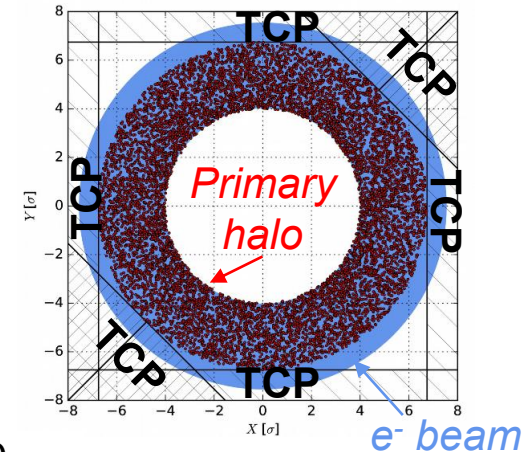
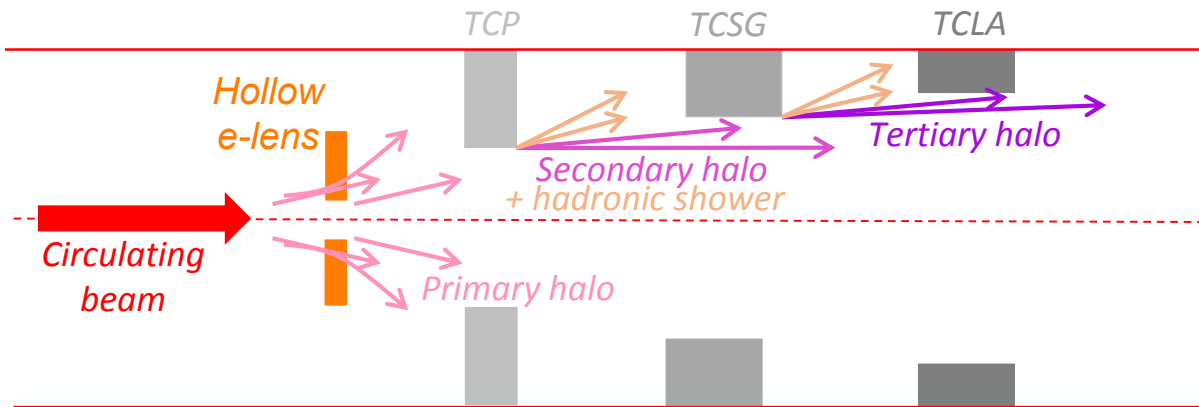


Preliminary study of the electron beam impedance

N. Mounet

Acknowledgements: D. Mirarchi, C. Zannini.

- Since 2020, the hollow electron lens is in the HL-LHC baseline:



From **D. Mirarchi**, WP2 15/09/2020

- Regarding **impedance & stability**, several potential impacts:
 - ❑ **Depletion of transverse distribution tails**, reducing Landau damping
→ taken into account in all stability predictions (tails cut at **3.2σ**)
 - ❑ **Impedance of the physical device**
→ see previous talk by **C. Zannini** and **B. Salvant**
 - ❑ But what about the **impedance of the electron beam** itself?

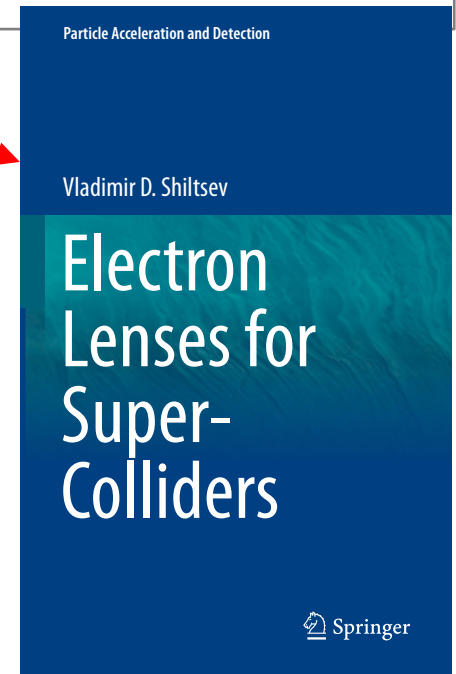
- The issue was studied by **A. Burov** et al in 1999:

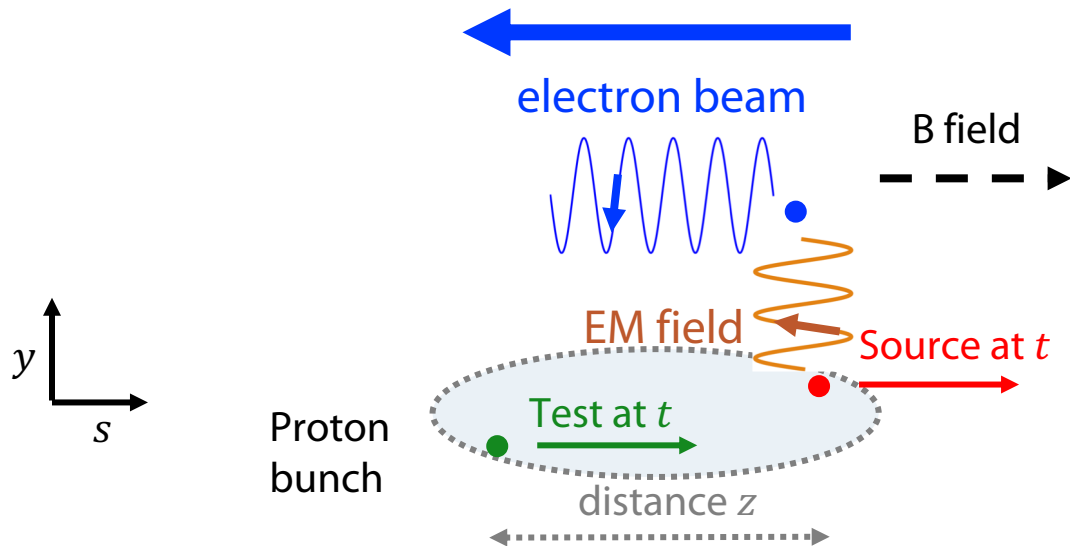
PHYSICAL REVIEW E	VOLUME 59, NUMBER 3	MARCH 1999
<p>Transverse beam stability with an “electron lens”</p> <p>A. Burov,[*] V. Danilov,[†] and V. Shiltsev <i>Fermi National Accelerator Laboratory, Batavia, Illinois 60510</i> (Received 29 July 1998)</p>		

- This study is partly reproduced in **V. Shiltsev's book**
- ... and used as reference in the conceptual design of hollow lens, by **G. Stancari** et al (**CERN-ACC-2014-0248**):

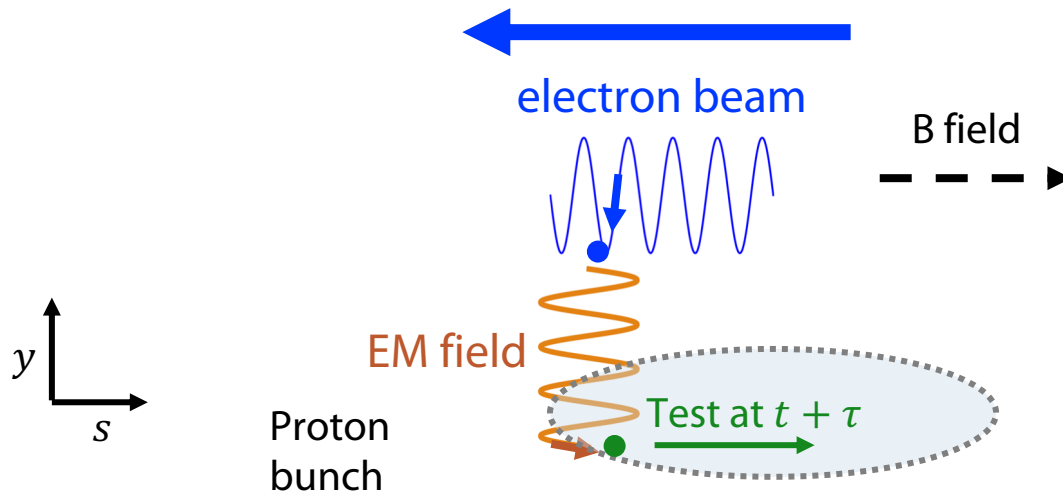
“An early concern on the use of electron lenses [...] was the **stability of the beams**. [...] In particular, a displaced head of the circulating bunch could distort the electron beam, whose electromagnetic fields could in turn act back on the bunch tail, causing oscillations in the electron trajectory and a **fast transverse mode coupling instability** [...].

The electron beam is made stiff **by increasing the axial solenoidal field, reducing its effective impedance** [...].”





- If a proton passes by with an offset, it **kicks the electrons** which are at the same longitudinal position.
- The electrons start **spiraling** under the action of the **solenoid** field.



- If a proton passes by with an offset, it **kicks the electrons** which are at the same longitudinal position.
- The electrons start **spiraling** under the action of the **solenoid** field.
- At a later time, a “test” proton behind the initial one, will in turn see offset electrons and receive a **kick** (in both the x and y directions \rightarrow coupling).

- Assuming both proton and electron beams are **uniform** and of **same radius** a_e , and looking only at the **first order perturbation** of the beams space-charge fields (i.e. keeping everything **linear**),
- **electrons** are **kicked** by a **displaced proton** slice, then move transversely under the sole effect of the **solenoid field** (**Larmor oscillations**), and finally **kick** the protons behind, at a later time.

⇒ we get **wake functions**: for a source proton displaced by Δx and Δy

Coupled terms

$$W_x = W \sin(kz) \Delta x - W(1 - \cos(kz)) \Delta y$$

$$W_y = W(1 - \cos(kz)) \Delta x + W \sin(kz) \Delta y$$

See **A. Burov** et al, PRE 59, 3 (1999) (converted to SI units and using PyHEADTAIL sign convention)

with
$$W = - \left(\frac{1}{4\pi\epsilon_0} \right)^2 \frac{4L_e I_e}{a_e^4 \beta_e c^2 B} \frac{(1 + \beta_p \beta_e)^2}{\beta_p + \beta_e},$$

$$k = \frac{\omega_L}{(\beta_p + \beta_e)c}, \quad \omega_L = \frac{eB}{m_e \gamma_e}$$

Absent from original formula

$$L_e = 3 \text{ m},$$

$$I_e = 5 \text{ A},$$

$$\gamma_e = (1 - \beta_e^2)^{-1/2} = 1.029 \text{ (15 keV } e^-),$$

$$\beta_p \approx 1 \text{ (top energy)}$$

$$B = 5 \text{ T}$$

$$a_e = 2.9 \text{ mm (9.4 } \sigma \text{ with 2.5 } \mu\text{m emittance).}$$

(Parameters from **D. Mirarchi**)

- In **absence of chromaticity** and of **any other kind of impedance**, a simple formula can be found for the **transverse mode coupling instability** threshold, in terms of **B** field:

$$B_{thr} \approx 39 \frac{eN_p \sqrt{\xi_x \xi_y}}{a_e^2 \sqrt{Q_s |Q_x - Q_y|}},$$

with the figures of merit of the e-lens defined as

$$\xi_x = \frac{\beta_x L_e r_p I_e (1 + \beta_e)}{2\pi a_e^2 \gamma_p e \beta_{ec}}, \quad \xi_y = \frac{\beta_y L_e r_p I_e (1 + \beta_e)}{2\pi a_e^2 \gamma_p e \beta_{ec}}$$

- In HL, one gets stability as soon as

$$B > B_{thr} \approx 0.07 \text{ T}$$

Even multiplying by the extra factor found $(1 + \beta_p \beta_e)^2$ (see previous slide), one gets $B_{thr} \approx 0.1 \text{ T}$.

⇒ It seems very far from the **nominal $B = 5 \text{ T}$** .

From **A. Burov** et al, PRE 59, 3 (1999) (converted to SI units)

$$\beta_x = \beta_y = 280 \text{ m},$$

$$Q_s = 2.1 \times 10^{-3},$$

$$|Q_x - Q_y| = 0.01,$$

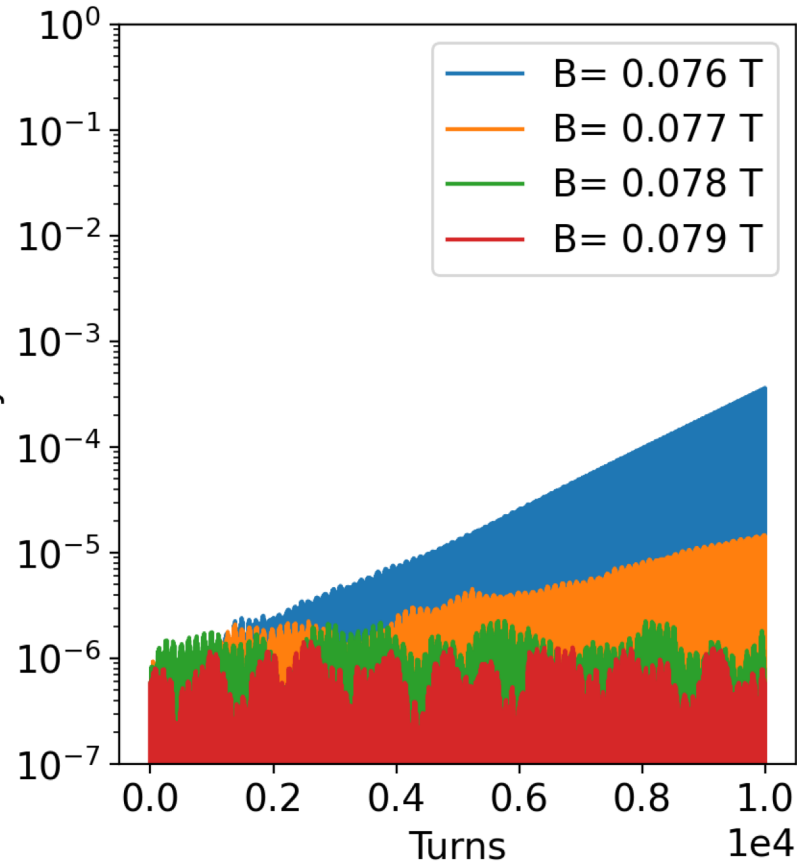
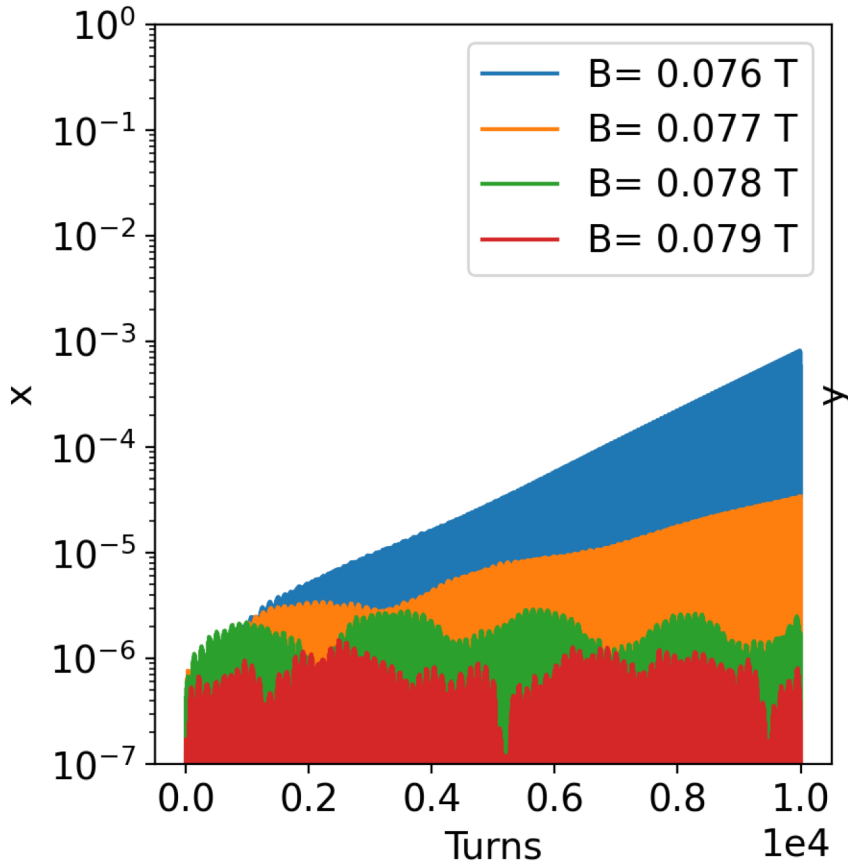
$$N_p = 2.3 \times 10^{11} \text{ p+ / bunch},$$

$$r_p = 1.535 \times 10^{-18} \text{ m},$$

$$\gamma_p = (1 - \beta_p^2)^{-\frac{1}{2}} = 7460.52,$$

$$\xi_x = \xi_y = 1.8 \times 10^{-3}.$$

➤ Checking the TMCI threshold with PyHEADTAIL:

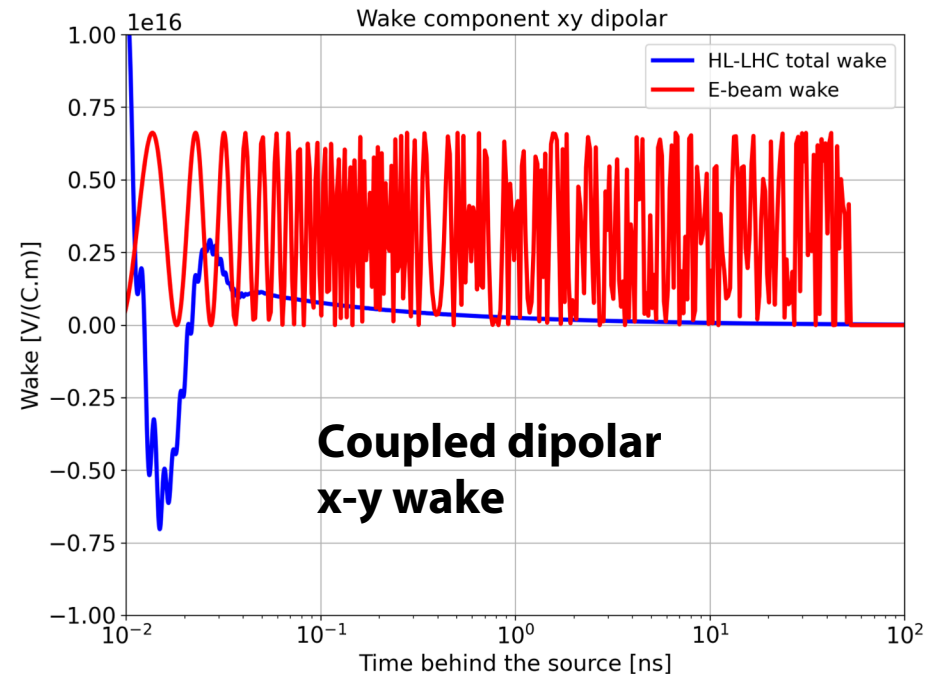
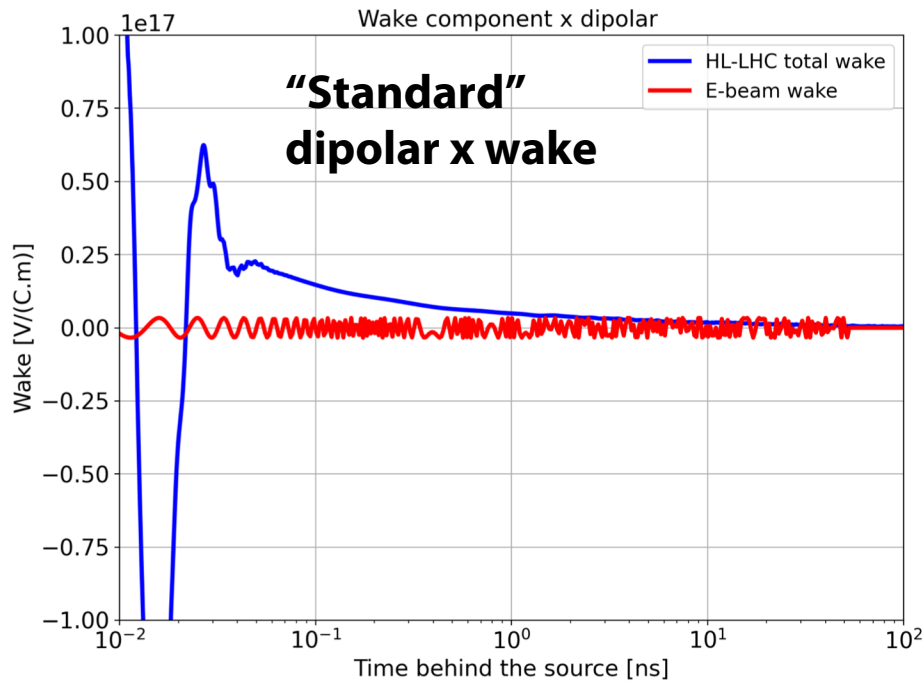


⇒ one gets $B_{thr} \approx 0.077$ T.

Many thanks to **Carlo Zannini** for providing an initial HL-LHC PyHEADTAIL script.

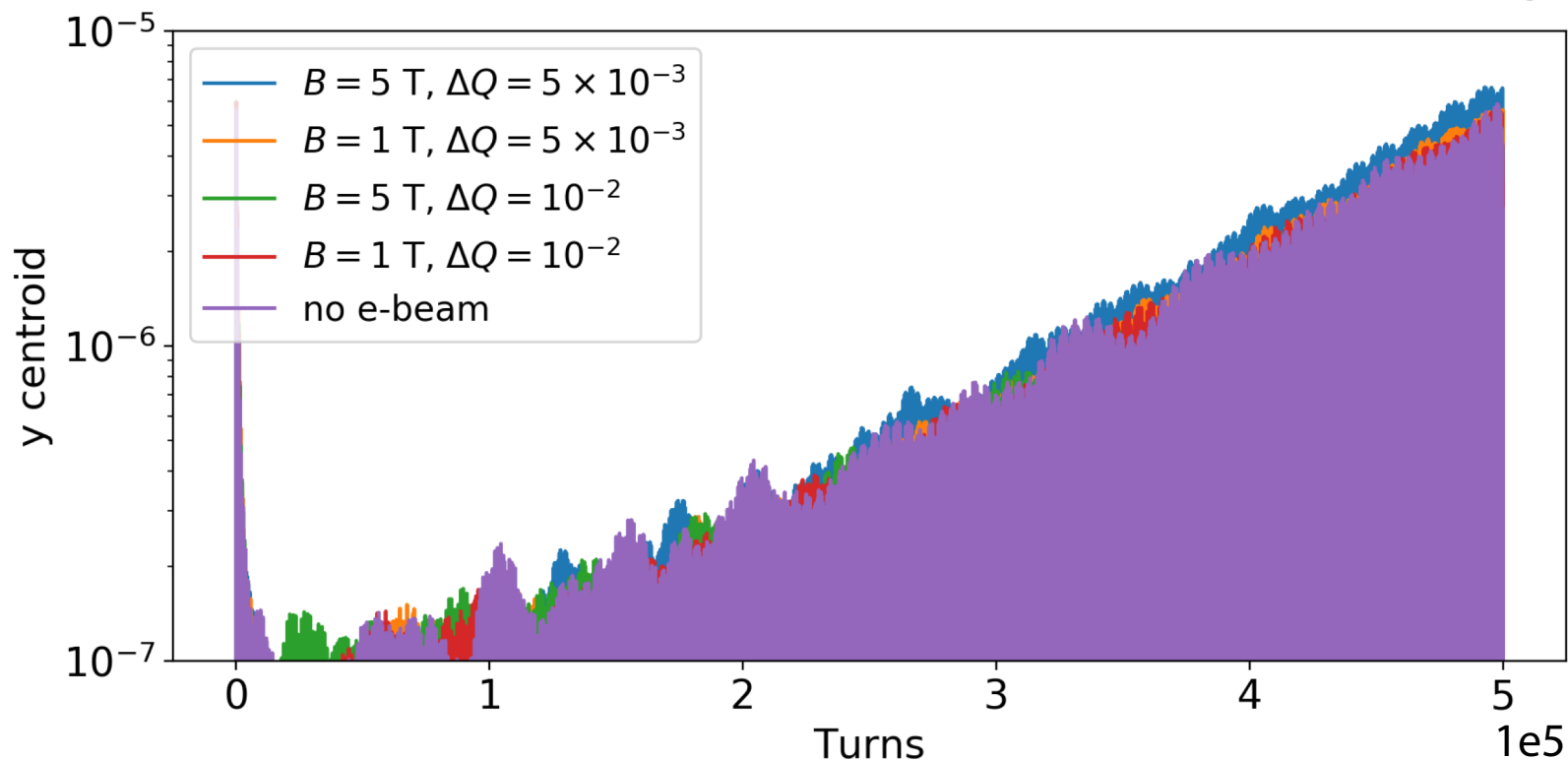
- Some **assumptions** of the model are **not true** in the case of the HL-LHC electron lens: the electron beam is hollow (so clearly not uniform), and the proton beam is Gaussian and typically much smaller.
- The wake extends up to **~ 52 ns** → potential **multibunch** effects.
- Only TMCI was checked, at **zero chromaticity**.
 - potential **weak headtail instabilities** at higher chromaticities ($Q' \sim 15$),
 - more generally, one should add the **electron beam impedance to the full model** and check the impact.

- Comparing the **total wake of the HL-LHC model** (latest update, retracted collimators, see WP2 28/07/2020) with the **electron-beam wake model**:



- ⇒ significantly smaller than the dipolar total wake (x or y) within the bunch,
- ⇒ but coupled terms much stronger than the rest of the wake.

- In a standard operational configuration ($Q'=15$, 100 turns damper), adding the wake from the electron beam on top of the full HL-LHC model (**both dipolar and coupled** terms):

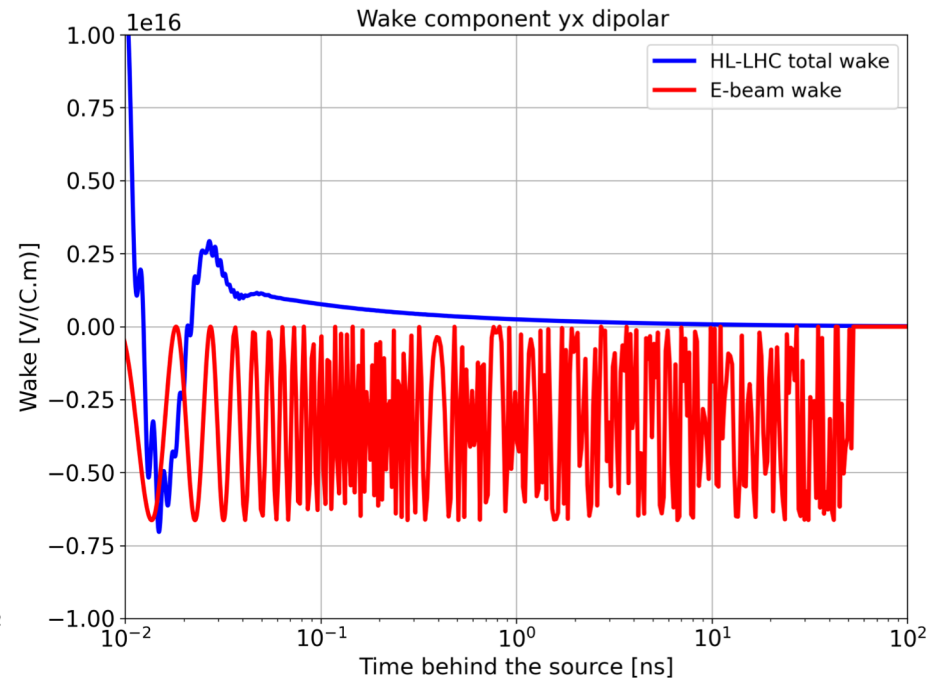
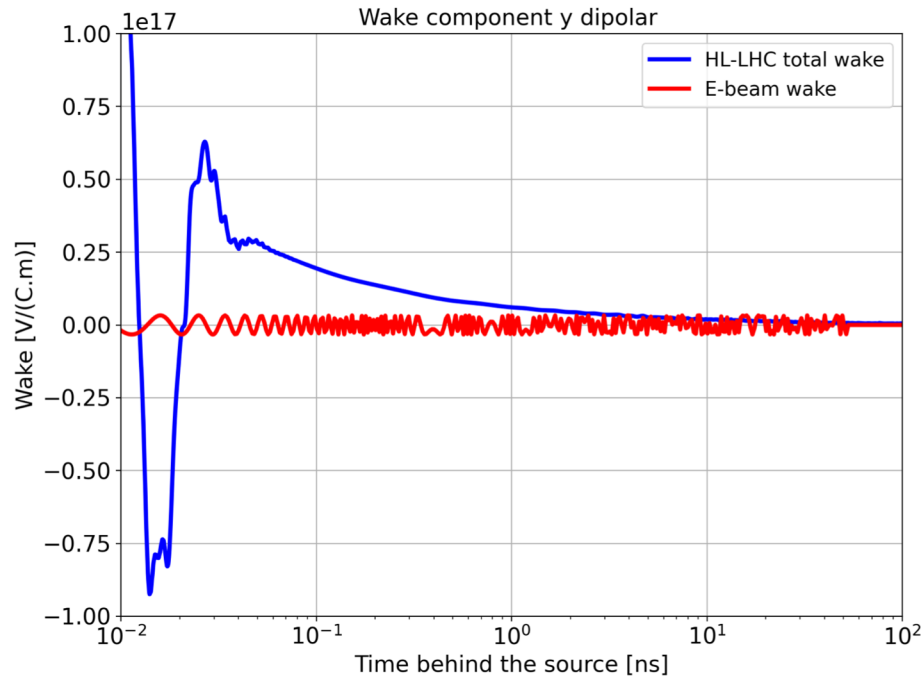


⇒ **no strong impact** on the most unstable plane (y), even at a lower B field and at a lower x-y tune difference (rise times stay within $\sim 10\%$).

- The electron beam of an electron lens can trigger coherent instabilities, in particular **TMCI-like**.
- These can be strongly mitigated with the solenoid field.
→ In the case of HL-LHC, as low as **B=0.1 T** should be enough to avoid TMCI **coming from the e-beam alone**.
- The simple wake model from A. Burov et al, shows nevertheless that **coupled terms in the e-beam wake** are **larger than those of the full HL-LHC model**.
- Still, **single-bunch instabilities** in a standard operational configuration ($Q'=15$, damper 100 turns) **seem not to be strongly affected**.
- Several **shortcomings** of the model need to be addressed:
 - ❑ the respective size and transverse shape of the protons and electron beams have to be modelled better,
 - ❑ potential **multibunch** effects to be looked at.

Appendix

- Comparing the **total wake of the HL-LHC model** (latest update, retracted collimators, see WP2 28/07/2020) with the **electron-beam wake model**:



- ⇒ significantly smaller than the dipolar total wake (x or y) within the bunch,
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