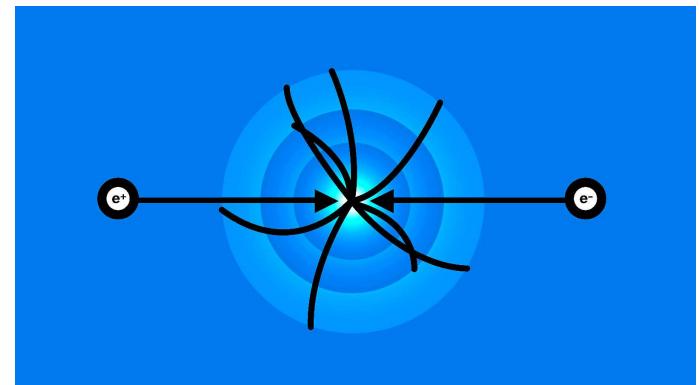


# Importance of Z-pole and WW running

A. Freitas

University of Pittsburgh

- Electroweak precision at  $Z$  pole &  $WW$
- Electroweak precision at  $\sqrt{s} = 250$  GeV
- $\alpha_s$  and  $\alpha(s)$  measurements

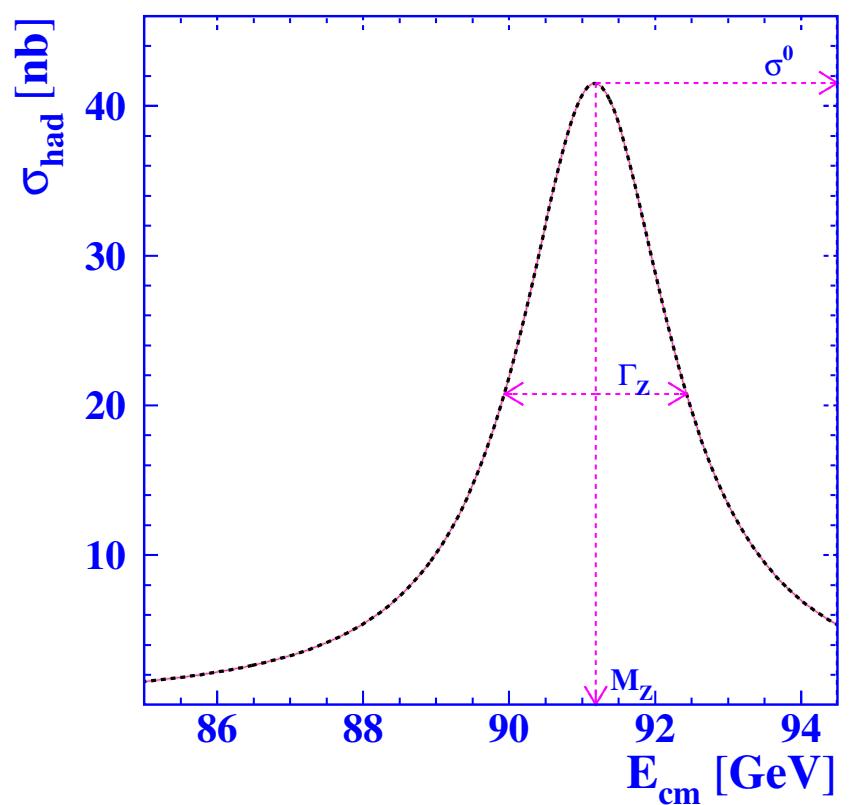
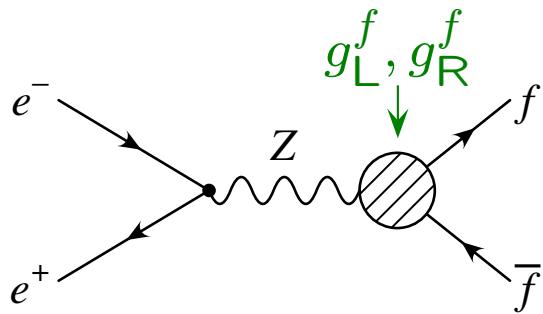


## $Z$ cross section and branching fractions

$e^+e^- \rightarrow f\bar{f}$  for  $E_{\text{CM}} \sim M_Z$ :

- Mass  $M_Z$
- Width  $\Gamma_Z = \sum_f \Gamma_{ff}$
- Braching ratio  $R_f = \Gamma_{ff}/\Gamma_Z$
- $\sigma^0 \approx \frac{12\pi \Gamma_{ee} \Gamma_{ff}}{(s - M_Z^2)^2 + M_Z^2 \Gamma_Z^2} = \frac{12\pi}{M_Z^2} R_e R_f$

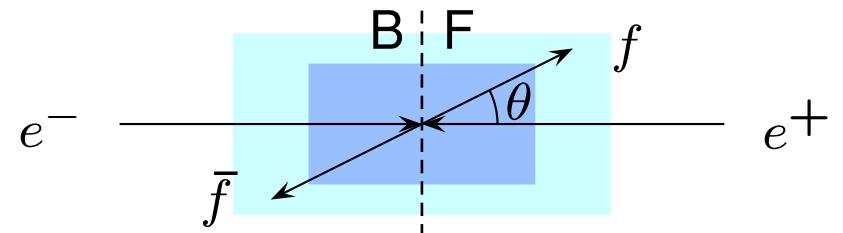
$$\Gamma_{ff} = C \left[ (g_L^f)^2 + (g_R^f)^2 \right]$$



Forward-backward asymmetry:

$$A_{FB} \equiv \frac{\sigma_F - \sigma_B}{\sigma_F + \sigma_B} = \frac{3}{4} \mathcal{A}_e \mathcal{A}_f$$

$$\mathcal{A}_f = \frac{2(1 - 4 \sin^2 \theta_{\text{eff}}^f)}{1 + (1 - 4 \sin^2 \theta_{\text{eff}}^f)^2}$$



$$\sin^2 \theta_{\text{eff}}^f = \frac{g_R^f}{2|Q_f|(g_R^f - g_L^f)}$$

Left-right asymmetry:

With polarized  $e^-$  beam:  $A_{LR} \equiv \frac{\sigma_L - \sigma_R}{\sigma_L + \sigma_R} = \mathcal{A}_e$

Polarization asymmetry:

Average  $\tau$  pol. in  $e^+ e^- \rightarrow \tau^+ \tau^-$ :  $\langle \mathcal{P}_\tau \rangle = -\mathcal{A}_\tau$

- Deconvolution of initial-state QED radiation:

$$\sigma[e^+e^- \rightarrow f\bar{f}] = \mathcal{R}_{\text{ini}}(s, s') \otimes \sigma_{\text{hard}}(s')$$

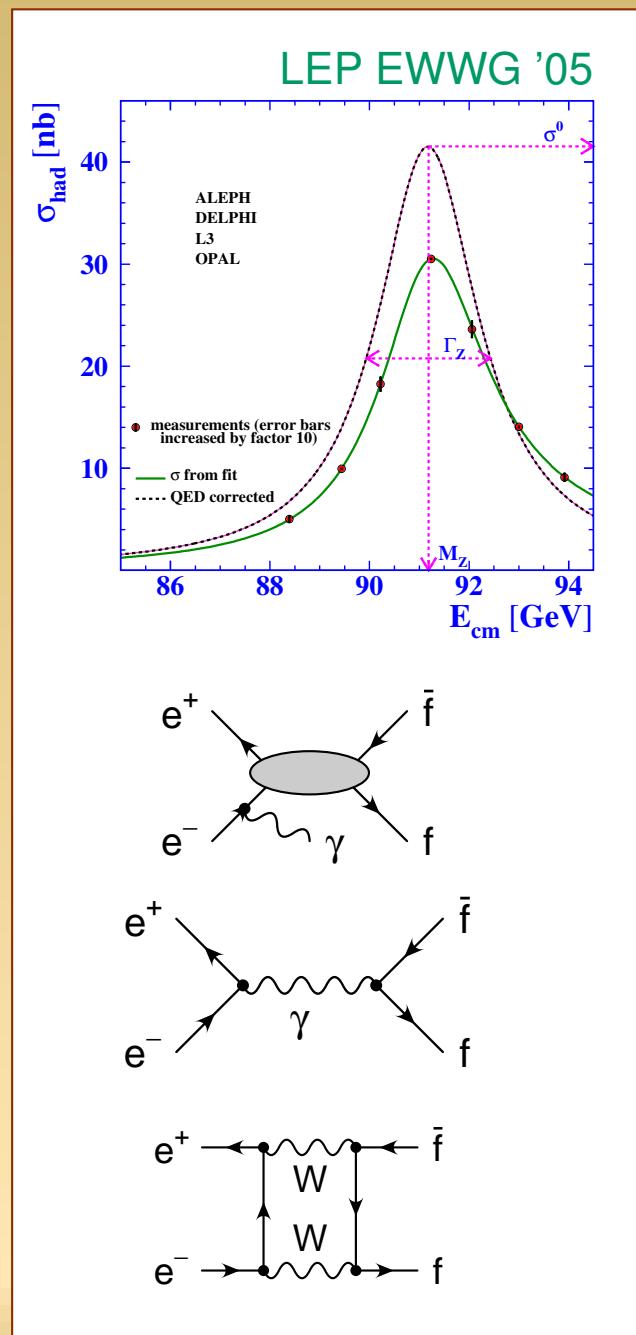
- Subtraction of  $\gamma$ -exchange,  $\gamma-Z$  interference, box contributions:

$$\sigma_{\text{hard}} = \sigma_Z + \sigma_\gamma + \sigma_{\gamma Z} + \sigma_{\text{box}}$$

- $Z$ -pole contribution:

$$\sigma_Z = \frac{R}{(s - M_Z^2)^2 + M_Z^2 \Gamma_Z^2} + \sigma_{\text{non-res}}$$

$\sigma_\gamma, \sigma_{\gamma Z}, \sigma_{\text{box}}, \sigma_{\text{non-res}}$  known at NLO  
 → need consistent pole expansion framework  
 → leading NNLO may be needed for future  $e^+e^-$



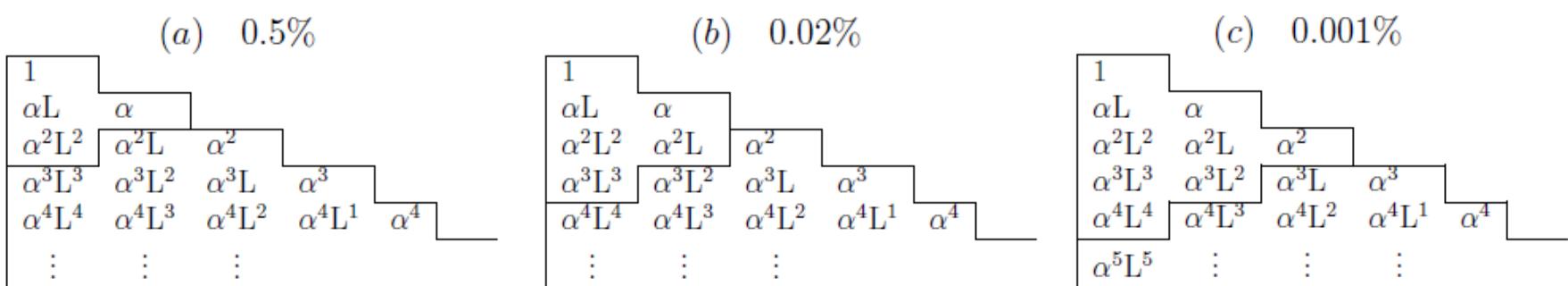
# Monte-Carlo methods for QED effects

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- Implementation in MC program to evaluate exp. efficiency and particle ID
  - Current state of art: e.g. KORALZ, KKMC Jadach, Ward, ...  
→  $\mathcal{O}(\alpha^2 L)$  accuracy [ $L = \ln(s/m_e^2)$ ]
  - One to two orders improvement needed:

Observable	Where from	Present (LEP)	FCC stat.	FCC syst	Now FCC
$M_Z$ [MeV]	Z linesh. [28]	$91187.5 \pm 2.1\{0.3\}$	0.005	0.1	3
$\Gamma_Z$ [MeV]	Z linesh. [28]	$2495.2 \pm 2.1\{0.2\}$	0.008	0.1	2
$R_l^Z = \Gamma_h/\Gamma_l$	$\sigma(M_Z)$ [33]	$20.767 \pm 0.025\{0.012\}$	$6 \cdot 10^{-5}$	$1 \cdot 10^{-3}$	12
$\sigma_{\text{had}}^0$ [nb]	$\sigma_{\text{had}}^0$ [28]	$41.541 \pm 0.037\{0.25\}$	$0.1 \cdot 10^{-3}$	$4 \cdot 10^{-3}$	6
$N_\nu$	$\sigma(M_Z)$ [28]	$2.984 \pm 0.008\{0.006\}$	$5 \cdot 10^{-6}$	$1 \cdot 10^{-3}$	6
$\sin^2 \theta_W^{eff} \times 10^5$	$A_{FB}^{lept.}$ [33]	$23099 \pm 53\{28\}$	0.3	0.5	55
$A_{FB,\mu}^{M_Z \pm 3.5 \text{ GeV}}$	$\frac{d\sigma}{d \cos \theta}$ [28]	$\pm 0.020\{0.001\}$	$1.0 \cdot 10^{-5}$	$0.3 \cdot 10^{-5}$	100

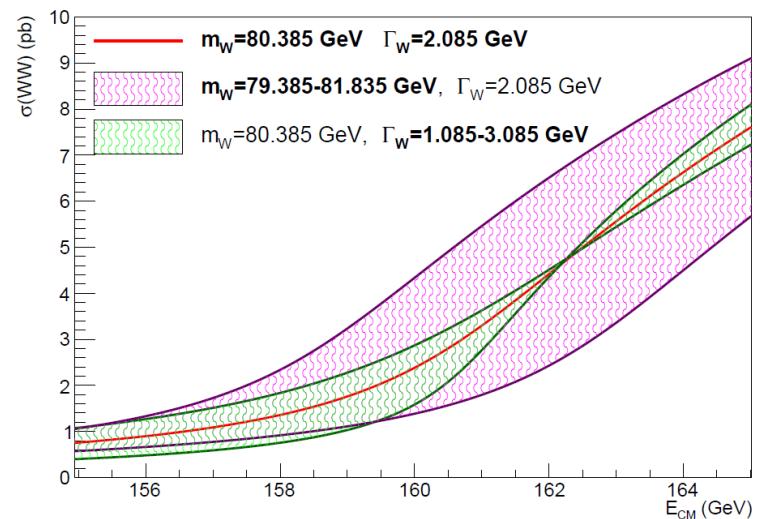
Jadach,  
Skrzypek '19



→ Need matching of h.o. matrix elements with QED parton shower  
(exclusive in all fs particles)

- High-precision measurement of  $M_W$  from  $e^+e^- \rightarrow W^+W^-$  at threshold
- a) Corrections near threshold enhanced by  $1/\beta$  and  $\ln \beta$   

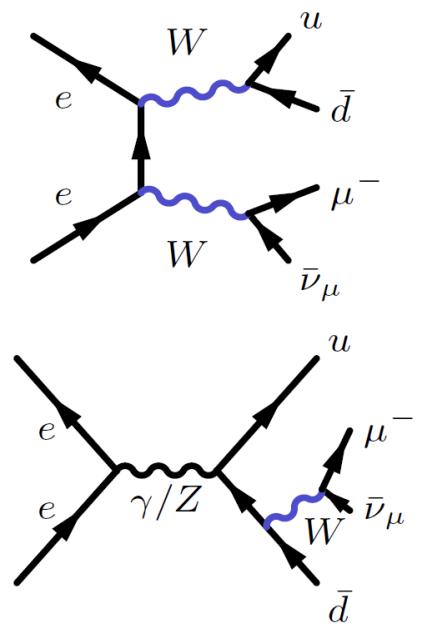
$$\beta \sim \sqrt{1 - 4 \frac{M_W^2 - i M_W \Gamma_W}{s}} \sim \sqrt{\Gamma_W/M_W}$$
- b) Non-resonant contributions are important



- Full  $\mathcal{O}(\alpha)$  calculation of  $e^+e^- \rightarrow 4f$   
Denner, Dittmaier, Roth, Wieders '05

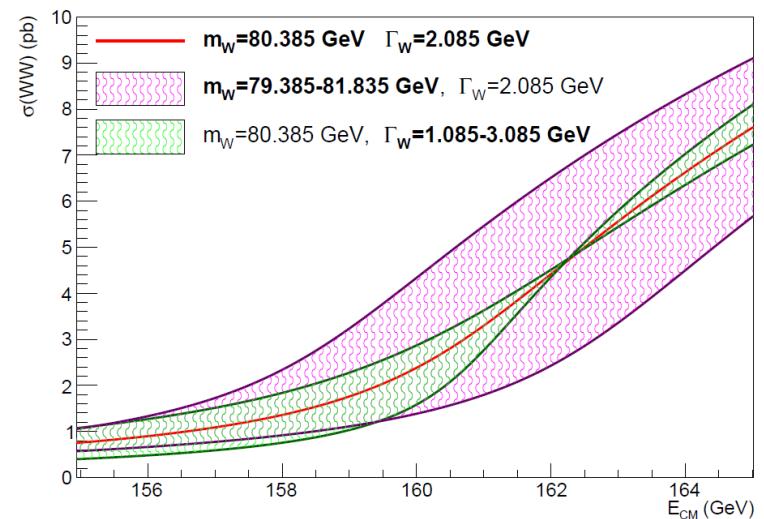
- EFT expansion in  $\alpha \sim \Gamma_W/M_W \sim \beta^2$   
Beneke, Falgari, Schwinn, Signer, Zanderighi '07

- NLO corrections with NNLO Coulomb correction ( $\propto 1/\beta^n$ ):  $\delta_{\text{th}} M_W \sim 3 \text{ MeV}$   
Actis, Beneke, Falgari, Schwinn '08
- Adding NNLO corrections to  $ee \rightarrow WW$  and  $W \rightarrow f\bar{f}$  and NNLO ISR:  $\delta_{\text{th}} M_W \lesssim 0.6 \text{ MeV}$



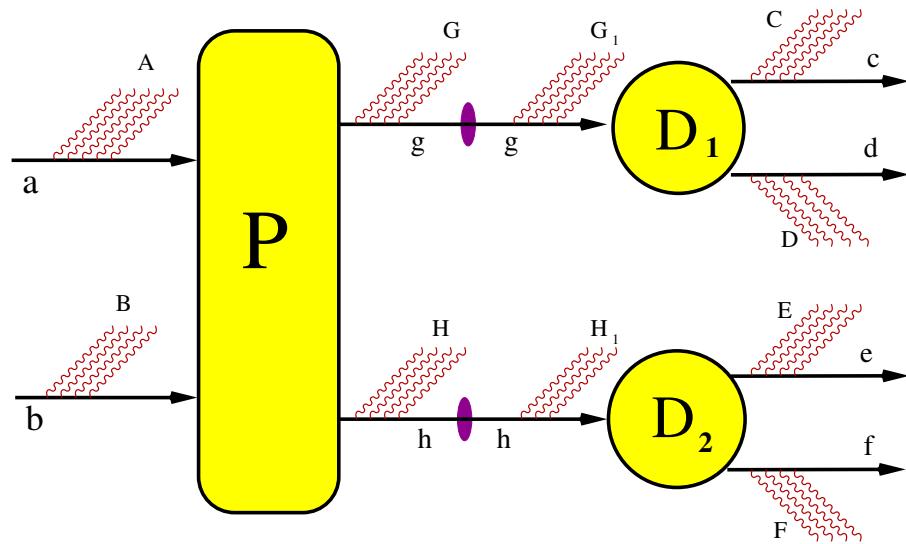
- High-precision measurement of  $M_W$  from  $e^+e^- \rightarrow W^+W^-$  at threshold
- a) Corrections near threshold enhanced by  $1/\beta$  and  $\ln \beta$ 

$$\beta \sim \sqrt{1 - 4 \frac{M_W^2 - i M_W \Gamma_W}{s}} \sim \sqrt{\Gamma_W/M_W}$$
- b) Non-resonant contributions are important



- Resummation of soft photon radiation

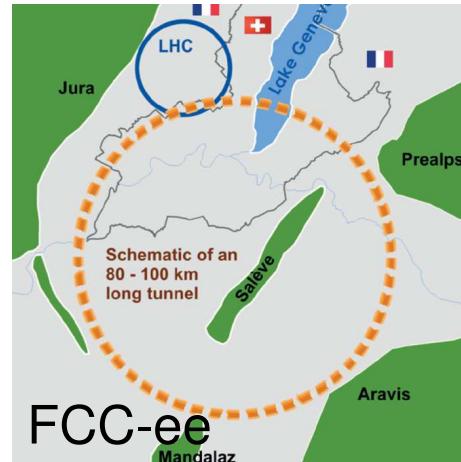
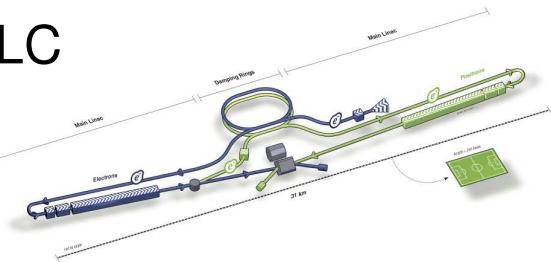
Jadach, Płaczek, Skrzypek '19



# Electroweak precision tests at future colliders

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ILC



$\sqrt{s}$	$M_Z$	$2M_W$
ILC/GigaZ	$100 \text{ fb}^{-1}$	$500 \text{ fb}^{-1}$ (6 pts.)
FCC-ee	$230 \text{ ab}^{-1}$	$10 \text{ ab}^{-1}$ (2 pts.)
CEPC	$45 \text{ ab}^{-1}$	$2.6 \text{ ab}^{-1}$ (3 pts.)

beam pol. ( $P_{e^-}=0.8$ ,  $P_{e^+}=0.3$ )

2 detectors

2 detectors

Anticipated precision for EWPOs:

	Current exp.	ILC/GigaZ	CEPC	FCC-ee
$M_W$ [MeV]	15	1–2 <sup>a,e</sup>	1 <sup>e</sup>	1 <sup>e</sup>
$M_Z$ [MeV]	2.1	–	0.5 <sup>e</sup>	0.1 <sup>e</sup>
$\Gamma_Z$ [MeV]	2.3	1 <sup>a</sup>	0.5 <sup>e</sup>	0.1 <sup>e</sup>
$R_\ell = \Gamma_Z^{\text{had}} / \Gamma_Z^\ell$ [ $10^{-3}$ ]	25	6 <sup>b</sup>	2 <sup>b</sup>	1 <sup>b</sup>
$R_b = \Gamma_Z^b / \Gamma_Z^{\text{had}}$ [ $10^{-5}$ ]	66	15 <sup>c</sup>	4.3 <sup>c</sup>	6 <sup>c</sup>
$\sin^2 \theta_{\text{eff}}^\ell$ [ $10^{-5}$ ]	16	1 <sup>d</sup>	<1 <sup>e</sup>	0.5 <sup>e</sup>

## Systematics:

<sup>a</sup> energy scale

<sup>b</sup> acceptance

<sup>c</sup> flavor tagging

<sup>d</sup> polarization

<sup>e</sup> beam energy calibration / beam-beam interactions

- To probe new physics, compare EWPOs with SM theory predictions
- Need to take theory error into account:

	Current exp.	Current th. <sup>†</sup>	CEPC	FCC-ee
$M_W$ [MeV]	15	4 *	1	1
$\Gamma_Z$ [MeV]	2.3	0.4	0.5	0.1
$R_\ell = \Gamma_Z^{\text{had}} / \Gamma_Z^\ell [10^{-3}]$	25	5	2	1
$R_b = \Gamma_Z^b / \Gamma_Z^{\text{had}} [10^{-5}]$	66	10	4.3	6
$\sin^2 \theta_{\text{eff}}^\ell [10^{-5}]$	16	4.5	<1	0.5

\* computed from  $G_\mu$

† full NNLO and leading NNNLO

- Theory error estimate is not well defined, ideally  $\Delta_{\text{th}} \ll \Delta_{\text{exp}}$
- Common methods:
  - Count prefactors ( $\alpha, N_c, N_f, \dots$ )
  - Extrapolation of perturbative series
  - Renormalization scale dependence
  - Renormalization scheme dependence

- **Electroweak precision tests** at future  $e^+e^-$  colliders require 1–2 orders improvement in SM theory calculations and tools
  - **Z-pole**: 3-loop & leading 4-loop EW + multi-loop/leg merging for QED MC
  - **off Z-pole / backgrounds**: ( $\geq 2$ )-loop EW
  - **WW** 2-loop EW for  $2 \rightarrow 2$  processes (+ 4-loop QCD)  
( $\geq 1$ )-loop for backgr. and non-resonant terms

**SMEFT:** Gauge-invariant operators with SU(2) Higgs doublet

$$\mathcal{L} = \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i + \mathcal{O}(\Lambda^{-3}) \quad (\Lambda \gg M_Z)$$

$$\mathcal{O}_{\phi 1} = (D_\mu \Phi)^\dagger \Phi \Phi^\dagger (D^\mu \Phi)$$

$$\alpha \Delta \textcolor{blue}{T} = -\frac{v^2}{2} \frac{c_{\phi 1}}{\Lambda^2}$$

$$\mathcal{O}_{BW} = \Phi^\dagger B_{\mu\nu} W^{\mu\nu} \Phi$$

$$\alpha \Delta S = -e^2 v^2 \frac{c_{BW}}{\Lambda^2}$$

$$\mathcal{O}_{LL}^{(3)e} = (\bar{L}_L^e \sigma^a \gamma_\mu L_L^e)(\bar{L}_L^e \sigma^a \gamma^\mu L_L^e)$$

$$\Delta G_{\textcolor{blue}{F}} = -\sqrt{2} \frac{c_{LL}^{(3)e}}{\Lambda^2}$$

$$\mathcal{O}_R^f = i(\Phi^\dagger \overset{\leftrightarrow}{D}_\mu \Phi)(\bar{f}_R \gamma^\mu f_R)$$

$$f = e, \mu, \tau, b, lq$$

$$\mathcal{O}_L^F = i(\Phi^\dagger \overset{\leftrightarrow}{D}_\mu \Phi)(\bar{F}_L \gamma^\mu F_L)$$

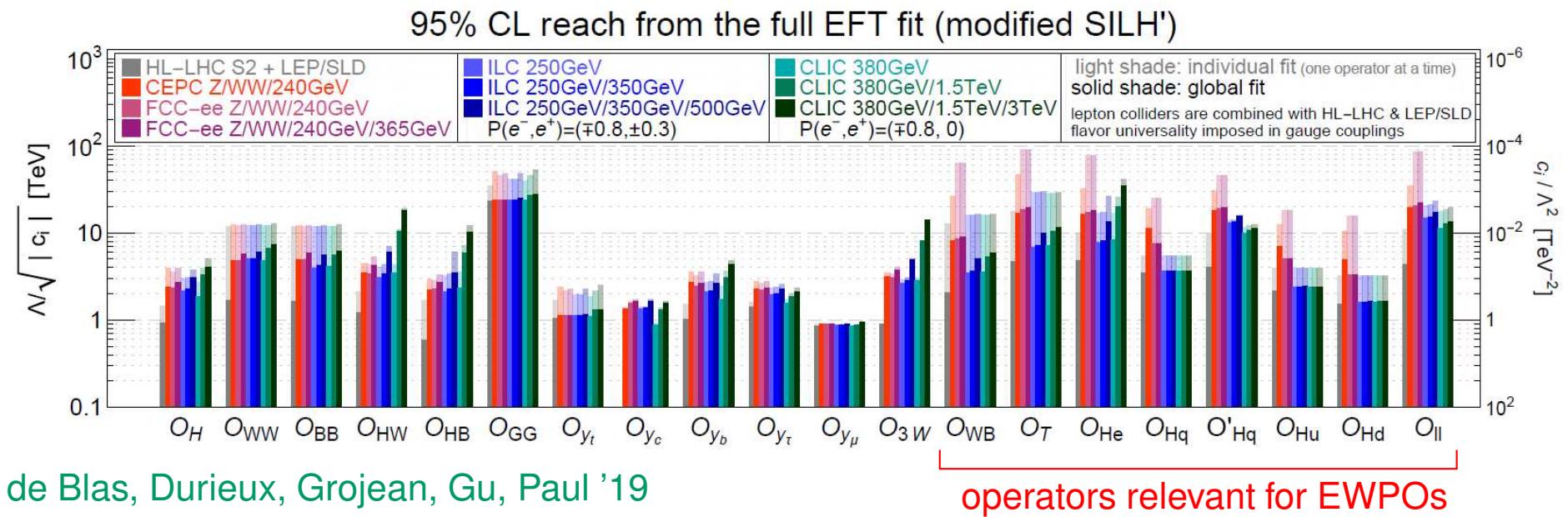
$$F = \binom{\nu_e}{e}, \binom{\nu_\mu}{\mu}, \binom{\nu_\tau}{\tau}, \binom{u, c}{d, s}, \binom{t}{b}$$

$$\mathcal{O}_L^{(3)F} = i(\Phi^\dagger \overset{\leftrightarrow}{D}_\mu^a \Phi)(\bar{F}_L \sigma_a \gamma^\mu F_L)$$

More operators than EWPOs

- Need to make flavor assumptions and/or  
use other obs. (e.g.  $W$  production and decay)

Projected reach assuming Minimal Flavor Violation:



de Blas, Durieux, Grojean, Gu, Paul '19

EWPOs accessible through **radiative return**  $e^+ e^- \rightarrow \gamma Z$

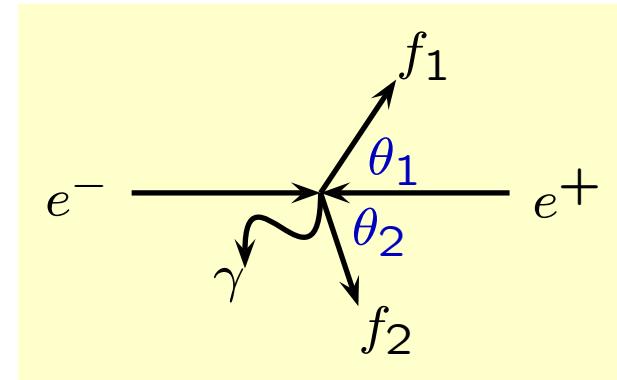
- $\gamma$  mostly collinear with beam
- Reduction in cross-section by

$$\sim \frac{\alpha}{\pi} \ln \frac{s}{m_e^2} \sim 0.06$$

- Precise det. of  $m_{ff}$  from measured angles:

$$m_{ff}^2 = s \frac{1 - \beta}{1 + \beta}, \quad \beta = \frac{|\sin(\theta_1 + \theta_2)|}{\sin \theta_1 + \sin \theta_2}$$

- Additional backgrounds from  $e^+ e^- \rightarrow WW, ZZ$  that are not flat in  $m_{ff}$

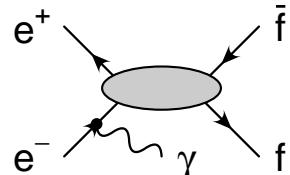


Ueno '19

Fujii et al. '19

- $A_{LR} \rightarrow \sin^2 \theta_{\text{eff}}^\ell$  (limited by sys. err. on beam polarization)
- $A_{FB}^{\mu,\tau,b}$  (statistics limited)
- $R_\ell, R_c, R_b$  (limited by sys. err. on flavor tag)
- No competitive measurements on  $M_Z, \Gamma_Z, \sigma^0$  (need to use LEP values)

## Leading effect: Soft+collinear multi- $\gamma$ ISR



$$\sigma[e^+e^- \rightarrow f\bar{f}] = \mathcal{R}_{\text{ini}}(s, s') \otimes \sigma_{\text{hard}}(s')$$

Kureav, Fadin '85

Berends, Burgers, v. Neerven '88

Kniehl, Krawczyk, Kühn, Stuart '88

Beenakker, Berends, v. Neerven '89

Bardin et al. '91; Skrzypek '92

Montagna, Nicrosini, Piccinini '97

$$\mathcal{R}_{\text{ini}} = \sum_n \left(\frac{\alpha}{\pi}\right)^n \sum_{m=0}^n h_{nm} \ln^m\left(\frac{s}{m_e^2}\right)$$

Universal ( $m=n$ ) logs known to  $n=6$ ,  
also some sub-leading terms

Ablinger, Blümlein, De Freitas, Schönwald '20

Exclusive description: MC tools

## Subleading effects:

Radiative corrections to  
 $e^+e^- \rightarrow f\bar{f}\gamma (+n\gamma)$

- Some corrections cancel for  $A_{\text{LR}}, A_{\text{FB}}, \text{BRs}$
- NLO for  $ee \rightarrow f\bar{f}\gamma$   
+ NNLO for  $ee \rightarrow Z\gamma$ ,  
 $Z \rightarrow f\bar{f}$  could be sufficient

**W mass measurement** from  $e^+e^- \rightarrow WW$ :

Baak et al. '13

- $\ell\nu_\ell\ell'\nu_{\ell'}$ : Endpoints of  $E_\ell$  or other distributions
- $\ell\nu_\ell jj$ : Kinematic reconstruction
- $jjjj$ : Systematic uncertainty from color reconnection

Expected precision with  $\mathcal{L}_{\text{int}} = 2 \text{ ab}^{-1}$  at  $\sqrt{s} = 250 \text{ GeV}$ :  $\Delta M_W \approx 2.5 \text{ MeV}$

**Theory needs:** Small impact of loop corrections, but accurate description of FSR  
QED effects needed

Anticipated precision for EWPOs:

Fujii et al. '19

	ILC-250*	ILC/GigaZ	CEPC	FCC-ee
$M_W$ [MeV]	$2.5^a$	$1\text{--}2^{\,a,e}$	$1^{\,e}$	$1^{\,e}$
$M_Z$ [MeV]	—	—	$0.5^{\,e}$	$0.1^{\,e}$
$\Gamma_Z$ [MeV]	—	$1^{\,a}$	$0.5^{\,e}$	$0.1^{\,e}$
$R_\ell = \Gamma_Z^{\text{had}}/\Gamma_Z^\ell [10^{-3}]$	$16^{\,c}$	$6^{\,b}$	$2^{\,b}$	$1^{\,b}$
$R_b = \Gamma_Z^b/\Gamma_Z^{\text{had}} [10^{-5}]$	$23^{\,c}$	$15^{\,c}$	$4.3^{\,c}$	$6^{\,c}$
$\sin^2 \theta_{\text{eff}}^\ell [10^{-5}]$	$2^{\,d}$	$1^{\,d}$	$<1^{\,e}$	$0.5^{\,e}$

## Systematics:

$a$  energy scale

$b$  acceptance

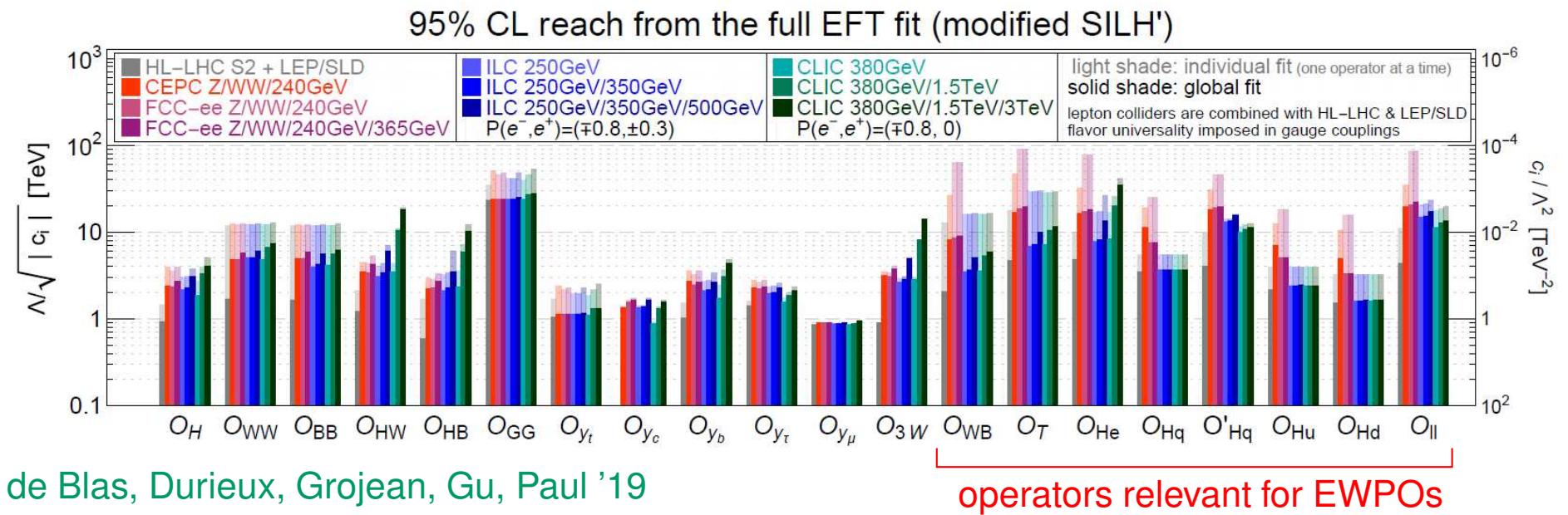
$c$  flavor tagging

$d$  polarization

$e$  beam energy calibration / beam-beam interactions

\*  $\sqrt{s} = 250$  GeV,  $\mathcal{L}_{\text{int}} = 2 \text{ ab}^{-1}$

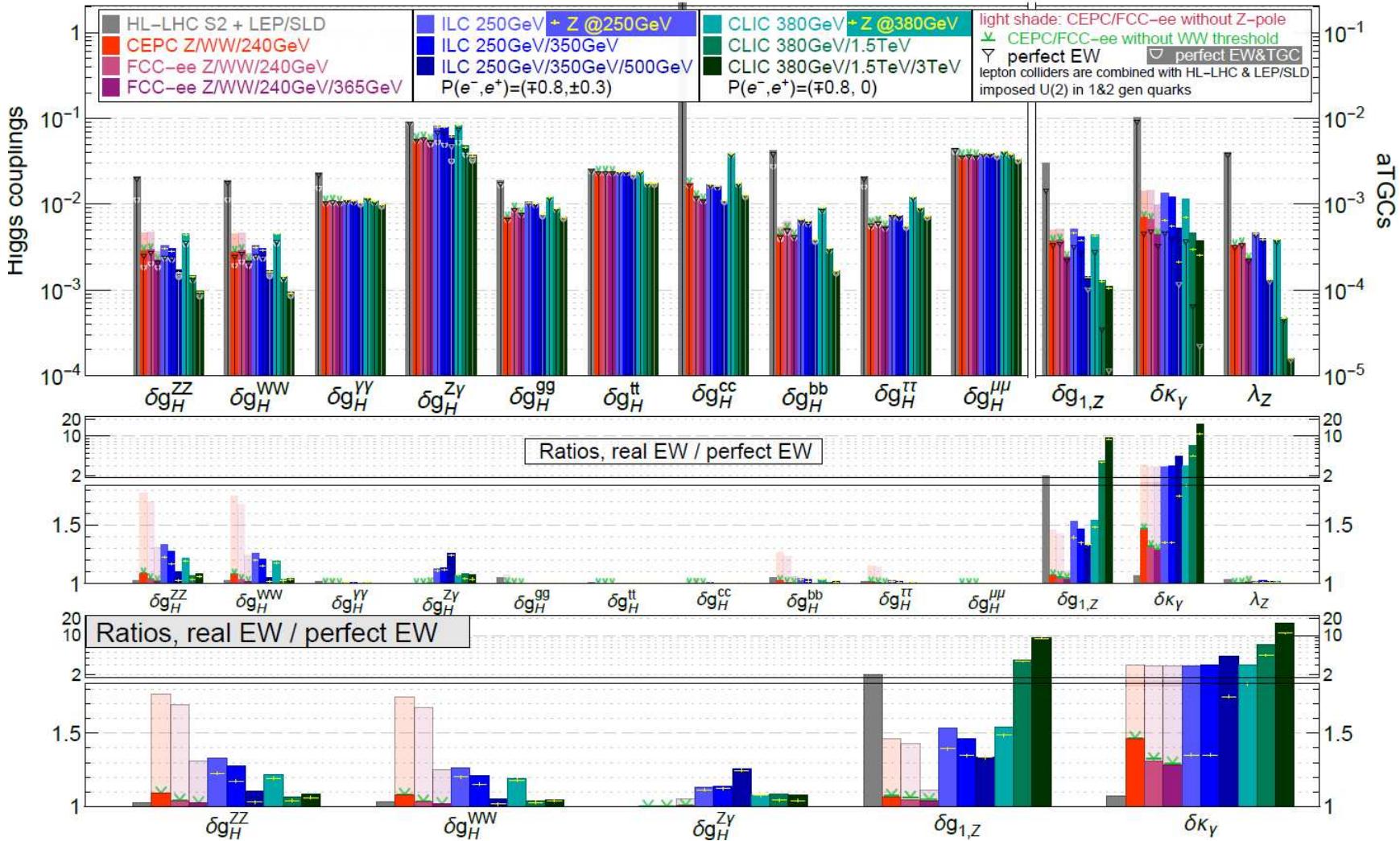
Projected reach assuming Minimal Flavor Violation:



# Correlation between EW and Higgs physics

18/21

- Improvement of  $Z$ -pole data important for Higgs physics and aGC:  
precision reach on effective couplings from full EFT global fit



## Strong coupling

- Electroweak precision ( $R_\ell = \Gamma_Z^{\text{had}} / \Gamma_Z^\ell$ ):

$$\alpha_s = 0.120 \pm 0.003 \quad \text{PDG '18}$$

→ No (negligible) non-perturbative QCD effects

$$\text{FCC-ee: } \delta R_\ell \sim 0.001$$

$$\Rightarrow \delta \alpha_s < 0.0002 \text{ (subj. to theory error)}$$

**Caviat:**  $R_\ell$  could be affected by new physics

- $R = \frac{\sigma[ee \rightarrow \text{had.}]}{\sigma[ee \rightarrow \mu\mu]}$  at lower  $\sqrt{s}$

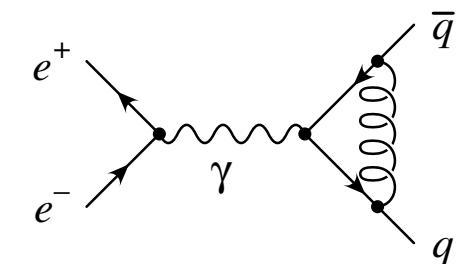
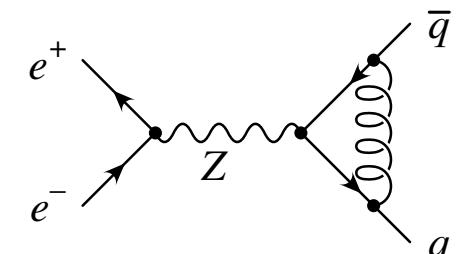
$$\text{e.g. CLEO } (\sqrt{s} \sim 9 \text{ GeV}): \alpha_s = 0.110 \pm 0.015$$

Kühn, Steinhauser, Teubner '07

→ dominated by  $s$ -channel photon, less room for new physics

→ QCD still perturbative

naive scaling to  $50 \text{ ab}^{-1}$  (BELLE-II):  $\delta \alpha_s \sim 0.0001$



- $\Delta\alpha_{\text{had}}$ : Could be limiting factor

a) From  $e^+e^- \rightarrow \text{had}$ . using dispersion relation

Current:  $\delta(\Delta\alpha_{\text{had}}) \sim 10^{-4}$

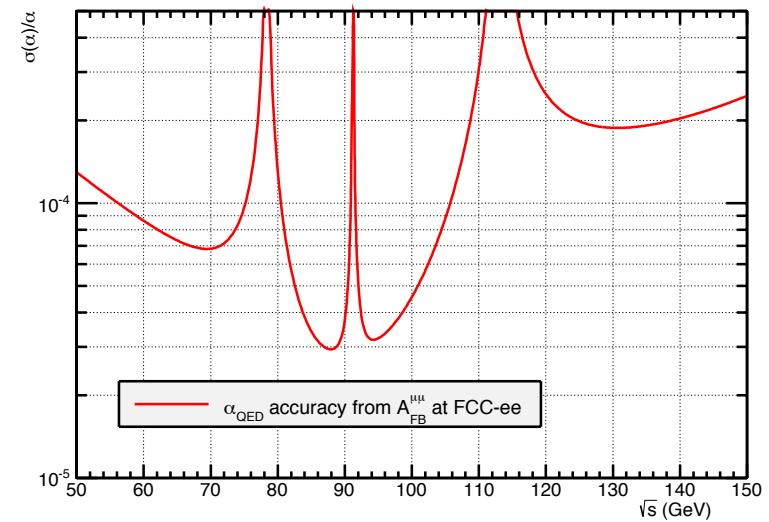
Improvement to  $\delta(\Delta\alpha_{\text{had}}) \sim 5 \times 10^{-5}$  likely

b) Direct determination at FCC-ee from  $e^+e^- \rightarrow \mu^+\mu^-$  off the Z peak  
(i.e.  $A_{\text{FB}}^{\mu\mu}$  at  $\sqrt{s} \sim 88$  GeV and  $\sqrt{s} \sim 95$  GeV)

$\rightarrow \delta(\Delta\alpha_{\text{had}}) \sim 3 \times 10^{-5}$  with  $\mathcal{L}_{\text{tot}} = 85 \text{ ab}^{-1}$

Janot '15

Requires high-precision theory prediction for  $e^+e^- \rightarrow \mu^+\mu^-$  including 2/3-loop corrections for  $\gamma$ -exchange and box contributions



- **Electroweak precision tests** at future  $e^+e^-$  colliders require allow to probe multi-TeV BSM physics (or feebly coupled lighter physics)
- Improved measurements in **several sectors** important:  
EW masses and couplings, Higgs couplings, top mass,  $\alpha_s$  and  $\alpha(s)$
- Theory progress needed both for **fixed-order loop corrections** as well as **MC tools**
- **ILC-250** can deliver similar physics goals as GigaZ, but with reduced precision  
→ **Open question:** Direct determination of  $\alpha_s$  and  $\alpha(s)$
- Unique theory challenges for description of  $ee \rightarrow Z\gamma$  vs.  $ee \rightarrow Z$

## Backup slides

# Z lineshape

- Deconvolution of initial-state QED radiation:

$$\sigma[e^+e^- \rightarrow f\bar{f}] = \mathcal{R}_{\text{ini}}(s, s') \otimes \sigma_{\text{hard}}(s')$$

- Subtraction of  $\gamma$ -exchange,  $\gamma-Z$  interference, box contributions:

$$\sigma_{\text{hard}} = \sigma_Z + \sigma_\gamma + \sigma_{\gamma Z} + \sigma_{\text{box}}$$

- Z-pole contribution:

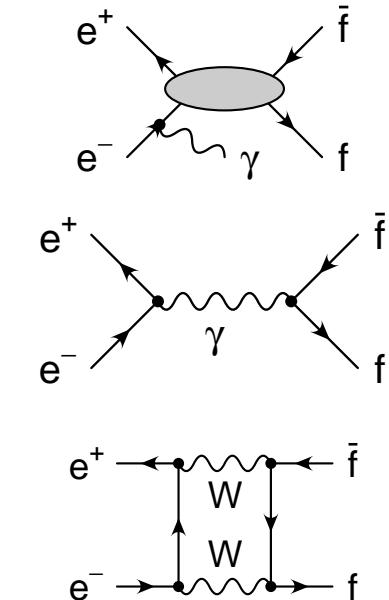
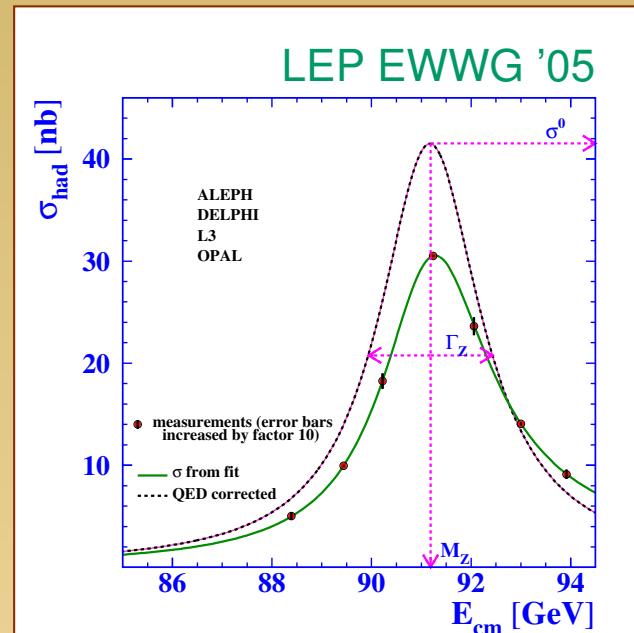
$$\sigma_Z = \frac{R}{(s - M_Z^2)^2 + M_Z^2 \Gamma_Z^2} + \sigma_{\text{non-res}}$$

- In experimental analyses:

$$\sigma \sim \frac{1}{(s - M_Z^2)^2 + s^2 \Gamma_Z^2 / M_Z^2}$$

$$M_Z = M_Z / \sqrt{1 + \Gamma_Z^2 / M_Z^2} \approx M_Z - 34 \text{ MeV}$$

$$\Gamma_Z = \Gamma_Z / \sqrt{1 + \Gamma_Z^2 / M_Z^2} \approx \Gamma_Z - 0.9 \text{ MeV}$$



## “Hard” matrix element

Consistent (gauge-invariant) theory setup:

Expansion of  $\mathcal{A}[e^+ e^- \rightarrow \mu^+ \mu^-]$  about  $s_0 = M_Z^2 - iM_Z\Gamma_Z$ :

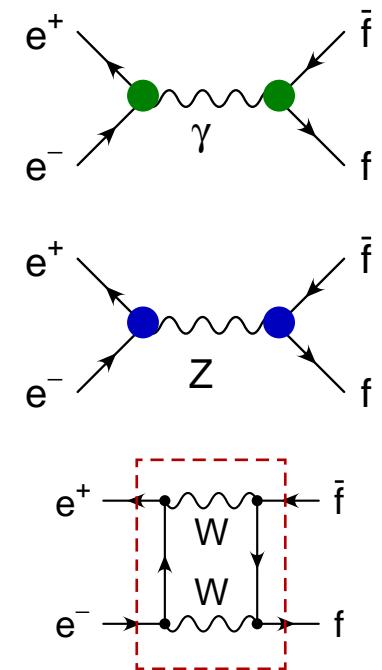
$$\mathcal{A}[e^+ e^- \rightarrow f\bar{f}] = \frac{R}{s - s_0} + S + (s - s_0)T + \dots$$

$$R = g_Z^e(s_0)g_Z^f(s_0)$$

$$S = \left[ \frac{1}{M_Z^2} g_\gamma^e g_\gamma^f + g_Z^e g_Z^{f'} + g_Z^{e'} g_Z^f + S_{\text{box}} \right]_{s=s_0}$$

$g_V^f(s)$  : effective  $V f\bar{f}$  couplings

At NNLO: Need  $R$  at  $\mathcal{O}(\alpha^2)$ ,  $S$  at  $\mathcal{O}(\alpha)$ , etc.



## Z-pole asymmetries

Blondel scheme:

(if  $e^-$  and  $e^+$  polarization available)

Blondel '88

Four independent measurements for  $P_{e^+}/P_{e^-} = ++, +-, -+, --$

$$A_{LR} = \sqrt{\frac{(\sigma_{++} + \sigma_{-+} - \sigma_{+-} - \sigma_{--})(-\sigma_{++} + \sigma_{-+} - \sigma_{+-} + \sigma_{--})}{(\sigma_{++} + \sigma_{-+} + \sigma_{+-} + \sigma_{--})(-\sigma_{++} + \sigma_{-+} + \sigma_{+-} - \sigma_{--})}}$$

**Note:** No need to know  $|P_{e^\pm}|$  !

Main systematic uncertainties:

- Difference of  $|P|$  for  $P > 0$  and  $P < 0$
- Difference of  $\mathcal{L}$  for  $P > 0$  and  $P < 0$

$$\delta A_{LR} \approx 10^{-4} \quad \Rightarrow \quad \delta \sin^2 \theta_{\text{eff}}^\ell \approx 1.3 \times 10^{-5}$$

Mönig, Hawkings '99

## Theory calculations: Status

- Many seminal works on 1-loop and leading 2-loop corrections

Veltman, Passarino, Sirlin, Marciano, Bardin, Hollik, Riemann, Degrassi, Kniehl, ...

- Full 2-loop results for  $M_W$ ,  $Z$ -pole observables

Freitas, Hollik, Walter, Weiglein '00

Hollik, Meier, Uccirati '05,07

Awramik, Czakon '02

Awramik, Czakon, Freitas, Kniehl '08

Onishchenko, Veretin '02

Freitas '14

Awramik, Czakon, Freitas, Weiglein '04

Dubovsky, Freitas, Gluza, Riemann, Usovitsch '16,18

Awramik, Czakon, Freitas '06

- Approximate 3- and 4-loop results (enhance by  $Y_t$  and/or  $N_f$ )

Chetyrkin, Kühn, Steinhauser '95

Chetyrkin et al. '06

Faisst, Kühn, Seidensticker, Veretin '03

Boughezal, Czakon '06

Boughezal, Tausk, v. d. Bij '05

Chen, Freitas '20

Schröder, Steinhauser '05

## Theory and parametric uncertainties

	CEPC	perturb. error with 3-loop <sup>†</sup>	Param. error CEPC*	main source
$M_W$ [MeV]	1	1	2.1	$m_t, \Delta\alpha$
$\Gamma_Z$ [MeV]	0.5	0.15	0.15	$m_t, \alpha_s$
$R_b$ [ $10^{-5}$ ]	4.3	5	< 1	
$\sin^2 \theta_{\text{eff}}^\ell$ [ $10^{-5}$ ]	<1	1.5	2	$m_t, \Delta\alpha$

<sup>†</sup> **Theory scenario:**  $\mathcal{O}(\alpha\alpha_s^2)$ ,  $\mathcal{O}(N_f\alpha^2\alpha_s)$ ,  $\mathcal{O}(N_f^2\alpha^2\alpha_s)$ , leading 4-loop  
 $(N_f^n = \text{at least } n \text{ closed fermion loops})$

Parametric inputs:

**\*CEPC:**  $\delta m_t = 600$  MeV,  $\delta\alpha_s = 0.0002$ ,  $\delta M_Z = 0.5$  MeV,  
 $\delta(\Delta\alpha) = 5 \times 10^{-5}$

## Theory and parametric uncertainties

	CEPC	perturb. error with 3-loop <sup>†</sup>	Param. error CEPC*	main source
$M_W$ [MeV]	1	1	0.6	$\Delta\alpha$
$\Gamma_Z$ [MeV]	0.5	0.15	0.1	$\alpha_s$
$R_b$ [ $10^{-5}$ ]	4.3	5	< 1	
$\sin^2 \theta_w^\ell$ [10 $^{-5}$ ]	<1	1.5	1	$\Delta\alpha$

<sup>†</sup> **Theory scenario:**  $\mathcal{O}(\alpha\alpha_s^2)$ ,  $\mathcal{O}(N_f\alpha^2\alpha_s)$ ,  $\mathcal{O}(N_f^2\alpha^2\alpha_s)$ , leading 4-loop  
 $(N_f^n = \text{at least } n \text{ closed fermion loops})$

Parametric inputs:

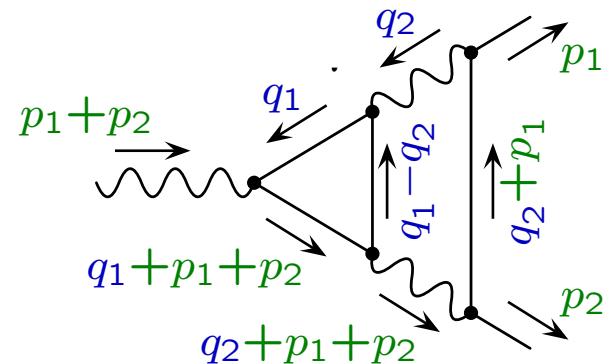
**\*FCC-ee:**  $\delta m_t = 50$  MeV,  $\delta\alpha_s = 0.0002$ ,  $\delta M_Z = 0.5$  MeV,  
 $\delta(\Delta\alpha) = 3 \times 10^{-5}$

# Calculational techniques

Experimental precision requires inclusion of **radiative corrections** in theory  
(1-loop, 2-loop, and partial 3-loop)

Integrals over loop momenta:

$$\int d^4 q_1 d^4 q_2 f(q_1, q_2, p_1, p_2, \dots, m_1, m_2, \dots)$$



Computer algebra tools:

- Generation of diagrams,  $\mathcal{O}(100) - \mathcal{O}(10000)$
- Lorentz and Dirac algebra
- Integral simplification (and expansion)

Evaluation of loop integrals:

- In general not possible analytically
- Numerical methods are more general, but computing intensive
- Special numerical techniques can balance precision and evaluation time

## Analytic calculations

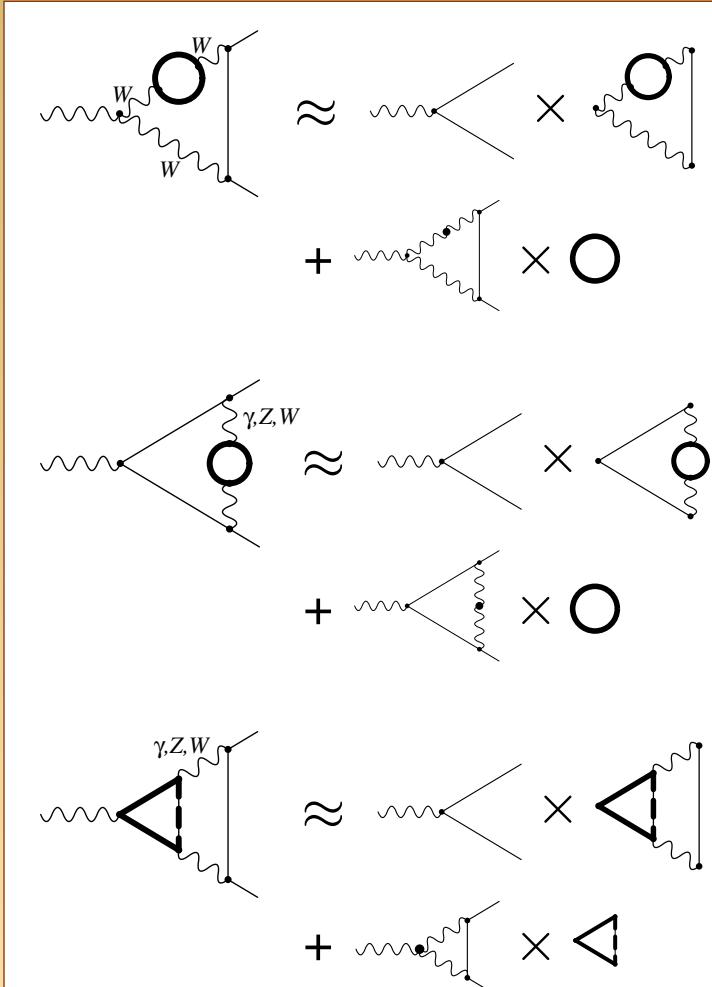
- Mostly used for diagrams with few mass scales
- Reduce to **master integrals** with integration-by-parts and other identities  
[Chetyrkin, Tkachov '81](#); [Gehrmann, Remiddi '00](#); [Laporta '00](#); ...

Public programs:	Reduze	<a href="#">von Manteuffel, Studerus '12</a>
	FIRE	<a href="#">Smirnov '13,14</a>
	LiteRed	<a href="#">Lee '13</a>
	KIRA	<a href="#">Maierhoefer, Usovitsch, Uwer '17</a>

- Large need for computing time and memory
- Evaluate master integrals with differential equations or Mellin-Barnes rep.  
[Kotikov '91](#); [Remiddi '97](#); [Smirnov '00,01](#); [Henn '13](#); ...
  - Result in terms of Goncharov polylogs / multiple polylogs
  - Some problems need iterated elliptic integrals / elliptic multiple polylogs  
[Broedel, Duhr, Dulat, Trancredi '17,18](#)  
[Ablinger et al. '17](#)
  - Even more classes of functions needed in future?

# Asymptotic expansions

- Exploit large mass ratios,  
e. g.  $M_Z^2/m_t^2 \approx 1/4$
  - Evaluate coeff. integrals analytically
  - Fast numerical evaluation
- Used in some 2/3-scale problems
- Public programs:  
exp      Harlander, Seidensticker, Steinhauser '97  
asy      Pak, Smirnov '10
- Possible limitations:
- Difficult coefficient integrals
  - bad convergence



# Numerical integration

Two general approaches:

- Automated treatment of UV/IR divergencies
- No restriction on number of loops or legs

## ■ Sector decomposition:

Public programs:	SecDec	Carter, Heinrich '10; Borowka et al. '12,15,17
	FIESA	Smirnov, Tentyukov '08; Smirnov '13,15

## ■ Mellin-Barnes representations:

Public programs:	MB/MBresolve	Czakon '06; Smirnov, Smirnov '09
	AMBRE/MBnumerics	Gluza, Kajda, Riemann '07 Dubovyk, Gluza, Riemann '15 Usovitsch, Dubovyk, Riemann '18

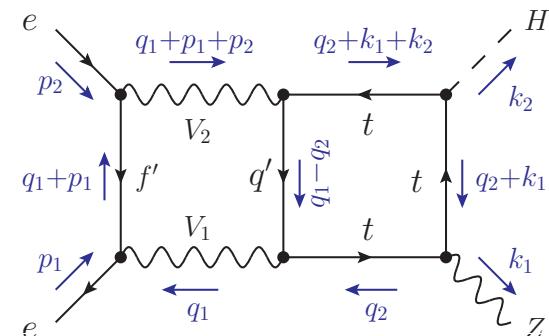
- Diagrams with internal thresholds can cause numerical instabilities
- Specialized techniques (for some type of diagrams) often improve computing time, robustness, precision (but not automated)

# Specialized numerical techniques

Example: HZ double boxes

Song, Freitas '21

- Introduce Feynman parameters and disp. rel.
- Expressions for second loop from, e.g., LoopTools  
Hahn, Perez-Victoria '98
- 3-dim. numerical integral with adaptive Gaussian integration
- $\mathcal{O}(0.1\%)$  precision in  $\mathcal{O}(\text{min.})$  on laptop



$$\int dx dy$$

Feynman diagram showing the loop variables for the HZ double box loop. The top-right diagonal gluon has a momentum  $k'_2 = xk_1 + (1-y)k_2$  and the bottom-right diagonal gluon has a momentum  $k'_1 = (1-x)k_1 + yk_2$ . The other momenta remain the same as in the previous diagram.

$$\int dx dy d\sigma$$

Feynman diagram showing the final result after numerical integration. A red vertical line labeled "mass  $\sigma$ " is inserted into the bottom-right diagonal gluon line. The top-right diagonal gluon has a momentum  $k'_2$  and the bottom-right diagonal gluon has a momentum  $q_1 + k'$ . The other momenta remain the same as in the previous diagram.