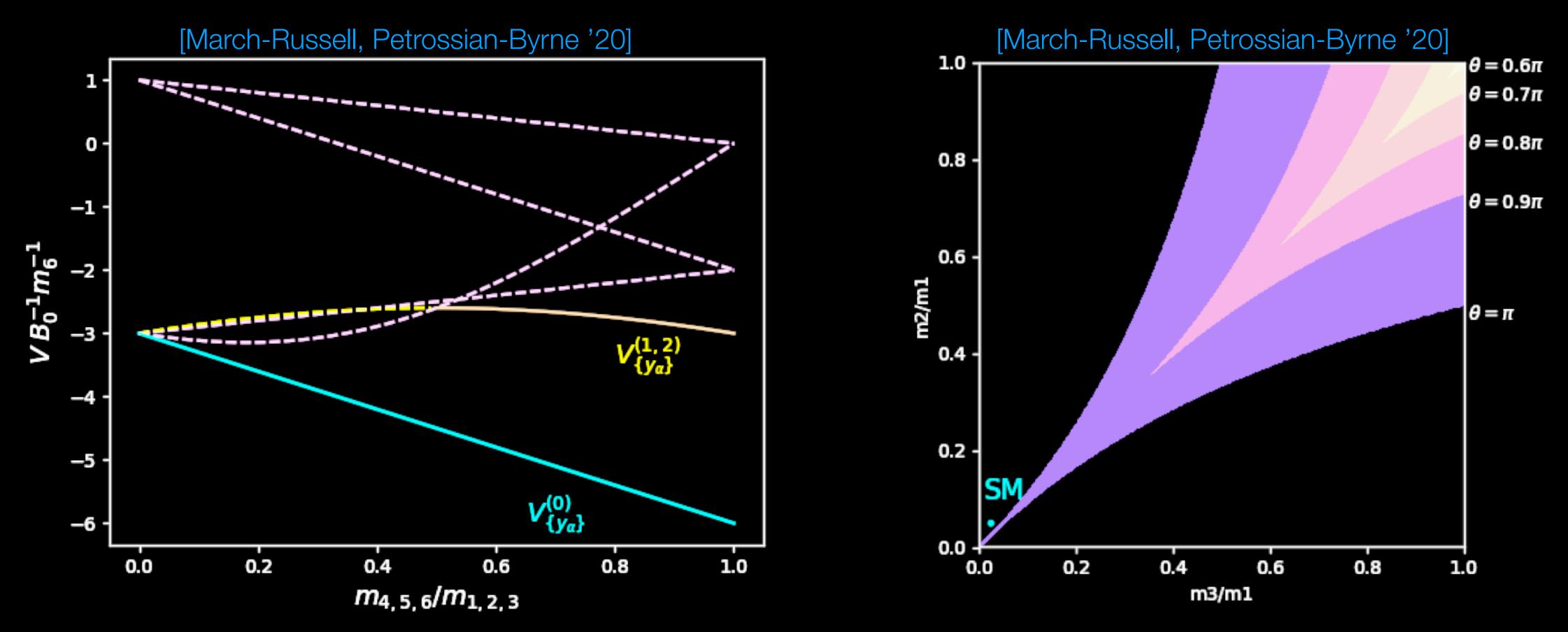
$$\begin{array}{ll} \text{venergy} \\ \text{onsistent} \\ \text{t satisfy} \\ \text{or} \end{array} & \begin{array}{l} \min(\nabla_i \nabla_j V) \leq -c' \frac{V}{M_{\mathrm{Pl}}^2} \\ \begin{array}{l} 0 \\ 0 \end{array} \\ \end{array} \\ \begin{array}{l} \text{[Ooguri, Palti, Shiu, Vafa '18]} \end{array} \end{array}$$

$$\frac{\min(\nabla_i \nabla_j V)}{V} \sim -\frac{10^{35}}{M_{\rm Pl}^2}$$

See also: [Ibañez, Martin-Lozano, Valenzuela '17,...]



The SM and the RdSC



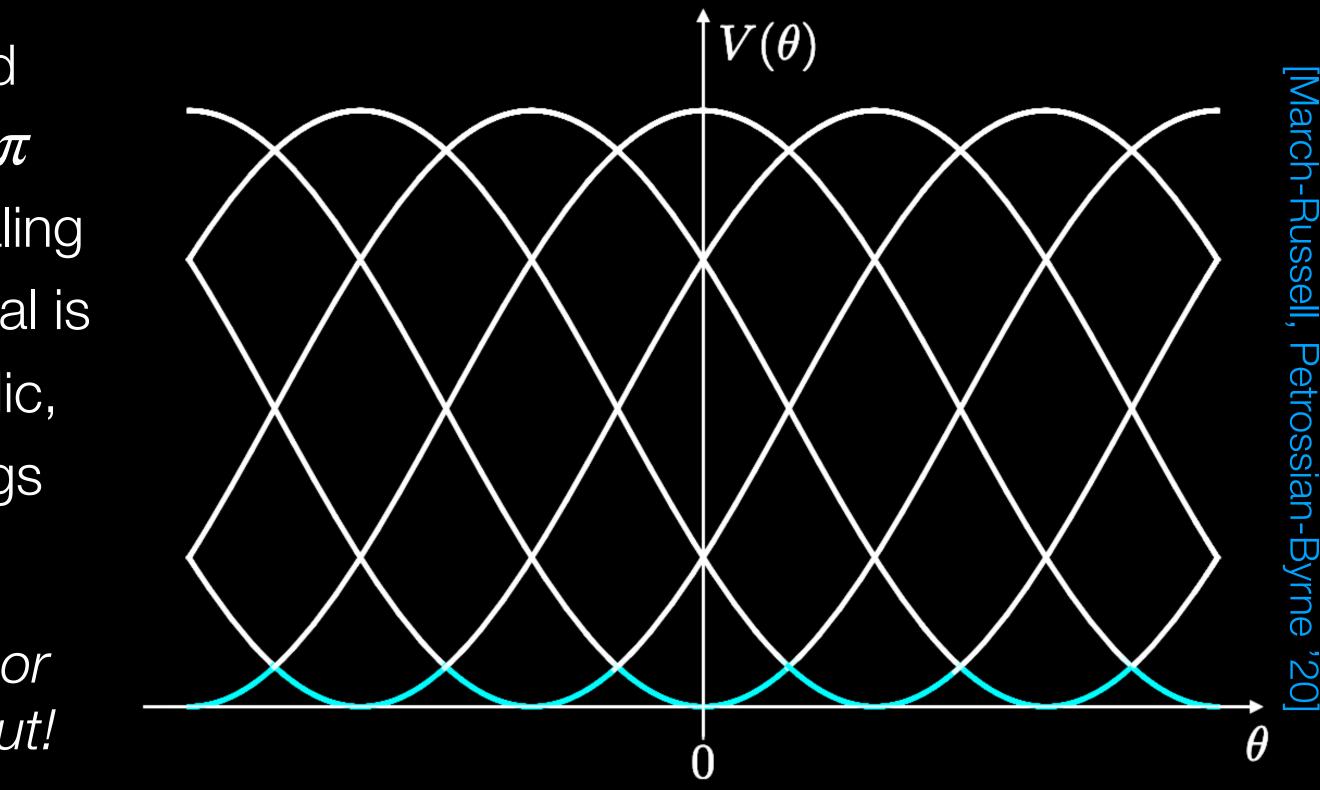
[March-Russell, Petrossian-Byrne '20] RdSC relevant to SM flavor! QCD has metastable vacua for certain values of yukawas, θ ; measured values consistent w/ RdSC

The Weak Scale & the RdSC

[March-Russell, Petrossian-Byrne '20] If pure SU(3) YM has metastable vacua, then for fixed yukawas the RdSC would bound the weak scale from above, $v \leq 50$ TeV.

[Witten '98, Shifman '98] Large-N YM should possess N-1 metastable vacua: reconciles 2π periodicity $V(\theta + 2\pi) = V(\theta)$ and large-N scaling $V(\theta) = N^2 f(\theta/N)$ for 2π periodic f(x): potential is multi-branched, with each branch $2\pi N$ periodic, ground state 2π periodic due to level crossings

Status at small N unknown, no lattice studies or reliable semiclassics, would be lovely to find out!



Not sure (personally) how much more there is to say about Higgs & hierarchy problem following traditional EFT logic.

• If the naturalness strategy fails, it should be for good reason. Failure of EFT logic due to gravitational effects is compelling, albeit hard to quantify. Swampland conjectures provide tools applicable to low-energy theory.

Surprising relevance to Higgs, SM, and its extensions, and visa versa.

 Applications to hierarchy problem convolve conjecture & supposition (e.g. applying WGC to EFT landscape at fixed dimensionless couplings), but follow familiar reasoning (anthropics) and have testable consequences.

Early days, but stand to learn much more about BSM from Swampland Conjectures and visa versa...

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Thank you!

