



Simulation of particle identification with cluster counting technique

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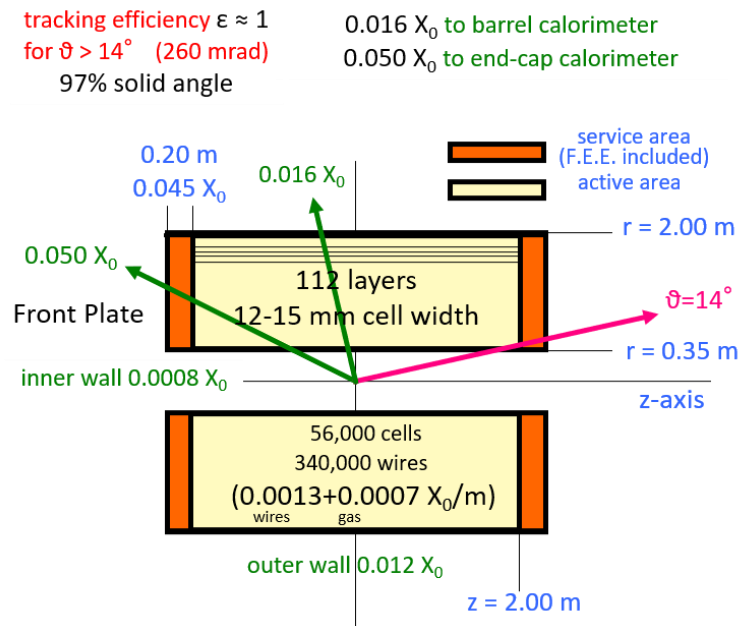
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^e Politecnico di Bari

IDEA DRIFT CHAMBER

The IDEA drift chamber (DCH) is the tracker of FCC-ee and CEPC.

It is designed to provide efficient tracking, high precision momentum measurement and excellent particle identification by exploiting the application of the cluster counting technique.



- **He based gas mixture** (90% He – 10% i-C₄H₁₀)
- **Full stereo configuration** with alternating sign stereo angles ranging from 50 to 250 mrad
- 12 ÷ 14.5 mm wide square cells 5 : 1 field to sense wires ratio
- 56,448 cells
- 14 co-axial super-layers, 8 layers each (112 total) in 24 equal azimuthal (15°) sectors

MAIN GOALS

- Gas containment – wire support functions separation: the total amount of material in radial direction, towards the barrel calorimeter, is of the order of 1.6% X_0 , whereas in the forward and backward directions it is equivalent to about 5.0% X_0 , including the endplates instrumented with front end electronics.
- Feed-through-less wiring: allows to increase chamber granularity and field/sense wire ratio to reduce multiple scattering and total tension on end plates due to wires by using thinner wires
- **Cluster timing:** allows to reach **spatial resolution < 100 μ m for 8 mm drift cells** in He based gas mixtures (such a technique is going to be implemented in the MEG-II drift chamber under construction)
- **Cluster counting:** allows to reach **dN_{cl}/dx resolution < 3%** for particle identification (a factor 2 better than dE/dx as measured in a beam test)

MORE INFORMATION:

https://indico.cern.ch/event/656491/contributions/2939121/attachments/1629781/2597342/IDEA-CDCH_FCCweek18.pdf

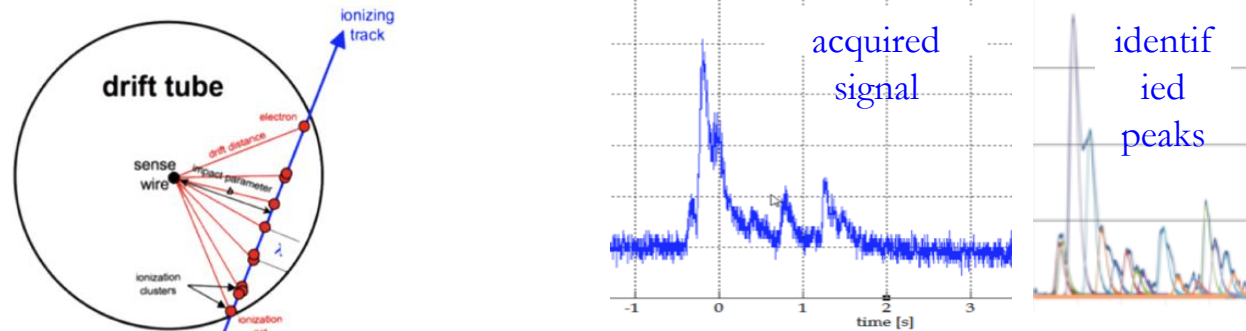
Cluster Counting/Timing and P.Id. expected performance

Gaseous counters provide signals with pulse height proportional to the numbers of electron liberated during the ionization process along the track length inside the sensitive volume and proportional to the energy deposit .

Using the information about energy deposit particle identification can be performed but, the large and inherent uncertainties in total energy deposition represent a limit to the particle separation capabilities.

Cluster counting technique can improve the particle separation capabilities.

The method consists in singling out, in ever recorded detector signal, the isolated structures related to the arrival on the anode wire of the electrons belonging to a single ionization act (dN/dx).



dE/dx

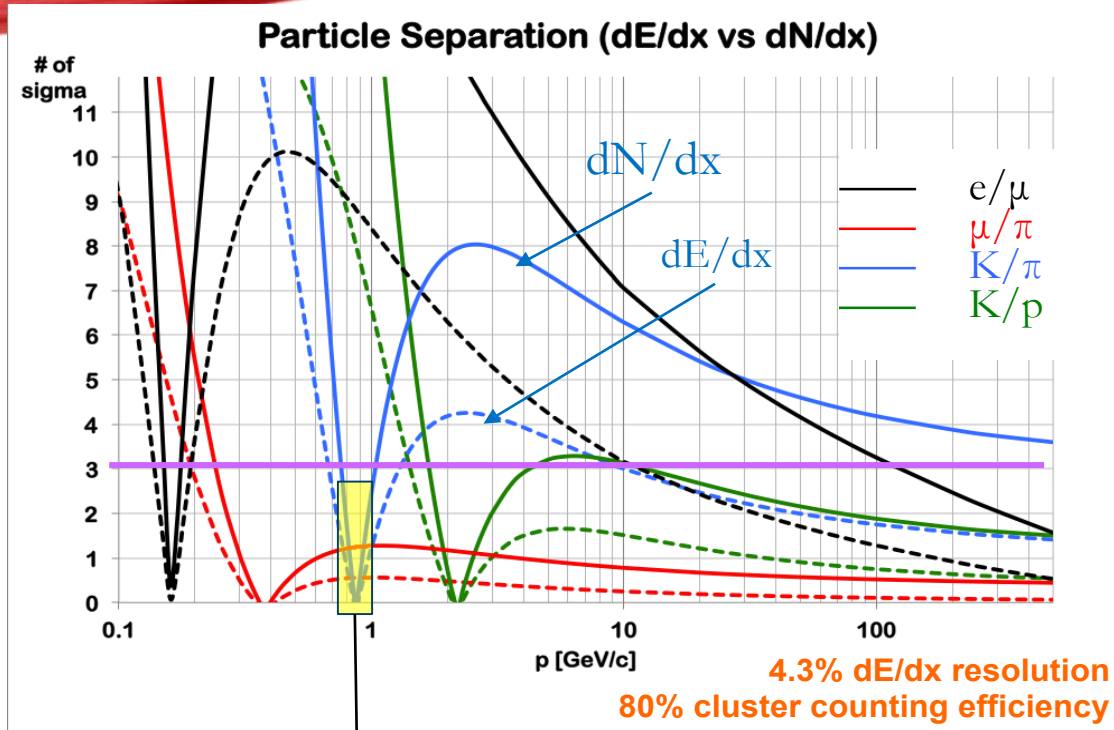
Truncated mean cut (70-80%) reduces the amount of collected information $n \approx 100$ and a 2m track at 1 atm give $\sigma \approx 4.3\%$

dN_{cl}/dx

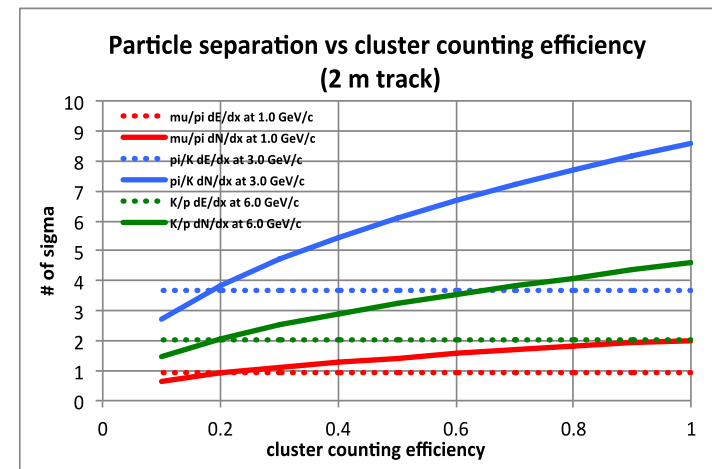
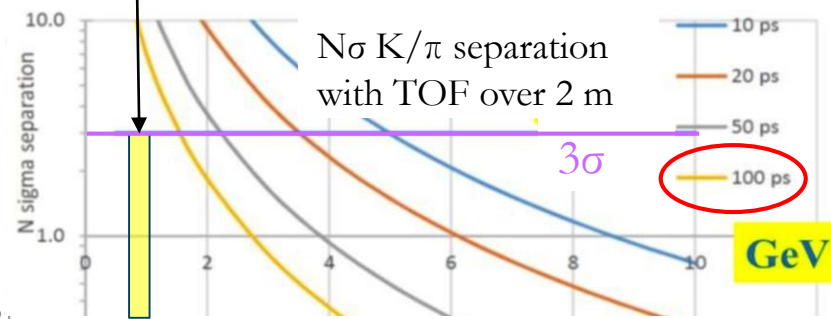
$\delta_{cl} = 12.5/cm$ for He/ iC_4H_{10} =90/10 and a 2m track give $\sigma \approx 2.0\%$

Moreover, C.C. may improve the spatial resolution $< 100 \mu m$ for 8 mm drift cells in He based gas mixtures

Cluster Counting/Timing and P.Id. expected performance



- 80% cluster counting efficiency.
- Expected excellent K/π separation over the entire range except 0.85 < p < 1.05 GeV (blue lines)
- Could recover with timing layer



Drift Chamber simulation - Cluster Counting/Timing simulation

To investigate the potential of the Cluster Counting technique (for He based drift chamber) on physics events a reasonable simulation/parameterization of the ionization clusters generation in Geant4 is needed.

Garfield/Garfield++:

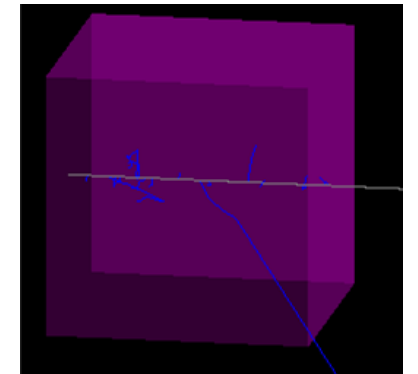
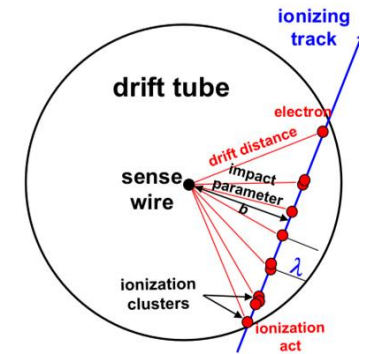
- (Heed) simulates the ionization process in the gasses (not only) in a detailed way.
- (Magboltz) computes the gas properties (drift and diffusion coefficients as function of the fields value)
- solves the electrostatic planar configuration and simulates the free charges movements and collections on the electrodes.

So Garfield can study and characterize the properties and performance of single cell or drift chamber with simple geometry, but it is not designed to simulate a full detector neither study collider events.

Geant4:

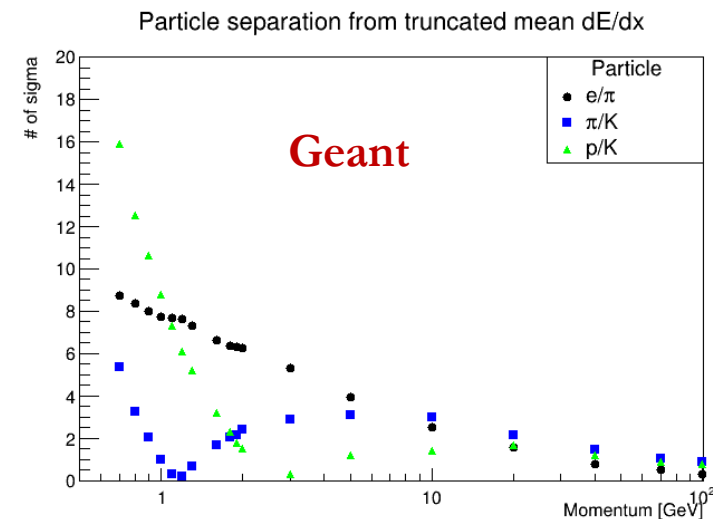
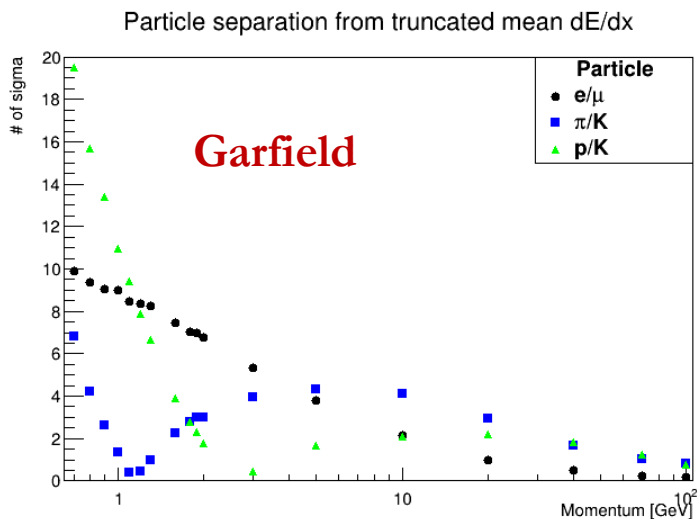
- Simulates the elementary particle interaction with material of a full detector
- Studies colliders events
- It doesn't simulate (normally) the ionization clustering process
- It doesn't simulate (normally) the free charges movements and collections on the electrodes.

It is very useful to simulate the elementary particle interaction with the material of a full (complex) detector and to study collider events. The fundamental properties and performance of the sensible elements (drift cells) have to be parametrized or ad-hoc physics models have to be defined.

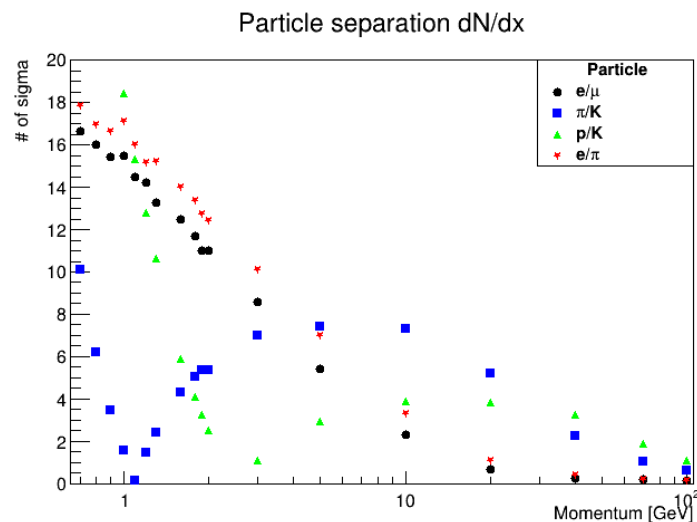


Drift Chamber simulation - Cluster Counting/Timing simulation-PID

We are simulating tracks crossing 200 cell (1 cm long side) of 90% He and 10% iC_4H_{10} with Garfield and Geant



Truncated mean at 70%



$$n_{\sigma} = \frac{\Delta_A - \Delta_B}{\langle \sigma_{A,B} \rangle} \quad \langle \sigma_{A,B} \rangle \text{ is the average of the two resolutions.}$$

Cluster counting leads to an **improvement** on particle separation power. As example, around 5 GeV the power separation of a pion from kaon obtained with traditional method is about 4, the one obtained with cluster counting is around 8.

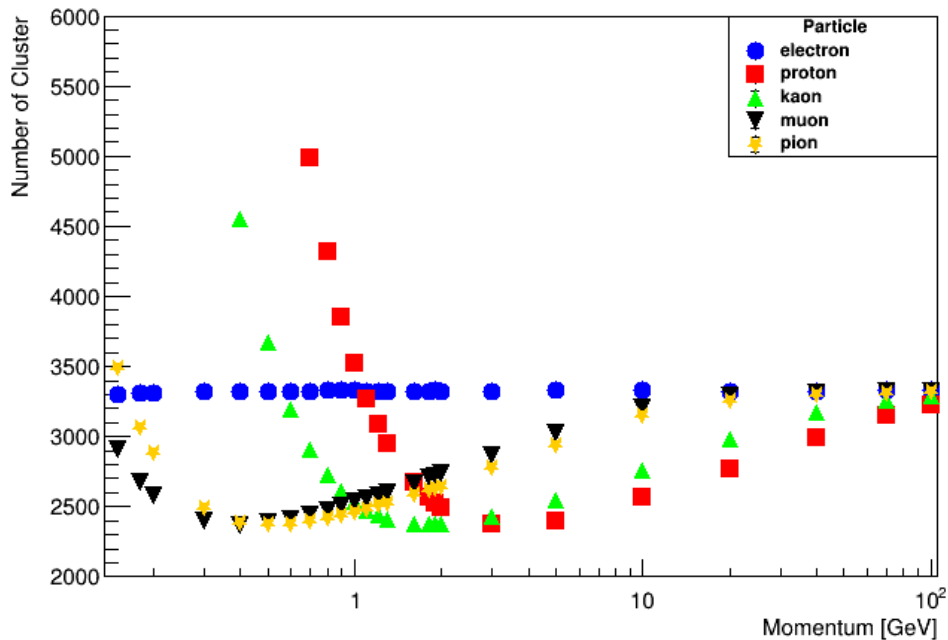
Drift Chamber simulation - Cluster Counting/Timing simulation-Garfield analysis

Studying the results from Garfield++ simulations, we can interpret correctly the results obtained from Geant4 simulations with the goal of reconstruct the number of clusters and the cluster size generated from different particles with different momenta passing through the tracker detector.

The goal is to extract from Garfield++ the relevant parameters to create models to convert the energy loss to cluster and then extract them as function of the primary particle $\beta\gamma$.

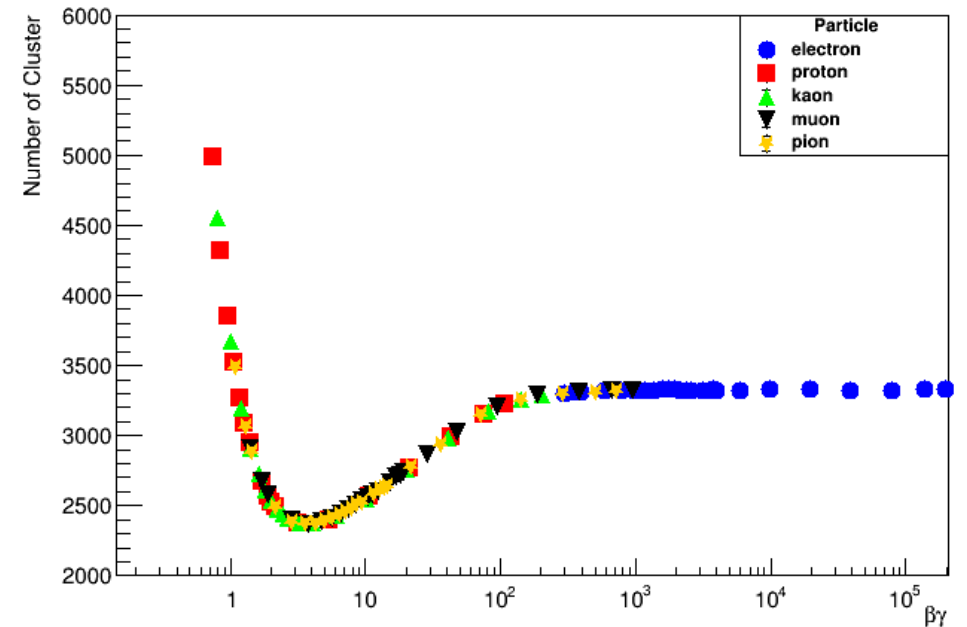
Number of cluster from Garfield++

Number of cluster for different particles vs momentum



Here the distribution of number of cluster produced by different particle at different momenta, obtained with Garfield++

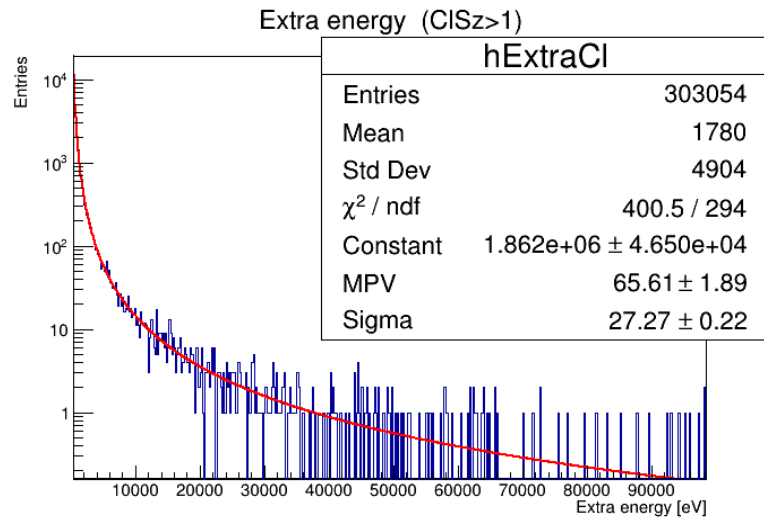
Number of cluster for different particles vs $\beta\gamma$



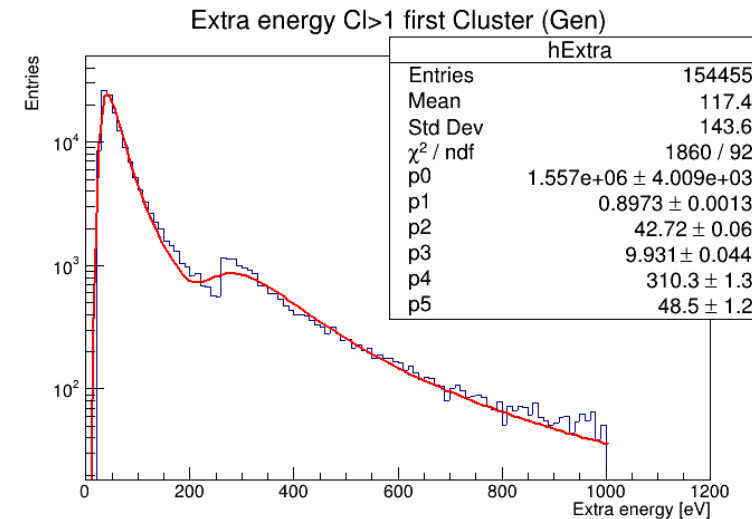
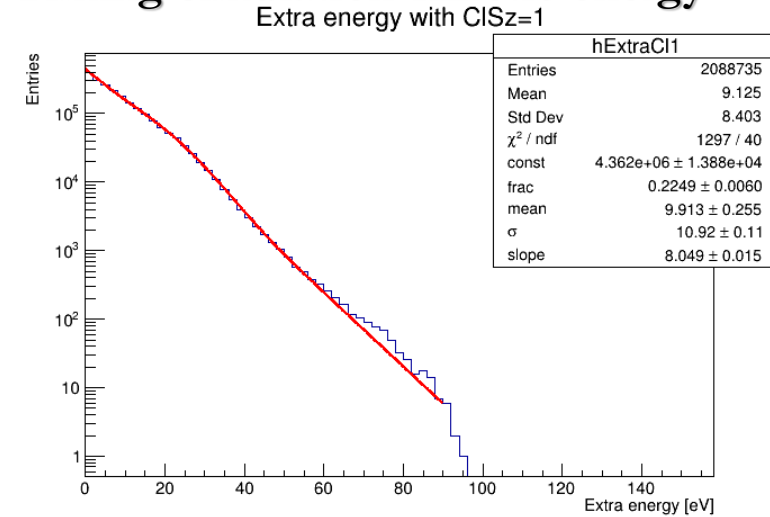
Drift Chamber simulation - Cluster Counting/Timing simulation-Kinetic energy

Kinetic energy distribution for cluster with cluster size equal to 1.
The fit is the sum of an exponential function plus a Gaussian function.

Kinetic energy distribution for cluster with cluster size higher than 1 (left)
and up to 1keV cut (right).
The fits are performed with a Landau functions.

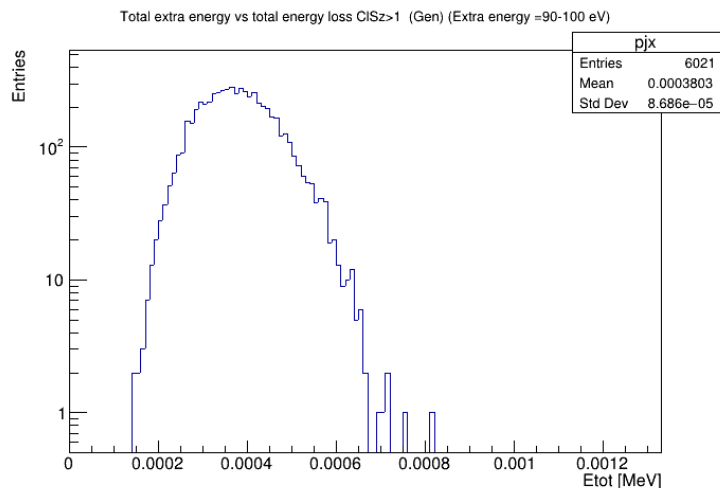
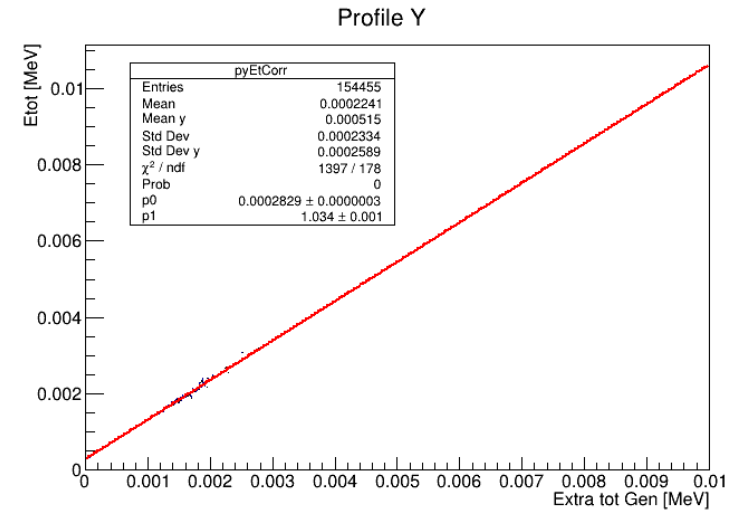
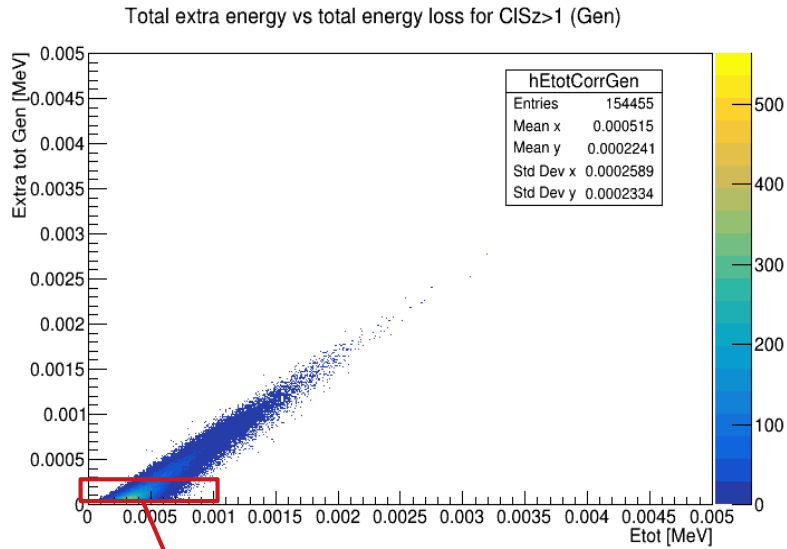


The **goal** is to create two different models for kinetic energy of clusters with cluster size equal to 1 and higher than 1, to interpret correctly the total energy loss by different particles in a single cell.



1keV cut is equivalent to the single interaction range cut set (by default) in Geant4

Drift Chamber simulation - Cluster Counting/Timing simulation-kinetic energy for cluster with cluster size higher than 1



maxExEcl = total kinetic energy spent to create cluster with cluster size higher than 1

MaxEx0 is the first parameter of the linear fit

MaxExSlp is the second parameter of linear fit

ExSgm is the average of the sigma of each point in the correlation trend.

$$\text{maxExEcl} = \frac{(E_{\text{tot}} - \text{maxEx0} + g\text{Random} \rightarrow \text{Gaus}(0, \text{ExSgm}))}{\text{maxExSlp}}$$

The figure shows an example of distribution of total energy loss for extra energy between 90 and 100 eV for cluster with cluster size higher than one.

Drift Chamber simulation - Cluster Counting/Timing simulation- The algorithms

We implemented seven different algorithms trying to reproduce the number of cluster and the cluster size.

The first step common to all algorithm is the evaluation of the total kinetic energy for cluster with cluster size higher than one (\maxExEcl) event by event.

1) The first algorithm uses a reference value of the ratio between clusters containing a single electron and clusters containing more than one electron (Rt). Using the Rt value, the algorithm chooses to create cluster with cluster size one or higher. Then, it assigns the kinetic energy to each cluster by using the proper distributions. If the cluster has more than one electron, a check on the total kinetic energy is performed and its cluster size is evaluated. The procedure is repeated until the sum of primary ionization energy and kinetic energy per cluster saturate the energy loss of the event.

2) **The second algorithm, if \maxExEcl is higher than zero, generates the kinetic energy for clusters with cluster size higher than one by using its distribution and evaluates cluster size. This procedure is repeated until the sum of primary ionization energy and kinetic energy per cluster saturate the \maxExEcl of the event.**

Then, using the remaining energy ($E_{loss} - \maxExEcl$), the algorithm creates clusters with cluster size equal to one by assigning their kinetic energy according to the proper distribution.

The reconstruction of clusters with cluster size equal to one remains the same for all next algorithms.

3) **The third algorithm (similar to the previous), during the generation of cluster with cluster size higher than one, assigns the kinetic energy to them, choosing the best over five extractions that makes the total kinetic energy for cluster with cluster size higher than one approximating better the \maxExEcl .**

To correct a systematic underestimation of the mean number of clusters, an additional correction to the residual energy for generating cluster with cluster size equal to one can be used.

Drift Chamber simulation - Cluster Counting/Timing simulation- The algorithms

4)The fourth algorithm (similar to the previous), during the generation of cluster with cluster size higher than one, assigns (by extracting from the proper distribution) the kinetic energy to them, until the total kinetic energy better approximates the \maxExEcl .

5)The fifth algorithm is similar to the fourth with almost differences in the technical implementation.

6)The sixth algorithm follows a different methodology. Indeed it uses the total kinetic energy of the event to evaluate a priori the number of cluster, applying the most likelihood criterium.

7)The last algorithm is similar to the second algorithm but generates the kinetic energy for cluster with cluster size higher than one by using the fit of kinetic energy distribution.

List of variables

\maxExEcl : total kinetic energy spent to create clusters with cluster size higher than 1

$ExEcl$: kinetic energy generated per cluster

$Ncl1$: number of clusters with cluster size equal to one

$Nclp$: number of clusters with cluster size higher than one

\maxCut : energy value equivalent to the range cut set in Geant4

$totExEcl$: total kinetic energy reconstructed to create clusters with cluster size higher than one

$Eloss$: energy loss from a track passing through the cell

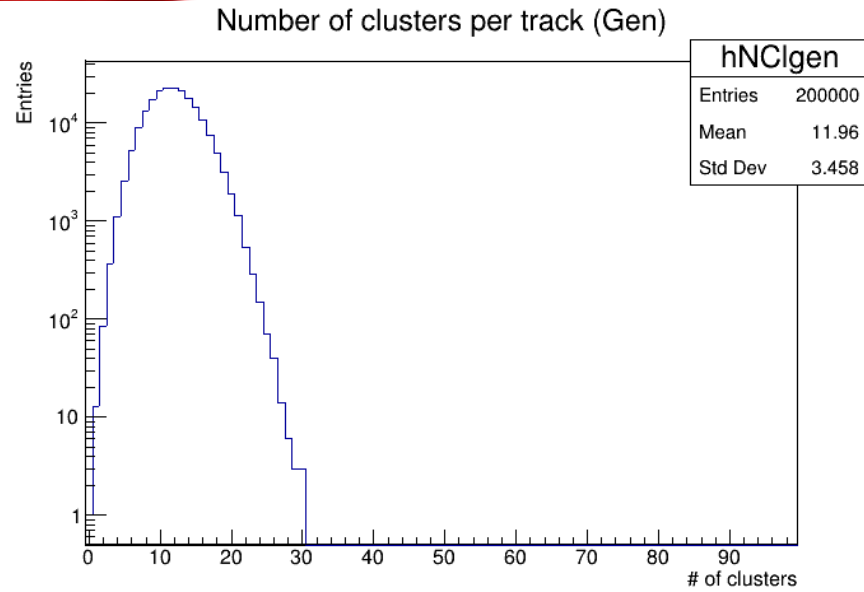
$ClSz$: cluster size

$Eizp$: primary ionization energy, 15.8 eV

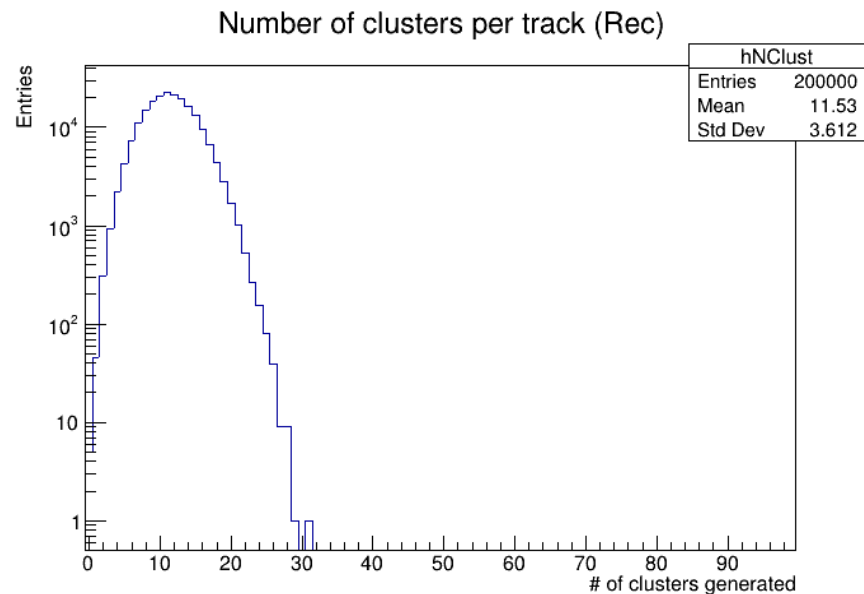
$Eizs$: secondary ionization energy, 25.6 eV

Drift Chamber simulation - Cluster Counting/Timing simulation-Results

MC Truth

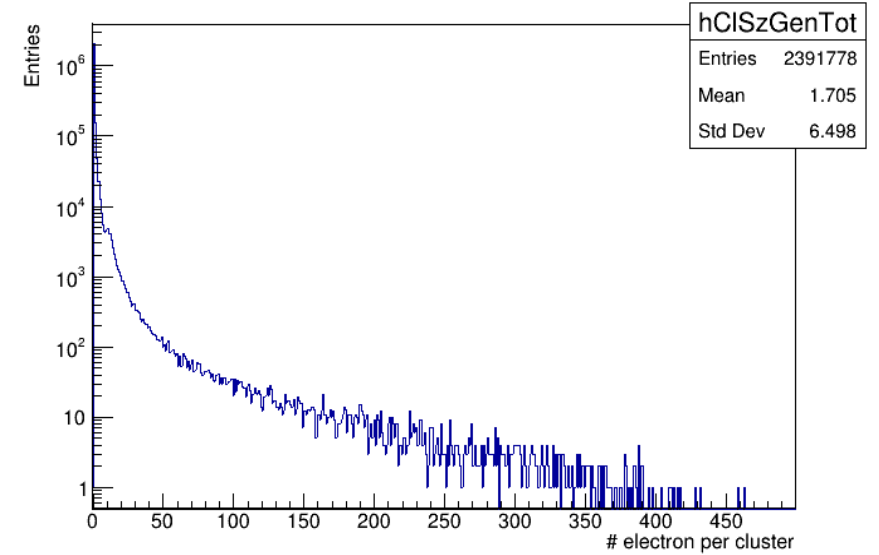


2st algorithm

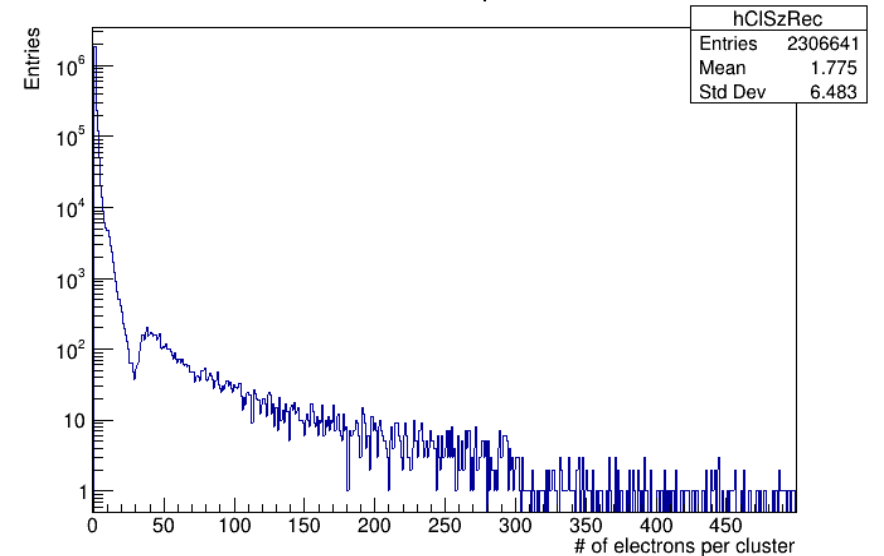


Case of study:
muon at 300 MeV

Cluster size (gen)

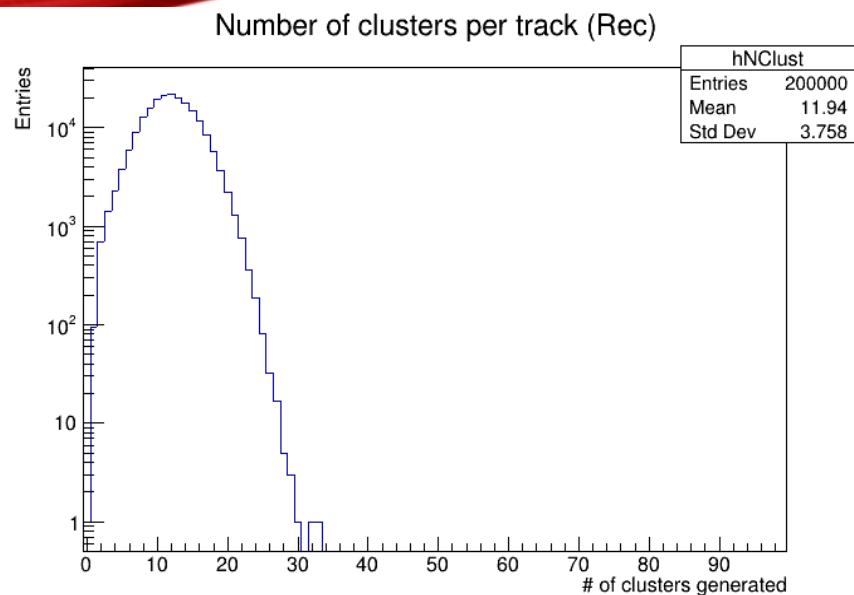


Number of electrons per cluster Rec

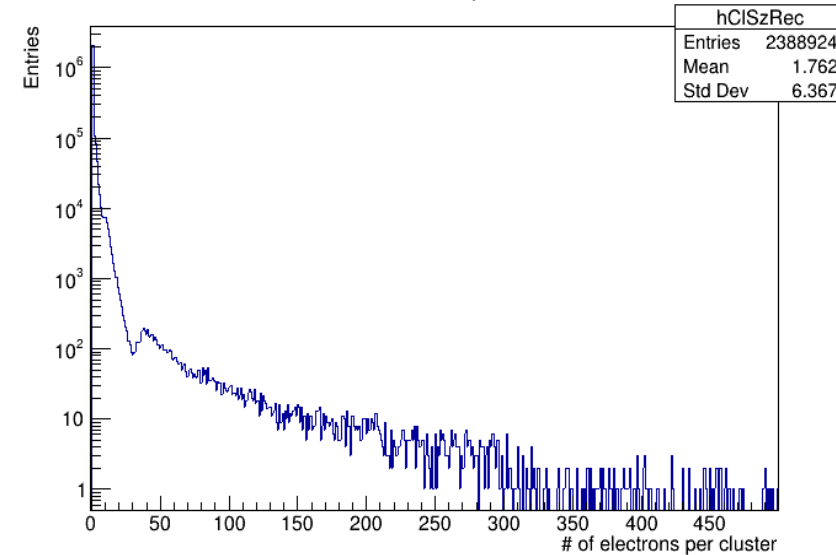


Drift Chamber simulation - Cluster Counting/Timing simulation-Results

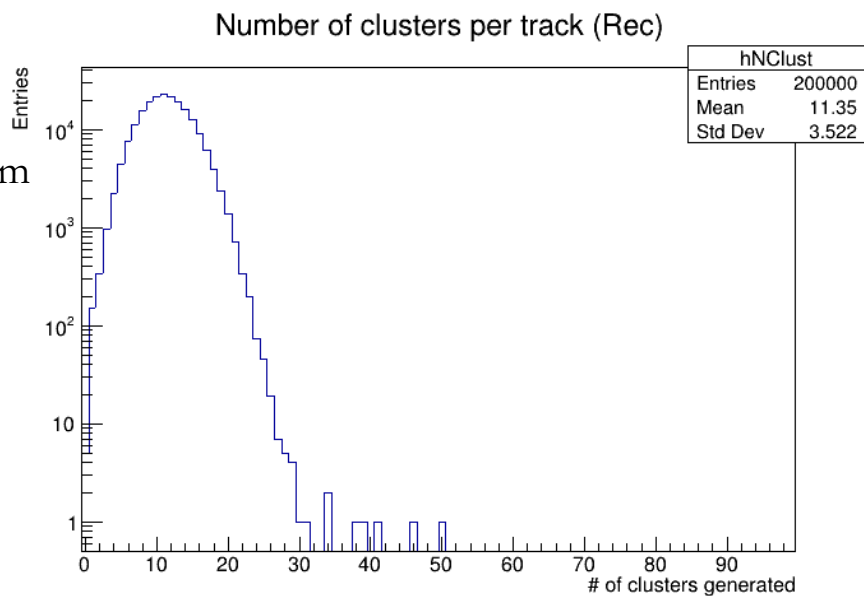
3st algorithm



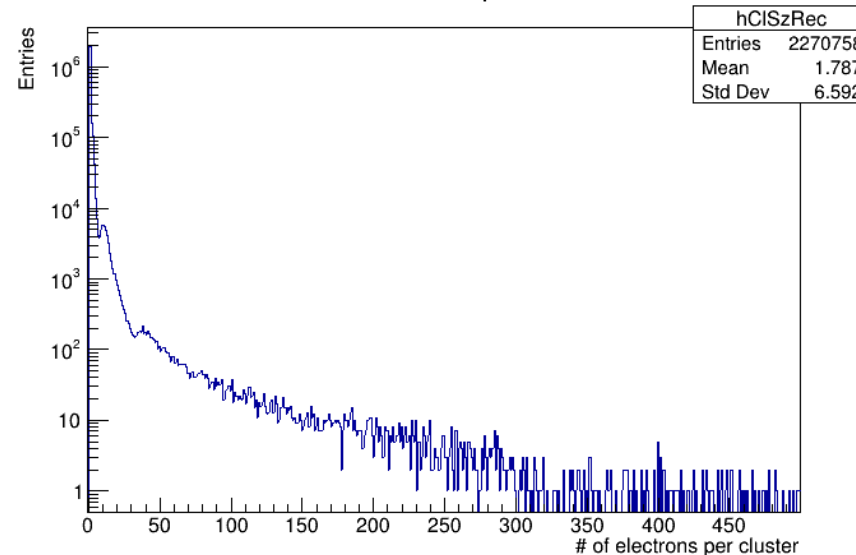
Number of electrons per cluster Rec



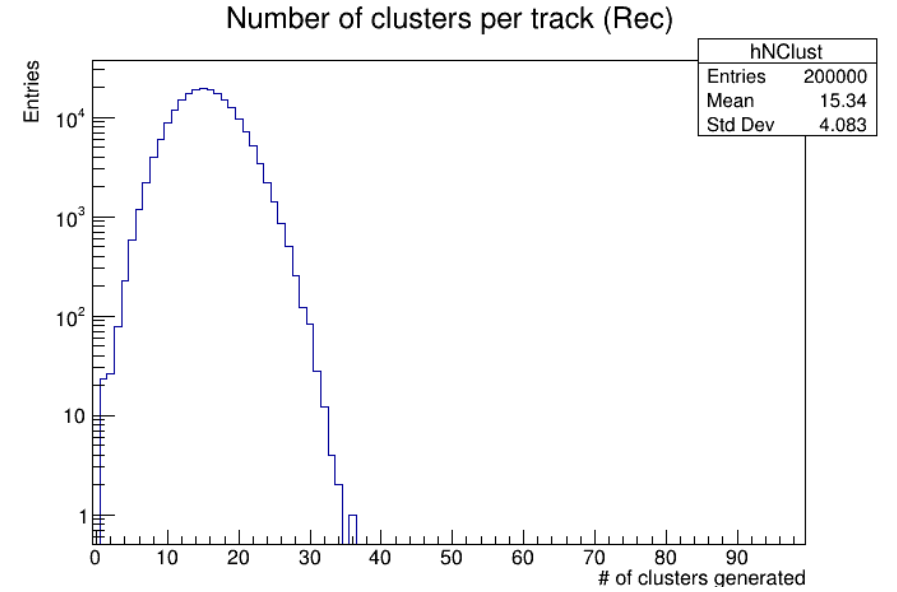
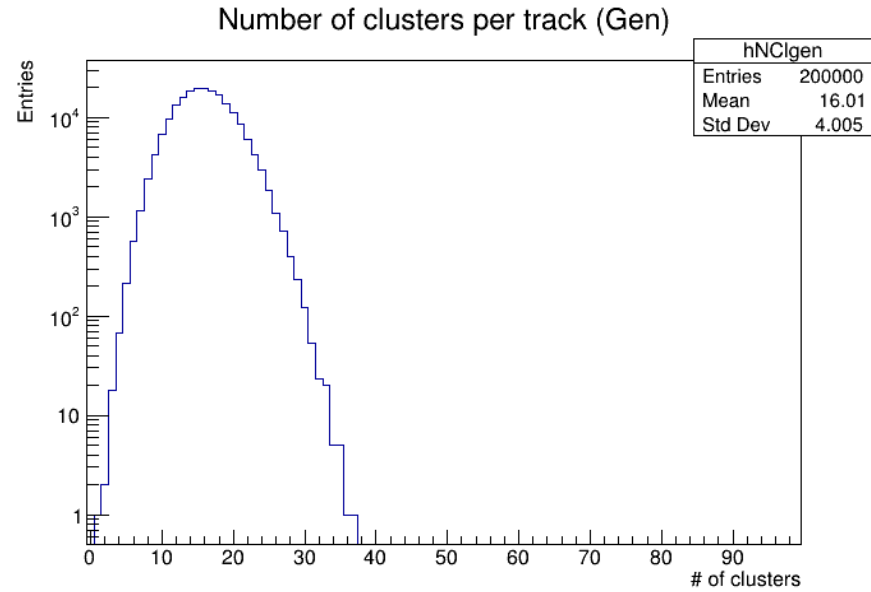
6st algorithm



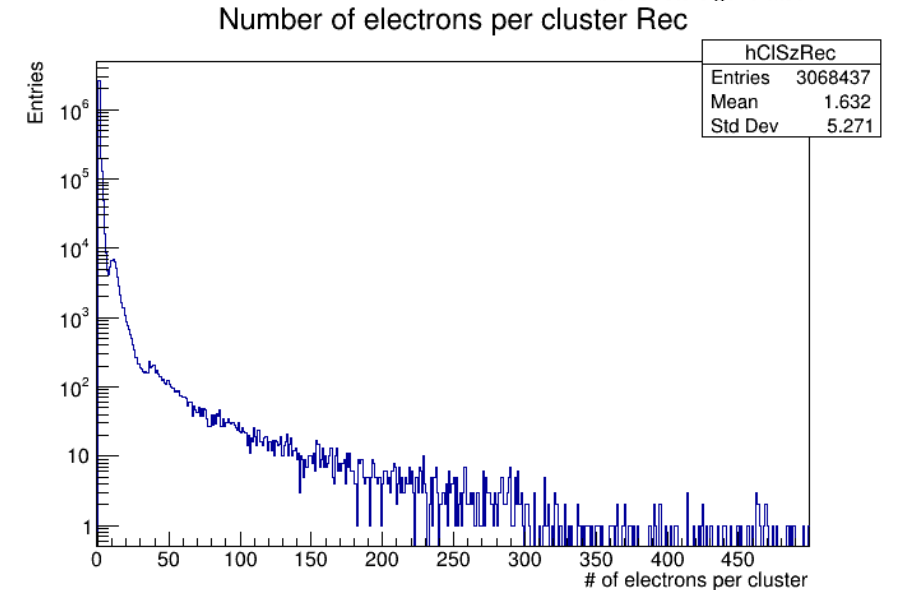
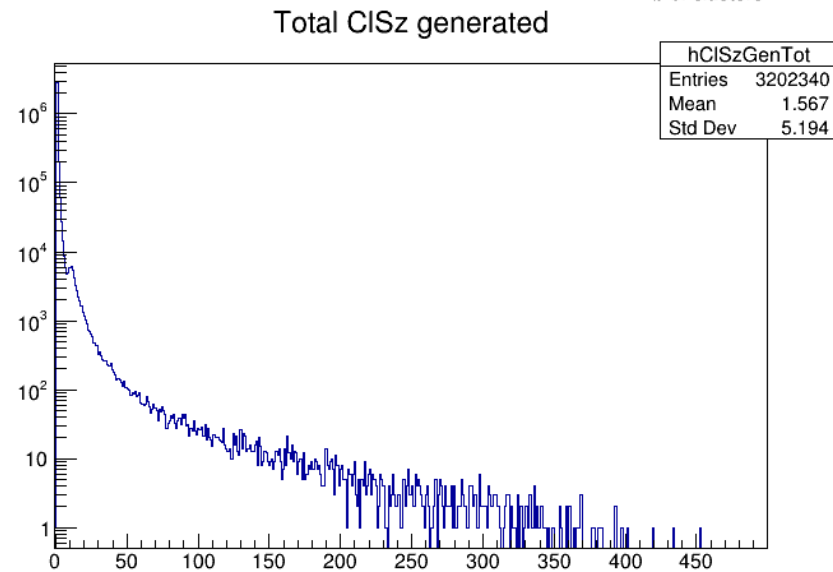
Number of electrons per cluster Rec



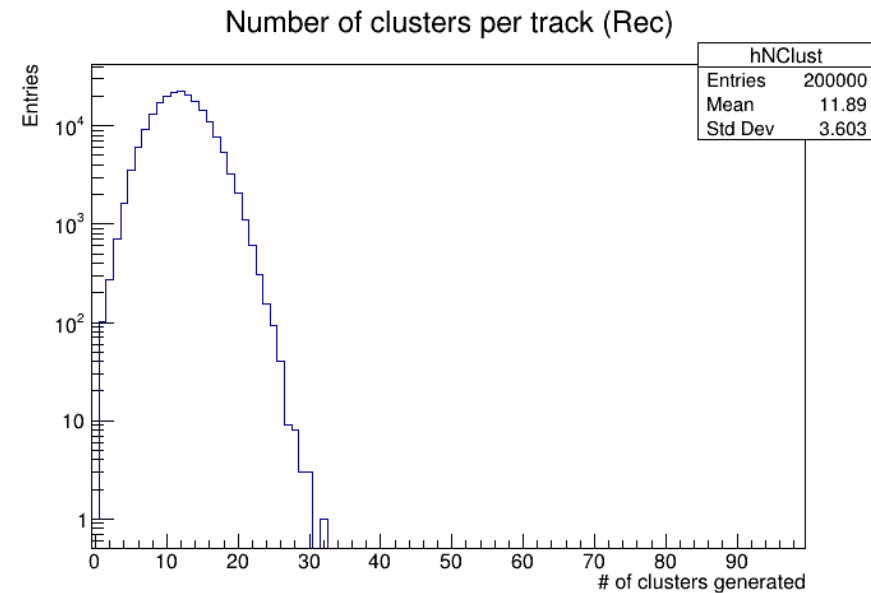
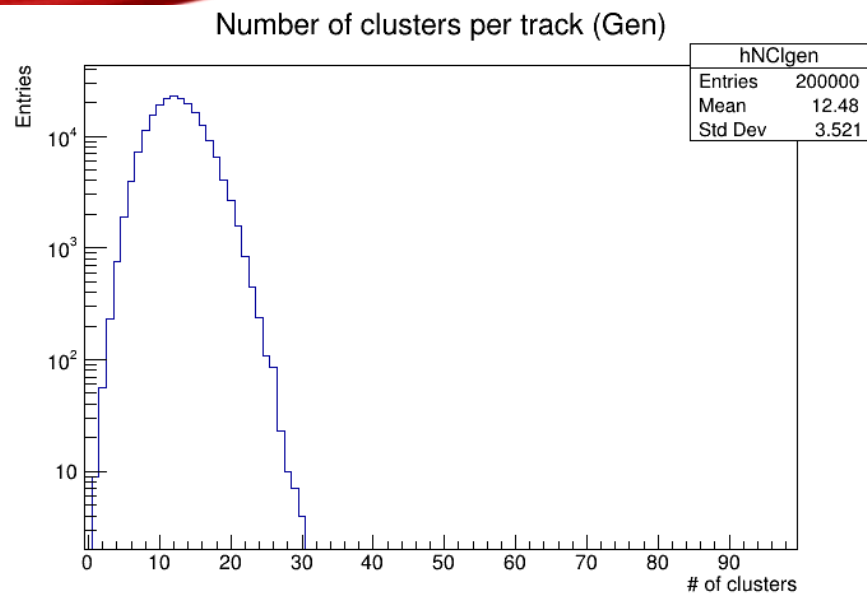
Drift Chamber simulation - Cluster Counting/Timing simulation-Test on different momenta and different particles



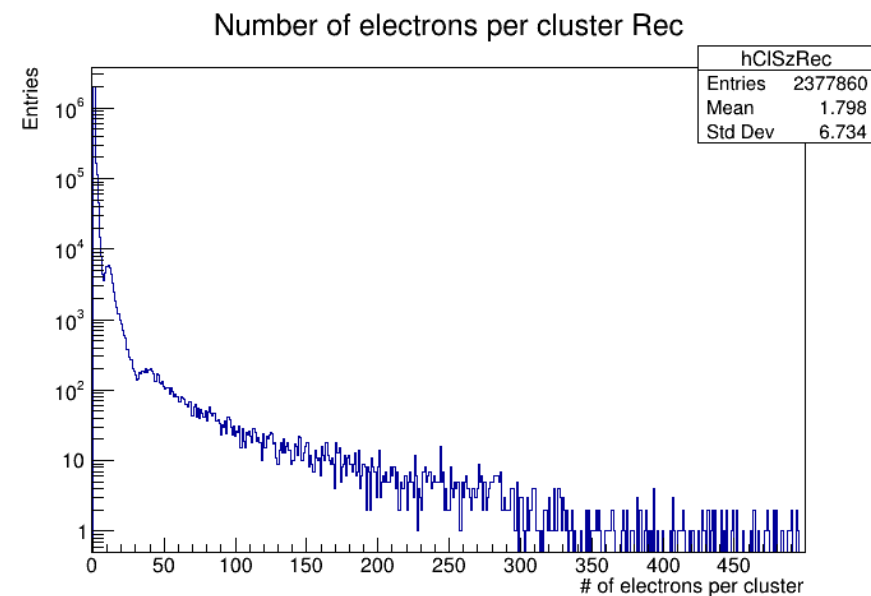
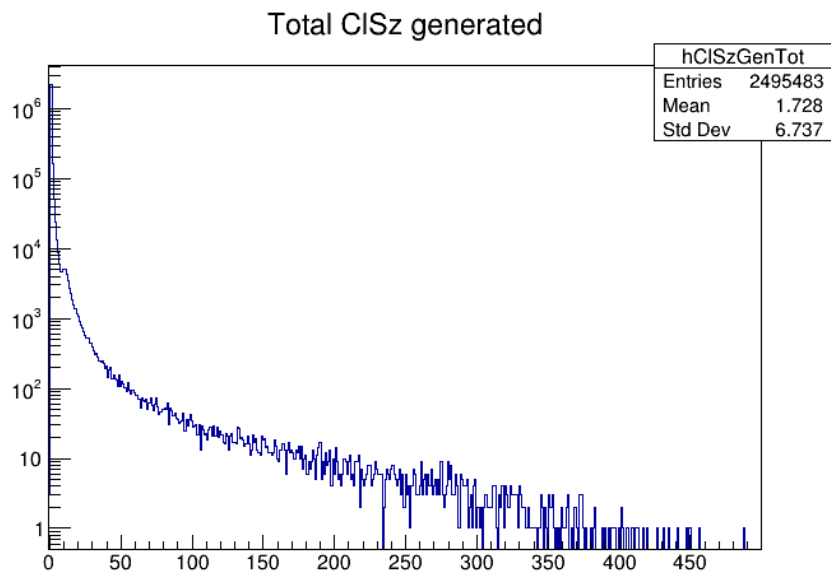
Case of study:
muon at 10 GeV



Drift Chamber simulation - Cluster Counting/Timing simulation-Test on different momenta and different particles



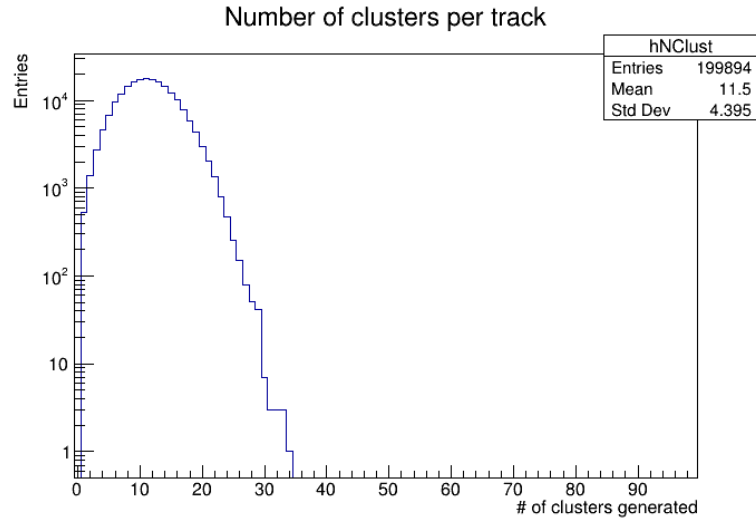
Case of study:
pion at 300 MeV



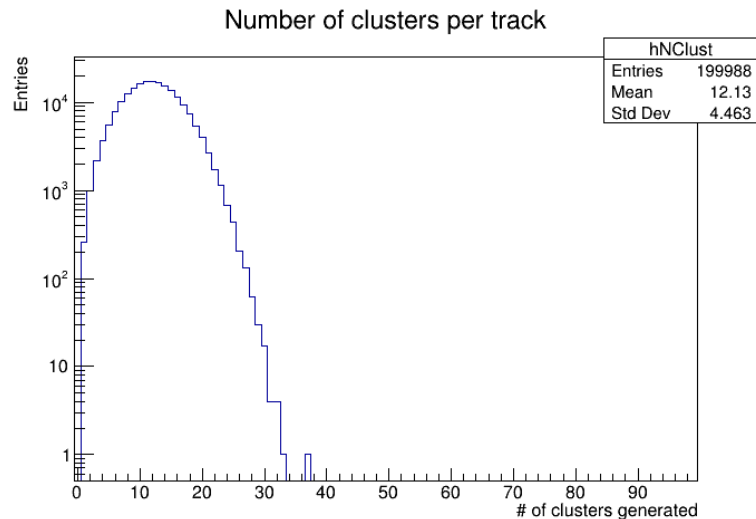
Drift Chamber simulation - Cluster Counting/Timing-Geant4

The algorithm is tested with Geant4 simulations and the results obtained are compatible with the ones obtained with Garfield++.

2st algorithm

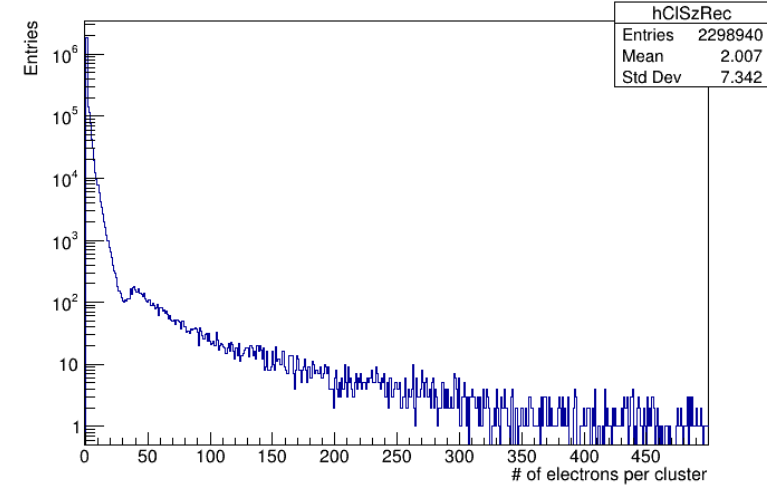


3st algorithm

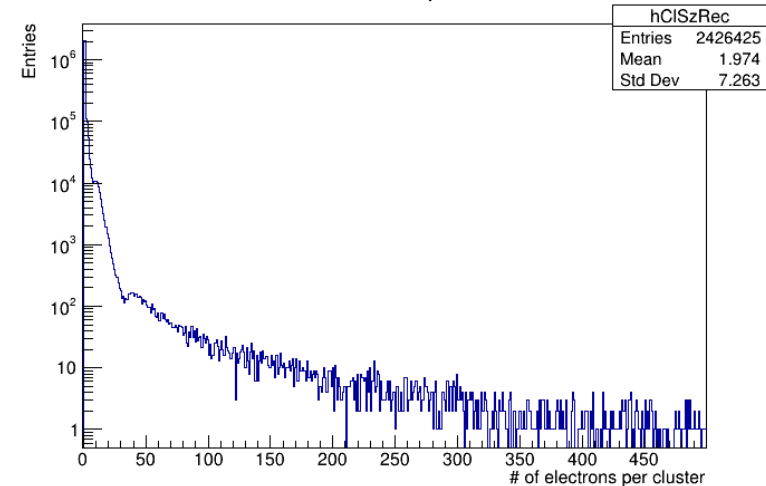


Case of study:
muon at 300 MeV

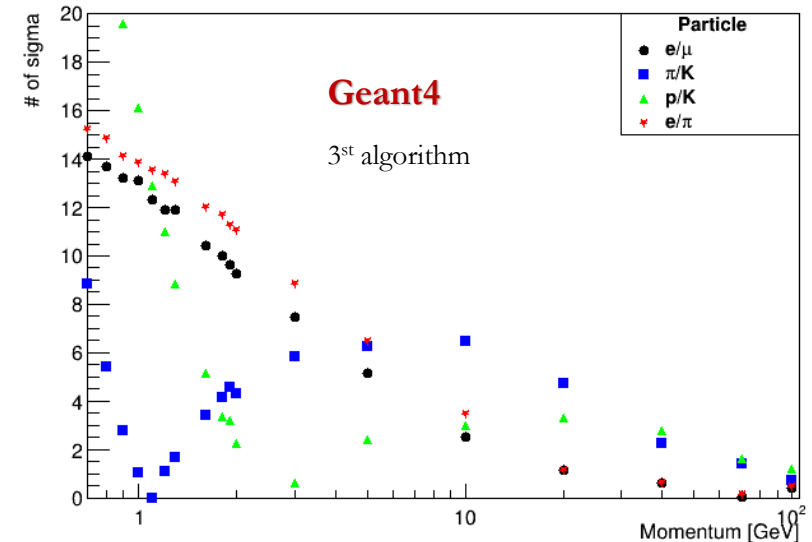
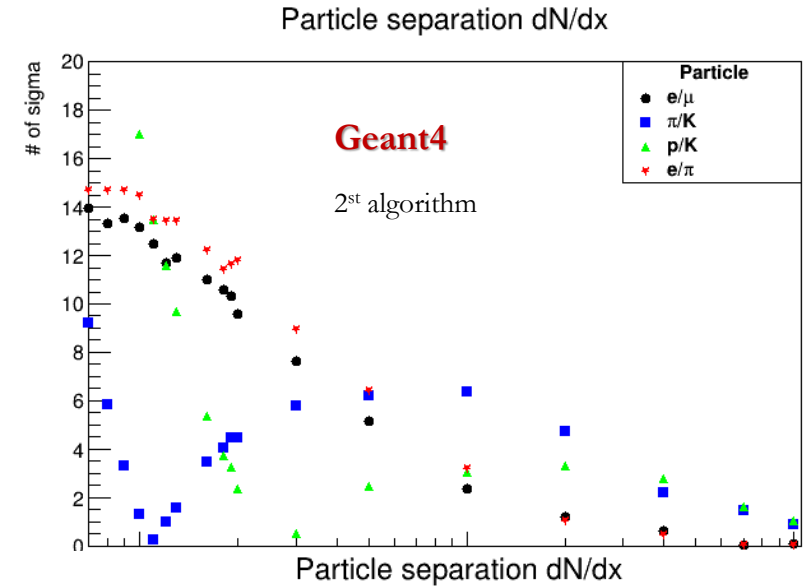
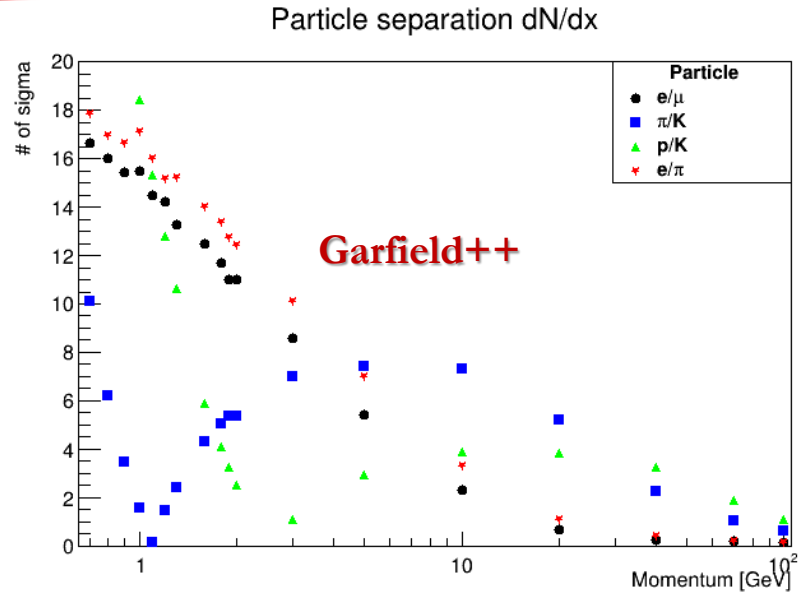
Number of electrons per cluster Rec



Number of electrons per cluster Rec



Drift Chamber simulation - Cluster Counting/Timing-PID in Geant4



We are simulating tracks crossing 200 cell (1 cm long side) of 90% He and 10% iC_4H_{10} with Garfield and Geant.
We are assuming a cluster counting efficiency of 100%.

Conclusions

- Cluster counting improves particle separation capabilities respect to the traditional dE/dx method.
- Reasonable algorithms to simulate the Ionization Clusters by using the Geant4 data are developed.
- The algorithm was tested on different cell size, providing consistent results.
- A first fast Cluster Finder algorithm was developed, implemented on an FPGA and tested on a test bench.
- We are working on importing the algorithm in the drift chamber full simulation.

To do list

- Continue to develop the full simulation and perform physics studies
- Perform PID studies with the full detector simulation
- Improve Clustering algorithm validation with measurements
- Continue to develop the DAQ prototype and test it
- Construct the monitor chamber and ad hoc prototypes

Thanks for your attention

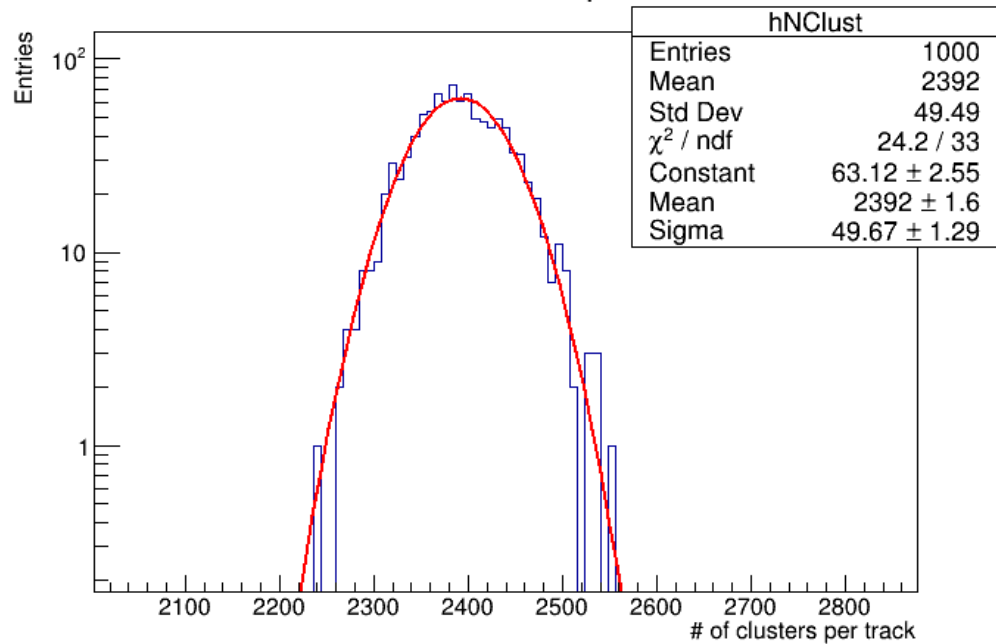


BACKUP

Number of cluster
for muon 300 MeV
in 200 cells

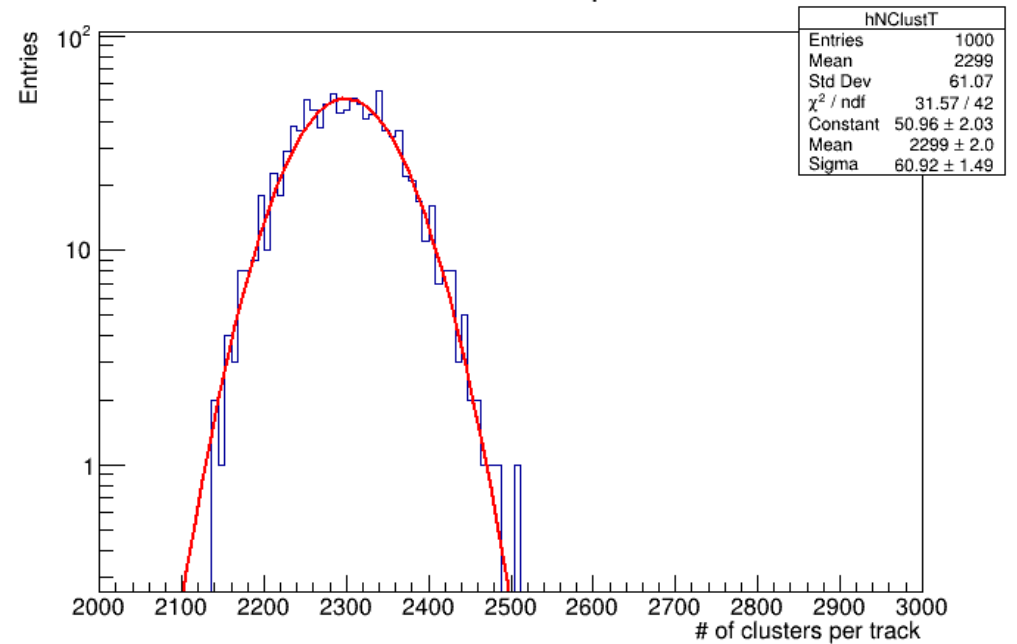
Geant4

Number of clusters per track

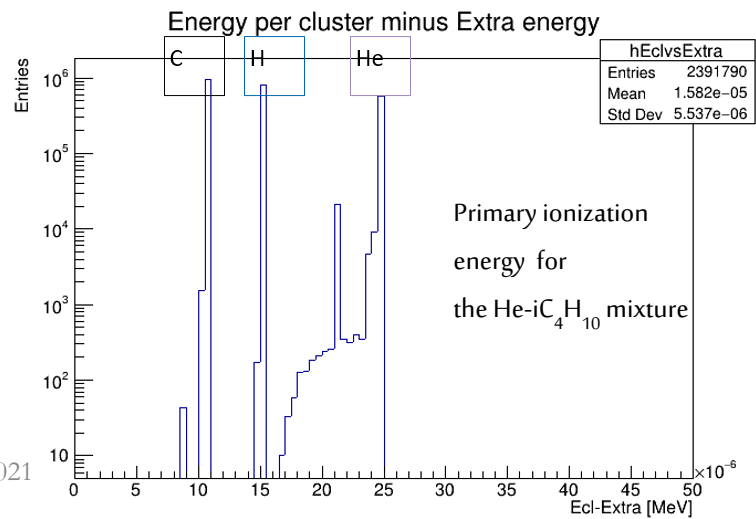
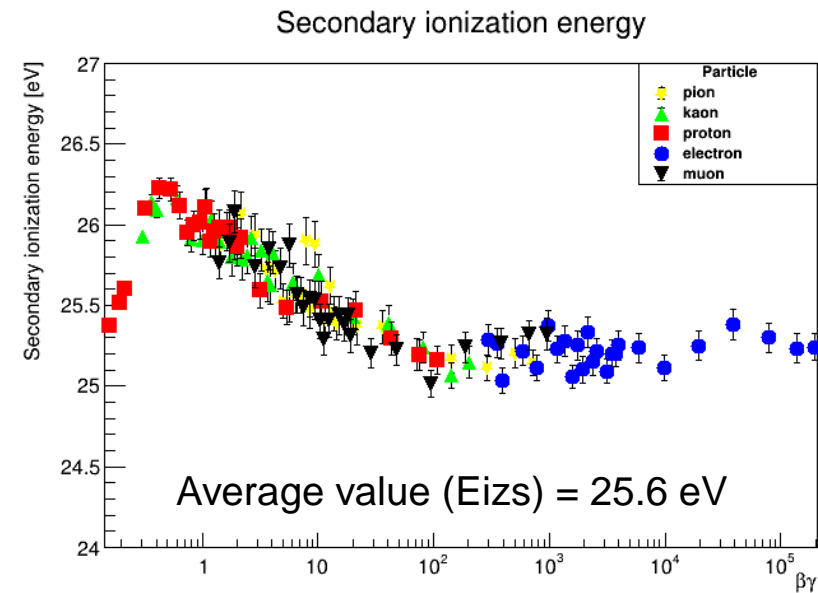
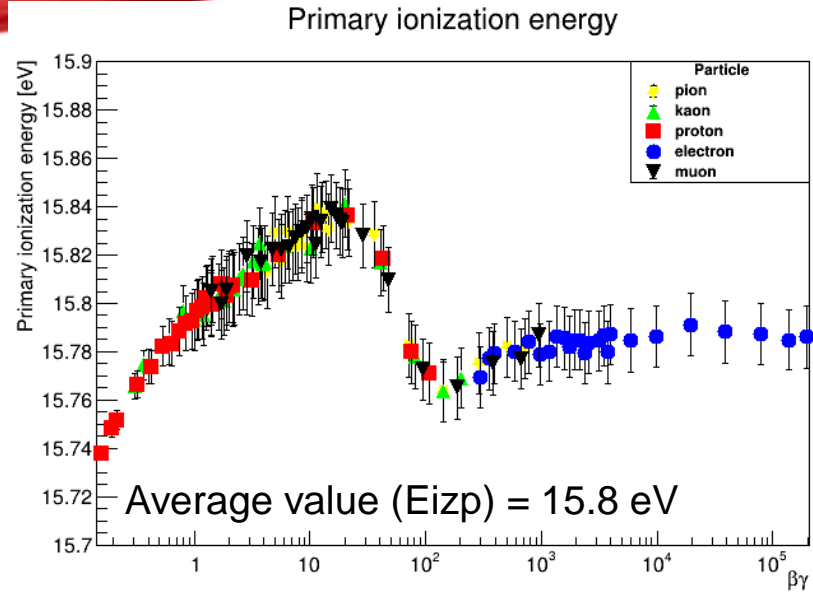


Garfield++

Number of clusters per track



Drift Chamber simulation - Cluster Counting/Timing simulation



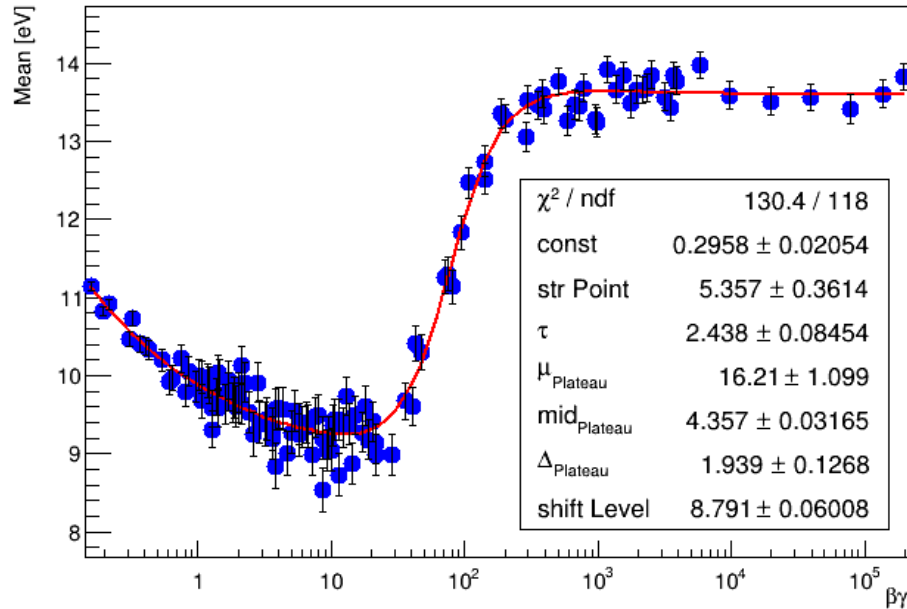
First step to the algorithm implementation is the evaluation of primary and secondary ionization energy

Fit parameters for kinetic energy distribution for cluster with cluster size equal to 1

Here the distribution of mean value (top left) and sigma value (bottom right) for different particles with different momenta. The distributions are fitted with an exponential function plus an efficiency function.

$$Eff = \frac{Eff_{plateau}}{1 + 81 \frac{V_{1/2} - V}{\Delta_{10\%}^{90\%}}}$$

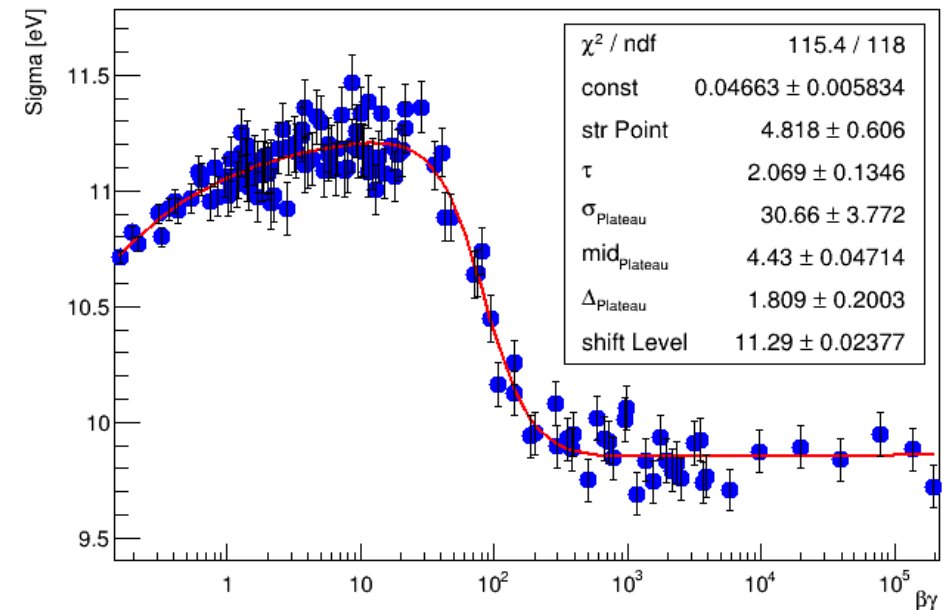
Mean from gaus+exp fit of Extra Energy with CISz=1



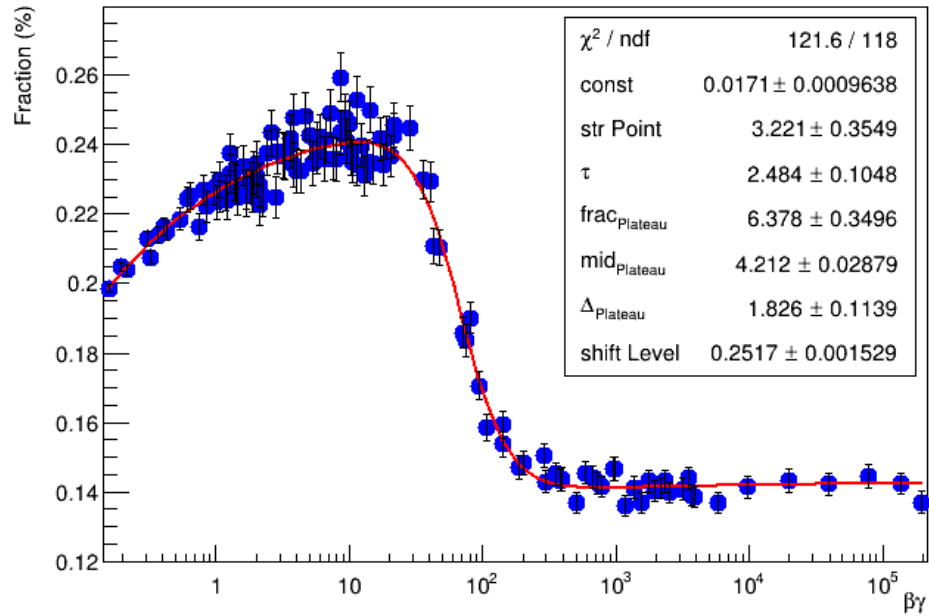
Other parameters studied are:

- Slope and fraction for exp+gaus fit
- Most probable value and sigma for landau fit
- The two MPV and sigma for double landau fit

Sigma from gaus+exp fit of Extra Energy with CISz=1



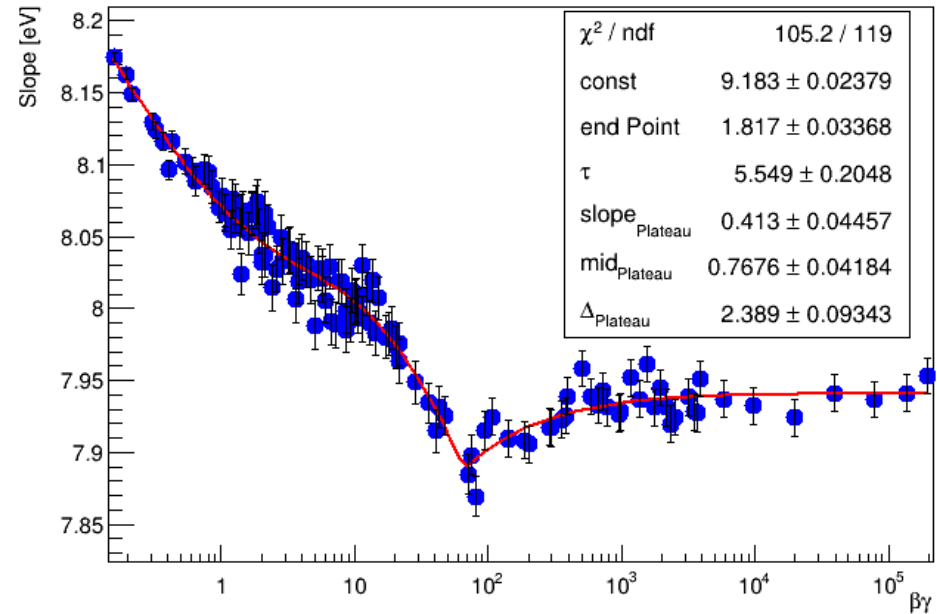
Fraction (%) gaus+exp fit of Extra Energy with CISz=1

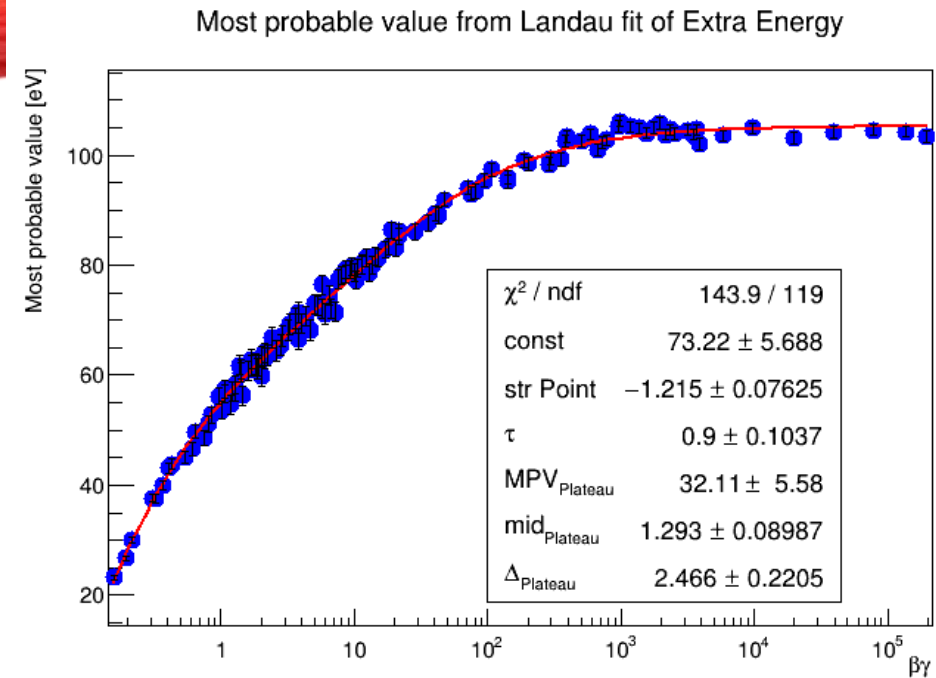


Here the distribution of the slope value of fit, fitted with a 1-exponential function plus an efficiency function.

Here the distribution of the fraction value from fit, fitted with a decreasing exponential function plus an efficiency function.

Slope from gaus+exp fit of Extra Energy with CISz=1

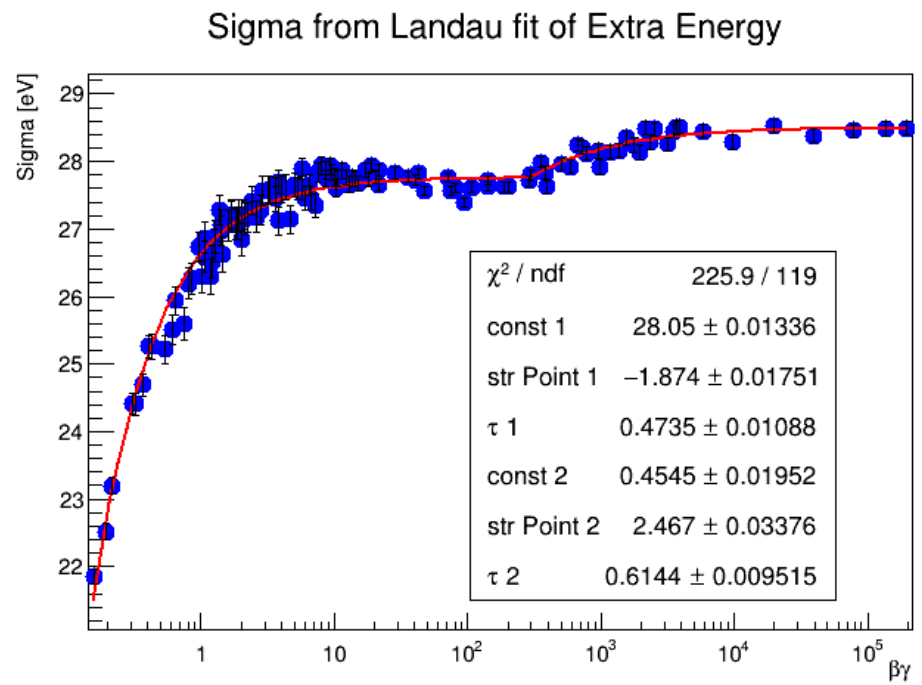




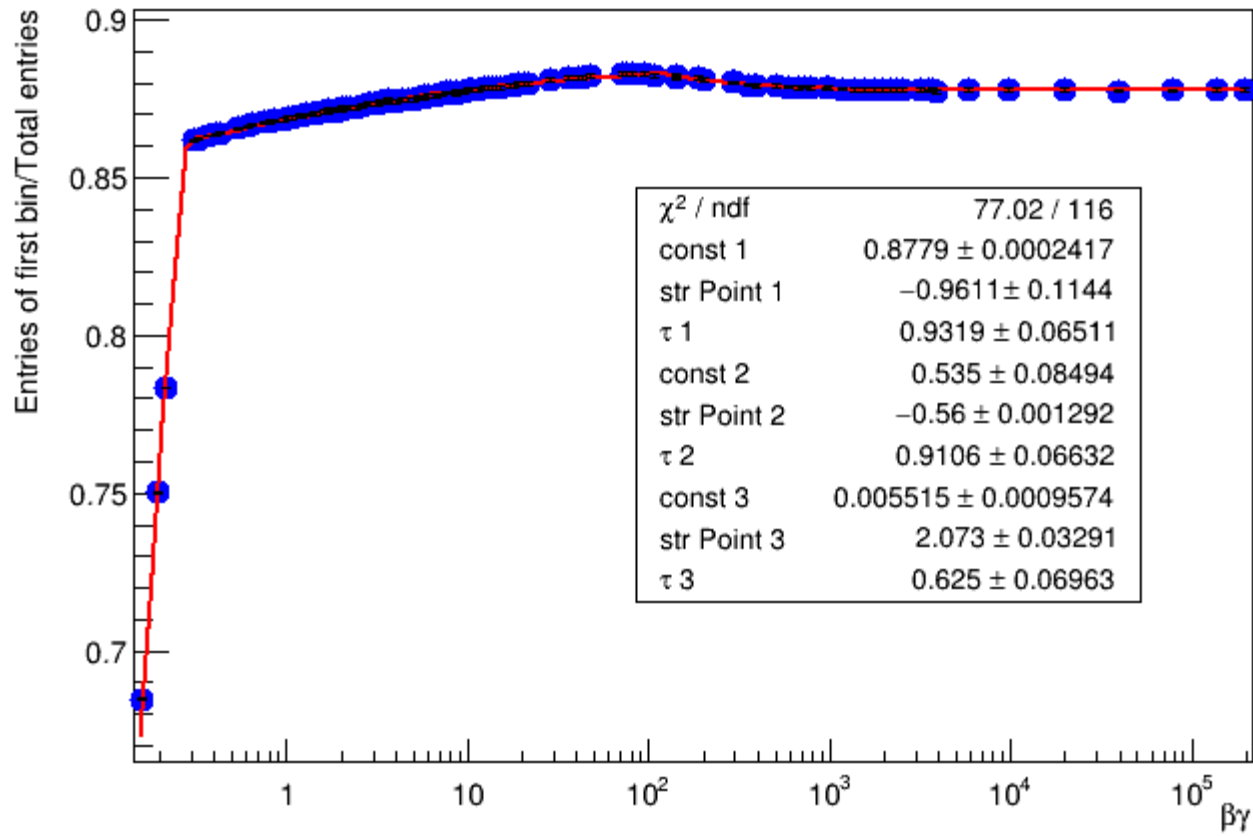
Here the distribution of the sigma value, fitted with 1- decreasing exponential function plus an efficiency function.

Fit parameters for kinetic energy distribution for cluster with cluster size higher than 1

Here the distribution of the most probable value, fitted with 1- decreasing exponential function plus an efficiency function.



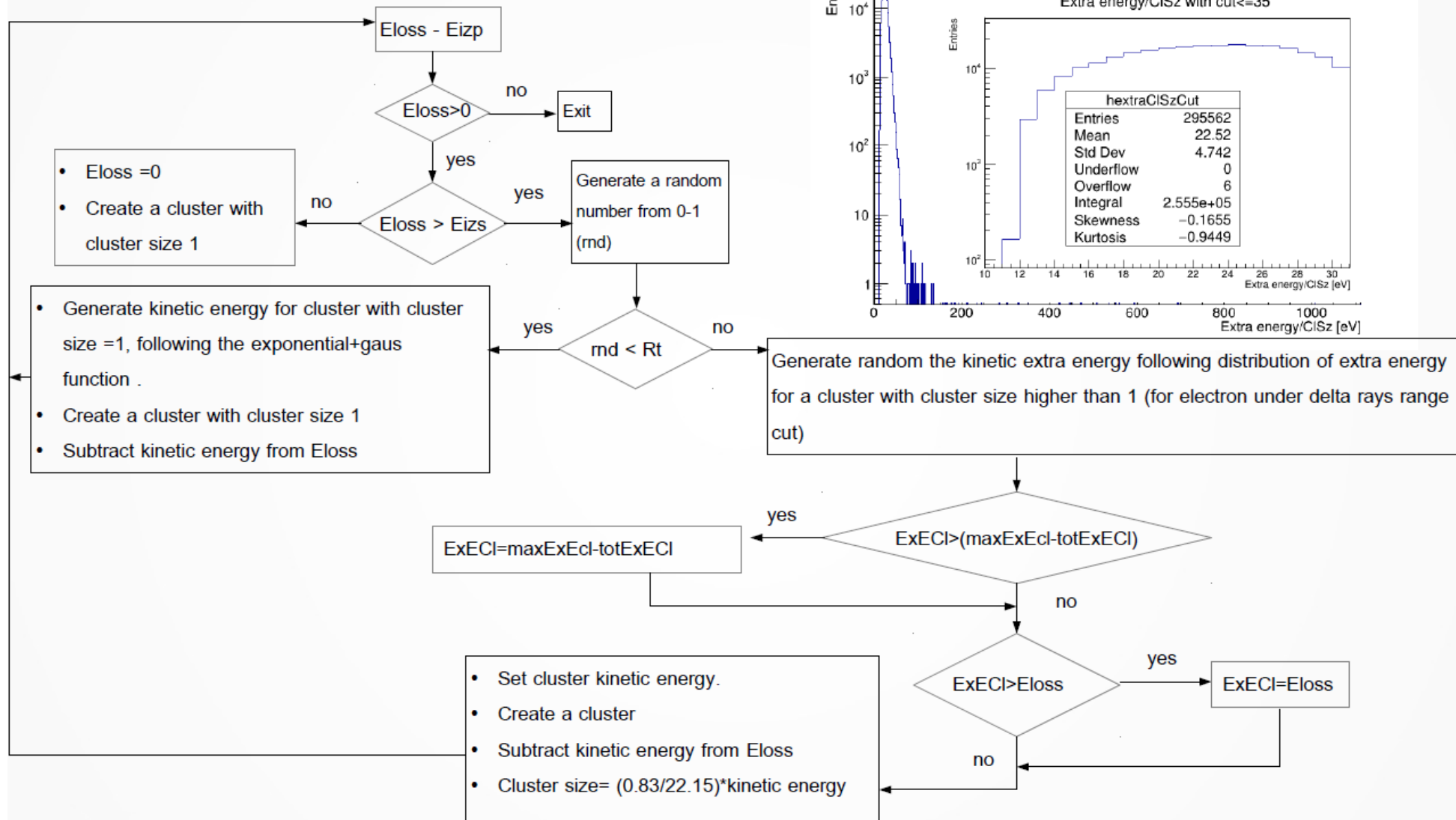
Entries of first bin/Total entries of CISz distribution



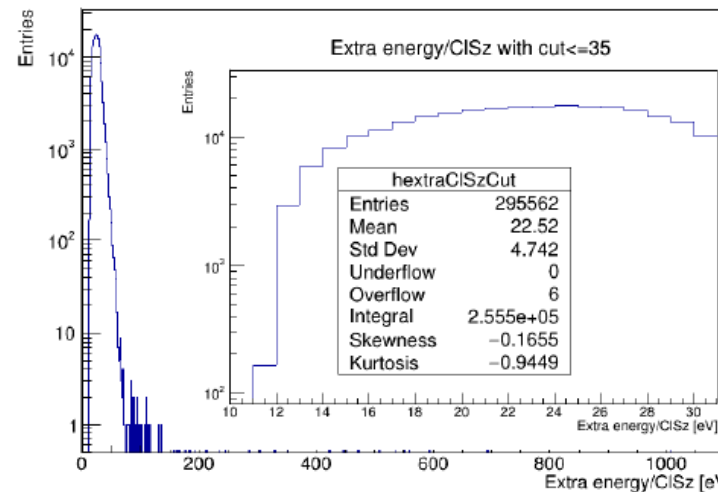
Here the ratio between the cluster containing a single electron and the cluster containing more than an electron.

The distribution is fitted with 1- decreasing exponential function plus two decreasing exponential functions.

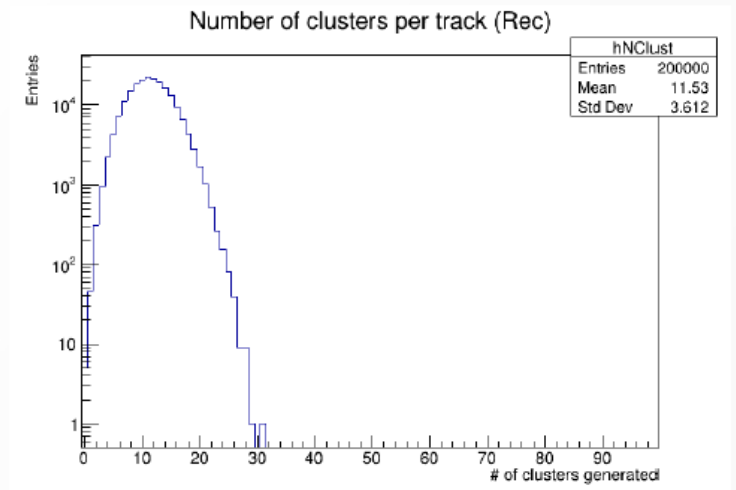
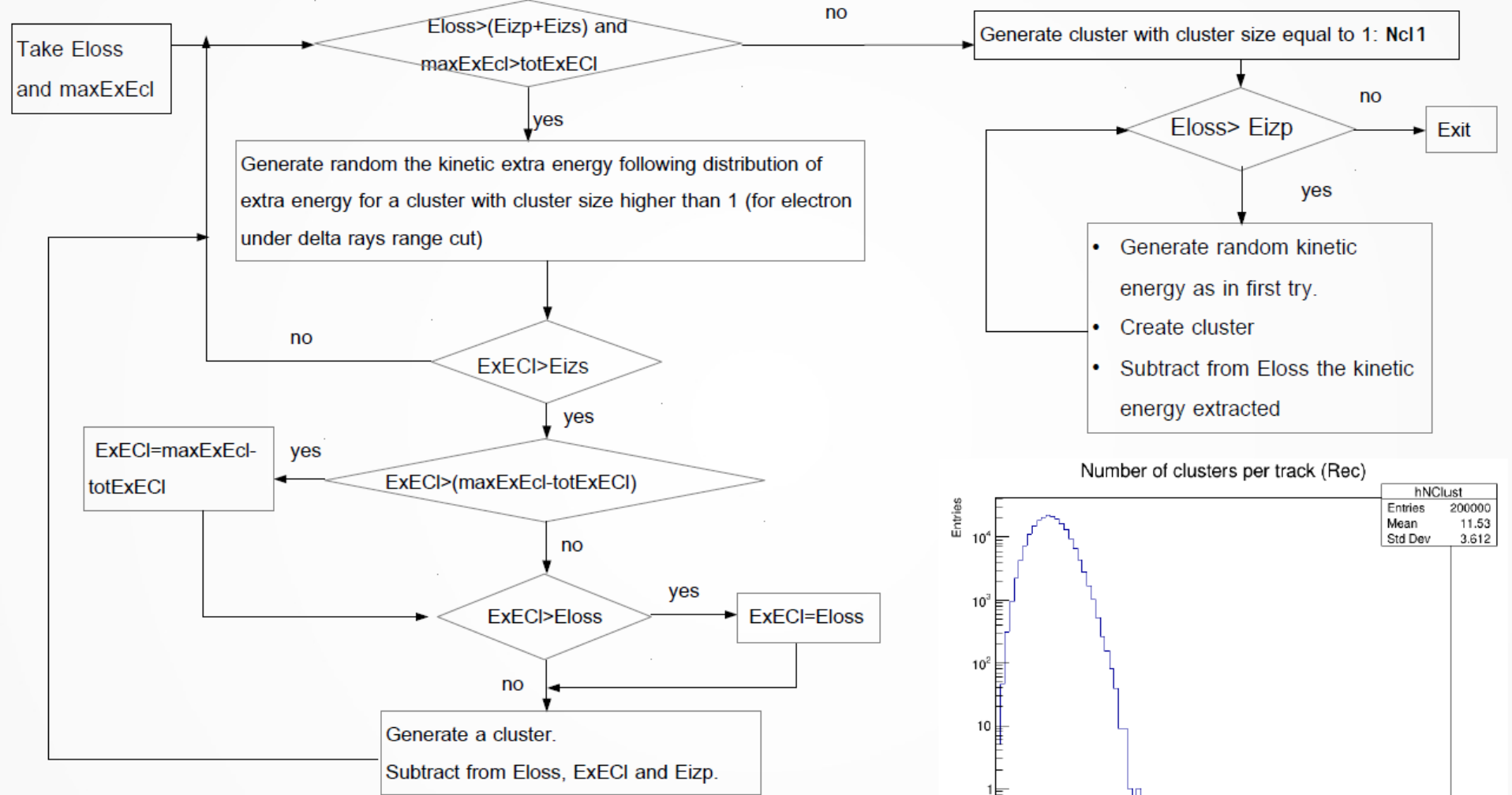
First algorithm



Extra energy/CISz with cut<=35

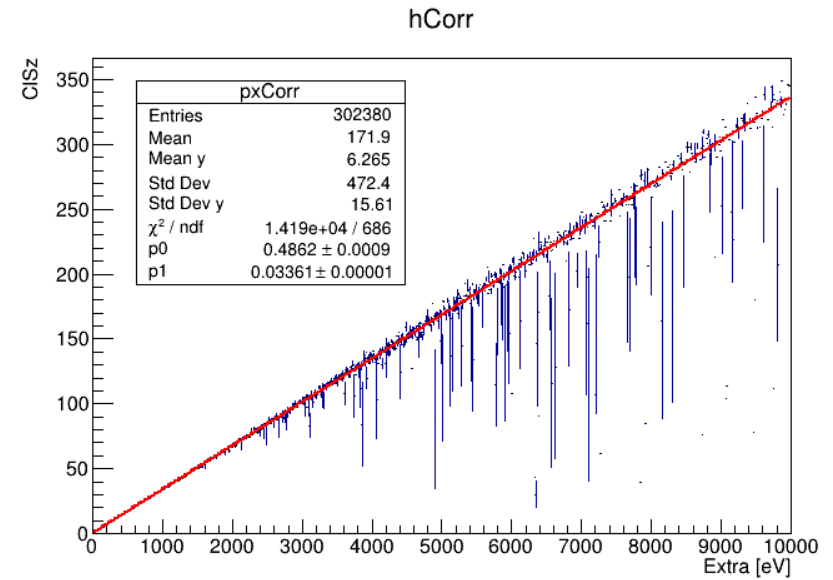
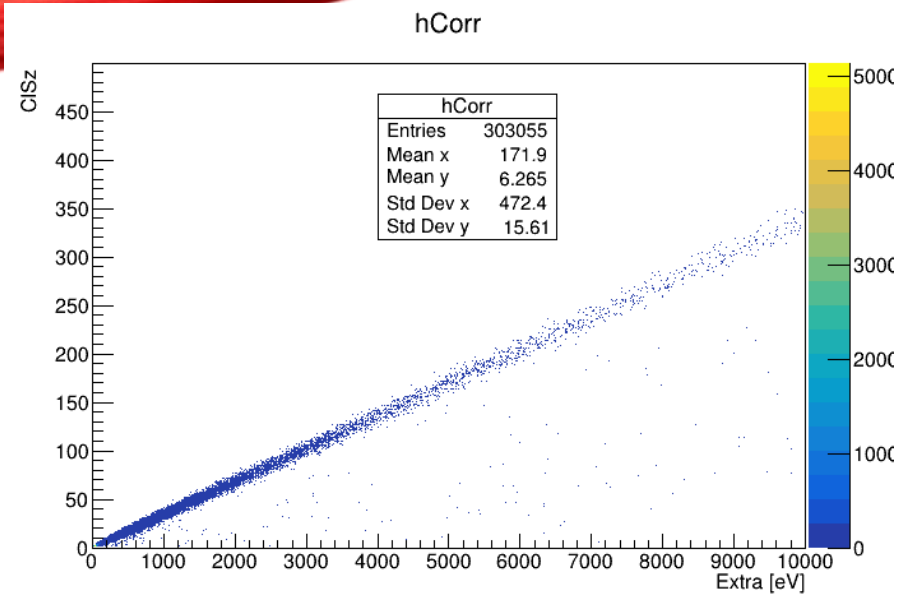


Second algorithm: reconstruction of number of cluster



Second try: reconstruction of cluster size

We studied the relation between extra energy and cluster size for cluster with cluster size higher than 1 (delta rays are included).



We fit the correlation trend with a first degree polynomial and save the parameter p0 and p1.

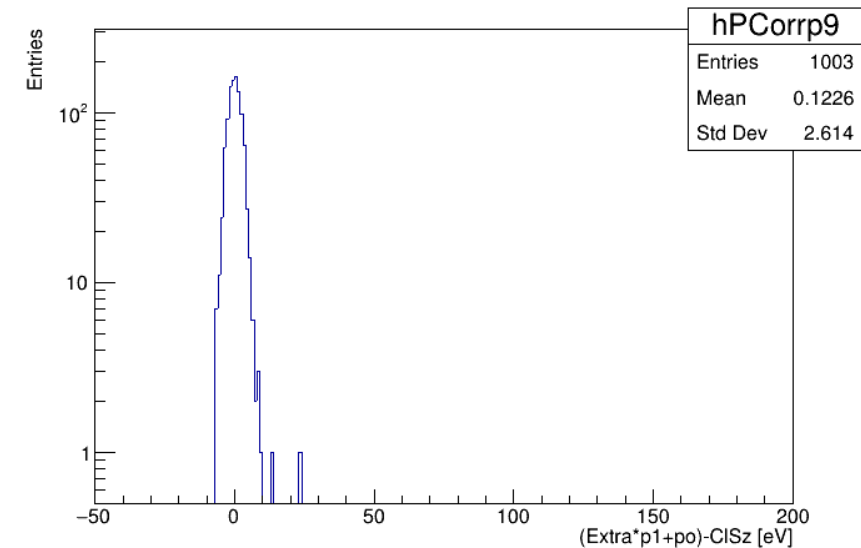
To evaluate the cluster size for cluster with more than one electron, we study the dispersion for different slice of extra energy up to delta rays cut range (1000 eV):

$$(\text{Extra energy} * p1 + p0) - \text{ClSz}$$

The figure shows an example for extra energy between 900 and 1000 eV.

At the end, the cluster size is evaluated as :

$$(\text{Extra energy} * p1 + p0) \text{ minus } \text{hPCorr} \rightarrow \text{GetRandom}()$$



Same evaluation is performed for cluster size generated by delta rays.

We study the dispersion for different slice of extra energy above the value of delta rays cut range (1000 eV):

$$(Extra\ energy * p1 + p0) - ClSz.$$

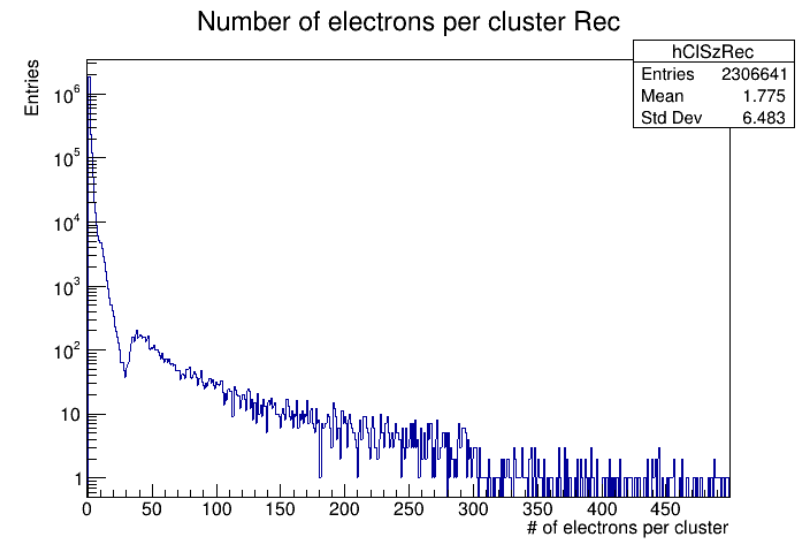
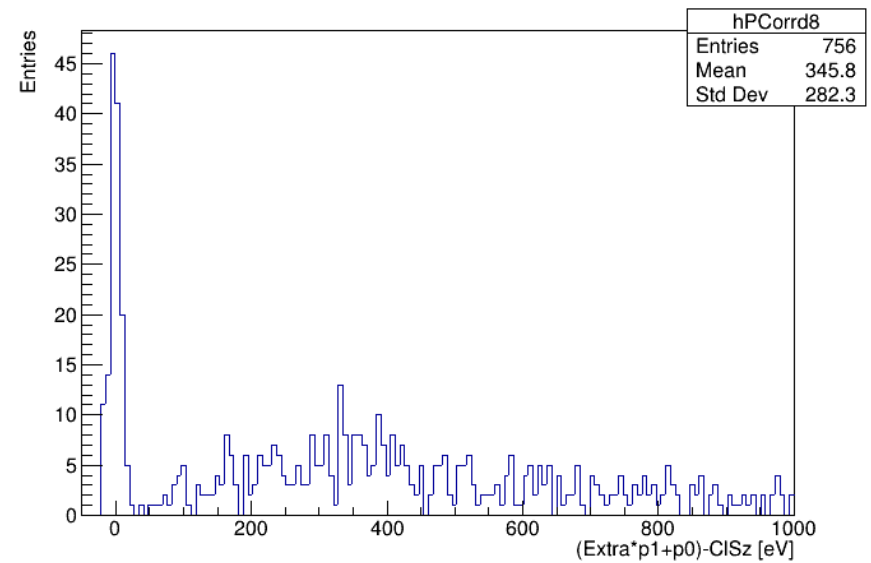
The figure shows a example for extra energy between 9 and 10 keV.

At the end, the cluster size is evaluated as :

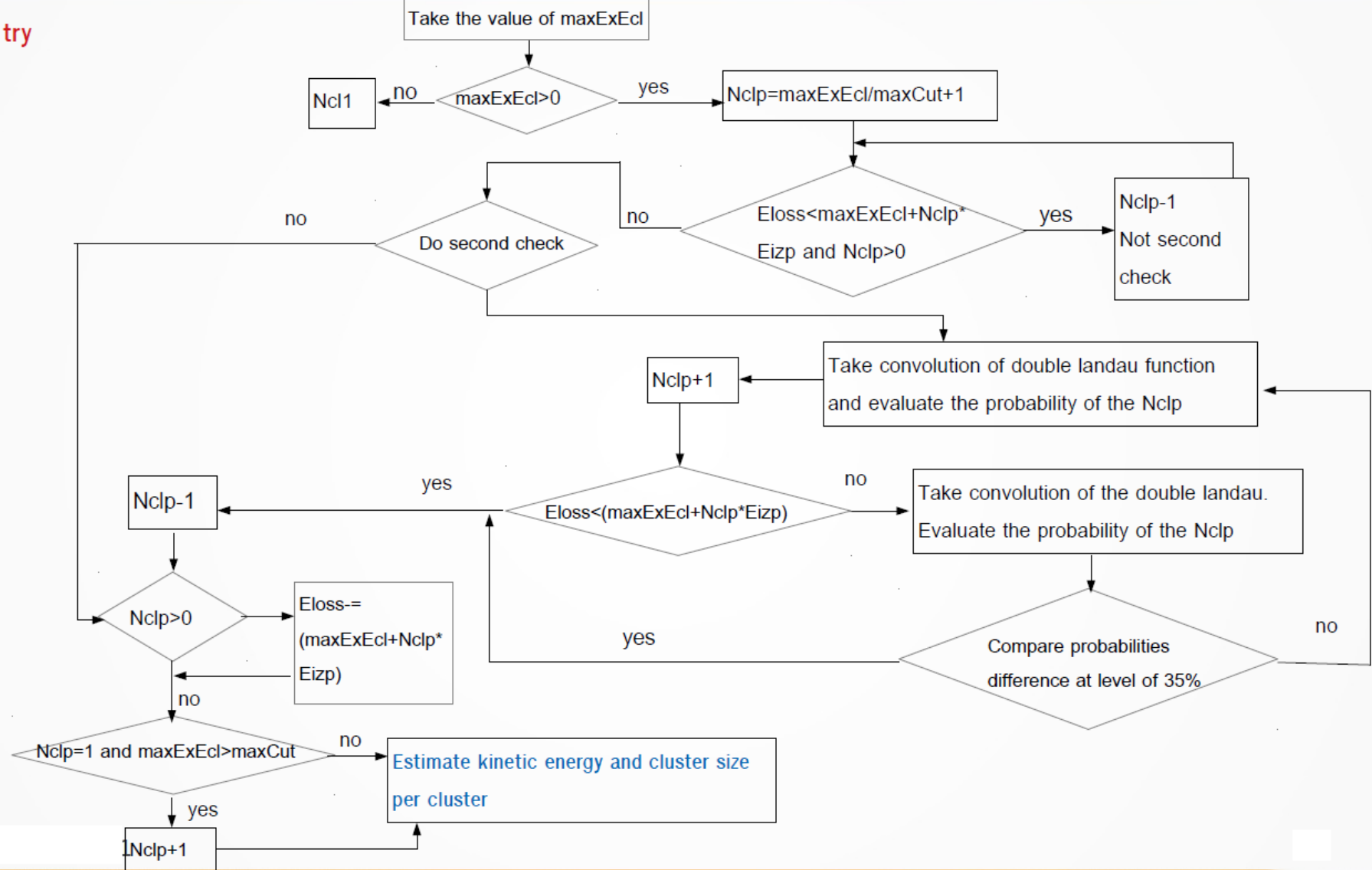
$$(Extra\ energy * p1 + p0) \text{ minus } hPCorrd \rightarrow GetRandom().$$

The result is better than the first try, besides the depression remains.

The evaluation of cluster size remains the same for all next attempts.

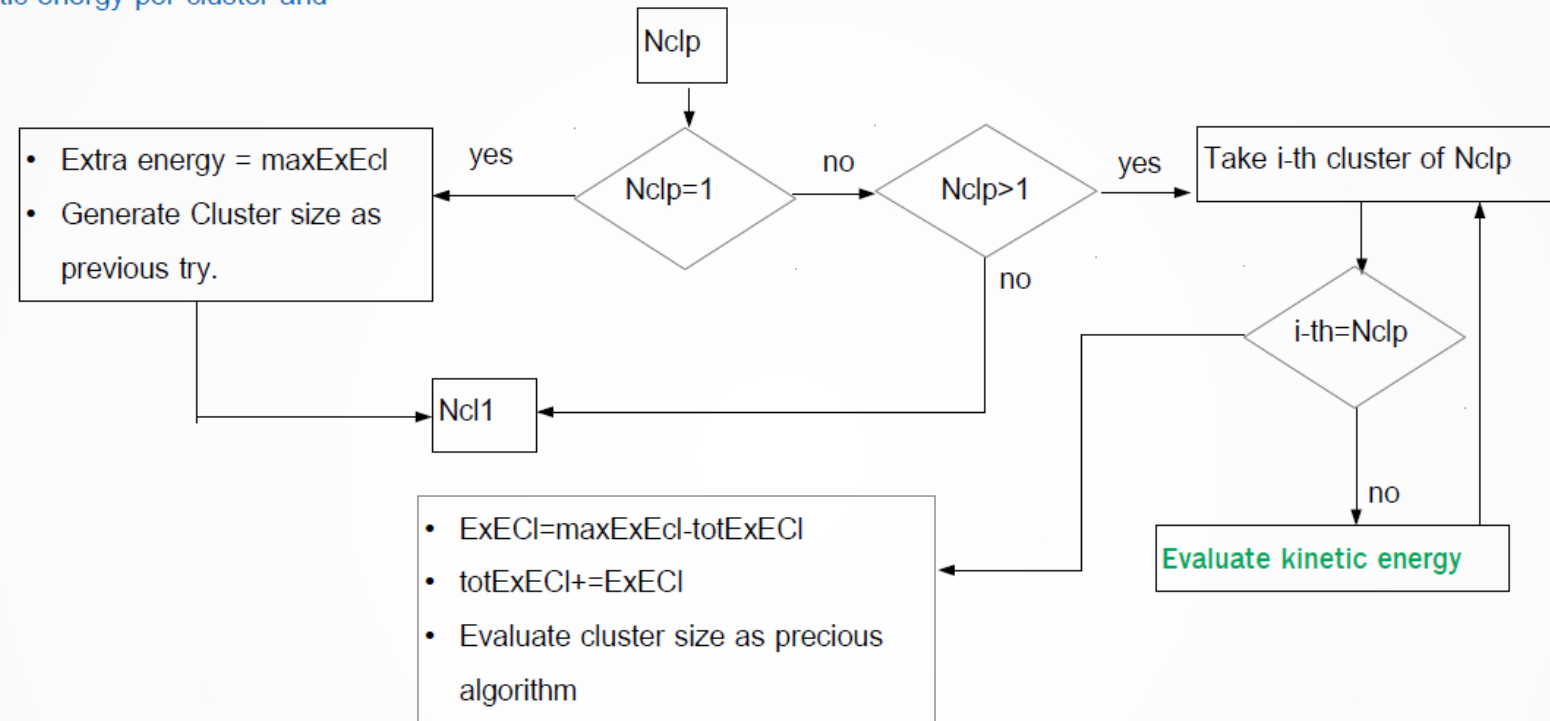


Sixth try



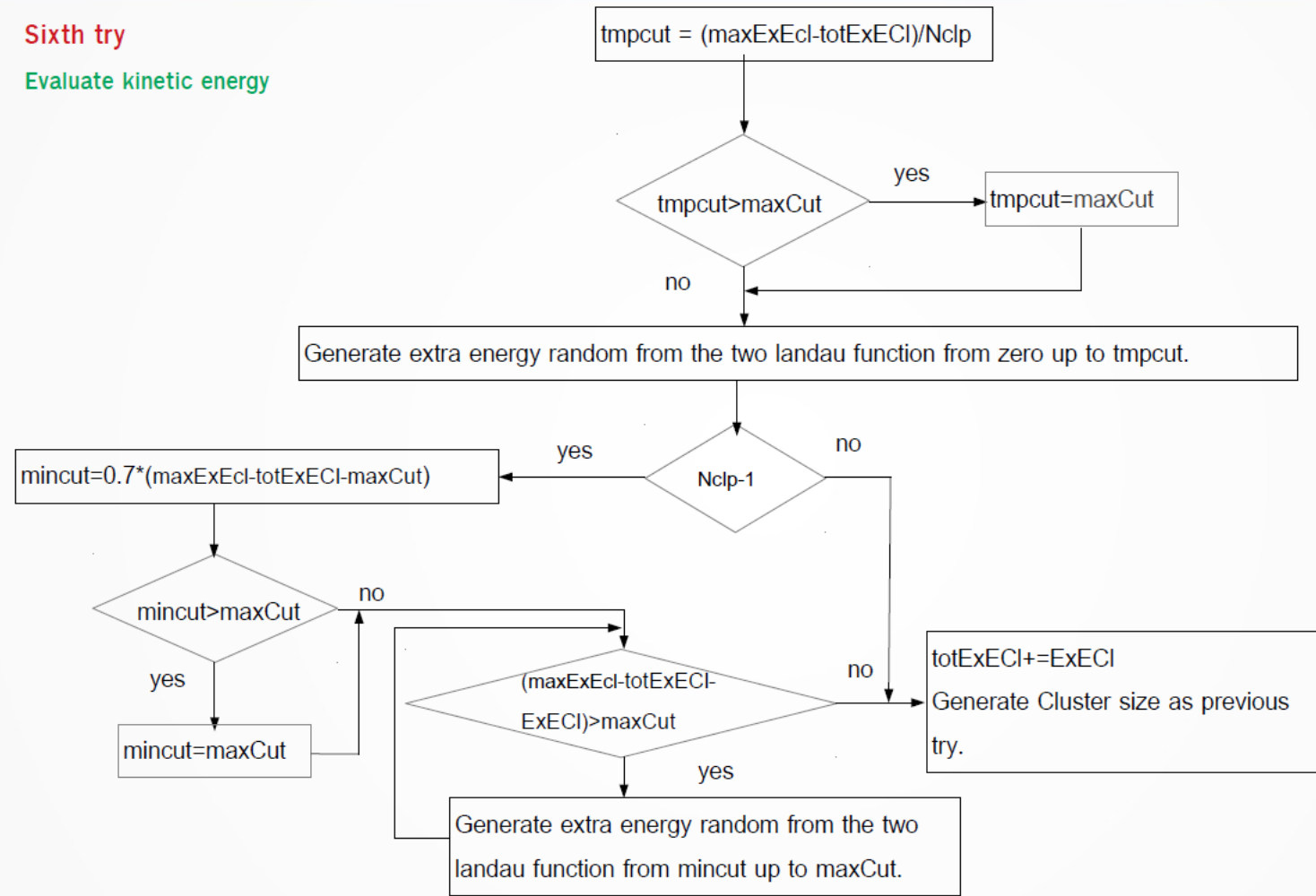
Sixth try

Estimate kinetic energy per cluster and cluster size



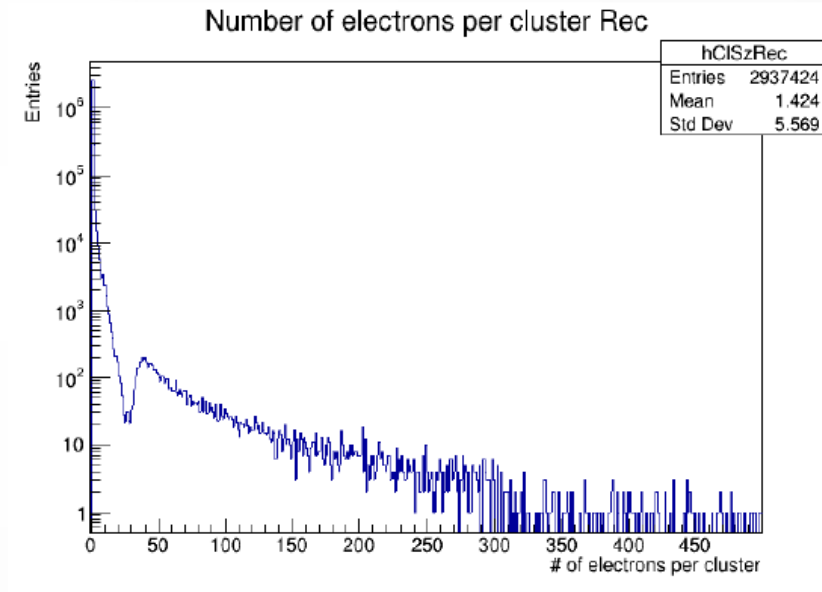
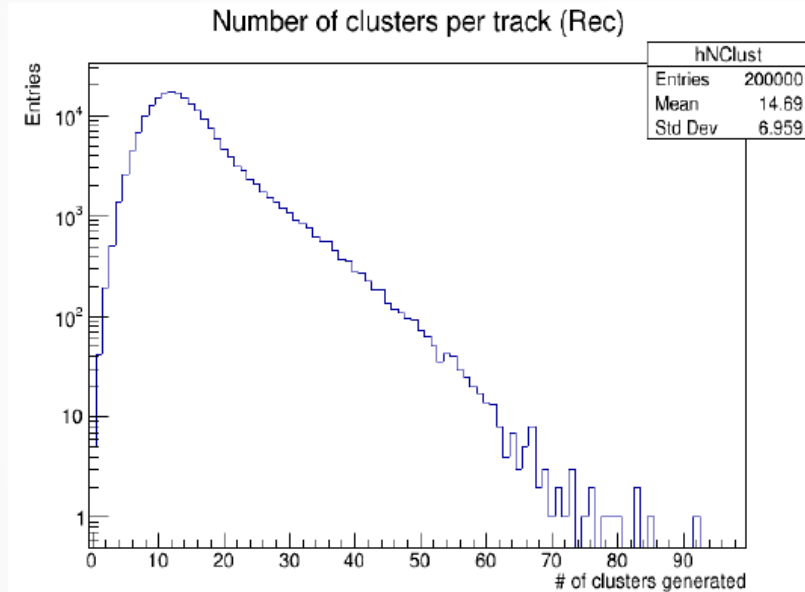
Sixth try

Evaluate kinetic energy



First algorithm

The first algorithm uses a reference value of the ratio between clusters containing a single electron and clusters containing more than one electron (R_t). Using the R_t value, the algorithm chooses to create cluster with cluster size one or higher. Then, it assigns the kinetic energy to each cluster by using the proper distributions. If the cluster has more than one electron, a check on the total kinetic energy is performed and its cluster size is evaluated. The procedure is repeated until the sum of primary ionization energy and kinetic energy per cluster saturate the energy loss of the event.



- The peak is almost in agreement with the expected value, but the distribution of number of clusters is not Poissonian, indeed presents a long fake tail.
- The cluster size presents a depression before the value of 50 electrons.

We change the method for algorithm implementation.

Summary table

	Ncl	σ Ncl	Ncl1	σ Ncl1	Nclp	σ Nclp	maxNclp	#entries of Nclp gen/ #entries of Nclp rec
expected value	11.96	3.458	10.44	3.228	1.912	1.04	10.05	
Algorithm								
1	14.69	6.959	12.85	6.426	2.157	1.25	13.5	1.082
2	11.53	3.612	9.225	3.633	3.448	2.602	25.5	0.899
3 (without correction)	10.99	3.72	9.339	3.608	2.428	1.321	14.5	0.886
3 (with correction)	11.94	3.758	10.25	3.69	2.429	1.317	12.5	0.889
4	11.63	3.642	9.388	3.633	3.349	2.675	24.5	0.889
5	12.11	3.808	9.533	3.935	4.186	2.972	24.5	0.820
6	11.36	3.525	9.501	3.511	2.724	1.311	12.5	0.886
7	7.012	4.026	7.593	3.862	2.286	1.258	12.5	1.295

	CISz	σ CISz
expected value	1.705	6.498
Algorithm		
1	1.424	5.569
2	1.775	6.483
3 (without correction)	1.828	6.695
3 (with correction)	1.762	6.367
4	1.753	6.434
5	1.698	6.231
6	1.787	6.67
7	2.485	9.012

Physics list for geant4

The physics list is “QGSP_BERT” with
G4EmStandardPhysics_option3(),
G4EmExtraPhysics(), G4DecayPhysics(),G4EmLowEPPysics().

From the ordered sequence of the electrons arrival times, considering the average time separation between clusters and their time spread due to diffusion, reconstruct the most probable sequence of clusters drift times: $\{t_i^{cl}\} \quad i = 1, N_{cl}$

$$dE/dx$$

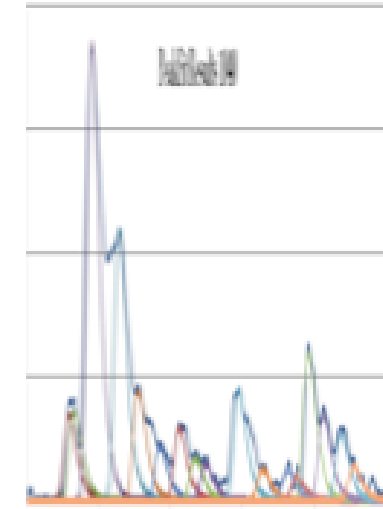
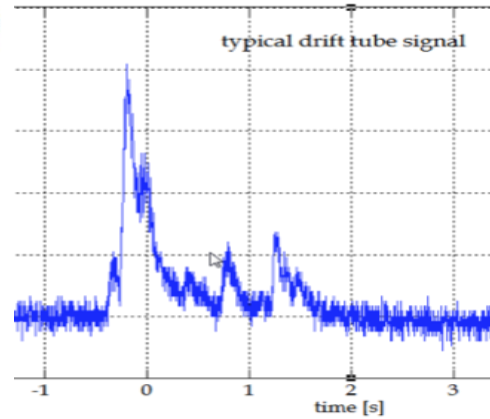
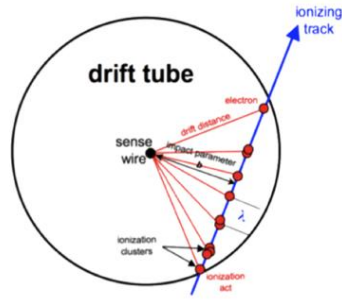
$$\frac{S_{dE/dx}}{(dE/dx)} = 0.41 \times n^{-0.43} \times (L_{track} [m] \times P [atm])^{-0.32}$$

from *Walenta parameterization (1980)*

truncated mean cut (70-80%) reduces the amount of collected information $n = 112$ and a 2m track at 1 atm give

$$\sigma \approx 4.3\%$$

Increasing P to 2 atm improves resolution by 20% ($\sigma \approx 3.4\%$) but at a considerable cost of multiple scattering contribution to momentum and angular resolutions.



$$dN_{cl}/dx$$

$$\frac{S_{dN_{cl}/dx}}{(dN_{cl}/dx)} = (d_{cl} \times L_{track})^{-1/2}$$

from *Poisson distribution*

n =number of sampling layers

$\delta_{cl} = 12.5/cm$ for He/iC₄H₁₀=90/10 and a 2m track give

$$\sigma \approx 2.0\%$$

A small increment of iC₄H₁₀ from 10% to 20% ($\delta_{cl} = 20/cm$) improves resolution by 20% ($\sigma \approx 1.6\%$) at only a reasonable cost of multiple scattering contribution to momentum and angular resolutions.