Top, flavour, and electroweak at a linear collider

A. Irles
IFIC (CSIC/UV)







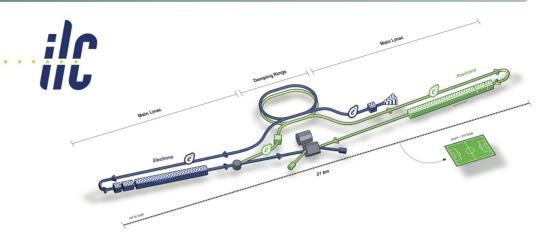


Thanks to J. Fuster, V. Miralles, R. Poeschl, F. Richard, A. Robson, M. Vos, the ILD conveners group, ... for the help preparing these slides

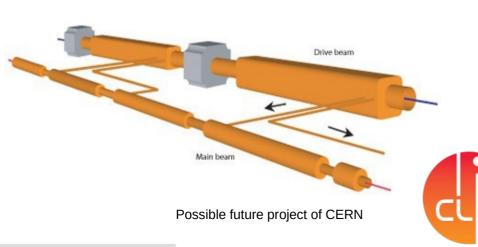
Apologies in advanced for the many interesting and cuttingedge studies not shown in the following slides.

Linear Colliders

- Energy: 0.1 1 TeV
- Electron and positron polarisation
- ► TDR in 2013
 - + DBD for detectors
- Footprint 31 km
- ► Initial Energy 250 GeV Footprint ~20km



Under discussion in Japanese Government and international community



- ► Energy: 0.4 3 TeV
- ► CDR in 2012
 - Project Implementation Plan in 2018
- Electron polarisation
- Footprint 50km
- ► Initial Energy 380 GeV Footprint ~11km

Higgs factories but also...

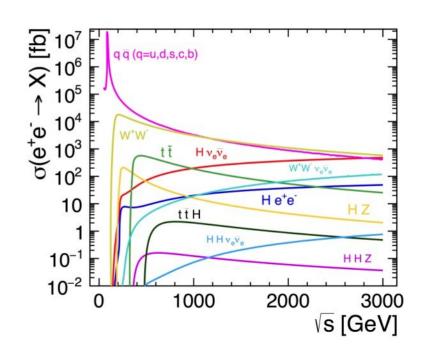


- ► All Standard Model particles within reach of planned LC projects
- ► High precision tests of Standard Model over wide range to detect onset of New Physics
- ► Machine settings can be "tailored" for specific processes
 - Centre-of-Mass energy & Beams polarisation (straightforward at linear colliders)

- ► Higgs factories but also...
 - "light" qq factory

(and Z-factory at Z-pole)

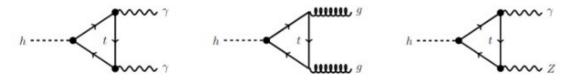
- WW factory
- Top-quark factory
- ttH facotry
- •



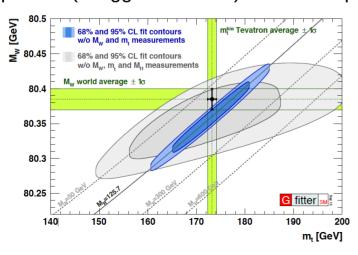
Higgs and top-quark

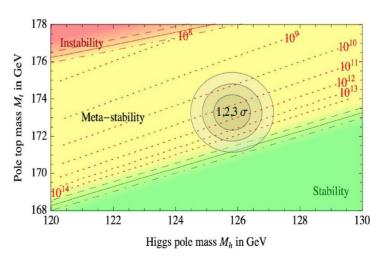


- "Recently" discovered particles still under scrutiny
 - The heaviest fundamental particles!
 - New physics by compositeness? Higgs and top composite objects?
- \blacktriangleright An enigmatic couple: the top quark has O(1) Yukawa coupling and rules the loop diagrams (gg \rightarrow H, H \rightarrow gg)



Top mass (&Higgs &W mass) are crucial parameters of the SM (and BSM)

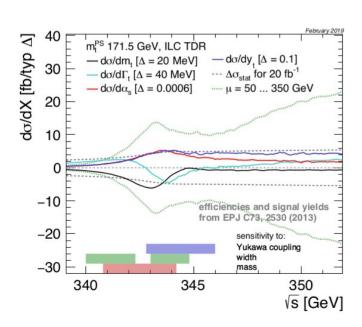




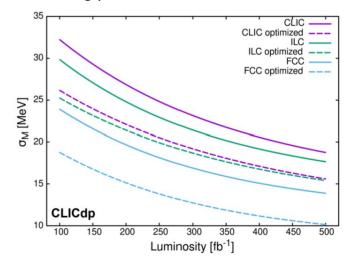
Top-quark mass: a key parameter



- ➤ The top threshold provides excellent sensitivity to the mass and other top quark properties
 - Measurement of the top quark mass in theoretically well-defined mass schemes
- Sensitivity to :
 - Top-quark mass, width, yukawa coupling, strong coupling constant



- ▶ Optimising top-quark threshold scan at CLIC & ILC using genetic algorithm
 - K. Novak et al [PD3]
 - Top-quark mass can be extracted with stat. unc. of the order of 25 MeV already for 100 fb-1
 - Without losing precision on width or Yukawa coupling



See also talks from:

Top-quark: A. Hoang, Harigaya. Pathak [PD1] b-quark mass: M. Vos [PD2] S. Tairafune [PD3/PD4]

...

LC: power of beam polarisation

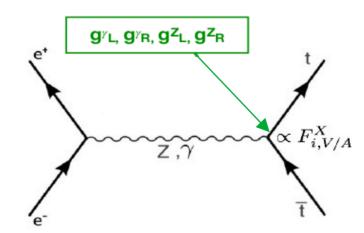


▶ Longitudinally polarised beams are a special feature of Linear e e Colliders:

- SLC: $P(e^{-}) = \pm 80\%$, $P(e^{+}) = 0\%$
- ILC: $P(e^{-}) = \pm 80\%$, $P(e^{+}) = \pm 30\%$ (upgrade 60%)
- CLIC: $P(e^{-}) = \pm 80\%$, $P(e^{+}) = 0\%$

$$P = \frac{N_R - N_L}{N_R + N_L}$$

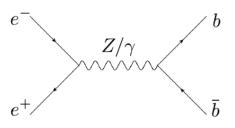
- \triangleright Electroweak interactions are highly sensitive to chirality of fermions: SU(2)L x U(1)
- Cross sections are sensitive to beam polarisation -> background suppression, signal enhancement, control of sytematics and...
- ▶ Beam polarisation allows the probe of the SM/BSM chiral structure
 - SM: Z and y differ in couplings to left- and right-handed fermions
 - BSM: unknown chiral structure!



Two fermions



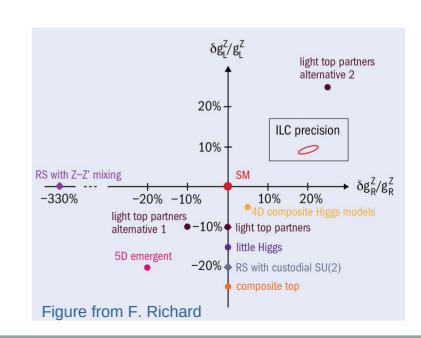
Differential cross section for (relativistic) di-fermion production



$$\frac{d\sigma}{d\cos\theta} (e_L^- e_R^+ \to f \, \overline{f}) = \Sigma_{LL} (1 + \cos\theta)^2 + \Sigma_{LR} (1 - \cos\theta)^2$$

$$\frac{d\sigma}{d\cos\theta} (e_R^- e_L^+ \to f \, \overline{f}) = \Sigma_{RR} (1 + \cos\theta)^2 + \Sigma_{RL} (1 - \cos\theta)^2$$

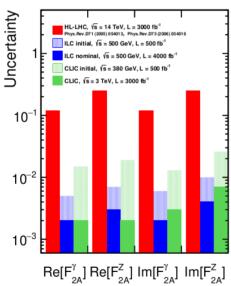
- The helicity amplitudes Σ_{IJ} , contain the couplings g_L/g_R (or Form factors or EFT factors)
- Left/right asymmetries (characteristic for each fermion)
- BSM in these topologies are mainly discussed in terms of new Z' bosons, coming from an extension of the SM gauge group
 - Most of these models modify the top-quark couplings



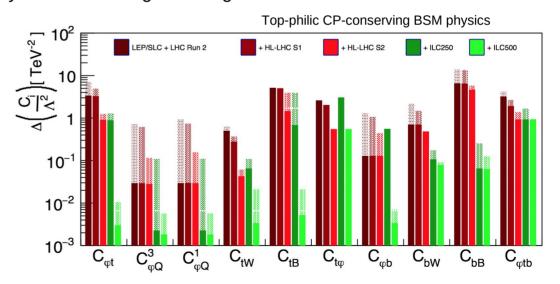
Top-Quark EW couplings



- ▶ Updated global fits with LEP/Tevatron/LHC data [for example JHEP12(2019)098]
 - Including ttZ, tty
- ▶ All studies show the extraordinary impact of adding the data of future LC
 - ILC 250GeV already helps for several operators (b quark observables)
 - Going above the top-quark pair production threshold is crucial for the top-operators
 - The determination of axial form factors highly benefit from higher energies



Top-philic CP-violating BSM physics,

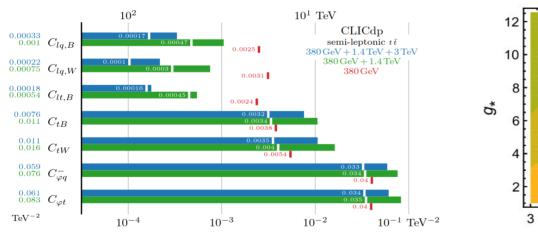


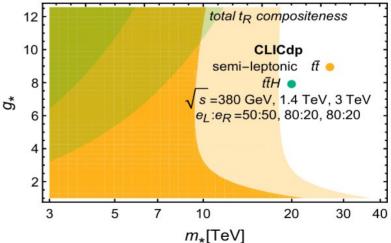
Plot: Top quark EW couplings and EFT fits [Miralles PD2] More by [Gu, Vries, Durieux, You, Najafabadi, Goldouzian... PD2]

Top-Quark EW couplings & compositeness



- ► CLIC studies at higher energies (boosted tops!)
- ▶ Potential to scrutinize the EFT operators with high precision at high energies
- Prospects on top compositeness including the study of ttH production (top Yukawa coupling) & tt global fits.
- Compositeness emerging in connection to the Naturalness Problem can be conclusively probed at a LC operating at high energies





Top-philic CP-conserving BSM physics,

JHEP11(2019)003

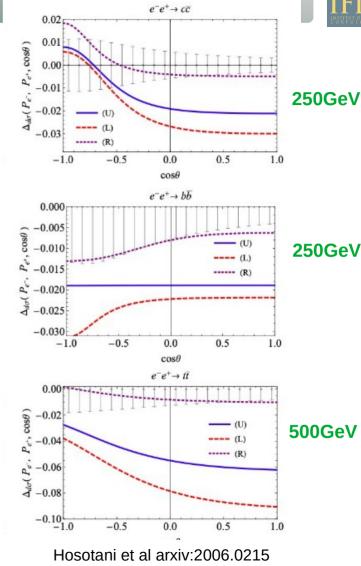
Only the top-quark?

► Randall-Sundrum, Grand Unification Higgs models (GUT):

- model new strong interactions by dynamics in a bounded 5th dimension
- strong-interaction resonances become Kaluza-Klein resonances
- Higgs is the 5th component of a 5D gauge boson

Modification of the EW couplings

- Hosotani et al arxiv:2006.0215
- These effects grow as, s/kR. (kR is the RS scale).
- Small effect at Z-pole (LEP/SLC) but visible already at 250 GeV (with polarized beams)
- Enhanced effects at large energies (500GeV-3TeV) & with beam polarization
- ► Expect particularly large effects for the 3rd generation, but also for ligther fermions



Access to Heavy Resonances at ILC250

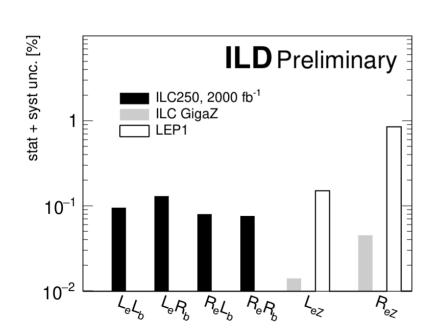


- ► At High Energies we are sensitive to interference effects of Z and photon!!
 - Extracted EW couplings of photon and Z can be influenced by new physics effects
- ► At Z-Pole: we are sensitive to Z/Z' mixing

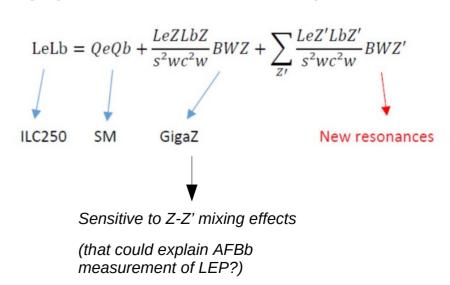
High importance of running at the Z-Pole A. Freitas [PD1], G. Durieux [PD2]

• Sensitivity to vector (and tensor) couplings of the Z (and not "disturbance" from photons)

e+e- → bb Full simulation studies (cc/ss in progress) [A.I. PD3]



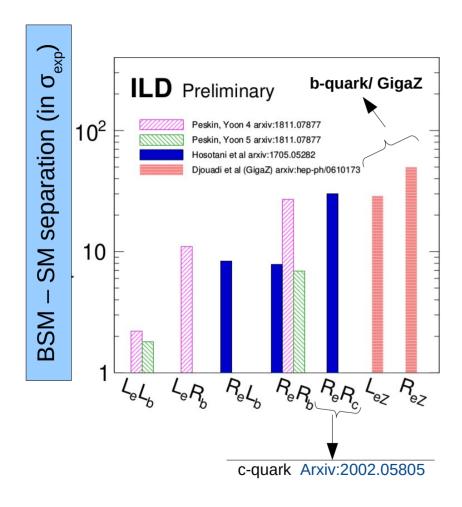
Couplings (notation for new resonances)



Arxiv:1709.04289, PoS(EPS-HEP2019)624

Access to Heavy Resonances at ILC250





▶ BEAM POLARISATION allows to distinguish between different models

Expected number of standard deviations for different **RS/compositeness BSM scenarios** when determining the different EW couplings to c- and b-quark at **ILC250** (with GigaZ input).

- Models that predict multi-TeV Z' resonances
- With or without mixing at Z-pole
- See backup for more details on the models

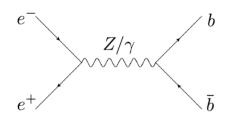
Potential for discovery of new resonances mZ' \sim O(10-20) TeV at ILC250

Arxiv:1709.04289, PoS(EPS-HEP2019)624

(few) Experimental challenges



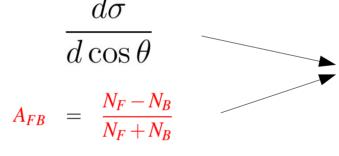
- **Example of observables used to extract the electroweak couplings**
- **cross section, Rq** and forward backward asymmetry **AFB** like observables



$$R_{q}^{0} = \Gamma_{q\bar{q}} / \Gamma_{had}(Z - pole)$$

$$\Rightarrow R_{q}^{cont.} = \sigma_{q\bar{q}} / \sigma_{had}(s > Z - pole)$$

Quark identification. No need to measure an angular distribution, (but possible)



Angular Distribution.

Quark ID + charge measurement (quark – antiquark disentangling)

Gives access to all left/right couplings.

(few) Experimental challenges



top-quark pairs... decay before

hadronizing

lepton ID

tracking

- ► C-quark pairs
- ► High efficient flavour tagging for cquarks expected at future colliders
- Charge measurement
 - **Primary method:** identification of Kaons produced D-meson decays → K-method (requires PID)
 - **Secondary method:** reconstruction of charged mesons → **Vtx-method**

PID is mandatory to reach competitive ac curacies

- s-quark pairs (in progress)
 - Check M. Basso's & Y. Okunawa's talk

- ► B-quark pairs
- ► High efficient flavour tagging for b-quarks expected at future colliders
- Charge Measurement
 - **Primary method:** reconstruction of charged mesons → Vtxmethod
 - **Secondary method:** identification of Kaons produced in b-hadron decays → K-method (requires PID)

positron electron flavor tagging 4 jets, isolated lepton

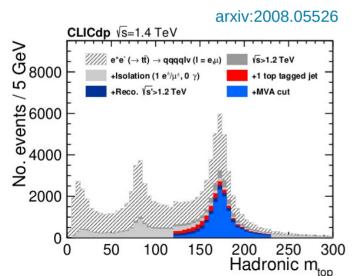
PID is very useful

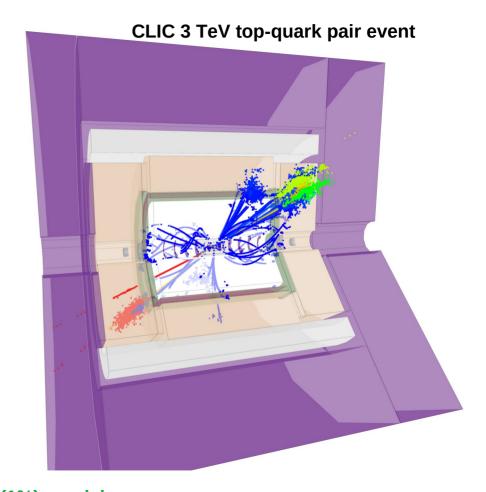
Boosted topologies



- Even Lepton/Linear Colliders have complicated event topologies
- Boosted jets: highly collimated jet environment above 1 TeV
- tops are reconstructed as "fat" jets with a rich substructure
 - Top tagger + MVA:

Maximizing full potential of PFA





Allows O(1%) precision on cross sections and AFB measurements

Flavor tagging and PID



Flavor tagging

Indispensable for any analysis with final state quarks

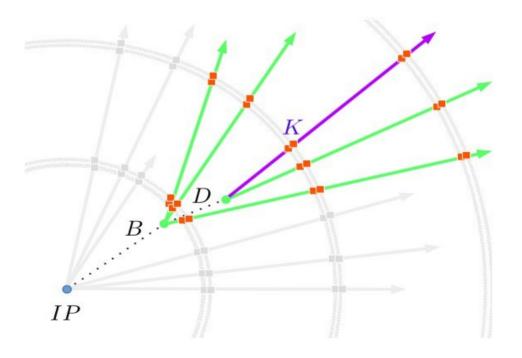
▶ Quark charge measurements

 Important for top-quark studies but Indispensable for ee → bb/cc/ss...

► Charge measurements:

- Vtx charge and Kaon Identification
- High efficiency (double tagging)
- High purity → control of the migrations
- ► Future detectors can base their entire measurements on double Tagging and vertex charge
 - LEP/SLC had to include single tags and semi-leptonic events

Exploit small beam spot of Linear Colliders Vtx detector close to beampipe + low material budget Tracking **efficiency** (>99%)

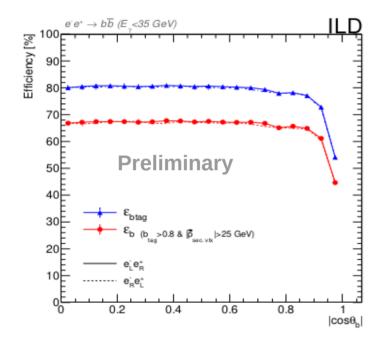


PhD thesis: S. Bilokin

Flavour tagging + PID



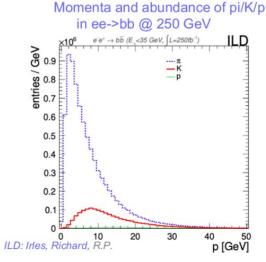
- Flavour tagging capabilities (b-quark example)
 - using LCFI+ tools from ILCsoft

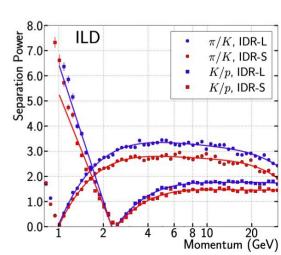


	<i>b</i> -quark		
Experiment	Eff. [%]	Pur. [%]	
DELPHI [19]	47%	86%	
ILD (this note)	80%	98.7%	

ILD note in progress → See S. Taraifune's talk

- Charged Hadron PID needed for a wide range of momentum
 - TPC as the main solution
- TOF as complement for low momentum (and/or forward tracks)
- Special need of K/pion separation





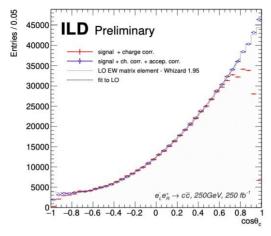
Improvable by clustering counting? [F. Cuna, PD4/PD5]

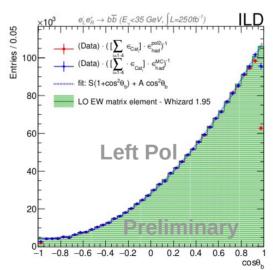
Realistic studies with full simulation

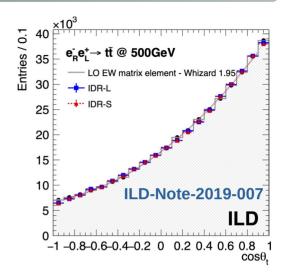


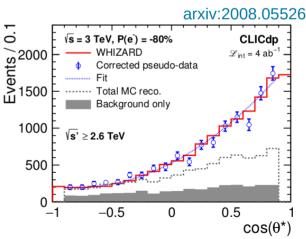
Probing the Chiral structure of the SM and BSM requires:

- high precision predictions & global fits
- ► High precision measurements (at the per mile level in some cases!!)
 - detailed studies with full simulations with realistic detectors
 - Optimization of detectors and reconstruction techniques
- ► EFT predictions will require input from differential distributions









FULL SIMULATION STUDIES

Summary / conclusions



- ► Linear colliders are not only Higgs Factories but EW factories.
 - Two fermion factory, Z factory, WW, ttH, etc
 - upgradability in energy (and also at Z/WW thresholds)
 - Polarization makes possible probing the full chiral structure of the SM / BSM
- ► Top (but also lighter fermions!!): excellent probes for BSM
 - Already at "low energies" of the Higgs Threshold
- Heavy activity improving the predictions to address the reach of future LC
 - EFT formalism, benchmark RS models...
- ► And also heavy activity on optimizing / understanding the limits of the future detectors

- Not covered in this talk...
 - EW at the Zpole
 - Higgs/Top interplay in EFT
 - b/t interplay in EFT
 - Top BSM bellow the threshold (FNCN)
 - etc

Back-up slides



Not only top-quark: LEP/SLC tension

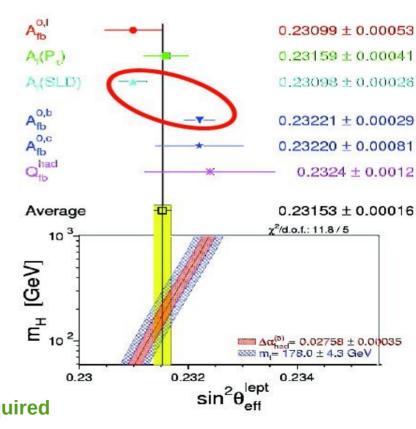


- Current LEP & SLC best $sin^2\theta^l_{eff}$ measurements show tension
 - This measurement is the one with largest tension with the SM fit.
 - Most precise single Individual determination of $sin^2\theta^l_{eff}$ from SLC \rightarrow Left-right asymmetry of leptons
 - Most precise single Individual determination of sin²θ¹
 _{eff}
 from LEP → forward backward assymetry (b-quark)
- ▶ Heavy quark effect, effect on all quarks/fermions, no effect at all?

The **resolution** of this issue requires improving the the measurements precission by an order of magnitude



Per mil level of experimental precision is required

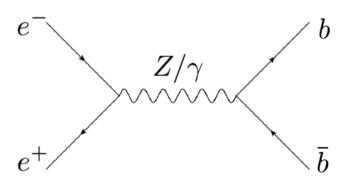


Cross sections



$$\sigma_{e^-e^+ \to q\overline{q}}$$

	Channel	$\sigma_{_{\mathrm{unpol}}}$ [fb]	σ _{.,+} [fb]	σ _{+,-} [fb]
	q=t	572	1564	724
500 GeV	q=b	372	1212	276
	q=u+d+s+c	2208	6032	2793
250 GeV	q=t			
	q=b	1756	5677	1283
	q=c	3020	8518	3565
	q=u+d+s	6750	18407	5463



▶ Beam polarisation also enhances the cross section values

Electroweak Physics also at the Z-pole



$$A_f = \frac{g_{Lf}^2 - g_{Rf}^2}{g_{Lf}^2 + g_{Rf}^2}$$

at an *un*polarised collider:

$$A_{FB}^{f} \equiv \frac{(\sigma_F - \sigma_B)}{(\sigma_F + \sigma_B)} = \frac{3}{4} A_e A_f$$

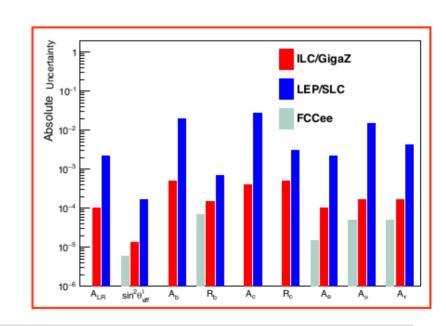
While at a *polarised* collider:

$$A_e = A_{LR} \equiv \frac{\sigma_L - \sigma_R}{(\sigma_L + \sigma_R)}$$

 $A_{FB,LR}^{f} \equiv \frac{(\sigma_F - \sigma_B)_L - (\sigma_F - \sigma_B)_R}{(\sigma_F + \sigma_B)_L + (\sigma_F + \sigma_B)_R} = \frac{3}{4} A_f$

studies by ILD:

- ▶ at least **factor 10**, often ~50 improvement **over LEP/SLC**
- ► Ac nearly 100 x better thanks to:
 - Excellent charm / anti-charm tagging:
 - excellent vertex detector
 - tiny ILC beam spot
 - Kaon-ID via dE/dx in ILD's TPC
- polarization buys:
 - a factor 100 in luminosity,

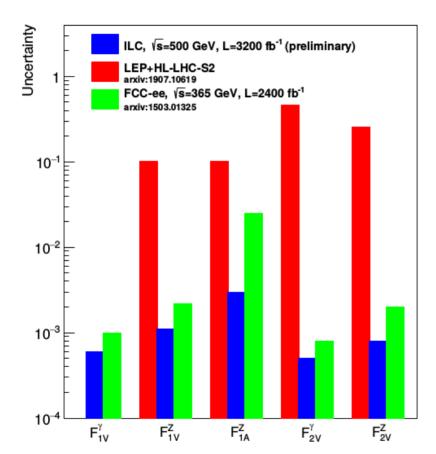


High importance of running at the **Z-Pole** A. Freitas [PD1], G. Durieux [PD2]

Electroweak Physics above the tt threshold



ILD-PHYS-PUB-2019-007



- ▶ e+e- collider way superior to LHC (\sqrt{s} = 14 TeV)
- ► Final state analysis at FCCee (polarisation)
 - Also possible at LC => Redundancy

$$\Gamma_{\mu}^{ttV}(k^{2},\,q,\,\bar{q}) = -ie \left\{ \gamma_{\mu} \left(F_{1V}^{V}(k^{2}) + \gamma_{5} F_{1A}^{V}(k^{2}) \right) + \frac{\sigma_{\mu\nu}}{2m_{t}} \left(q + \bar{q} \right)^{\nu} \left(F_{2V}^{V}(k^{2}) + \gamma_{5} F_{2A}^{V}(k^{2}) \right) \right\}$$

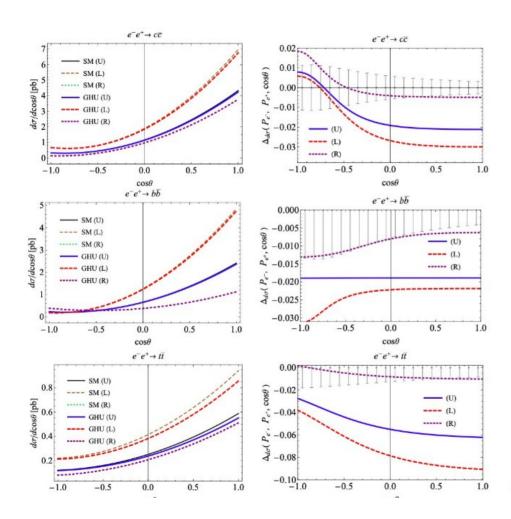
- Two remarks:
 - 500 GeV is nicely away from QCD Matching regime
 Less systematic uncertainties
 - The determination of axial form factors highly benefit from higher energies

Mapping between FF and EFT Coefficients

$$\begin{split} F_{1V}^{Z} &= \frac{\frac{1}{4} - \frac{2}{3}s_{W}^{2}}{s_{W}c_{W}} - \frac{m_{t}^{2}}{\Lambda^{2}} \frac{1}{2s_{W}c_{W}} \left[C_{\varphi q}^{V} = C_{\varphi u}^{(33)} + (C_{\varphi q}^{1(33)} - C_{\varphi q}^{3(33)}) \right], \\ F_{1A}^{Z} &= \frac{-\frac{1}{4}}{s_{W}c_{W}} - \frac{m_{t}^{2}}{\Lambda^{2}} \frac{1}{2s_{W}c_{W}} \left[C_{\varphi q}^{A} = C_{\varphi u}^{(33)} - (C_{\varphi q}^{1(33)} - C_{\varphi q}^{3(33)}) \right], \\ F_{2V}^{Z} &= 4 \frac{m_{t}^{2}}{\Lambda^{2}} \left[C_{uZ}^{R} = \text{Re}\{c_{W}^{2}C_{uW}^{(33)} - s_{W}^{2}C_{uB}^{(33)}\}/s_{W}c_{W} \right], \\ F_{2A}^{Z} &= 4 \frac{m_{t}^{2}}{\Lambda^{2}} i \left[C_{uZ}^{I} = \text{Im}\{c_{W}^{2}C_{uW}^{(33)} - s_{W}^{2}C_{uB}^{(33)}\}/s_{W}c_{W} \right], \end{split}$$

a BSM example: GUT Inspired Grand Higgs Unification Model





- Model parameter is Hosotani angle θ_{H} yielding the Higgs-Potential as consequence of Aharanov-Bohm Phase in 5^{th} dimension
- Model defined in Randall-Sundrum warped extra dimensions
 - KK excitations of gauge bosons and new bosons modify fermion couplings
- Predictions for ILC
 - $m_{\scriptscriptstyle KK}$ = 13 TeV and $\theta_{\scriptscriptstyle H}$ = 0.1
- Deviations from SM of the order of a few %
 - Effects measurable already at 250 GeV
 - · Effects amplified by beam polarisations
 - Effects for tt, bb and cc (and other light fermions)
- One concrete example for importance to measure full pattern of fermion couplings
- etingFull pattern only available with beam polarisation polarisation



Polarisation & Electroweak Physics at high energies



similarly, disentangle Z / γ exchange in $e^+e^- \rightarrow f\bar{f}$

g_{Lf}, g_{Rf}: helicity-dependent couplings of Z to fermions

$$\Rightarrow A_f = \frac{g_{Lf}^2 - g_{Rf}^2}{g_{Lf}^2 + g_{Rf}^2}$$

specifically for the electron:
$$A_e = \frac{(\frac{1}{2} - \sin^2 \theta_{eff})^2 - (\sin^2 \theta_{eff})^2}{(\frac{1}{2} - \sin^2 \theta_{eff})^2 + (\sin^2 \theta_{eff})^2} \approx 8(\frac{1}{4} - \sin^2 \theta_{eff})$$

at an *un*polarised collider:

$$A_{FB}^f \equiv rac{(\sigma_F - \sigma_B)}{(\sigma_F + \sigma_B)} \ = rac{3}{4} A_e A_f \quad ext{ => no direct access to Ae,}$$

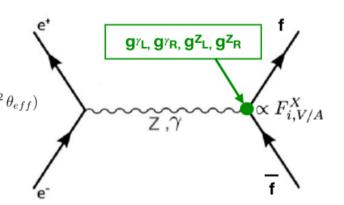
only via tau polarisation

While at a *polarised* collider:

$$A_e = A_{LR} \equiv \frac{\sigma_L - \sigma_R}{(\sigma_L + \sigma_R)} \quad \text{and} \quad A_{FB,LR}^f \equiv \frac{(\sigma_F - \sigma_B)_L - (\sigma_F - \sigma_B)_R}{(\sigma_F + \sigma_B)_L + (\sigma_F + \sigma_B)_R} = \frac{3}{4} A_f$$

trading theory uncertainy:

the **polarised** $A_{FB,LR}^{J}$ receives 7 x smaller radiative corrections than the unpolarised A_{FB}^{J} !





Detector Technologies

Vertex: CMOS, DEPFET, FPCCD, ...

Tracker:

TPC (GEM, micromegas, pixel) + silicon pixels/strips

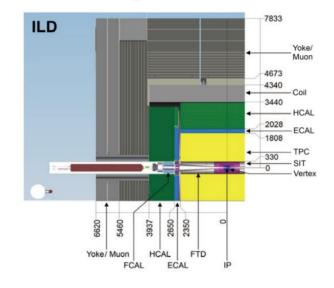
ECAL:

Silicon (5x5mm²) or Scintillator (5x45mm²) with Tungsten absorber

HCAL:

Scintillator tile (3x3 cm²) or Gas RPC (1x1 cm²) with Steel absorber

All inside solenoidal coil of 3-4 T



Detector R&D collaborations:







ILD Design Goals

Features of ILC:

low backgrounds, low radiation, low collision rate (5-10 Hz)

These allow us to pursue aggressive detector design:

	Detector Requirements	Physics
•	Impact parameter resolution $\sigma(d_0) < 5 \bigoplus 10 / (p[GeV] \sin^{3/2}\theta) \mu m$	H→bb,cc,gg,ττ
•	Transverse momentum resolution $\sigma(1/p_T) = 2 \times 10^{-5} \text{ GeV}^{-1} \bigoplus 1 \times 10^{-3} / (p_T \sin^{1/2}\theta)$	Total e+e-→ZH cross section
•	Jet energy resolution 3-4% (around E _{jet} ~100 GeV)	H→invisible
•	Hermeticity $\theta_{min} = 5 \text{ mrad}$	H→invisible; BSM

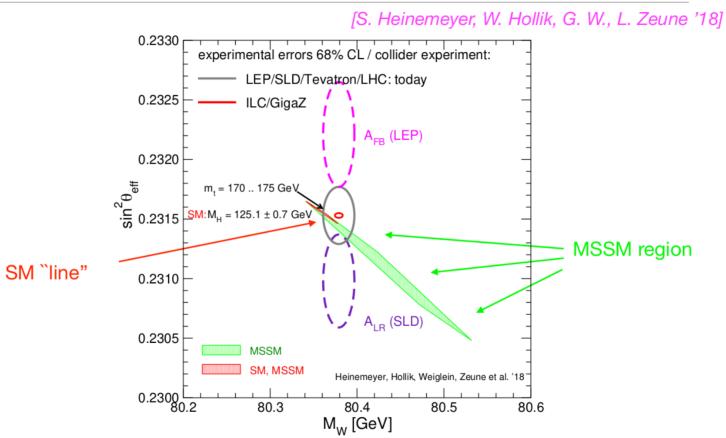
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 $\Rightarrow M_{\rm W}$ and $\sin^2\theta_{\rm eff}$ have high sensitivity for model discrimination

BSM opportunities at e+e- colliders, Georg Weiglein, LCWS2021, 03 / 2021

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