

Second HEP Graduate Workshop
University of Mohamed Boudiaf
M'sila – Algeria, 03-05 April 2021

Dark Matter Searches at Colliders

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Academia Sinica, Taiwan



中央研究院
ACADEMIA SINICA



Outline: part-II

- Production at Colliders
- DM Signatures and Searches
- Dark Sector at the LHC
- Future Prospects

Production at colliders

Proton-Proton (Heavy Ion) Accelerator

Circumference: 27 km

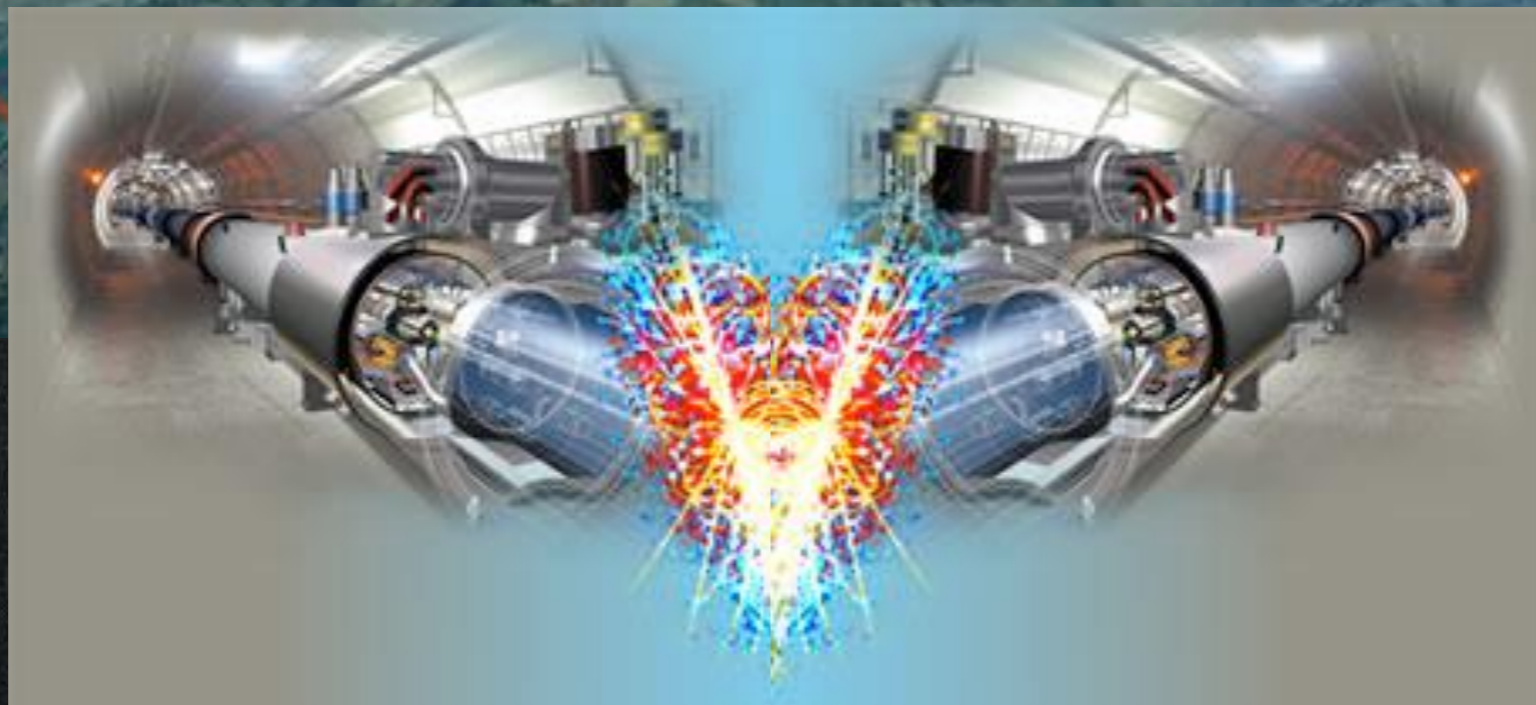
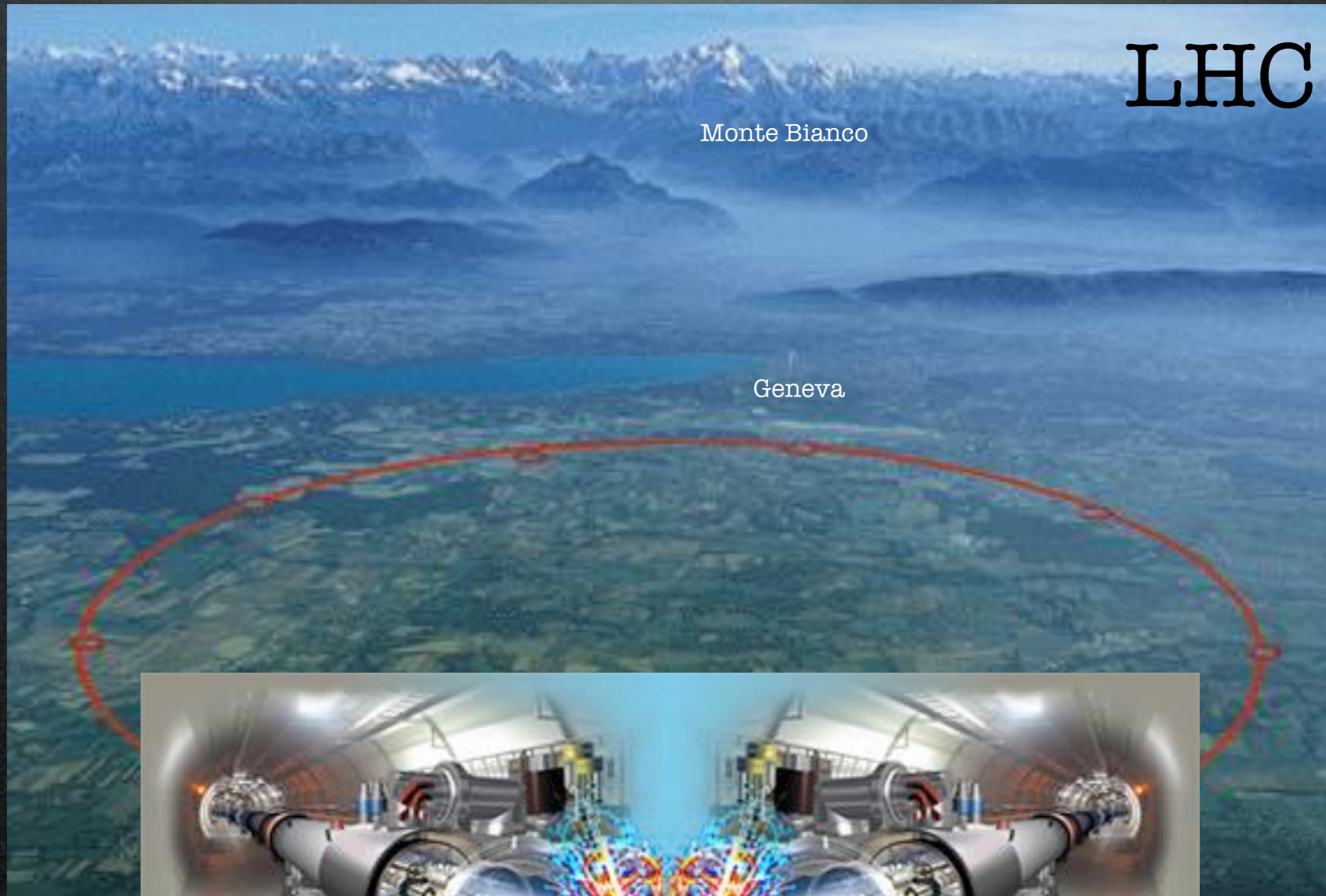
Energy: $\sqrt{s} = 7 \text{ TeV}$ (2010/2011)
 $\sqrt{s} = 8 \text{ TeV}$ (2012)
 $\sqrt{s} = 13 \text{ TeV}$ (2015-2018) Run2
 $\sqrt{s} = 14 \text{ TeV}$ (2021-2023) Run3



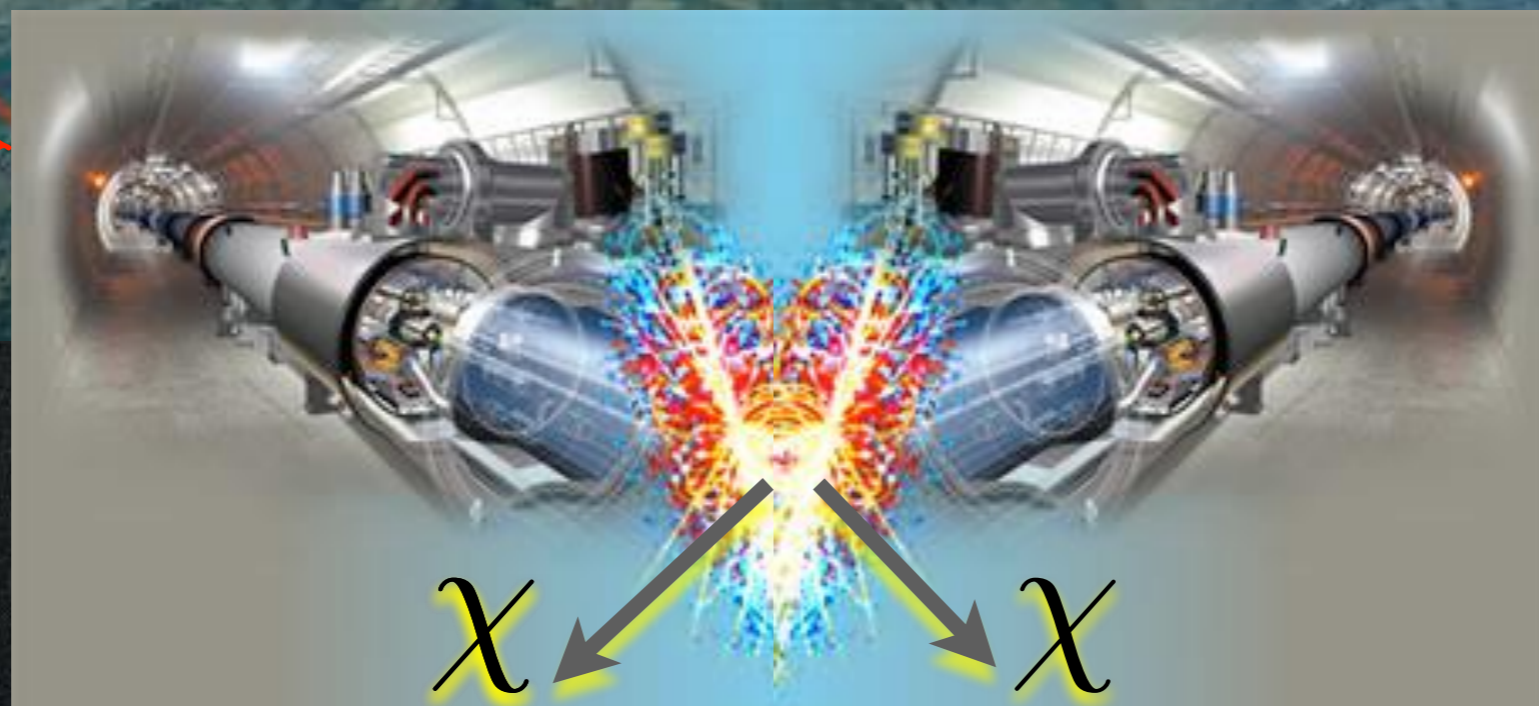
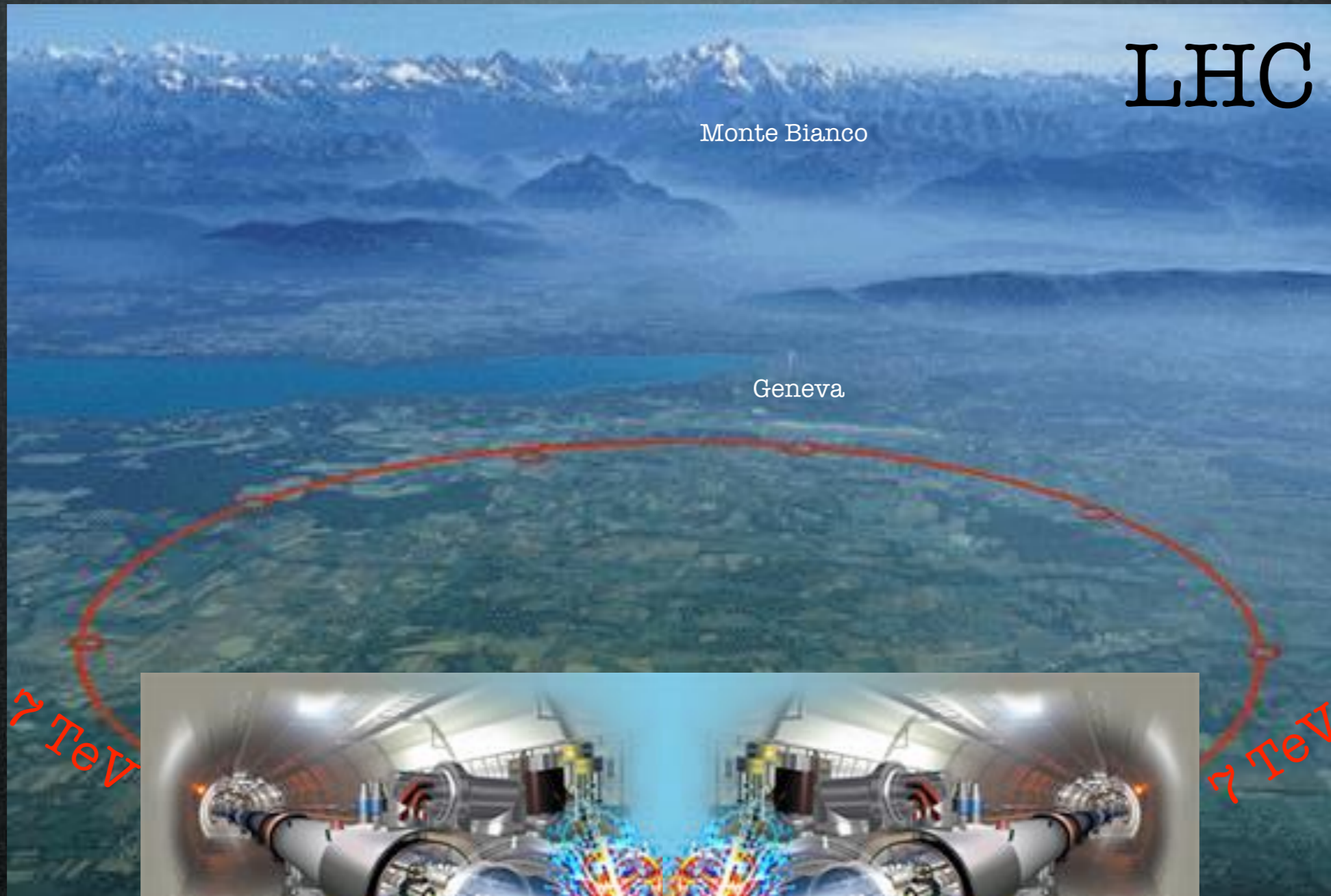
protons

protons

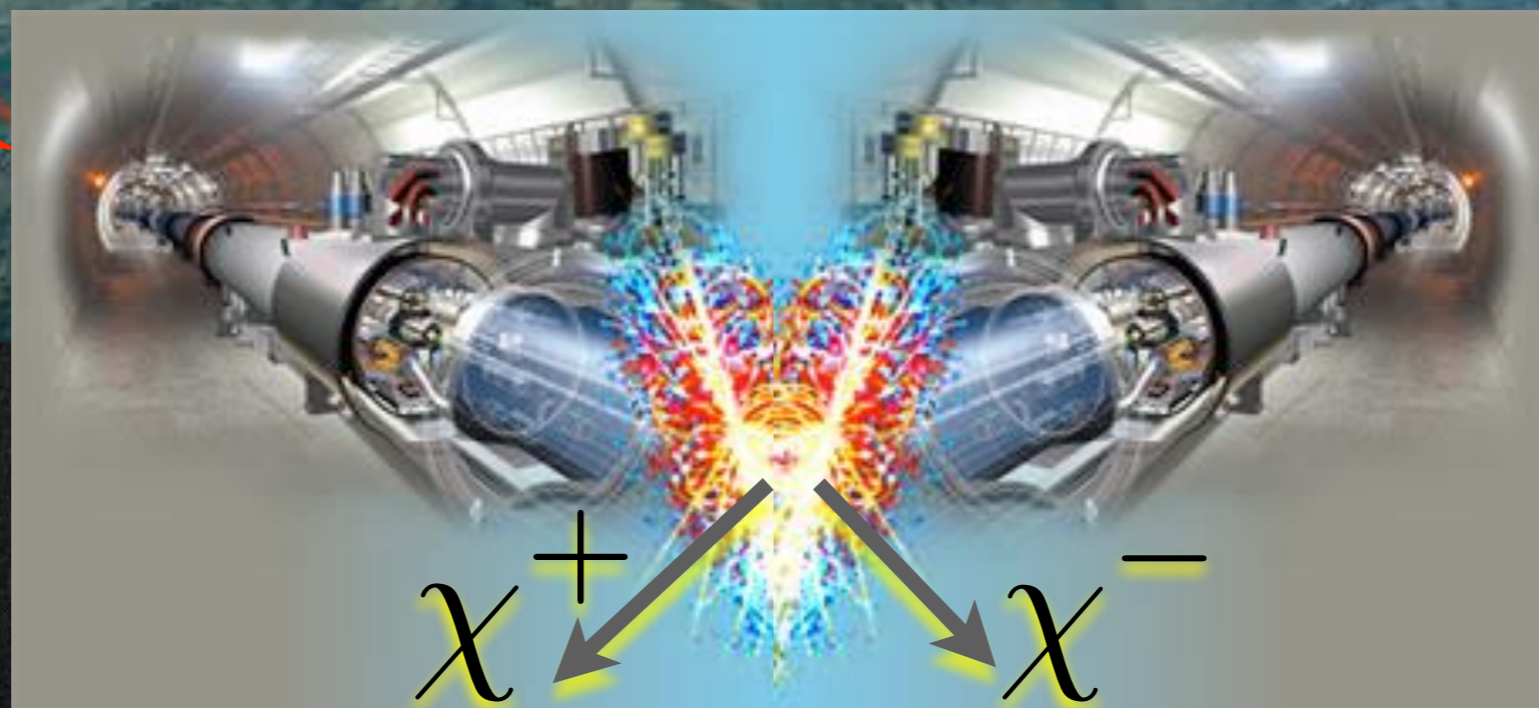
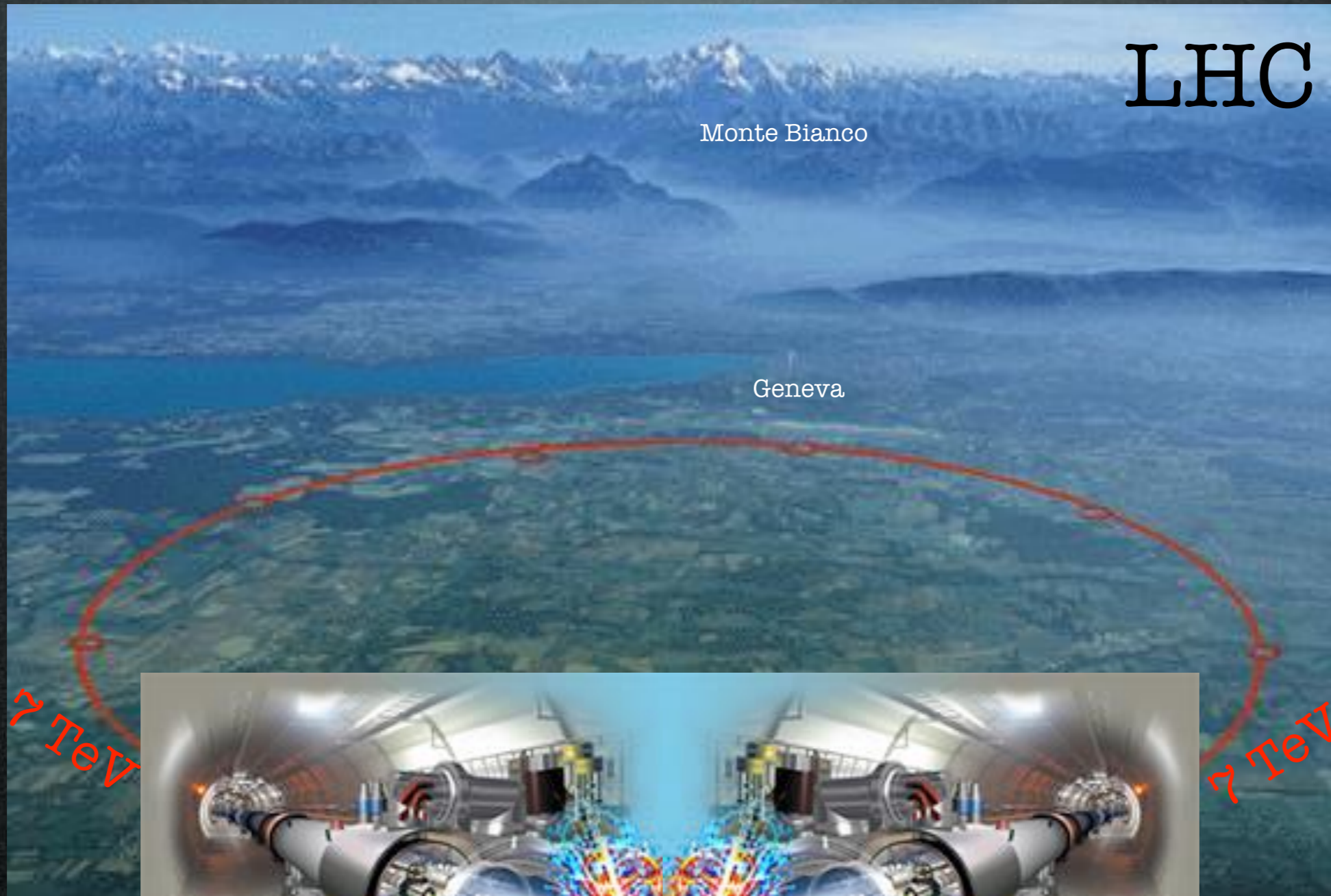
Production at colliders



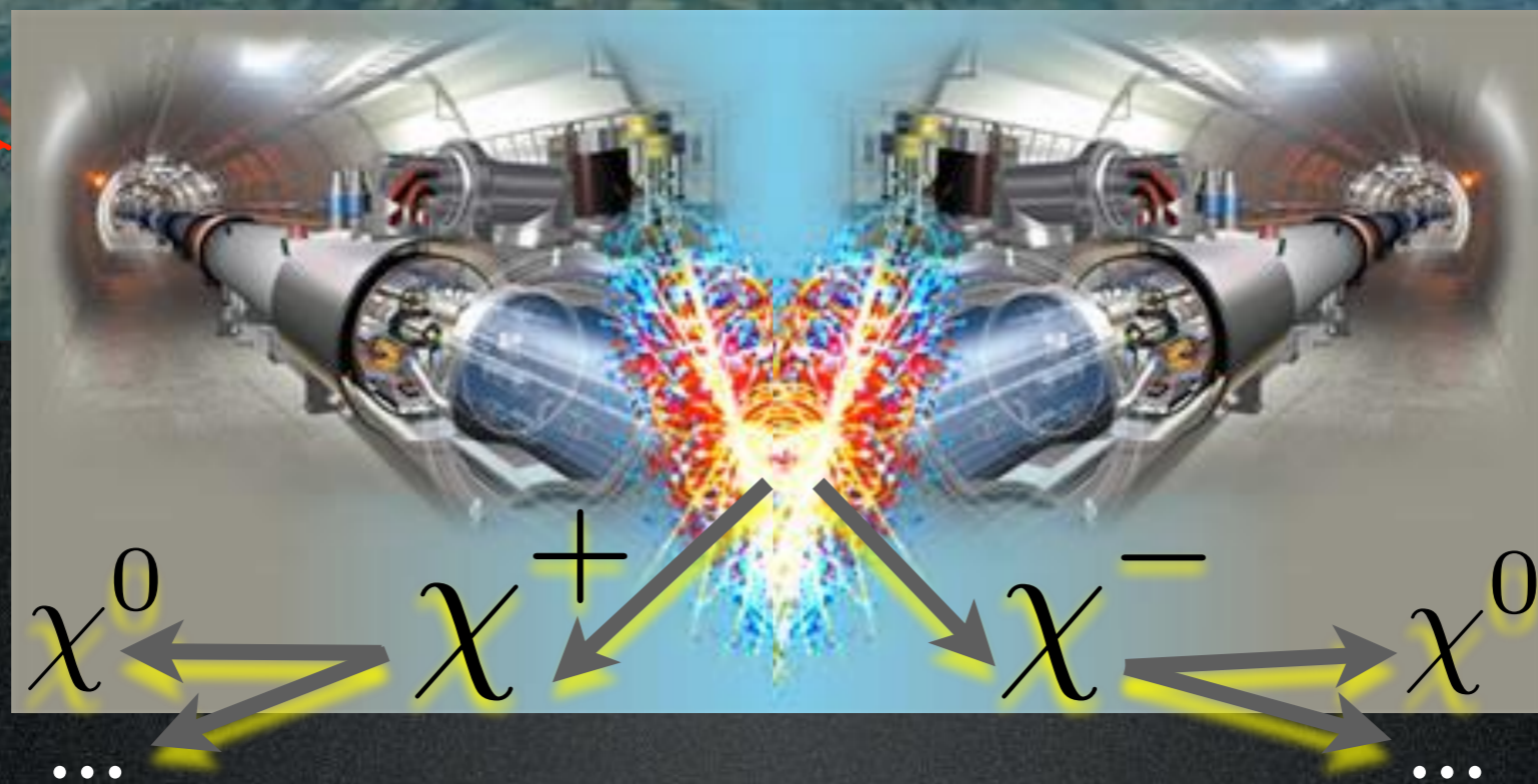
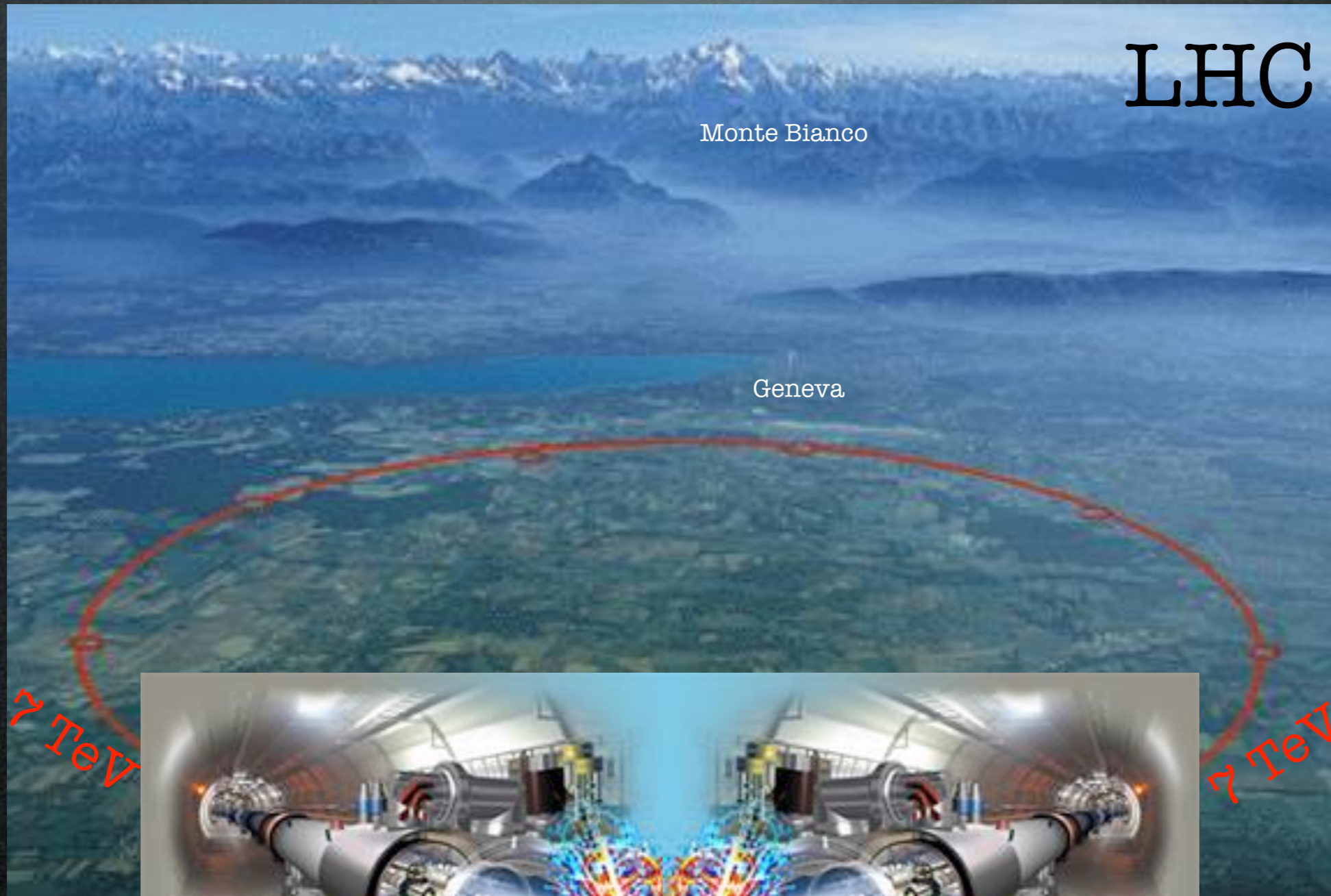
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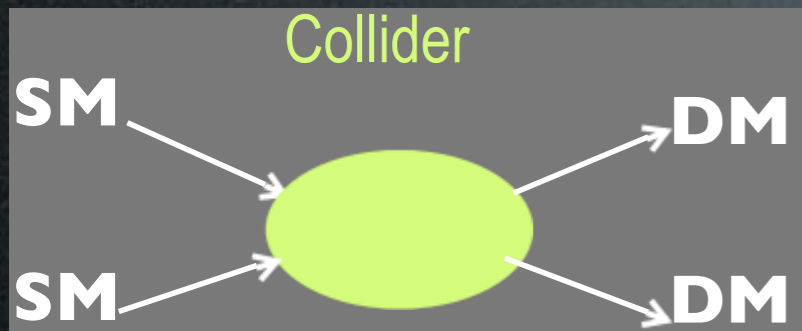
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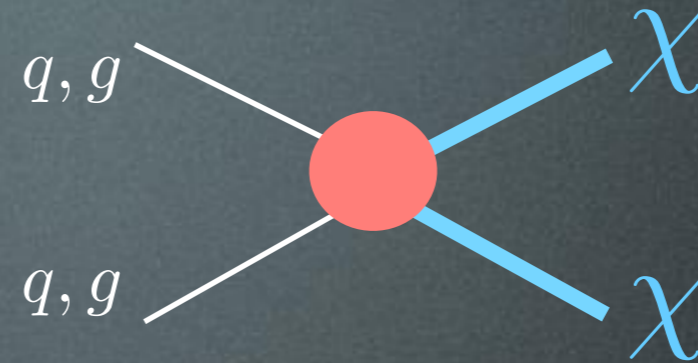
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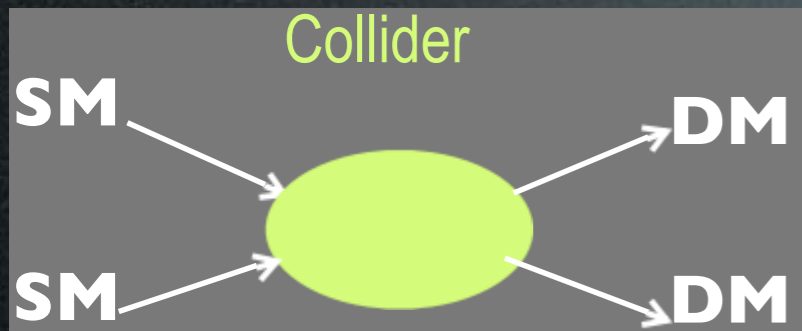
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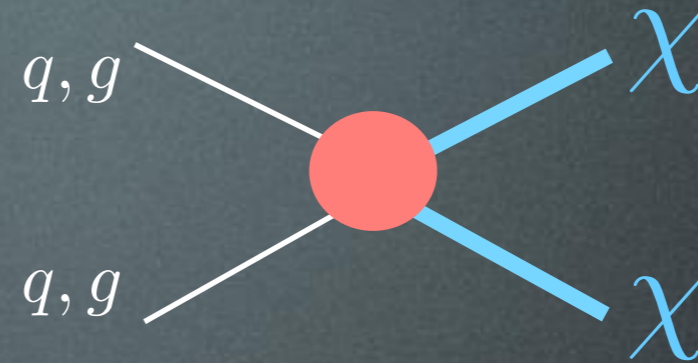
At LHC:



Production at colliders

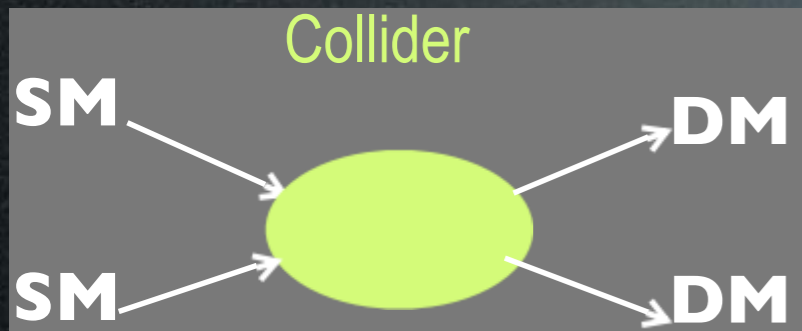


At LHC:

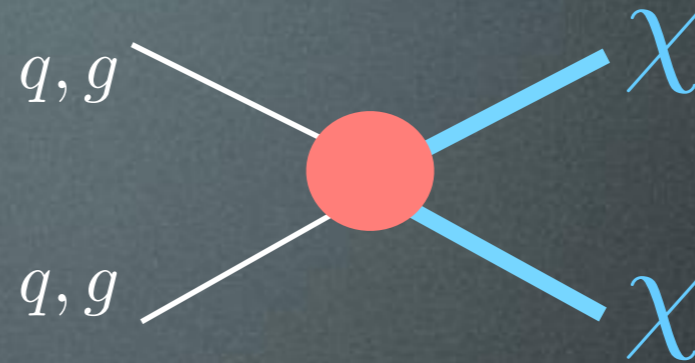


'Problem' is: DM flies away

Production at colliders



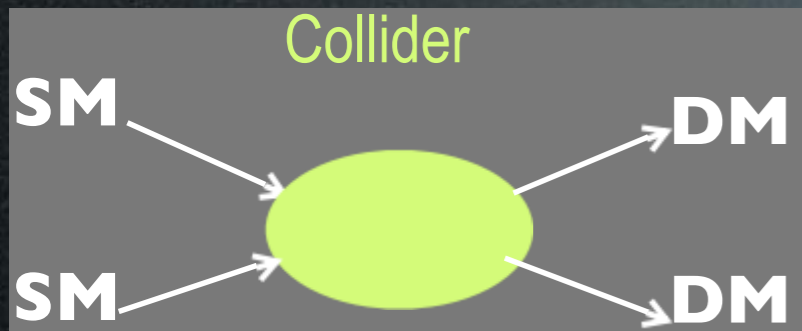
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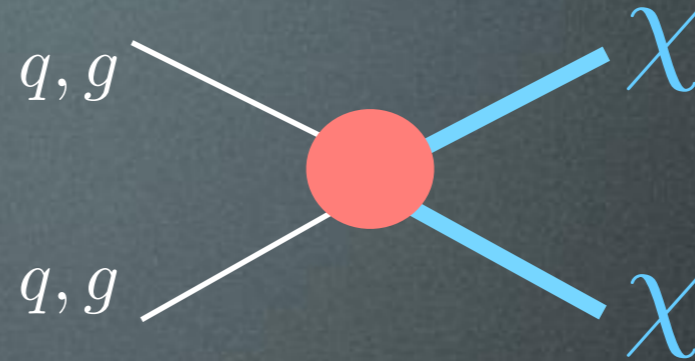
'Problem' is: DM flies away

Signature is: **missing energy**

Production at colliders



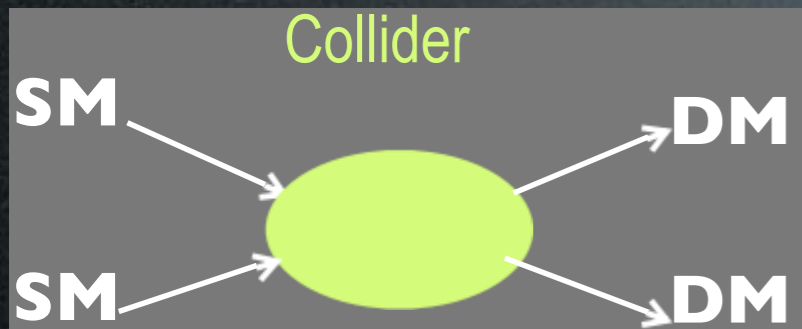
At LHC:



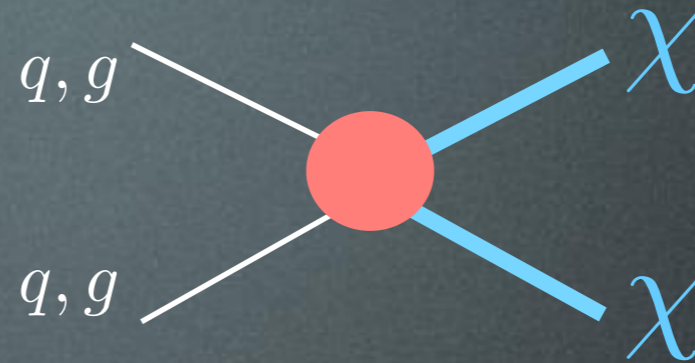
'Problem' is: DM flies away

Signature is: missing energy
transverse

Production at colliders

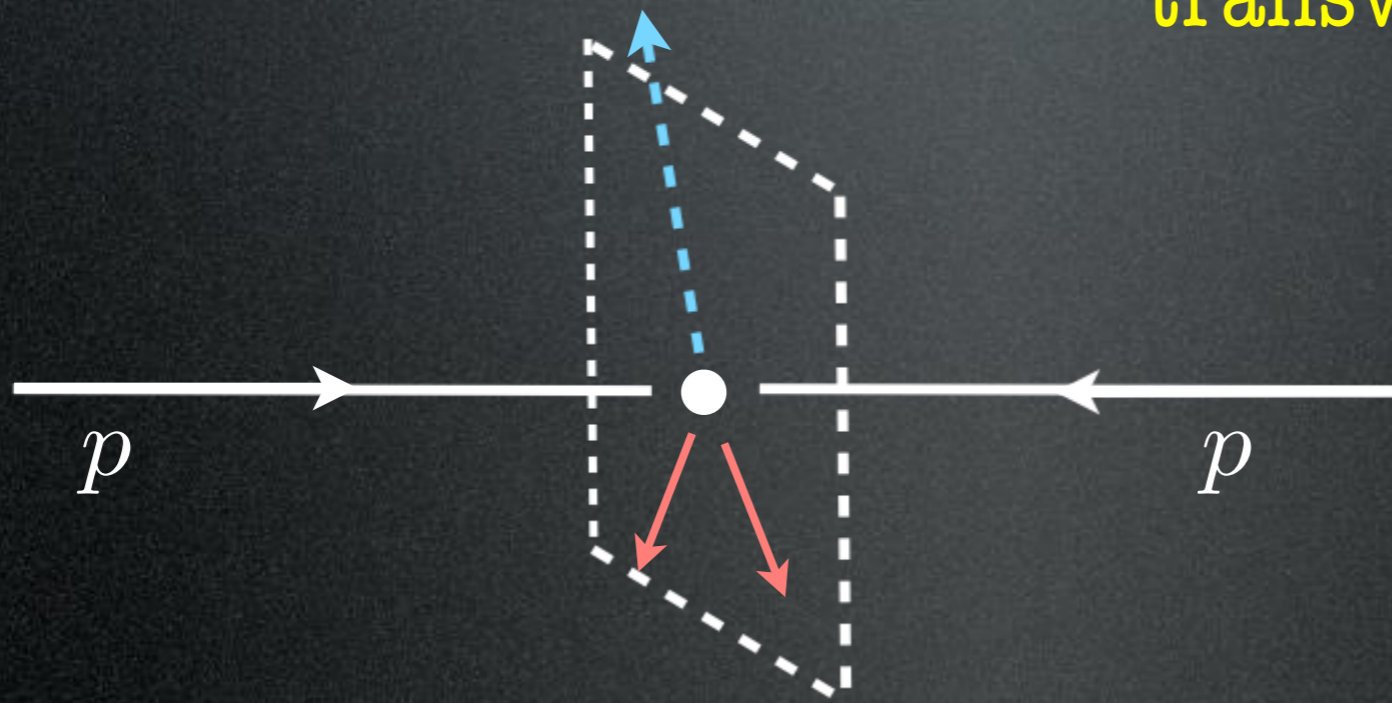


At LHC:



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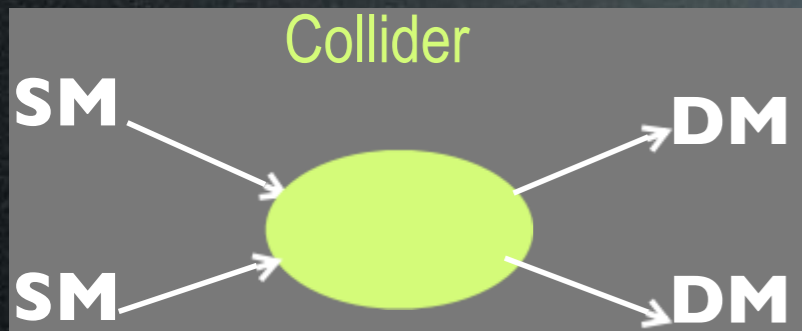


Before collision: $\vec{P}_T^{\text{tot}} \equiv 0$
(NB: $\vec{P}_L^{\text{tot}} \neq 0$ in general)

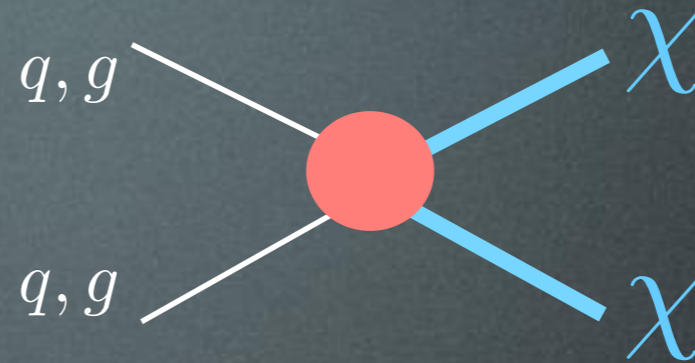
After collision: $\vec{P}_T^{\text{vis}} \stackrel{?}{=} 0$

If \neq , then **'MET'**

Production at colliders

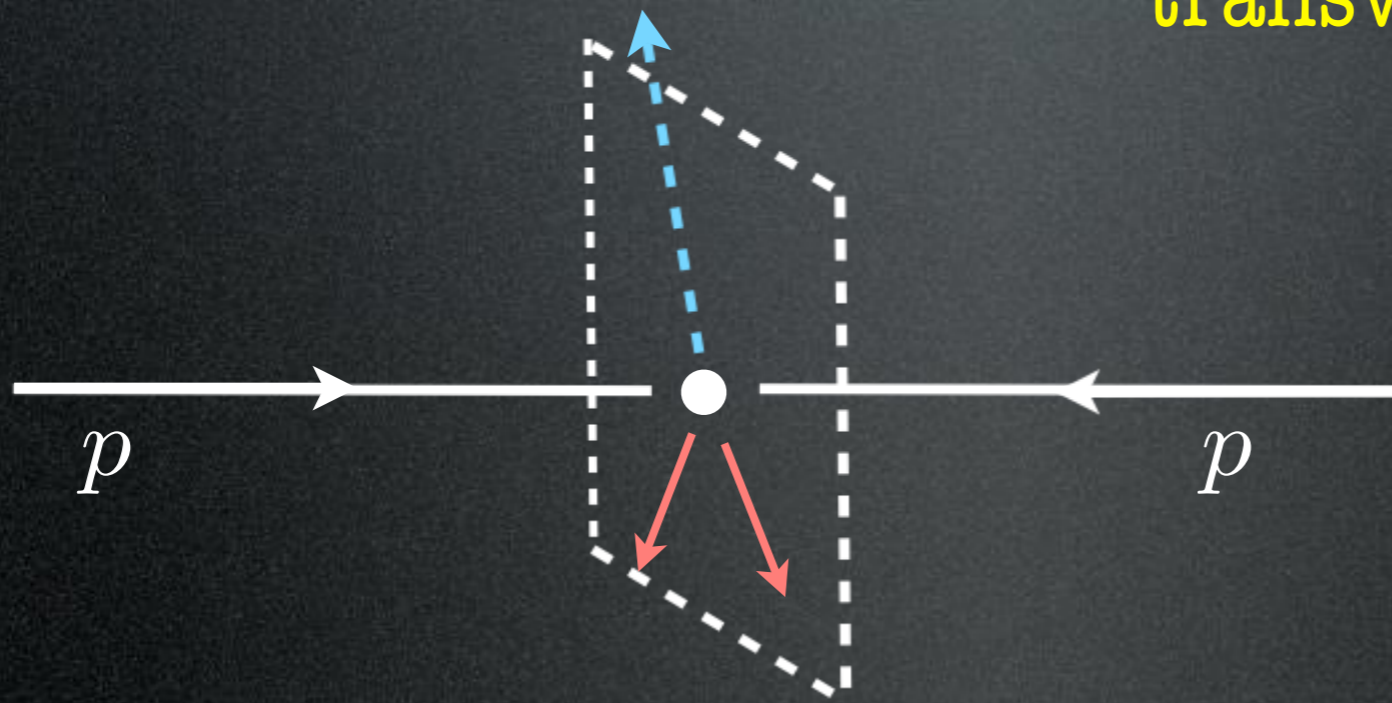


At LHC:



‘Problem’ is: DM flies away

Signature is: **missing energy**
transverse



Before collision: $\vec{P}_T^{\text{tot}} \equiv 0$

(NB: $\vec{P}_L^{\text{tot}} \neq 0$ in general)

After collision: $\vec{P}_T^{\text{vis}} \stackrel{?}{=} 0$

If \neq , then **MET**

Background: neutrinos (e.g. $W \rightarrow e\nu$)

- model your background and look for anomalies
- construct kinematic variables sensitive to χ mass

Production at colliders

OK, **MET** is the crucial signature. Then **what else?**

Production at colliders

OK, **MET** is the crucial signature. Then **what else?**

the bare minimum

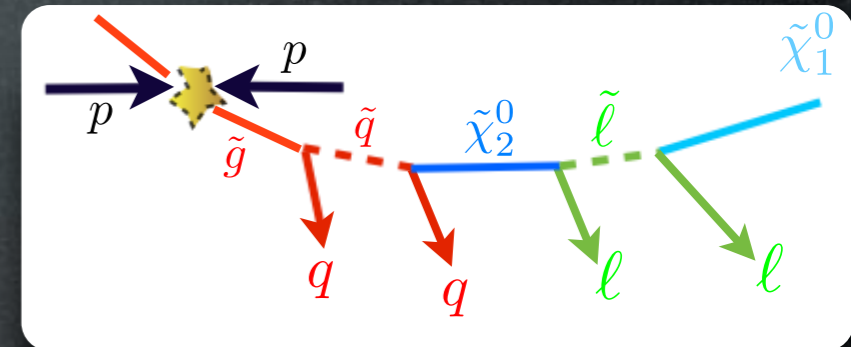
lots of things

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lots of things



'trigger on 4j+4l+MET...'

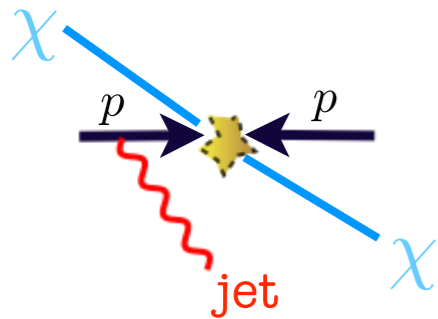
huge literature

- well studied (M_T^2 ...)
- model dependent

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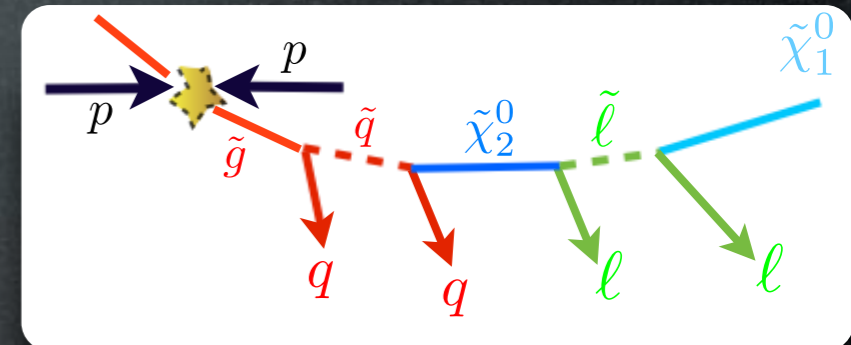


‘mono-jet’

‘started’ by
J. Goodman et al.,
1008.1783

- ‘new’
- model independent

lots of things



‘trigger on 4j+4l+MET...’

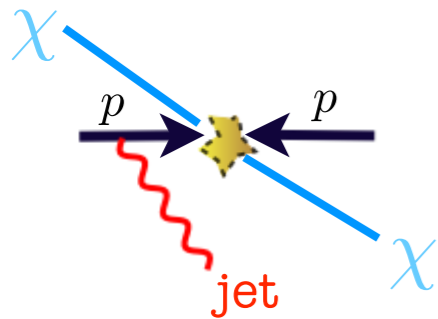
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‘mono-photon’

‘mono-Z/W’

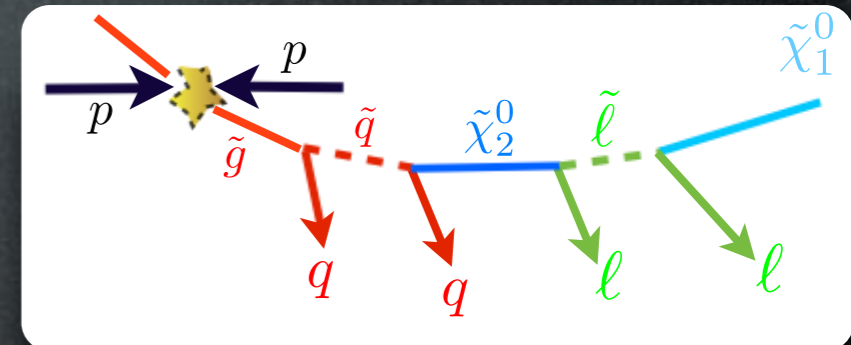
‘mono-top’...

‘mono-higgs’...



‘mono-X’

lots of things



‘trigger on 4j+4l+MET...’

huge literature

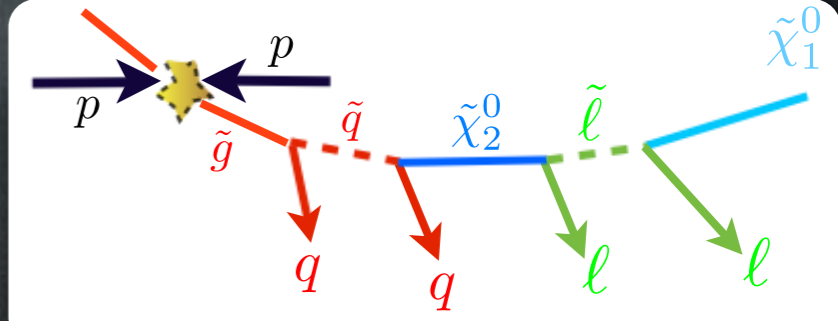
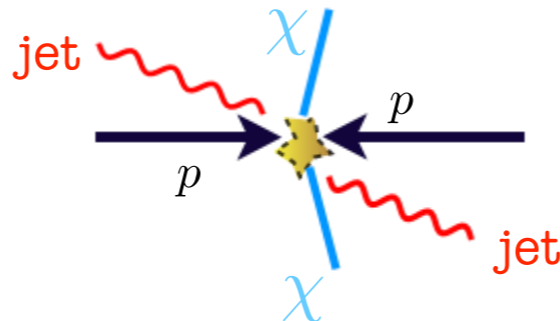
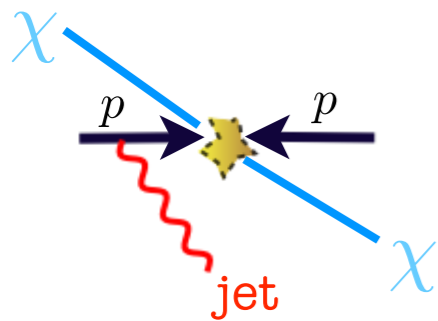
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‘(forward) di-jets’

‘trigger on 4j+4l+MET...’

huge literature

- ‘new’
- model independent

- ‘clean’ topology
- flexible interpretation
(see later)

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- ‘mono-photon’
- ‘mono-Z/W’
- ‘mono-top’...
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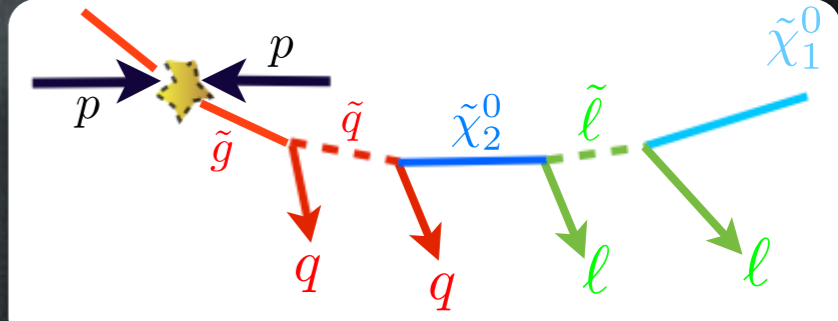
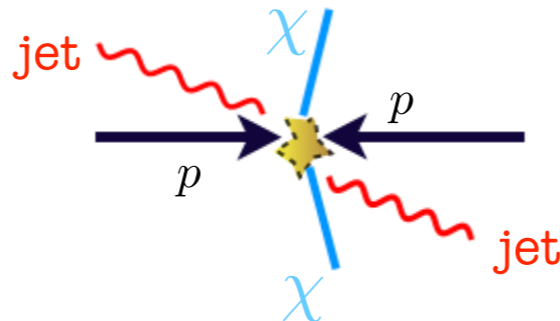
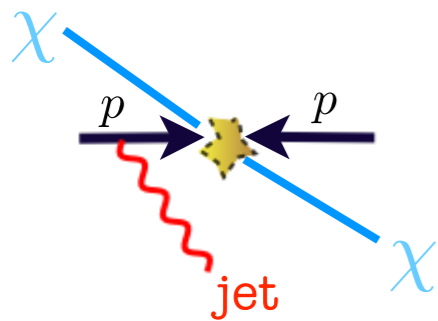
‘mono-X’

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- ‘mono-Z/W’
- ‘mono-top’...
- ‘mono-higgs’...



‘mono-X’

NB: not an exhaustive list

Production at colliders

OK, given what we (do not) see, how do we interpret?

effectively

simplified models

a simple model

a full theory

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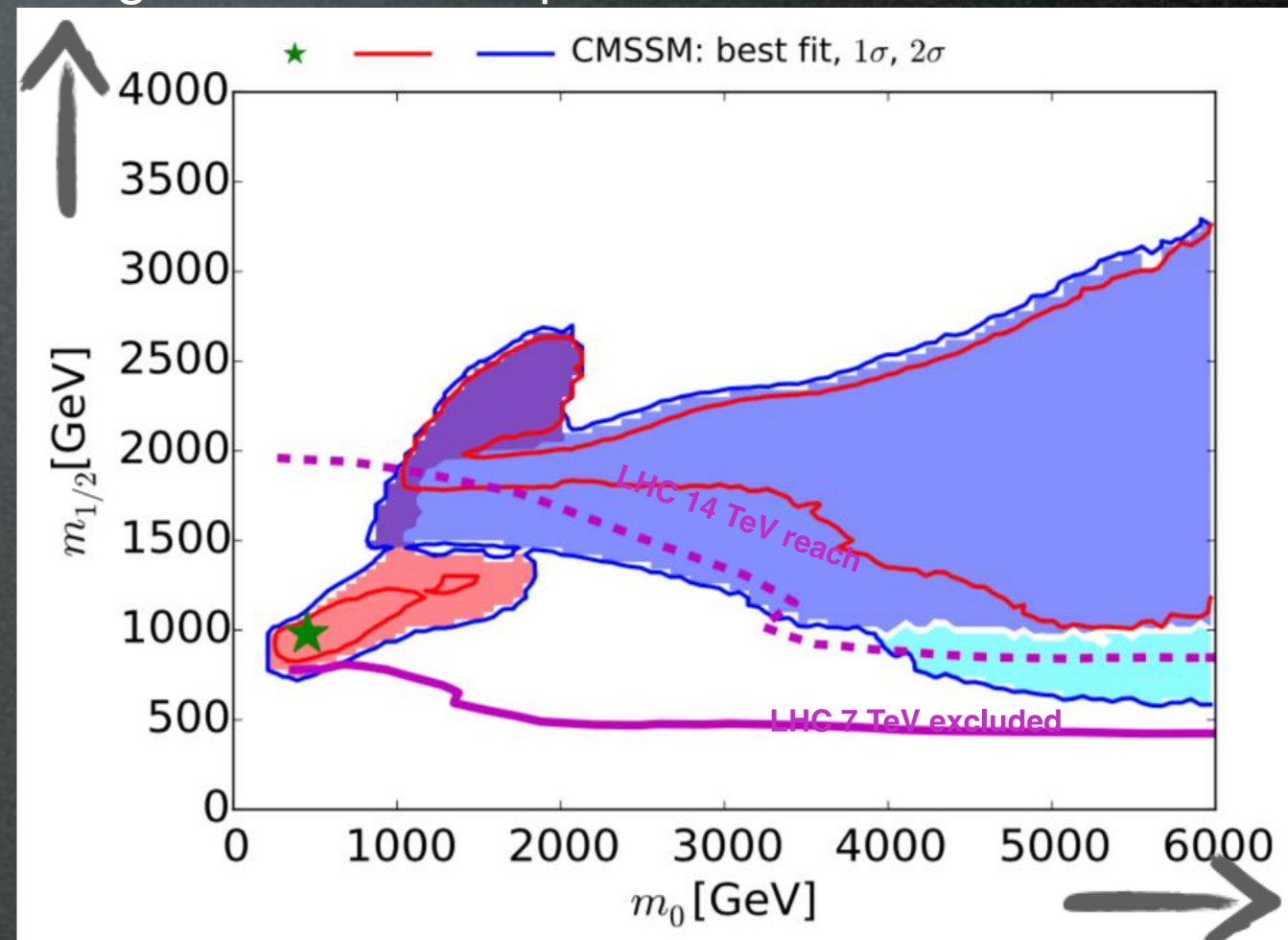
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E.g. SUSY: cMSSM, pMSSM...

parameters: **tens** (or even **hundreds!**)

the glorious $m_0/m_{1/2}$ plane:



filled areas get Ω_{DM} right due to:

- stau coann.
- hybrid
- stop coann.
- A/H funnel
- $\tilde{\chi}_1^\pm$ coann.
- focus point

Production at colliders

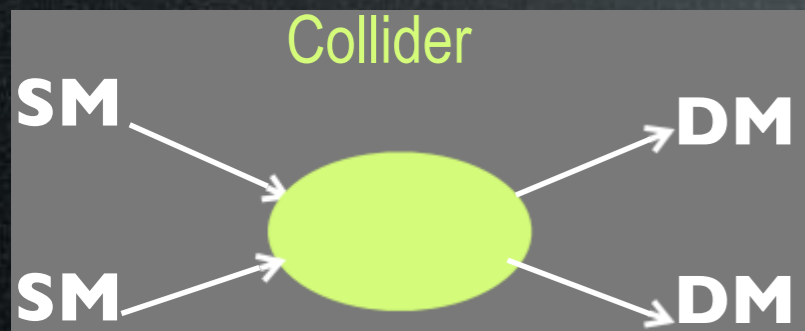
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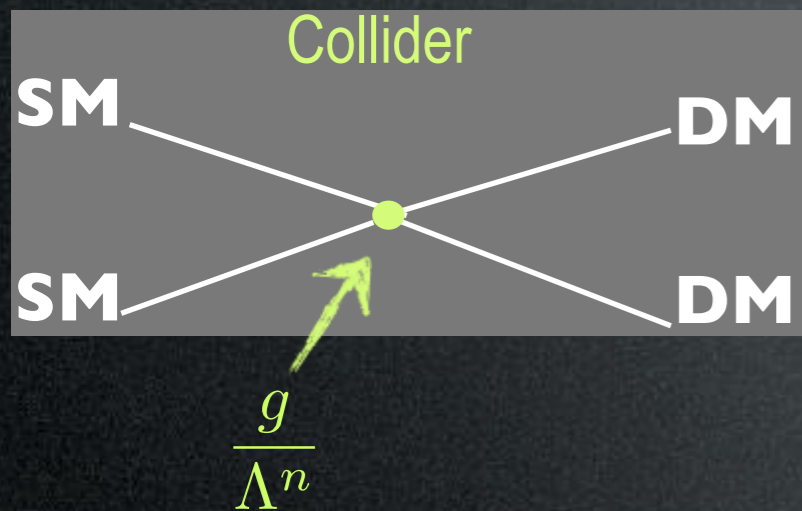
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$$\frac{1}{\Lambda_1^2} [q\bar{q}] [\chi\bar{\chi}]$$

$$\frac{1}{\Lambda_2^2} [q\gamma_\mu\bar{q}] [\chi\gamma^\mu\bar{\chi}]$$

parameters: Λ (assuming $g \simeq 1$)

Production at colliders

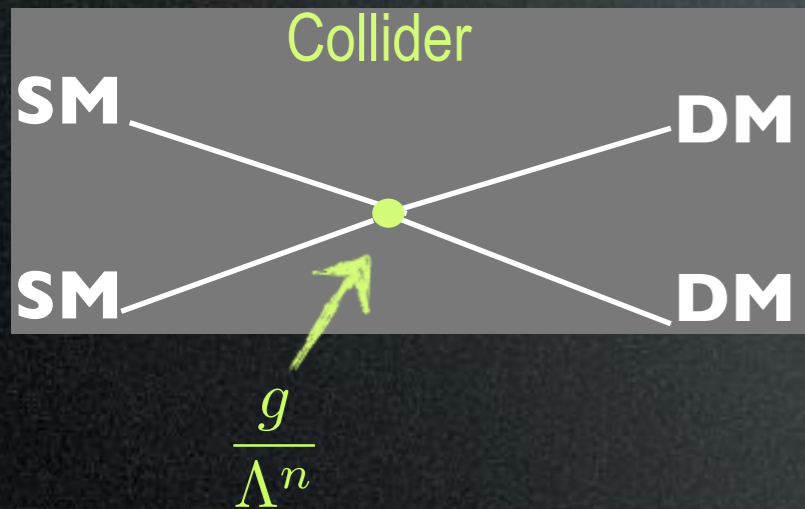
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(a) Operators for Dirac fermion DM

Name	Operator	Dimension	SI/SD
D1	$\frac{m_q}{\Lambda^3} \bar{\chi} \chi \bar{q} q$	7	SI
D2	$\frac{i m_q}{\Lambda^3} \bar{\chi} \gamma^5 \chi \bar{q} q$	7	N/A
D3	$\frac{i m_q}{\Lambda^3} \bar{\chi} \chi \bar{q} \gamma^5 q$	7	N/A
D4	$\frac{m_q}{\Lambda^3} \bar{\chi} \gamma^5 \chi \bar{q} \gamma^5 q$	7	N/A
D5	$\frac{1}{\Lambda^2} \bar{\chi} \gamma^\mu \chi \bar{q} \gamma_\mu q$	6	SI
D6	$\frac{1}{\Lambda^2} \bar{\chi} \gamma^\mu \gamma^5 \chi \bar{q} \gamma_\mu q$	6	N/A
D7	$\frac{1}{\Lambda^2} \bar{\chi} \gamma^\mu \chi \bar{q} \gamma_\mu \gamma^5 q$	6	N/A
D8	$\frac{1}{\Lambda^2} \bar{\chi} \gamma^\mu \gamma^5 \chi \bar{q} \gamma_\mu \gamma^5 q$	6	SD
D9	$\frac{1}{\Lambda^2} \bar{\chi} \sigma^{\mu\nu} \chi \bar{q} \sigma_{\mu\nu} q$	6	SD
D10	$\frac{i}{\Lambda^2} \bar{\chi} \sigma^{\mu\nu} \gamma^5 \chi \bar{q} \sigma_{\mu\nu} q$	6	N/A
D11	$\frac{\alpha_s}{\Lambda^3} \bar{\chi} \chi G^{\mu\nu} G_{\mu\nu}$	7	SI
D12	$\frac{\alpha_s}{\Lambda^3} \bar{\chi} \gamma^5 \chi G^{\mu\nu} G_{\mu\nu}$	7	N/A
D13	$\frac{\alpha_s}{\Lambda^3} \bar{\chi} \chi G^{\mu\nu} \tilde{G}_{\mu\nu}$	7	N/A
D14	$\frac{\alpha_s}{\Lambda^3} \bar{\chi} \gamma^5 \chi G^{\mu\nu} \tilde{G}_{\mu\nu}$	7	N/A

(b) Operators for Complex scalar DM

Name	Operator	Dimension	SI/SD
C1	$\frac{m_q}{\Lambda^2} \phi^\dagger \phi \bar{q} q$	6	SI
C2	$\frac{m_q}{\Lambda^2} \phi^\dagger \phi \bar{q} \gamma^5 q$	6	N/A
C3	$\frac{1}{\Lambda^2} \phi^\dagger \overleftrightarrow{\partial}_\mu \phi \bar{q} \gamma^\mu q$	6	SI
C4	$\frac{1}{\Lambda^2} \phi^\dagger \overleftrightarrow{\partial}_\mu \phi \bar{q} \gamma^\mu \gamma^5 q$	6	N/A
C5	$\frac{\alpha_s}{\Lambda^3} \phi^\dagger \phi G^{\mu\nu} G_{\mu\nu}$	6	SI
C6	$\frac{\alpha_s}{\Lambda^3} \phi^\dagger \phi G^{\mu\nu} \tilde{G}_{\mu\nu}$	6	N/A

Tim Tait, 2010+

and many many many others

Production at colliders

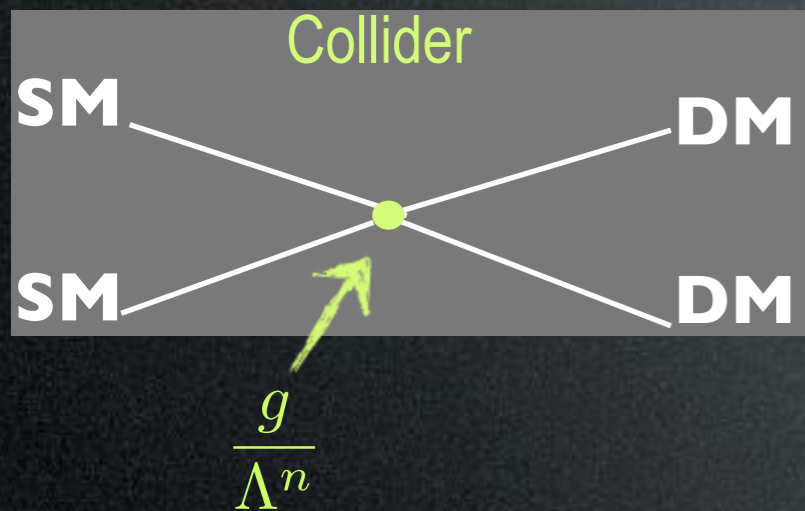
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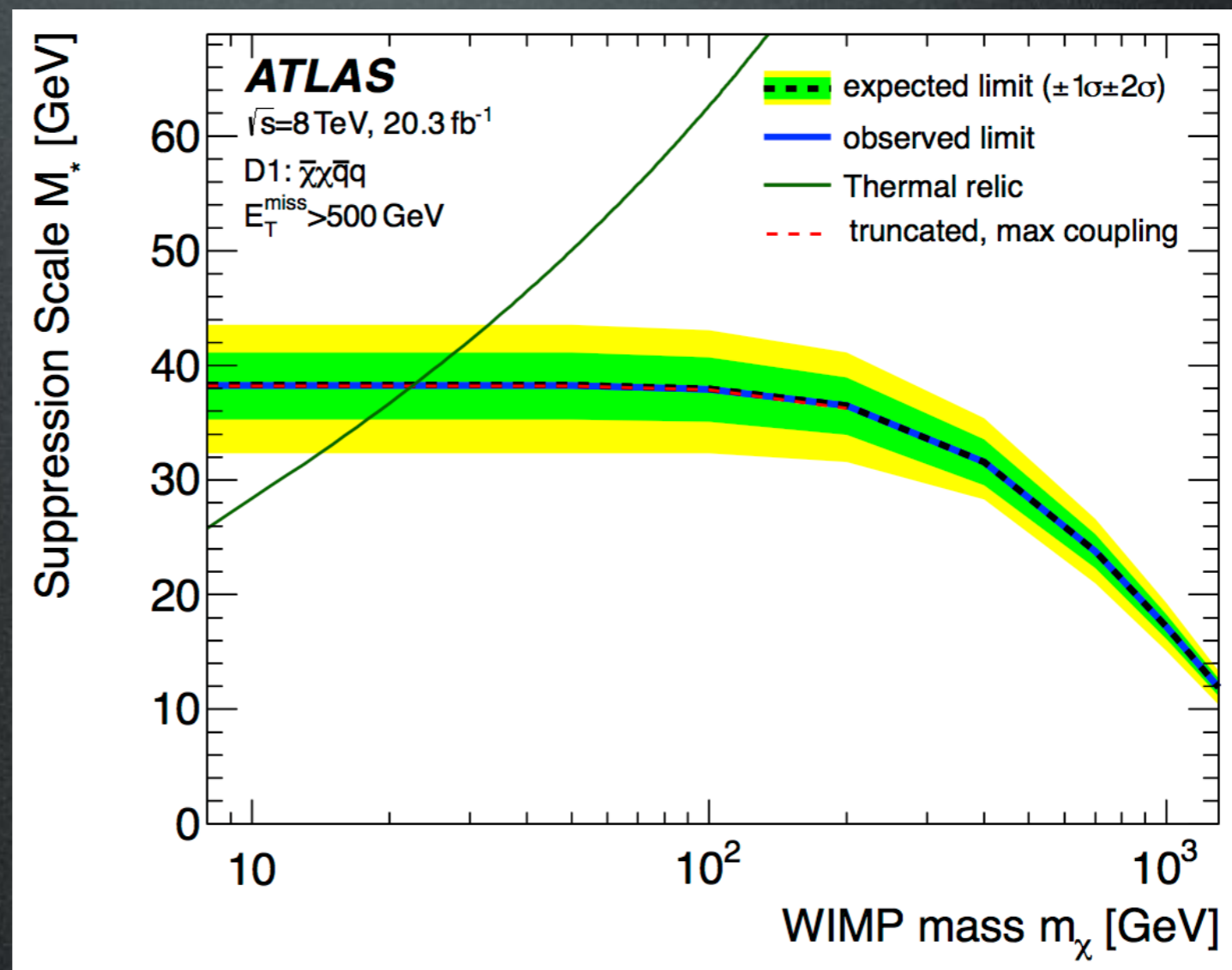
simplified models

a simple model

a full theory



limits on Δ



Production at colliders

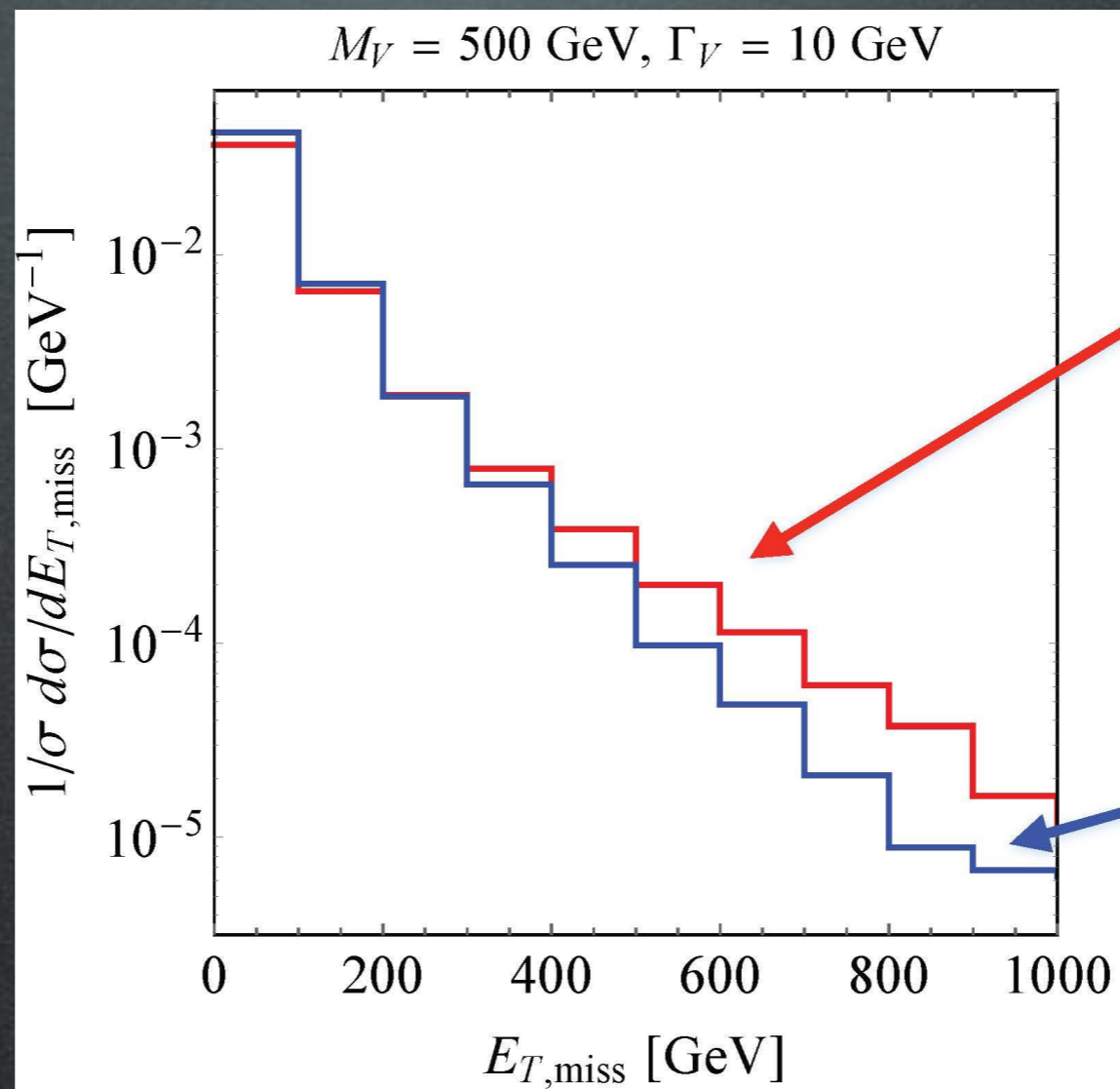
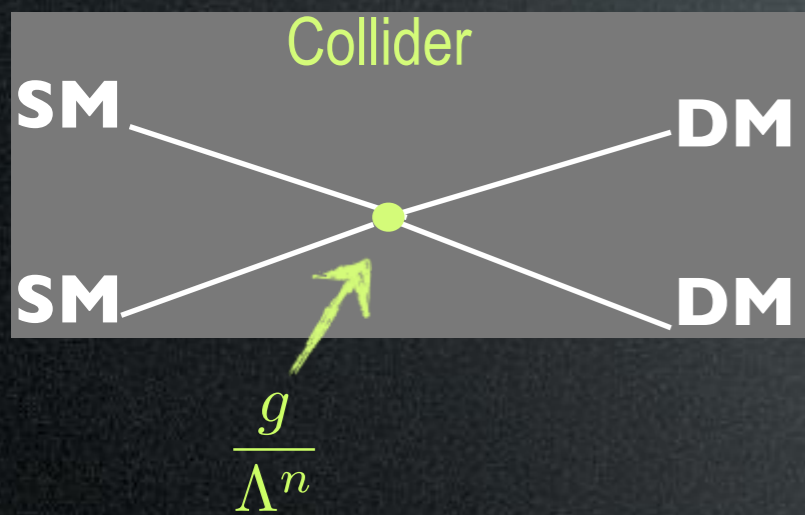
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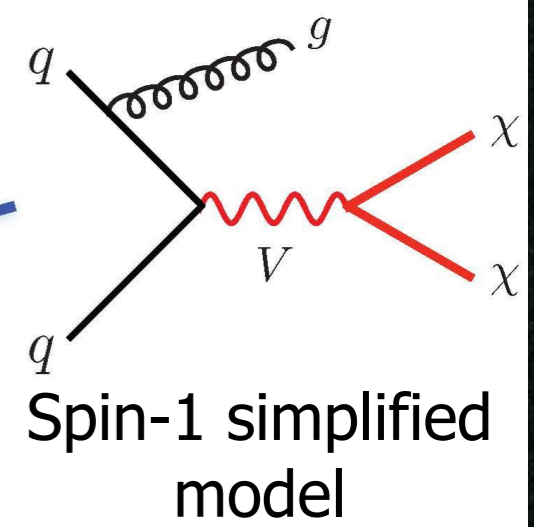
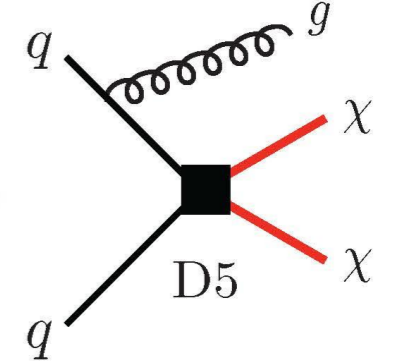
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Vector operator



Production at colliders

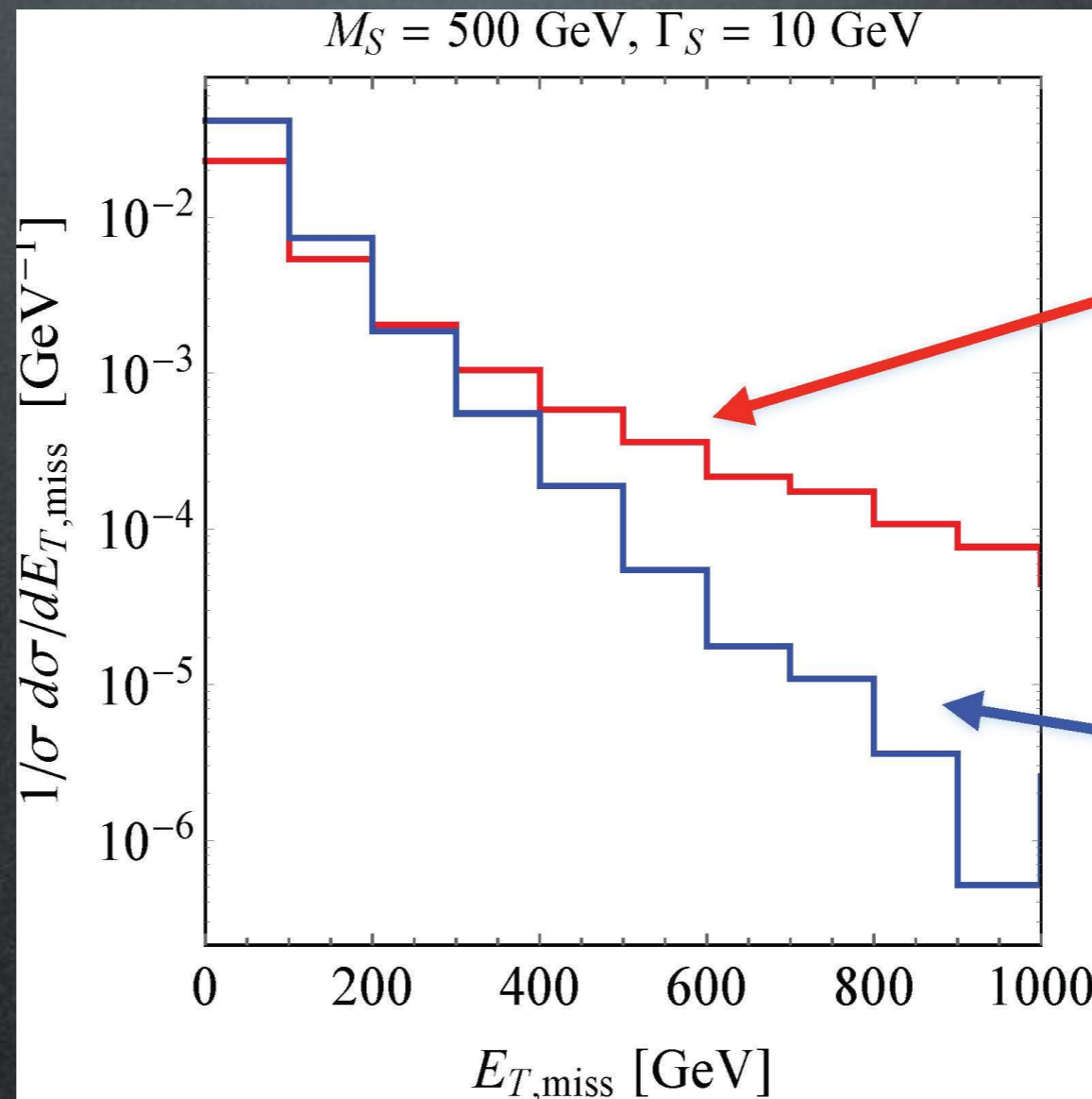
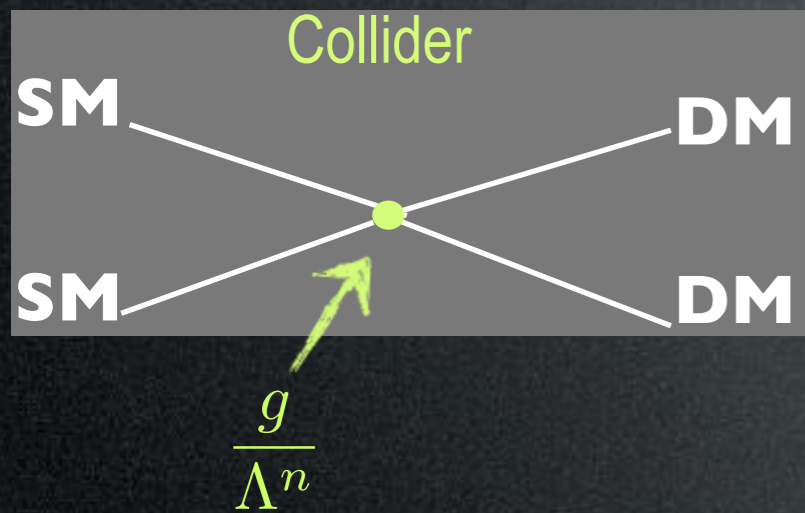
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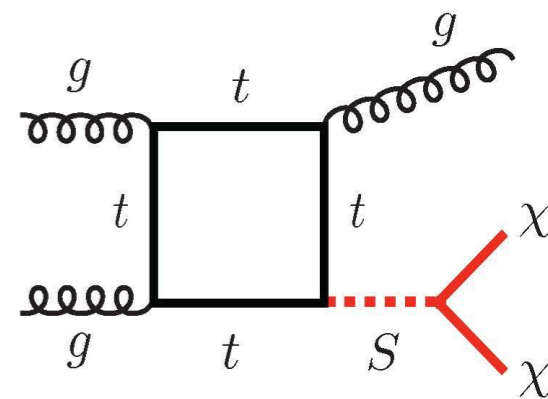
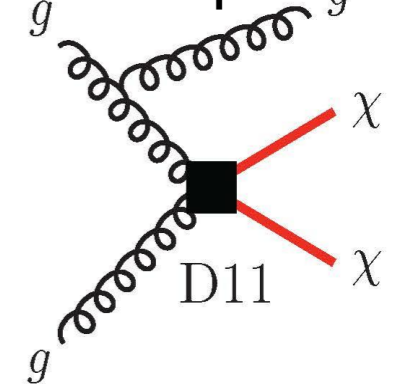
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Gluonic operator



Spin-0 simplified model

Production at colliders

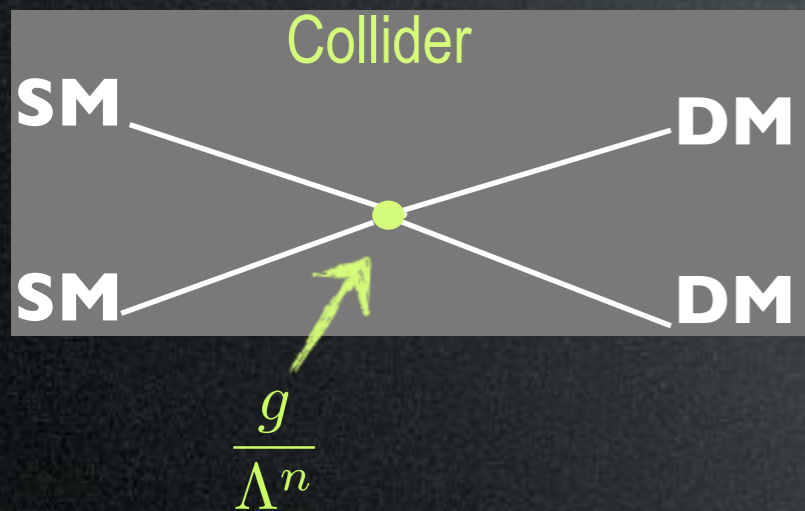
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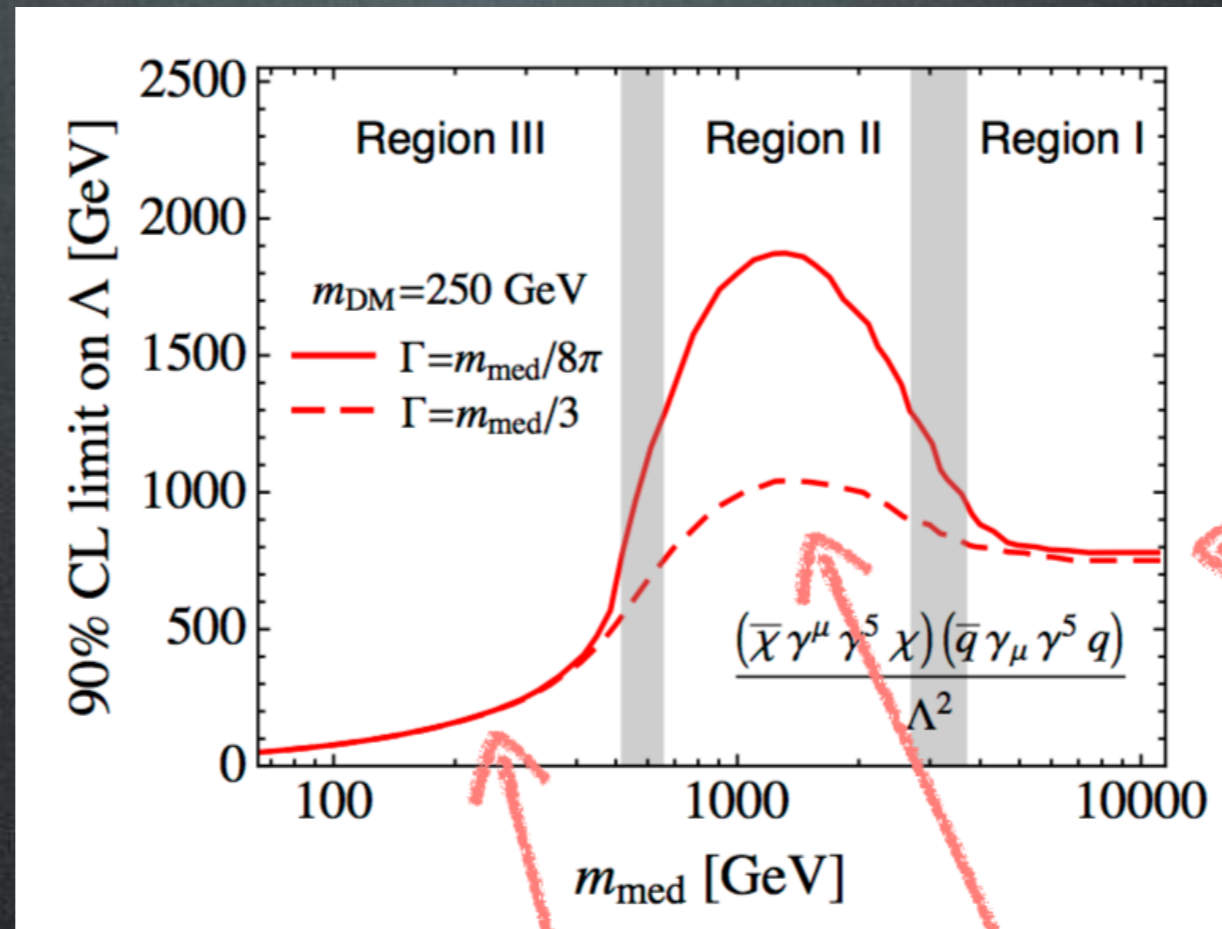
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but how 'sensible' are they?

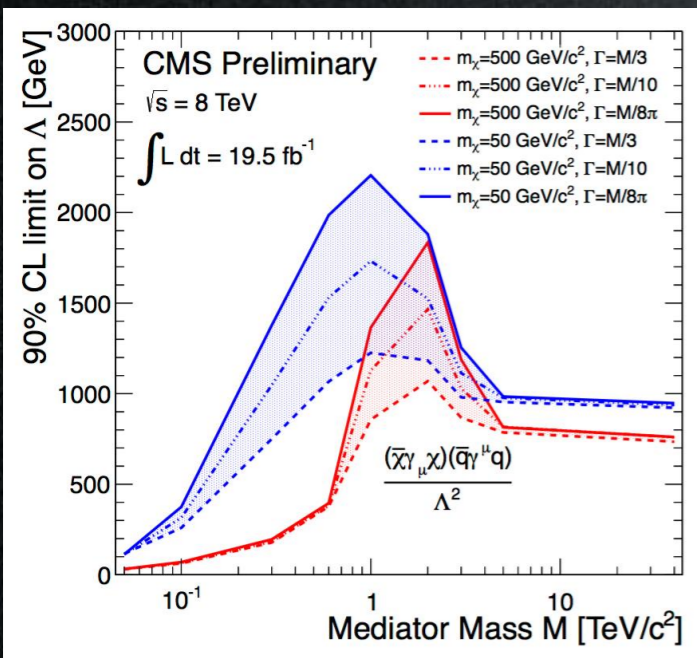


EFT puts bounds on Λ

but Λ is actually $\Lambda = \frac{m_{med}}{\sqrt{g_q g_\chi}}$
at large m_{med} , the two agree

resolving, at small m_{med} ,
what is the discrepancy?
i.e. what is the limit deduced on Λ ,
defined as above?

limits agree



Buchmuller, Dolan,
McCabe 1308.6799

EFT bound is
much larger

EFT underestimates bound
by some % up to factor 2 or more

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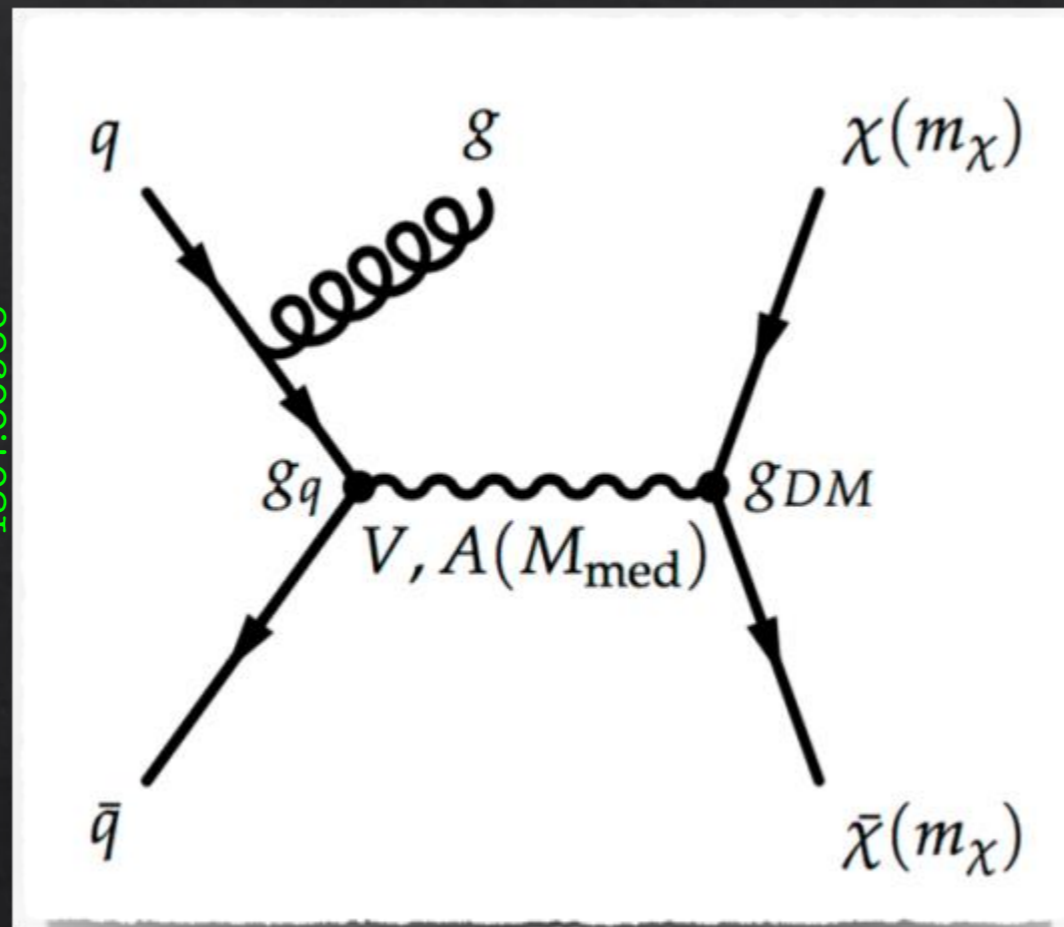
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ATLAS+CMS DM forum,
1507.00966



parameters:

$m_\chi, m_{\text{med}}, g_q, g_{\text{DM}}$

$$\mathcal{L}_{\text{vector}} = g_q \sum_{q=u,d,s,c,b,t} Z'_\mu \bar{q} \gamma^\mu q + g_\chi Z'_\mu \bar{\chi} \gamma^\mu \chi$$

$$\mathcal{L}_{\text{axial-vector}} = g_q \sum_{q=u,d,s,c,b,t} Z'_\mu \bar{q} \gamma^\mu \gamma^5 q + g_\chi Z'_\mu \bar{\chi} \gamma^\mu \gamma^5 \chi.$$

(and similarly for t-channel, scalar mediator, scalar DM etc...)

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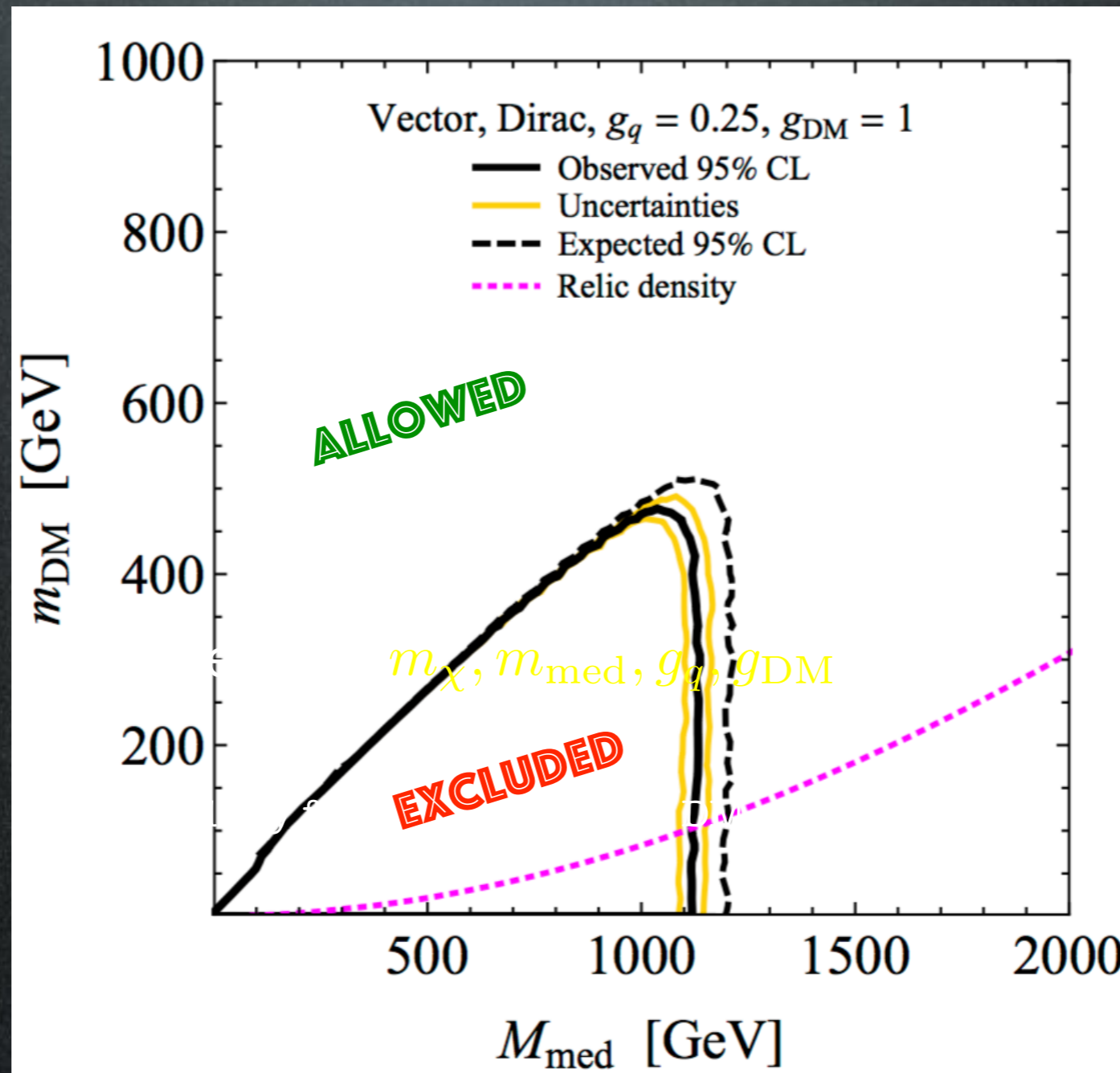
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mock plot:



Production at colliders

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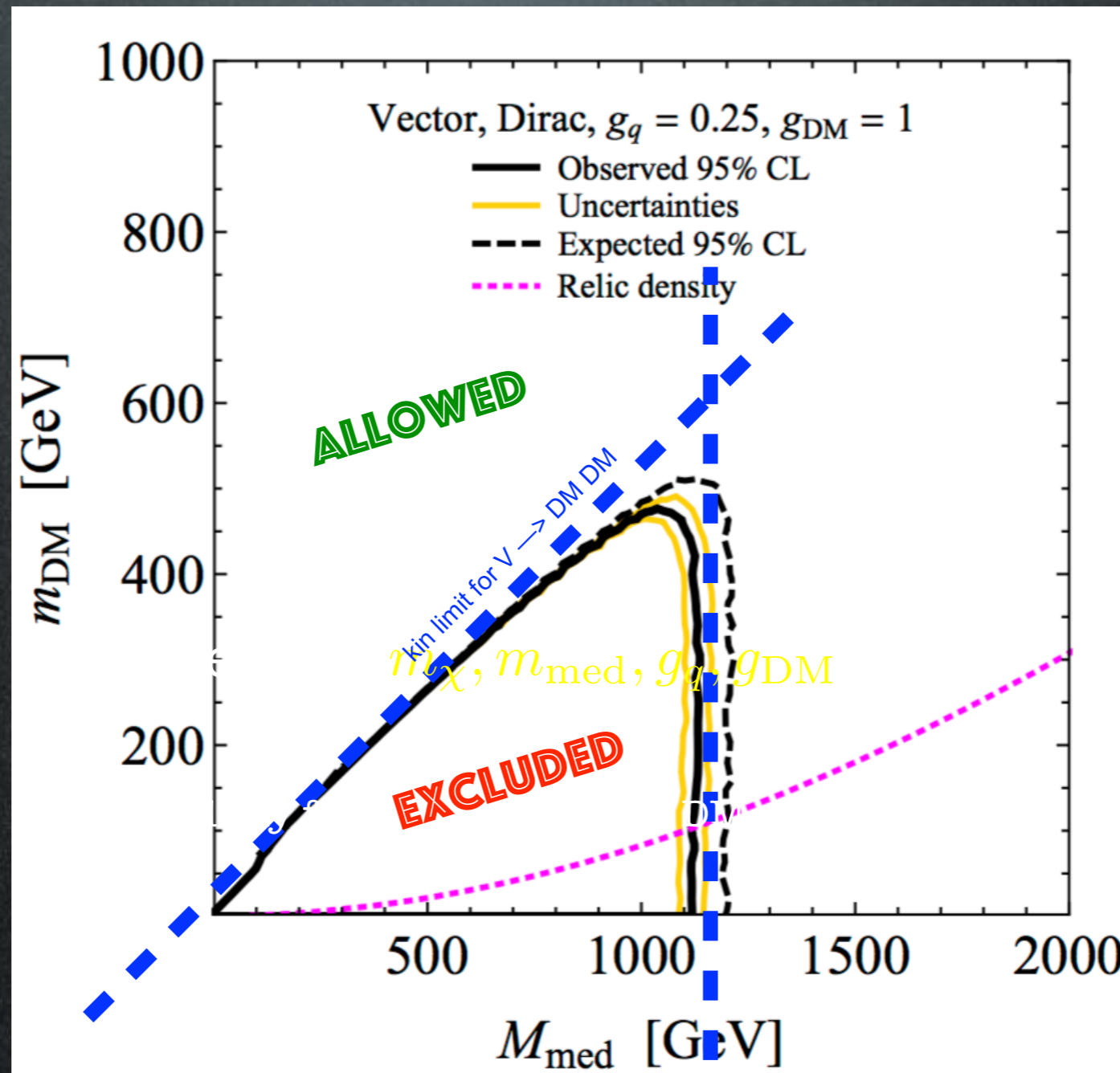
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mock plot:



LHCDDMWG, 1603.04156

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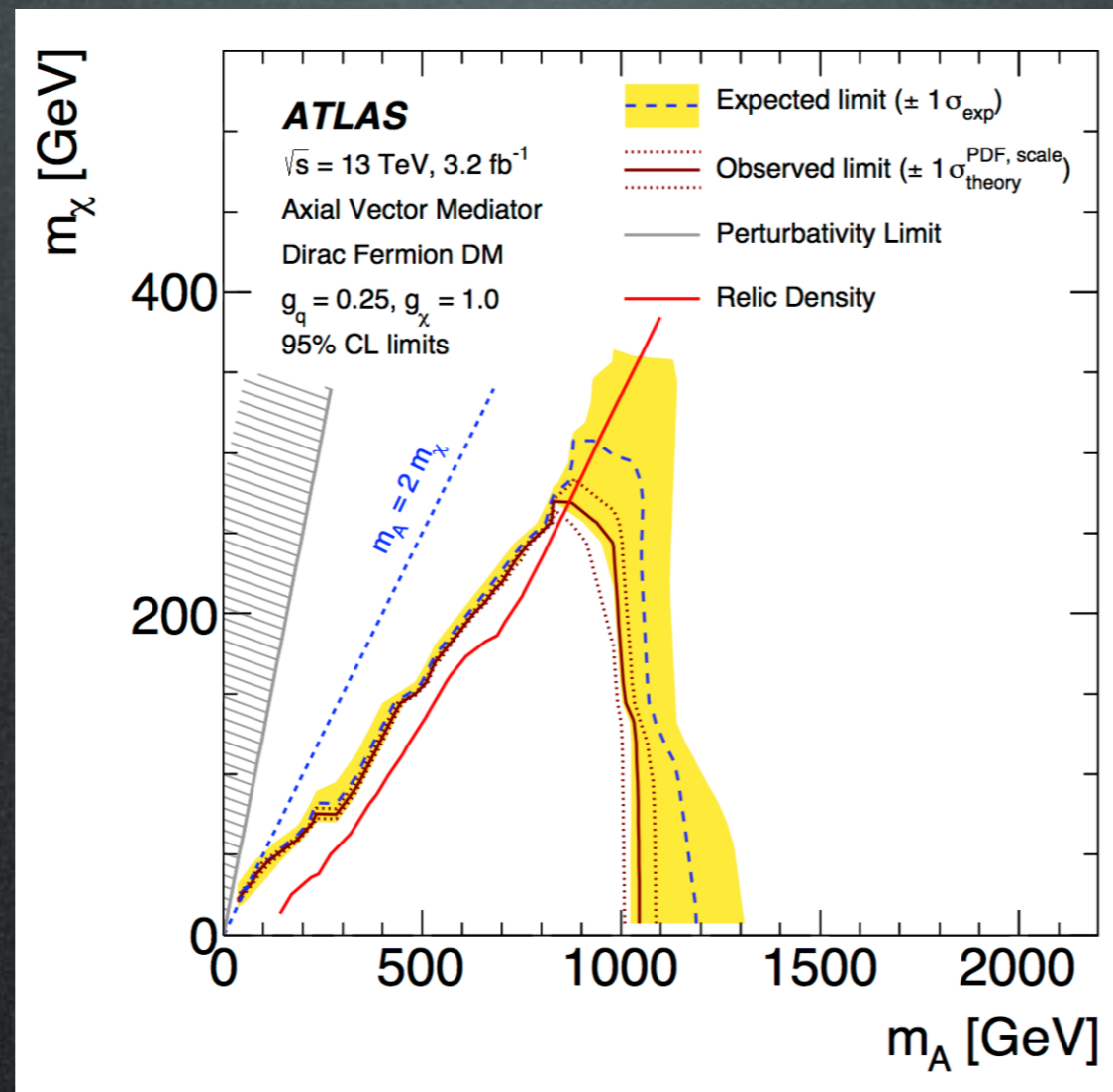
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actual plots:



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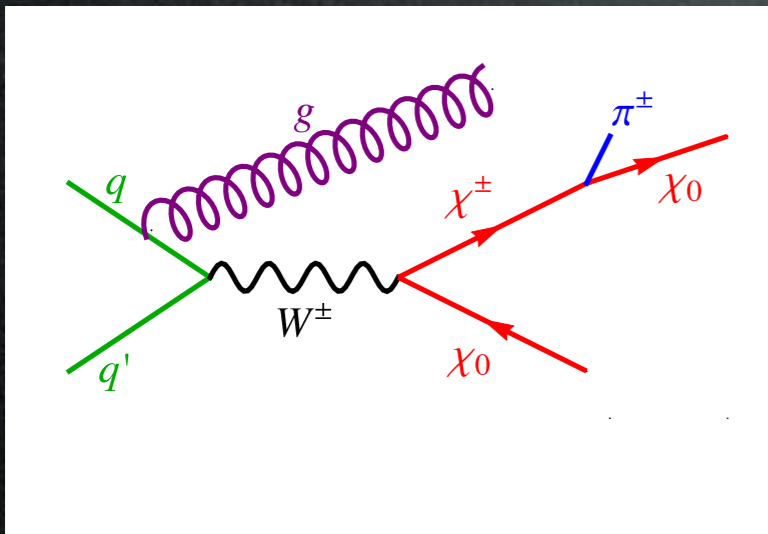
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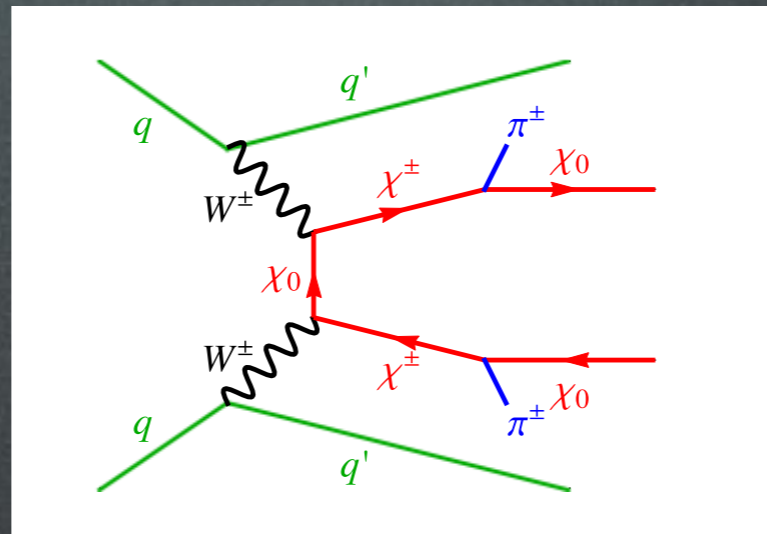
a simple model

a full theory

e.g. **pure WIMP** model (a.k.a. 'Minimal DM')



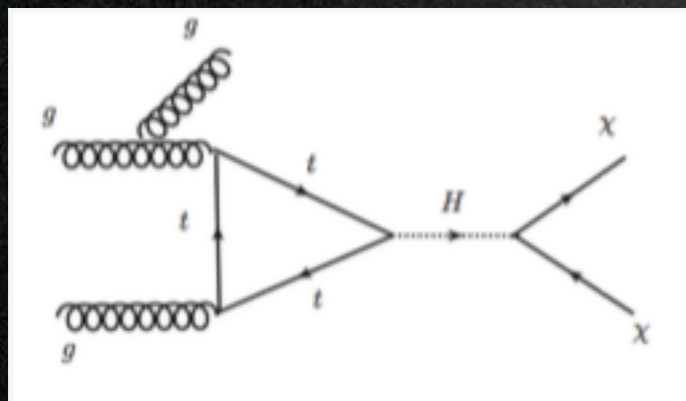
mono-jet



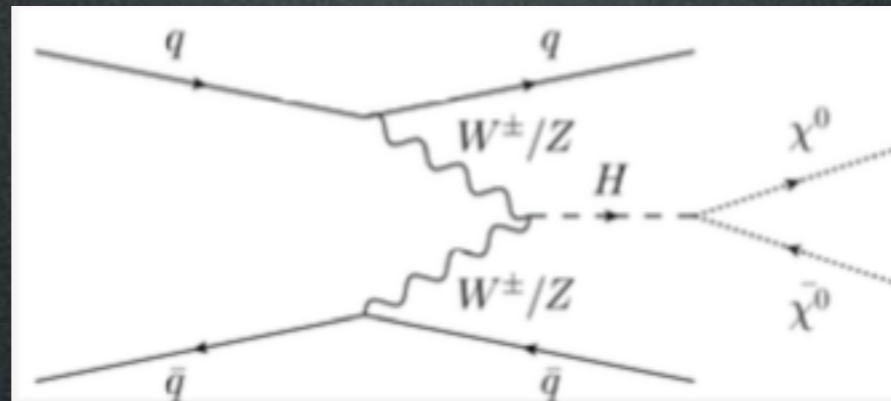
Vector Boson Fusion (VBF)

parameters: M_{DM}

e.g. **higgs portal** models



mono-jet

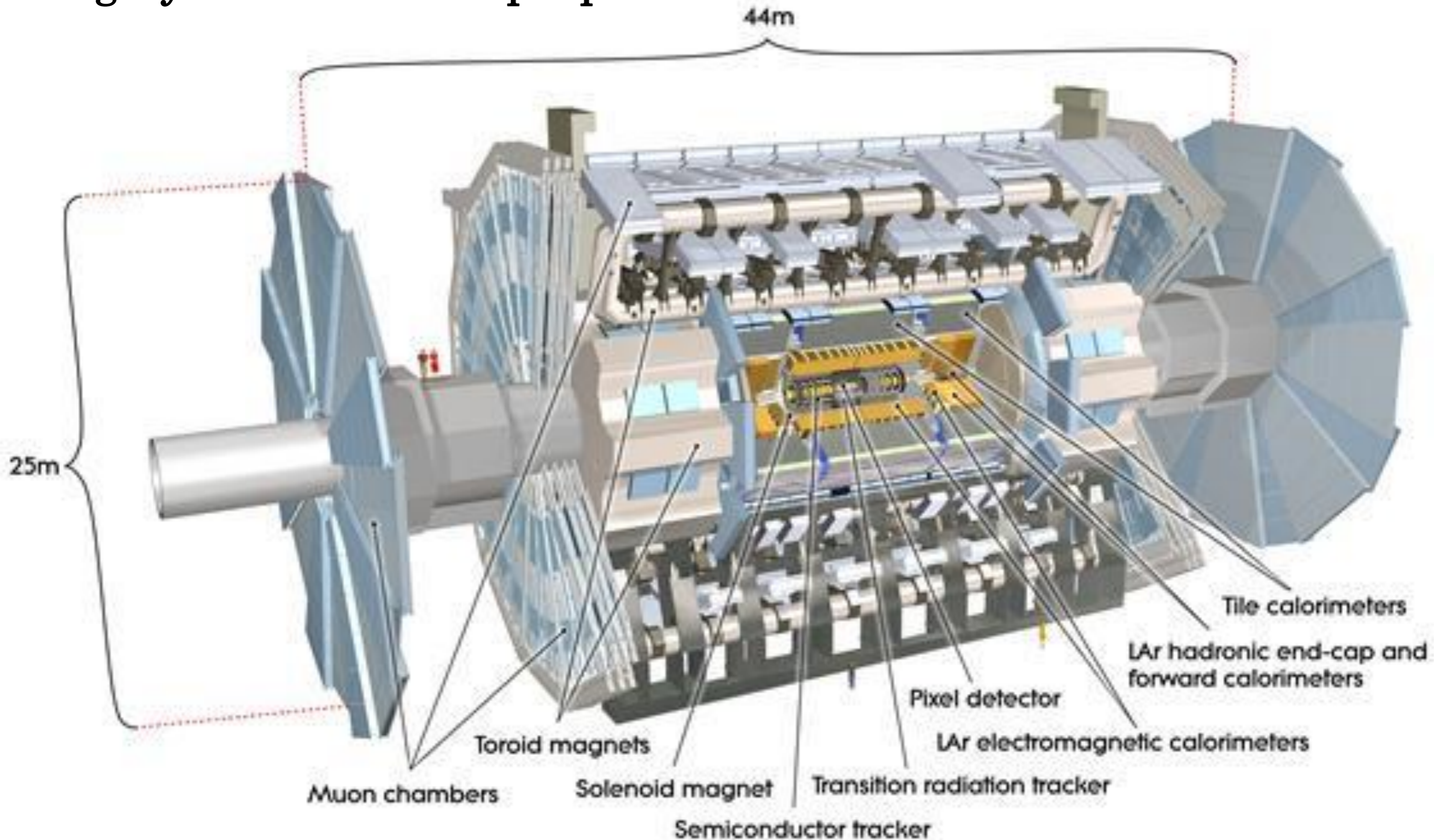


Vector Boson Fusion (VBF)

parameters: $g_{h\chi}$

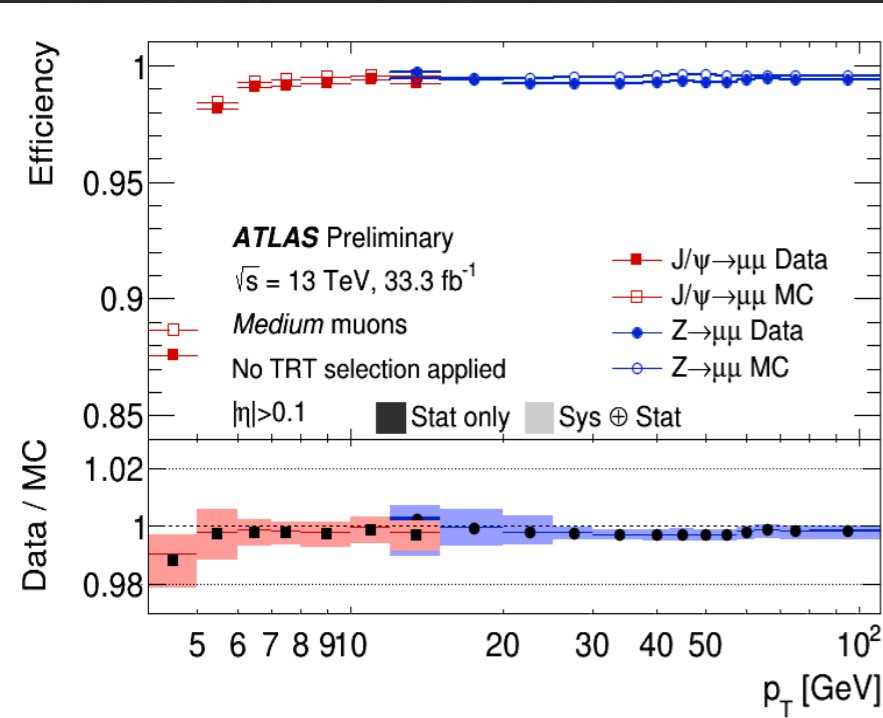
The ATLAS Detector

A highly efficient multi-purpose detector

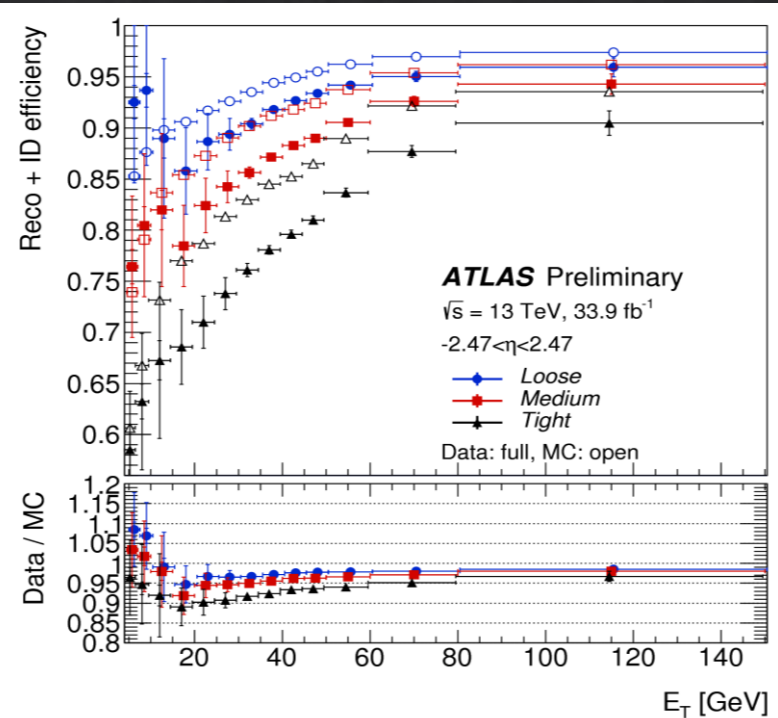


The ATLAS Detector

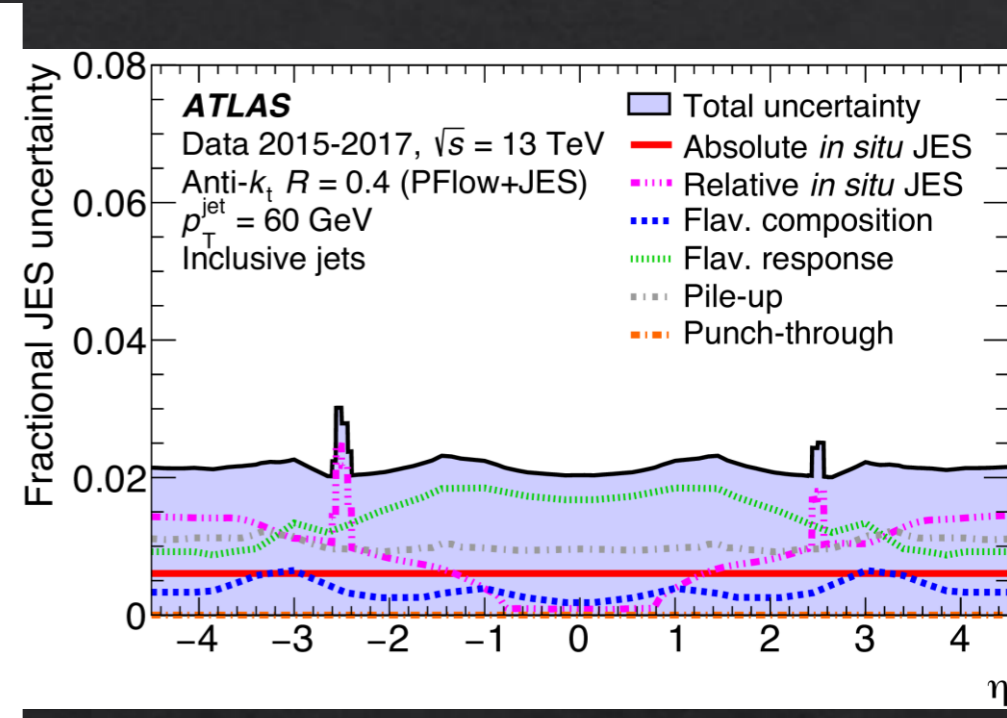
Example of measurements with the ATLAS detector



Electrons



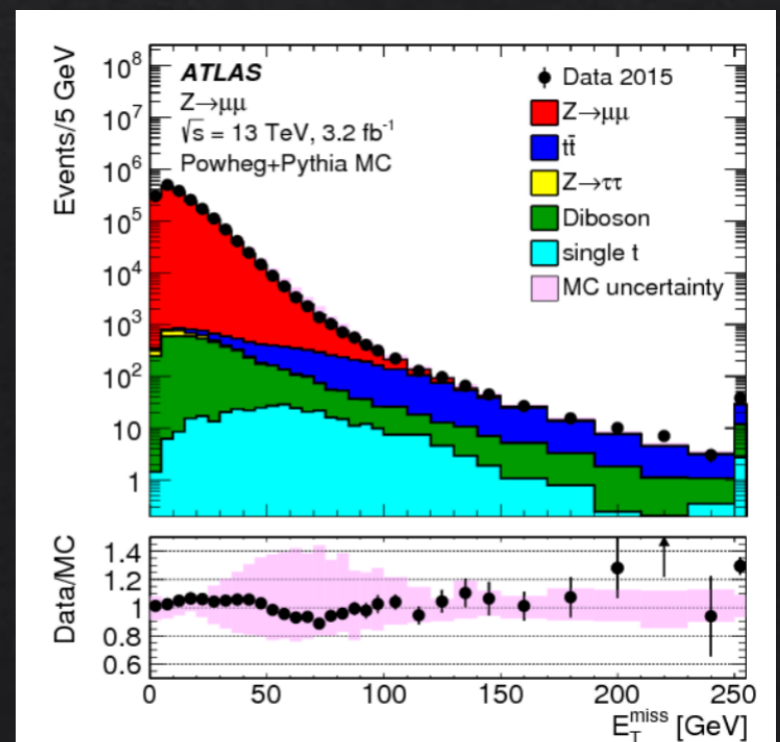
Muons



Jets

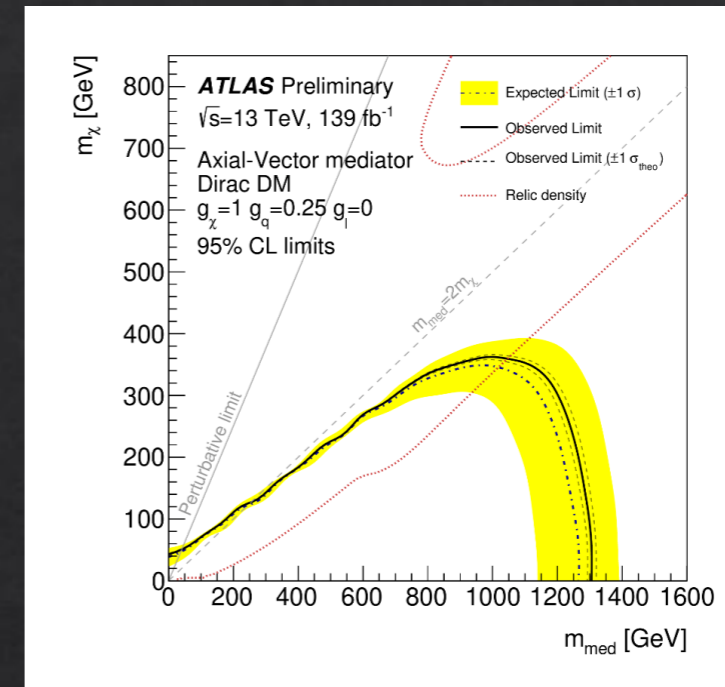
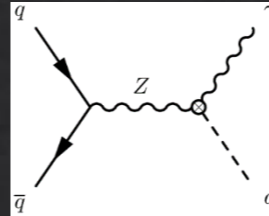
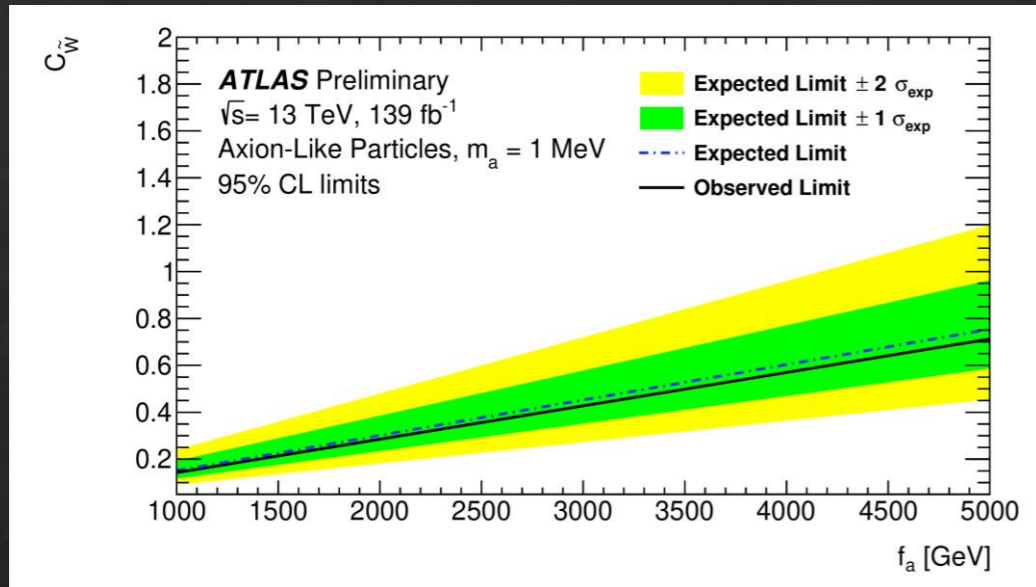
Missing Transverse Momentum

$$E_{x(y)}^{\text{miss}} = - \sum_{i \in \{\text{hard objects}\}} p_{x(y),i} - \sum_{j \in \{\text{soft signals}\}} p_{x(y),j}$$

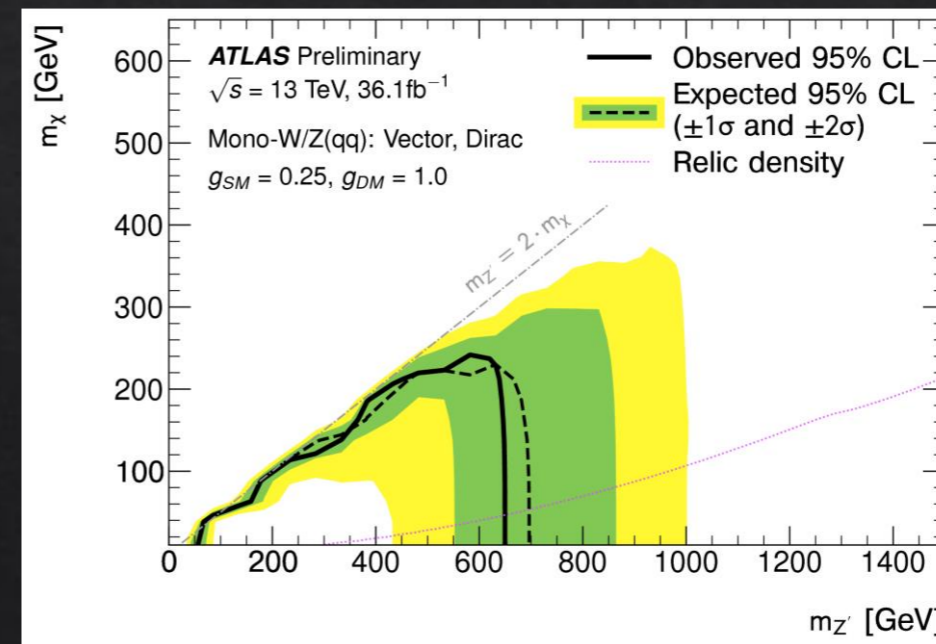
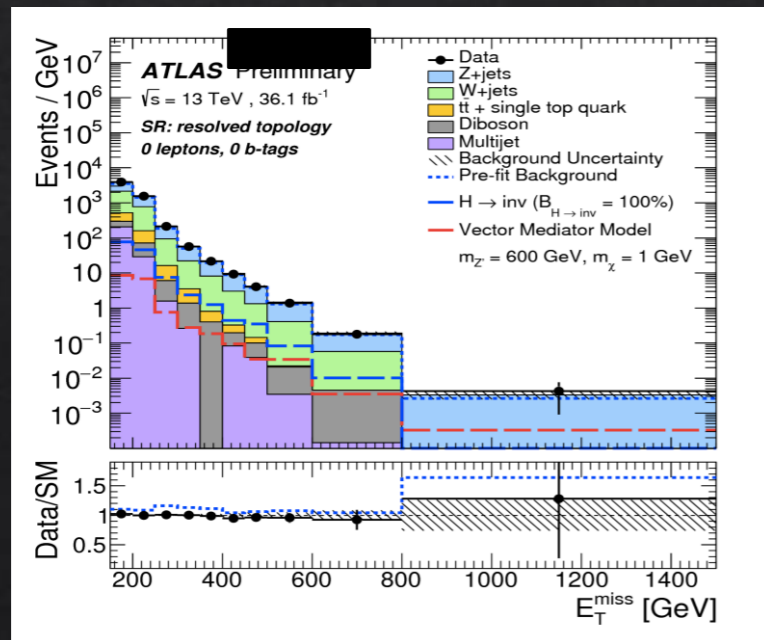


Mono-X searches

◇ Mono-photon searches:



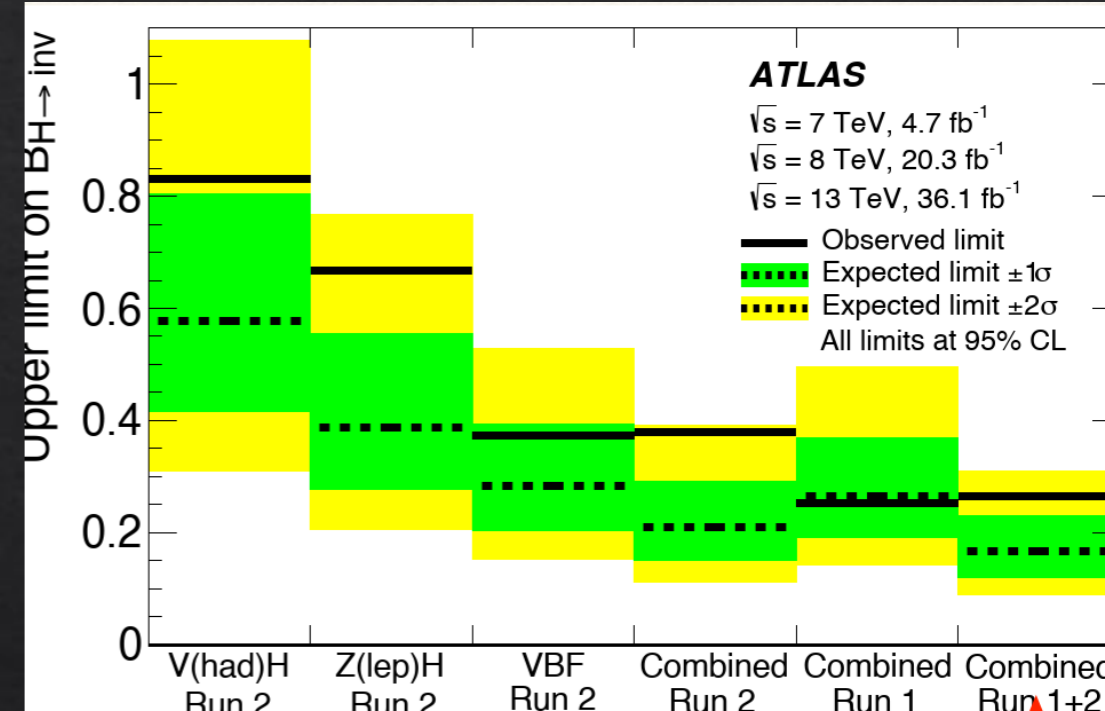
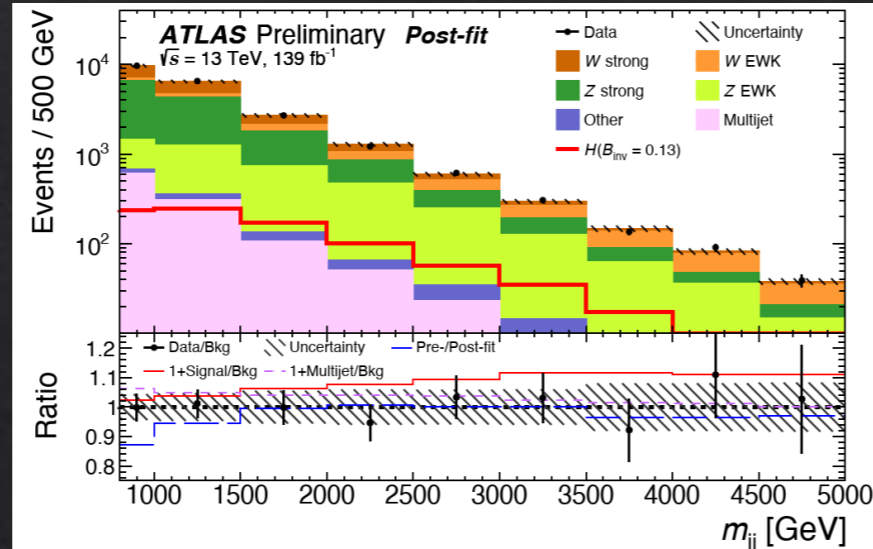
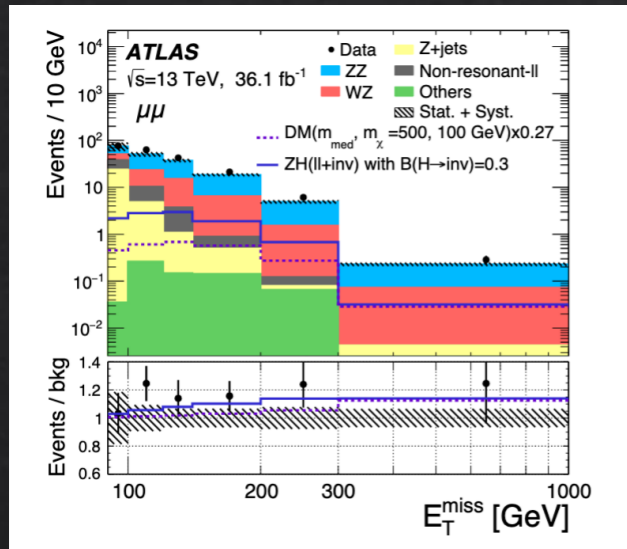
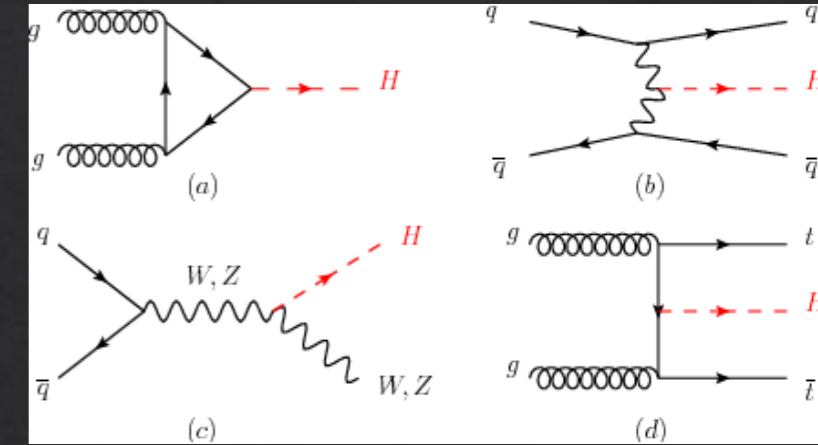
◇ Mono-W/Z searches:



◇ Many other Mono-X signatures being searched at the LHC

Higgs Portal

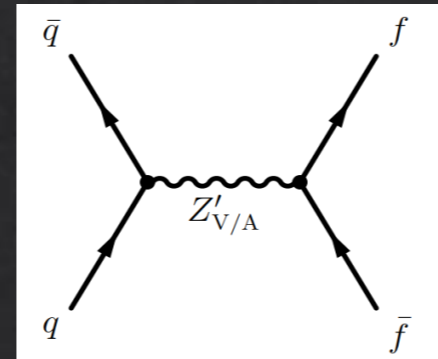
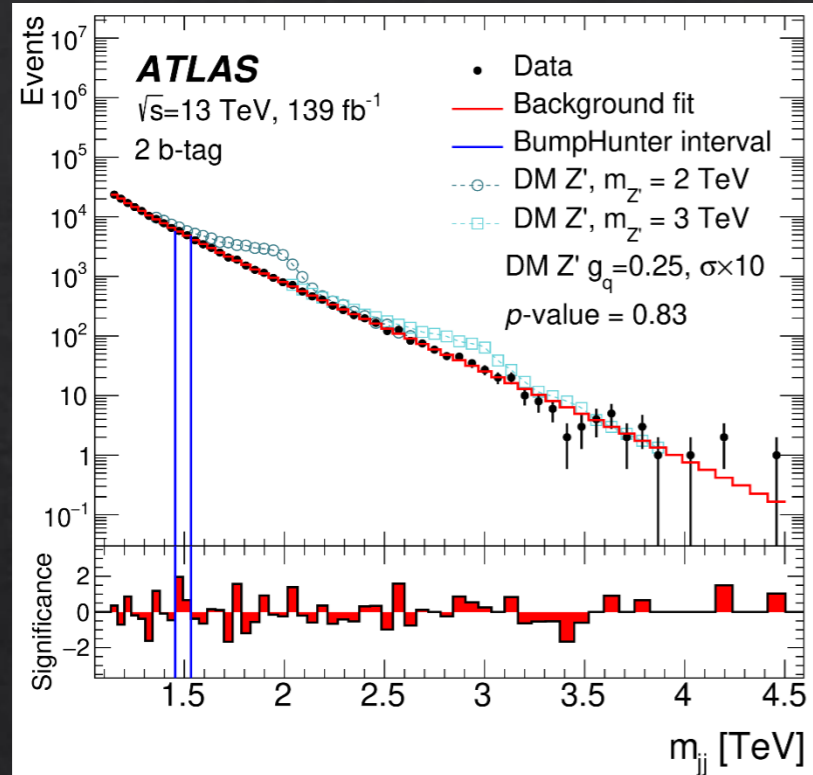
- ◇ Up-to-date measurements at the LHC
 - ◇ 20% BR possible for new physics
 - ◇ We can investigate unknown (invisible) decays
- ◇ Searches done in ZH, ttH and VBF modes:



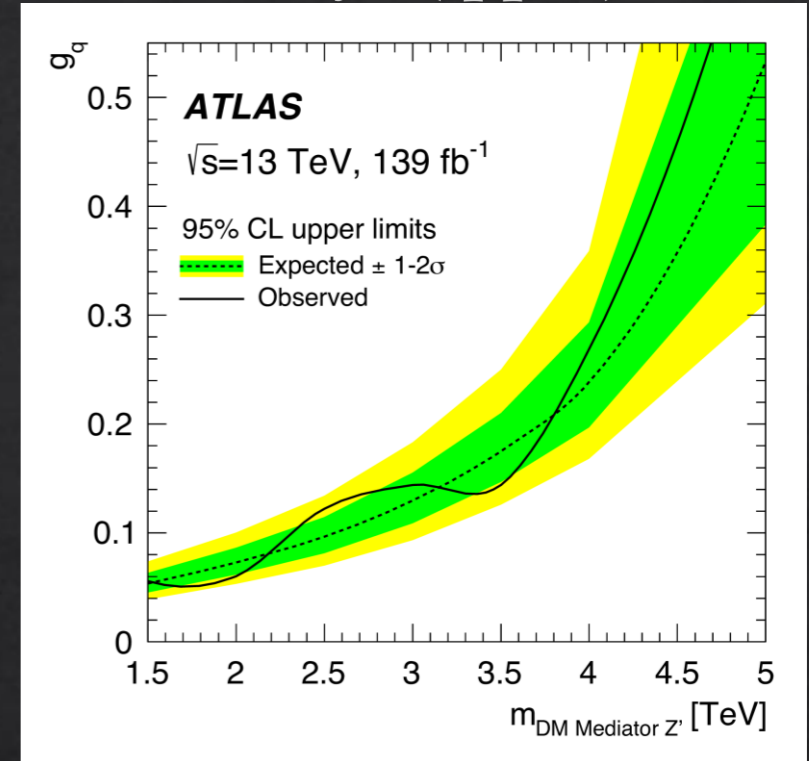
- ◇ Best limit achieved with ATLAS VBF
 - ◇ $BR(H \rightarrow \chi\chi) < 13\%$

Other Signatures

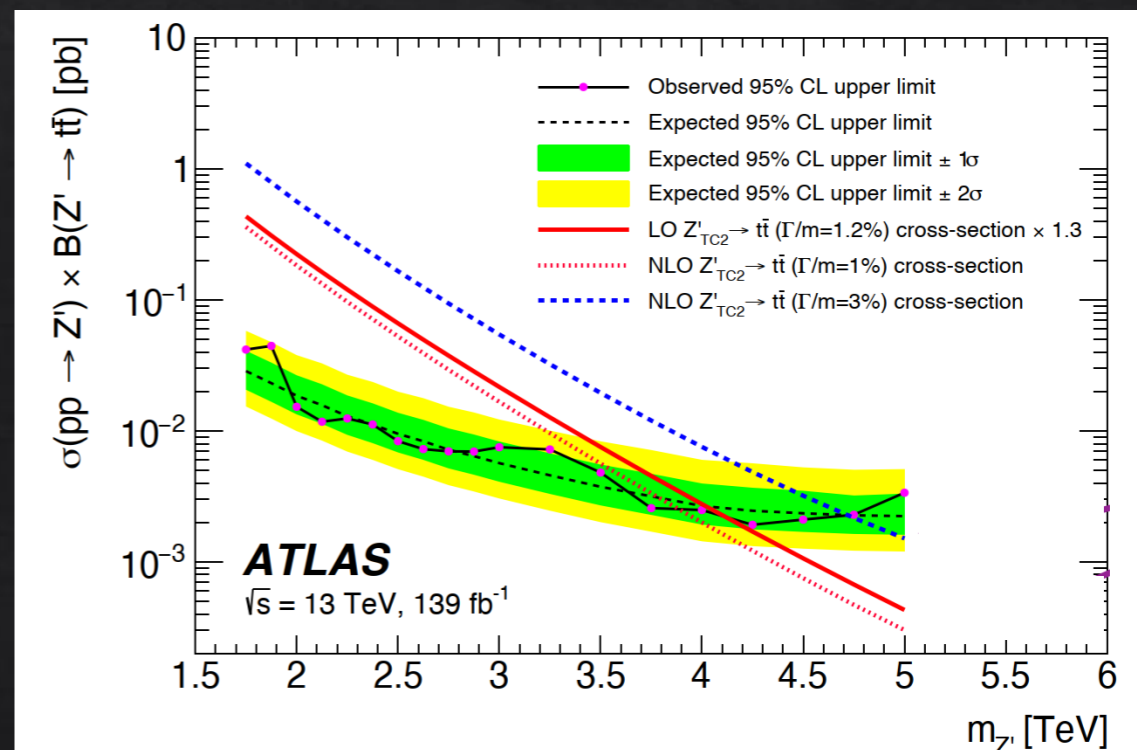
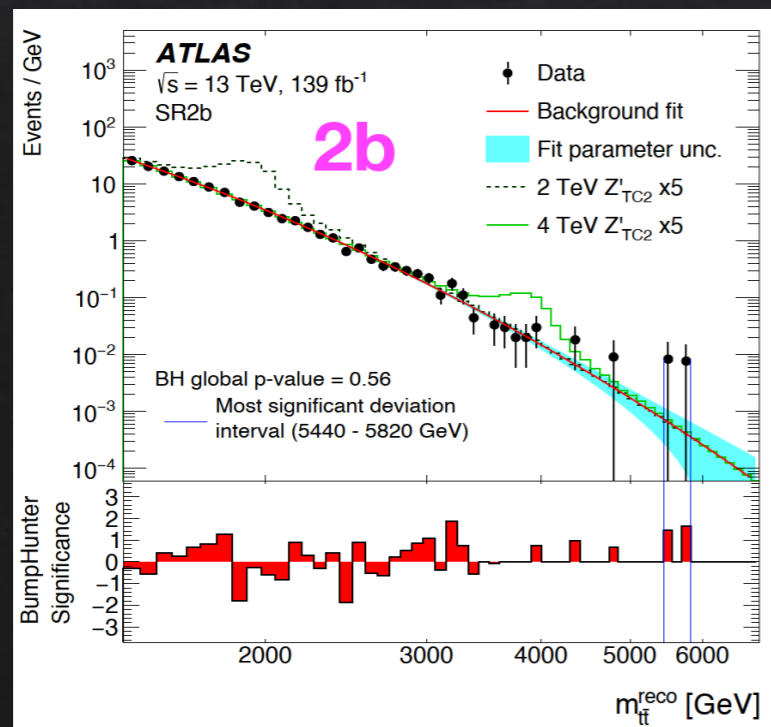
◆ Mediator searches: look for resonances in visible decays (qq, ll)



Leptophilic or
 Leptophobic Z' ,
 Low mass resonances

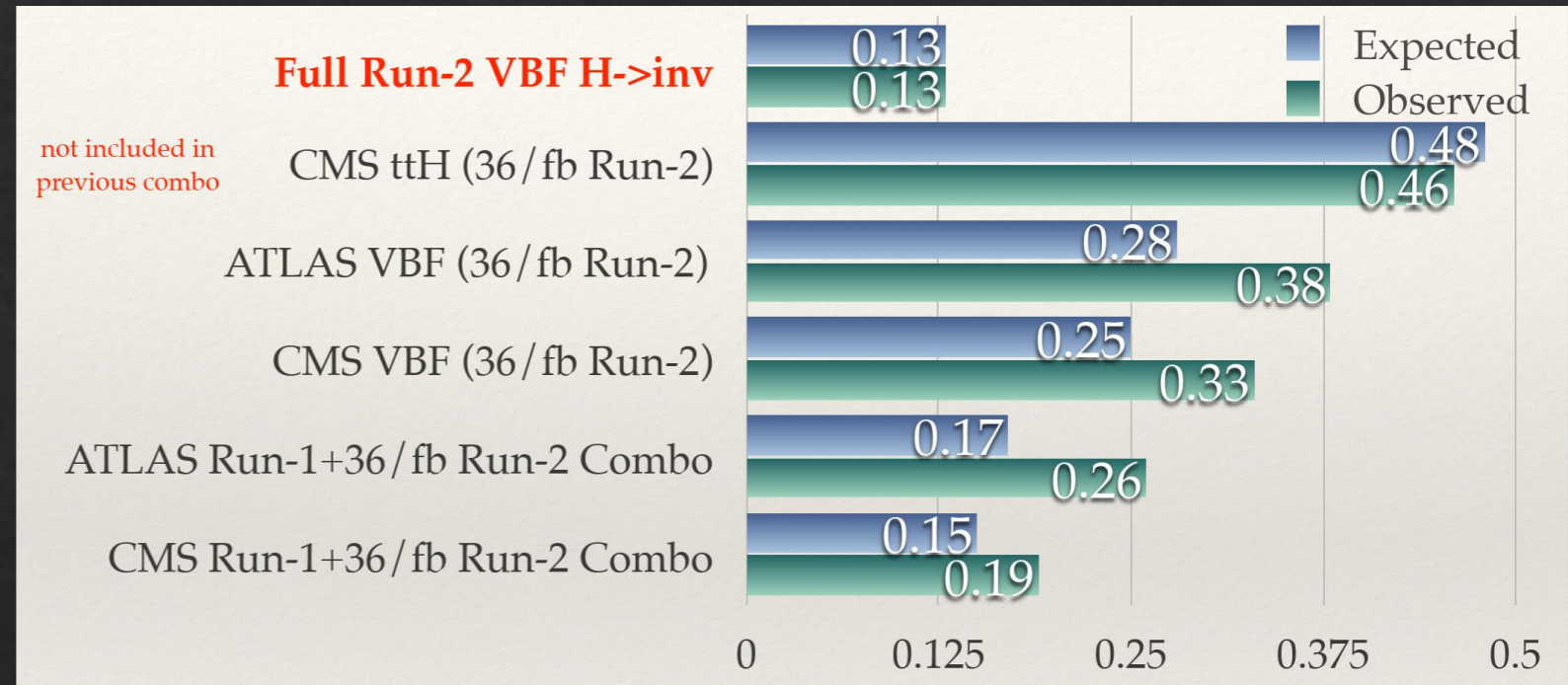


◆ $t\bar{t}$ resonances

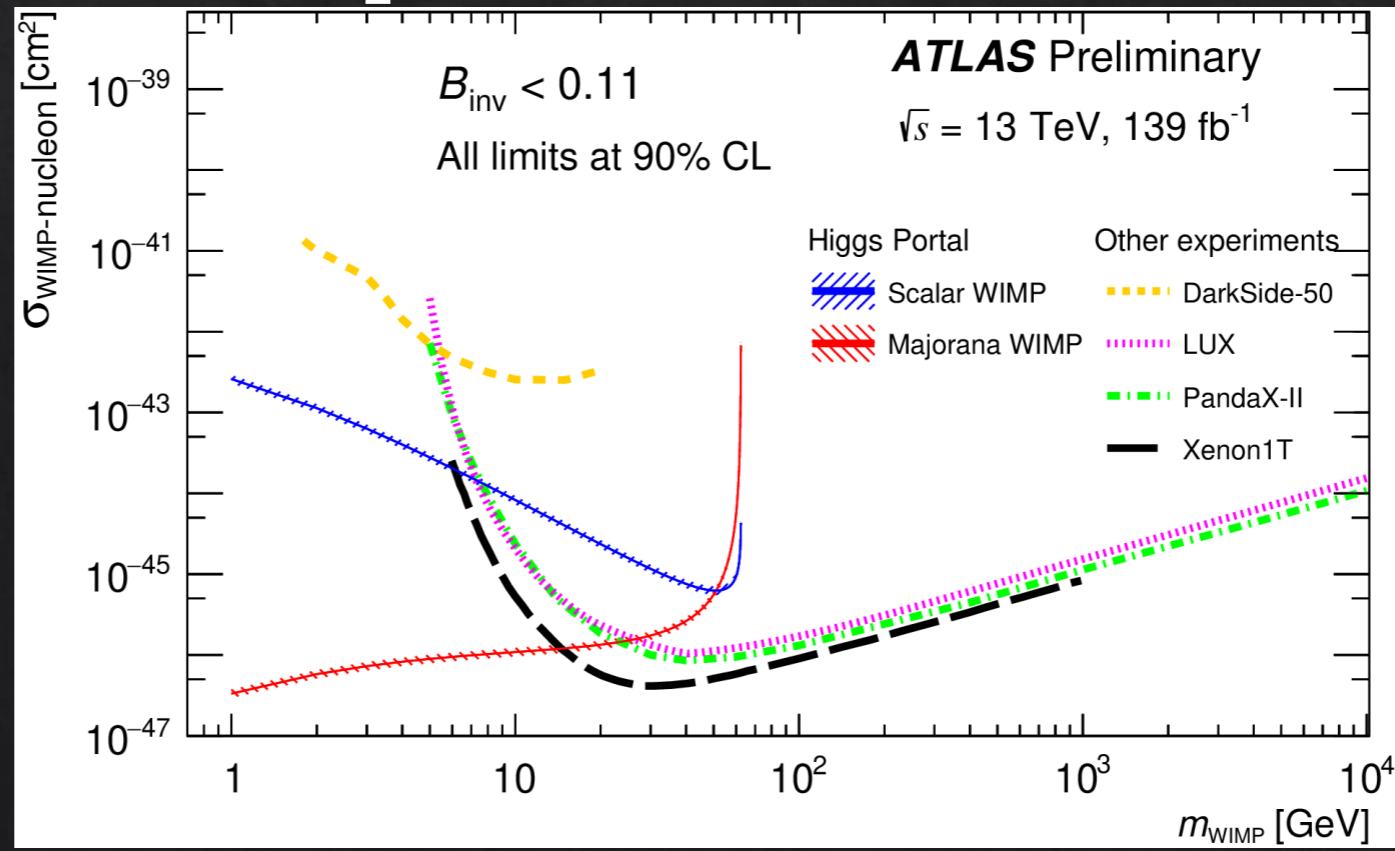


Higgs Portal

◆ Combined results:

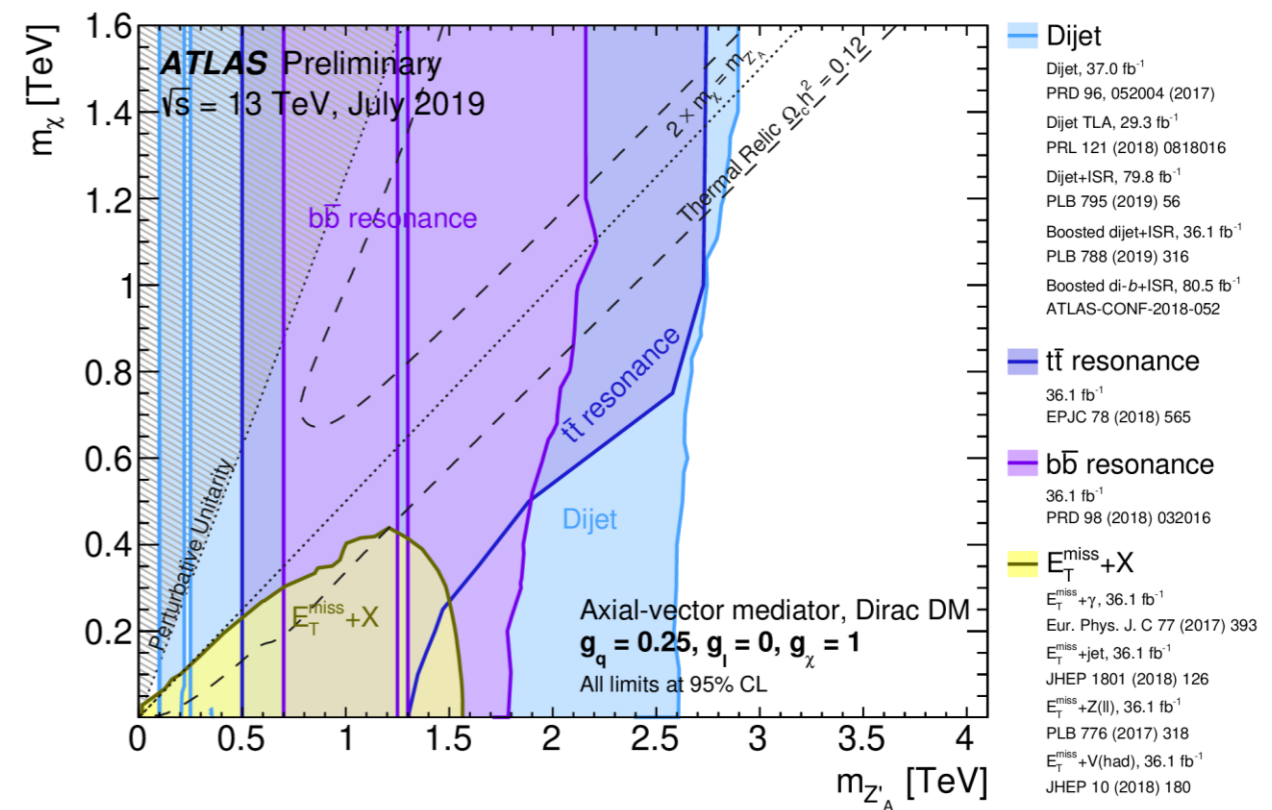
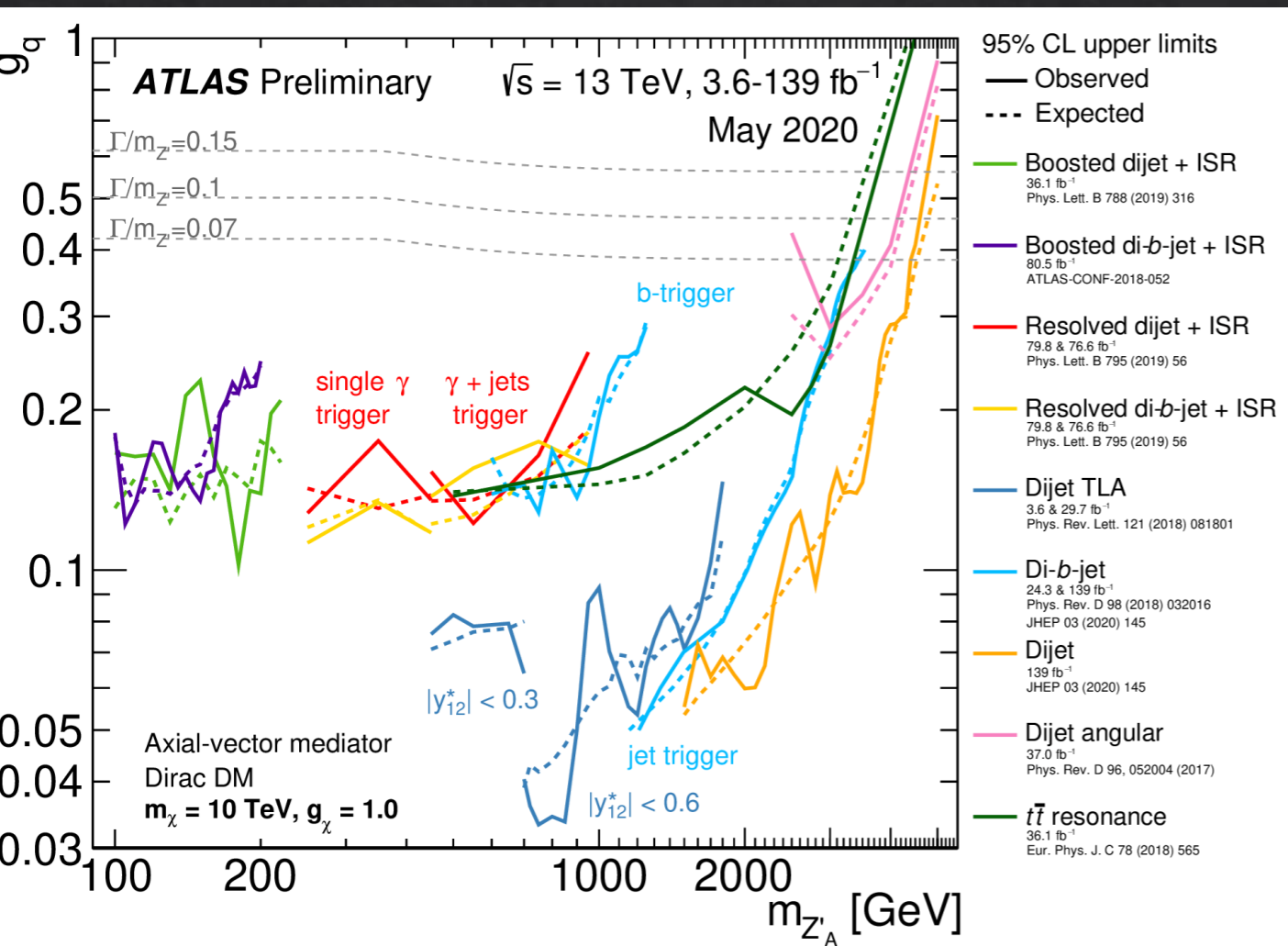


◆ Comparison to DD experiments:



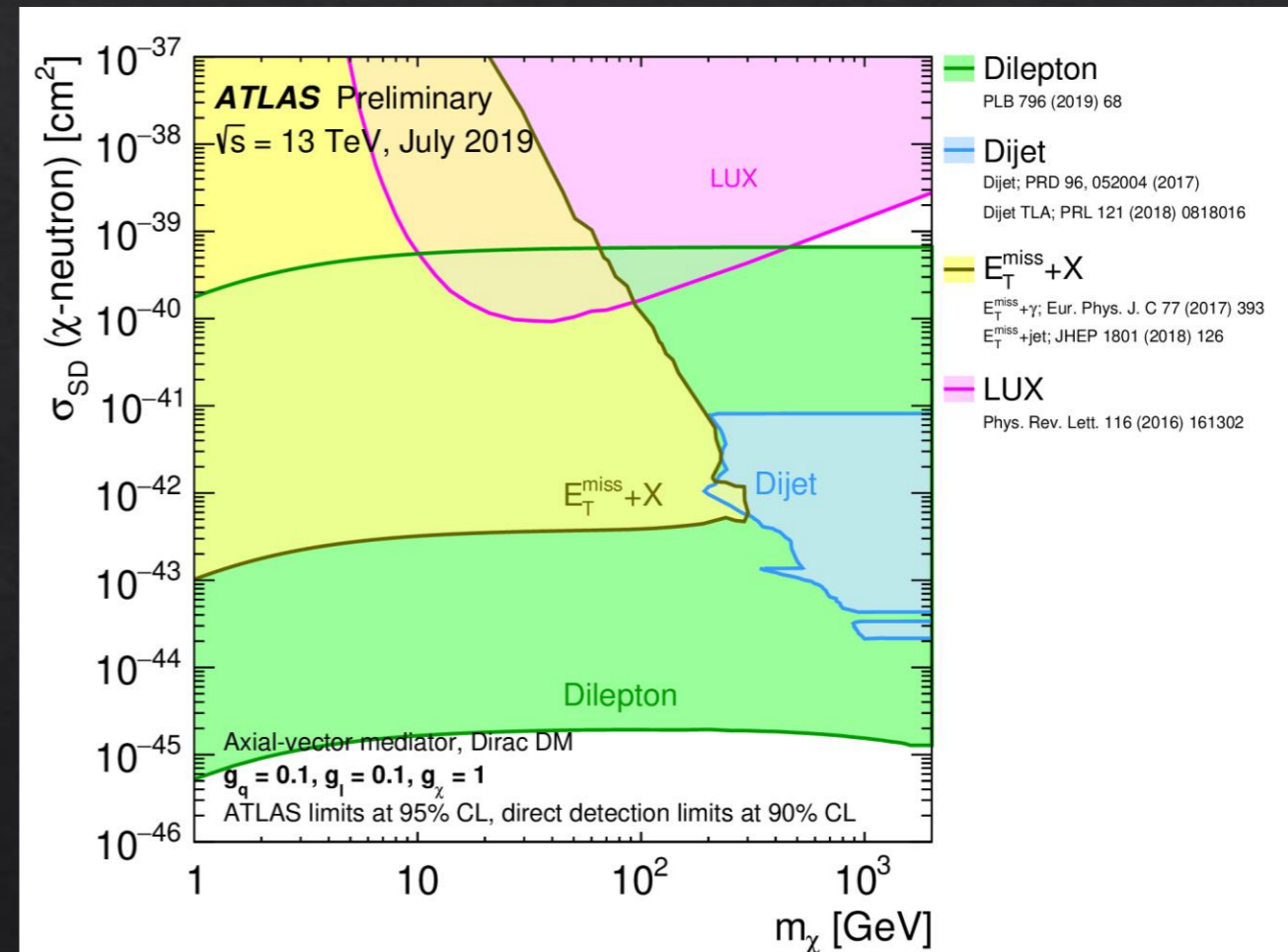
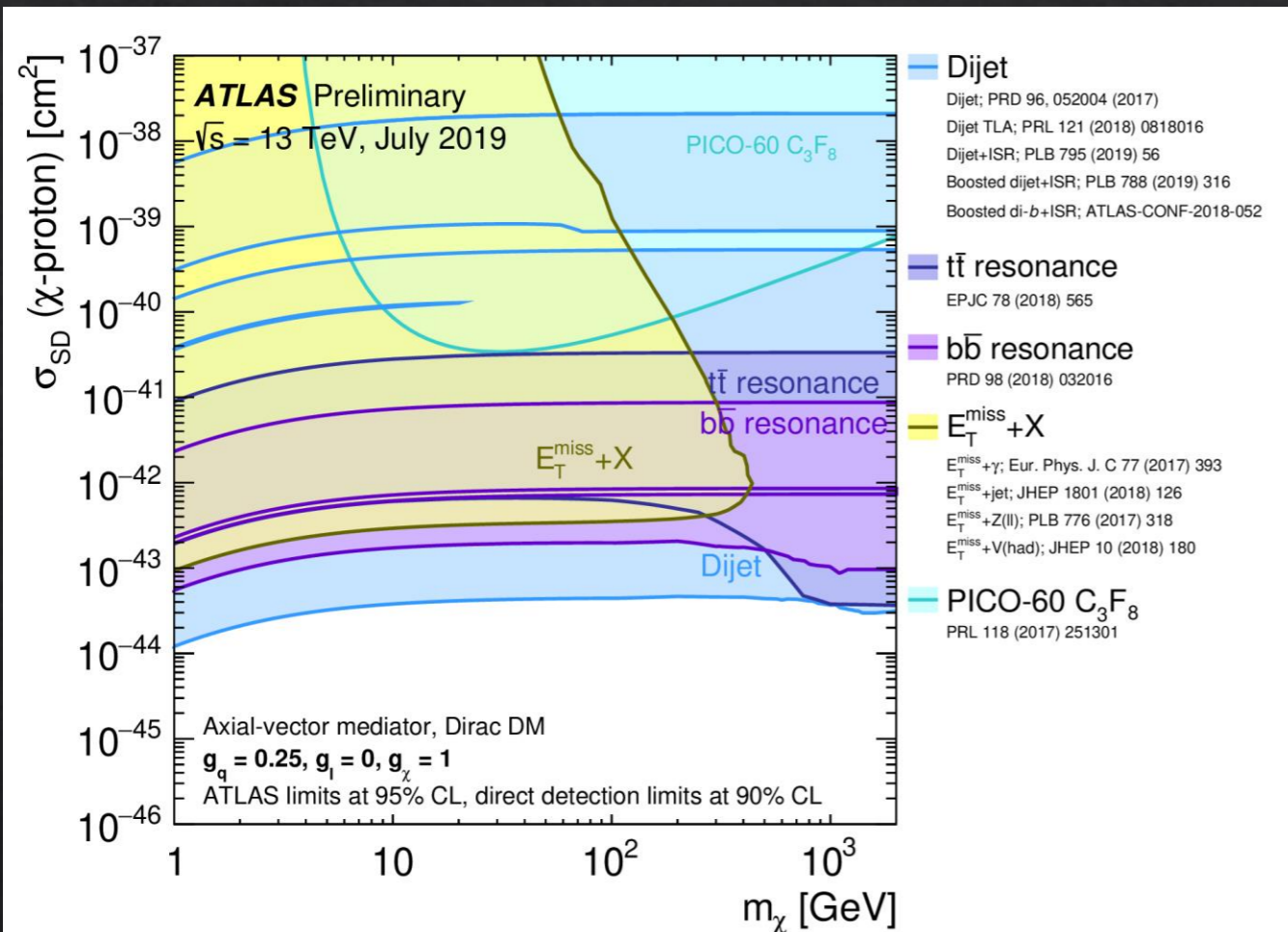
Combinations

◆ Extensive work to combine ATLAS results within Simplified Models interpretation



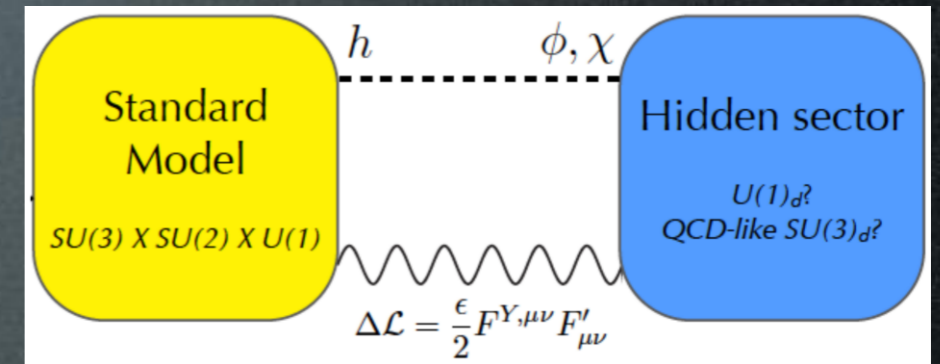
Combinations

- ◆ Extensive work to combine ATLAS results within Simplified Models interpretation and comparison to DD experiments:



What about the Dark Sector

- The “dark sector” consists of particles that do not couple to known SM fields, but interact through a mediator:
 - Dark photons (vector portal), dark scalars (Higgs portal), ALPs (axion), sterile neutrinos...
- Mediators can provide “portal” to DM candidates or be candidates themselves.
- Exotics decays of Higgs boson are predicted in many BSM models to explain g-2 discrepancy or positron excesses



But there physics laws or constraints that SM – Dark Sector interactions must respect

$$\mathcal{L} \supset -\frac{\epsilon}{2} B^{\mu\nu} A'_{\mu\nu} - H^\dagger H (AS + \lambda S^2) - Y_N^{ij} \bar{L}_i H N_j$$

Vector Portal:
 Min. Lagrangian =
 SM Lagrangian
 + Dark QED
 + “Kinetic Mixing (e)”

Vector portal
 [Okun; Galison & Manohar; Holdom; Foot et al]

Higgs portal
 [Patt & Wilczek]

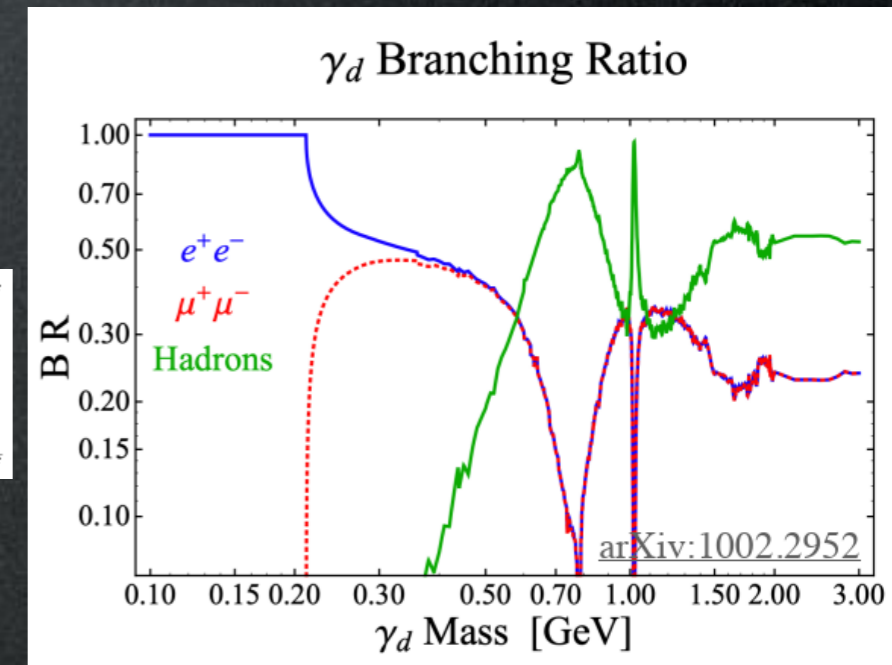
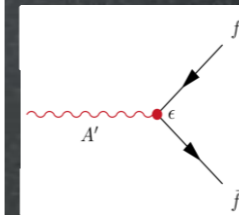
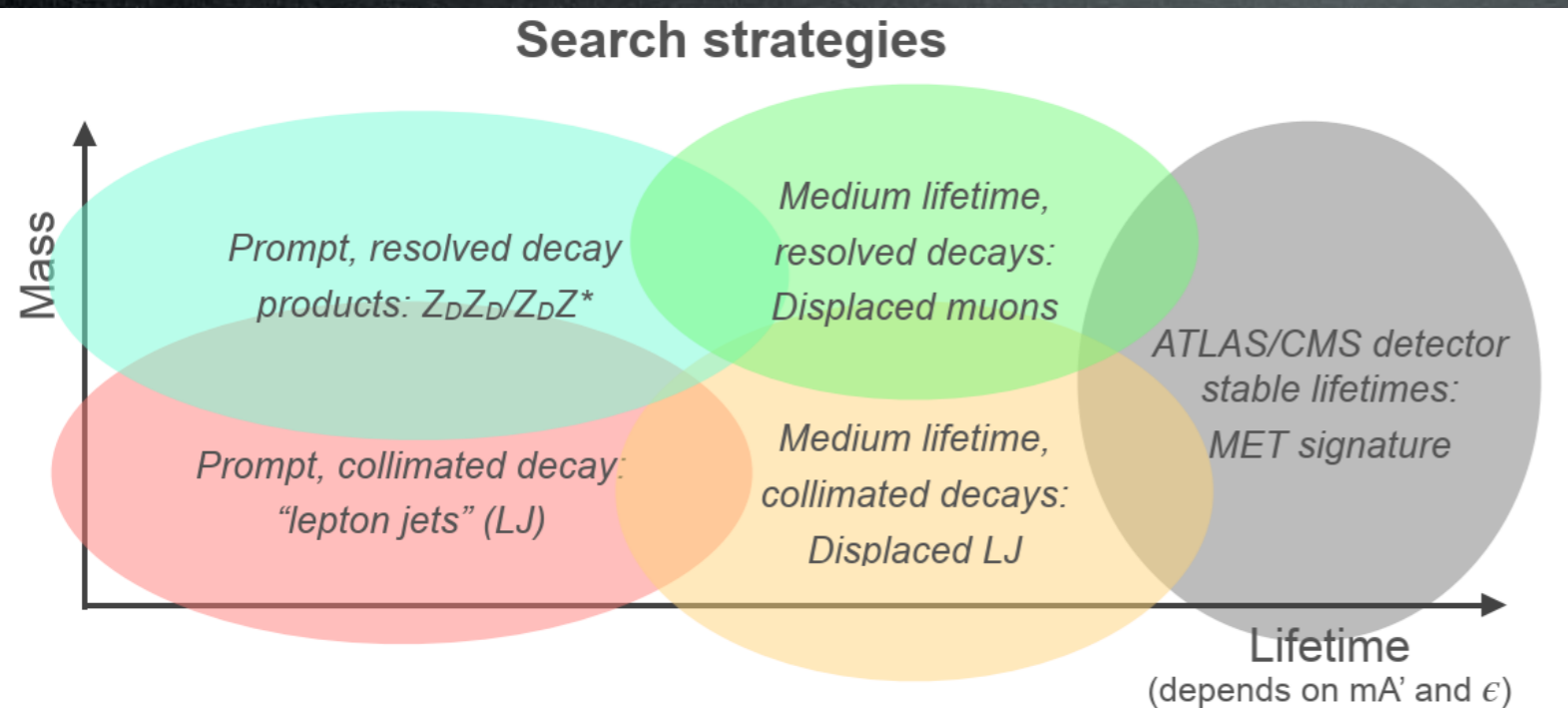
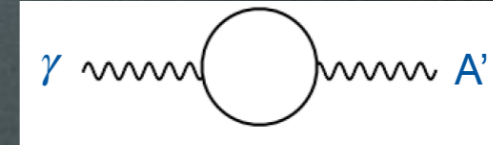
Neutrino portal

$$+ \frac{1}{f_a} \left(\text{tr}(G\tilde{G}) + c_F F\tilde{F} \right) a + \mathcal{O}(\text{dim} \geq 5)$$

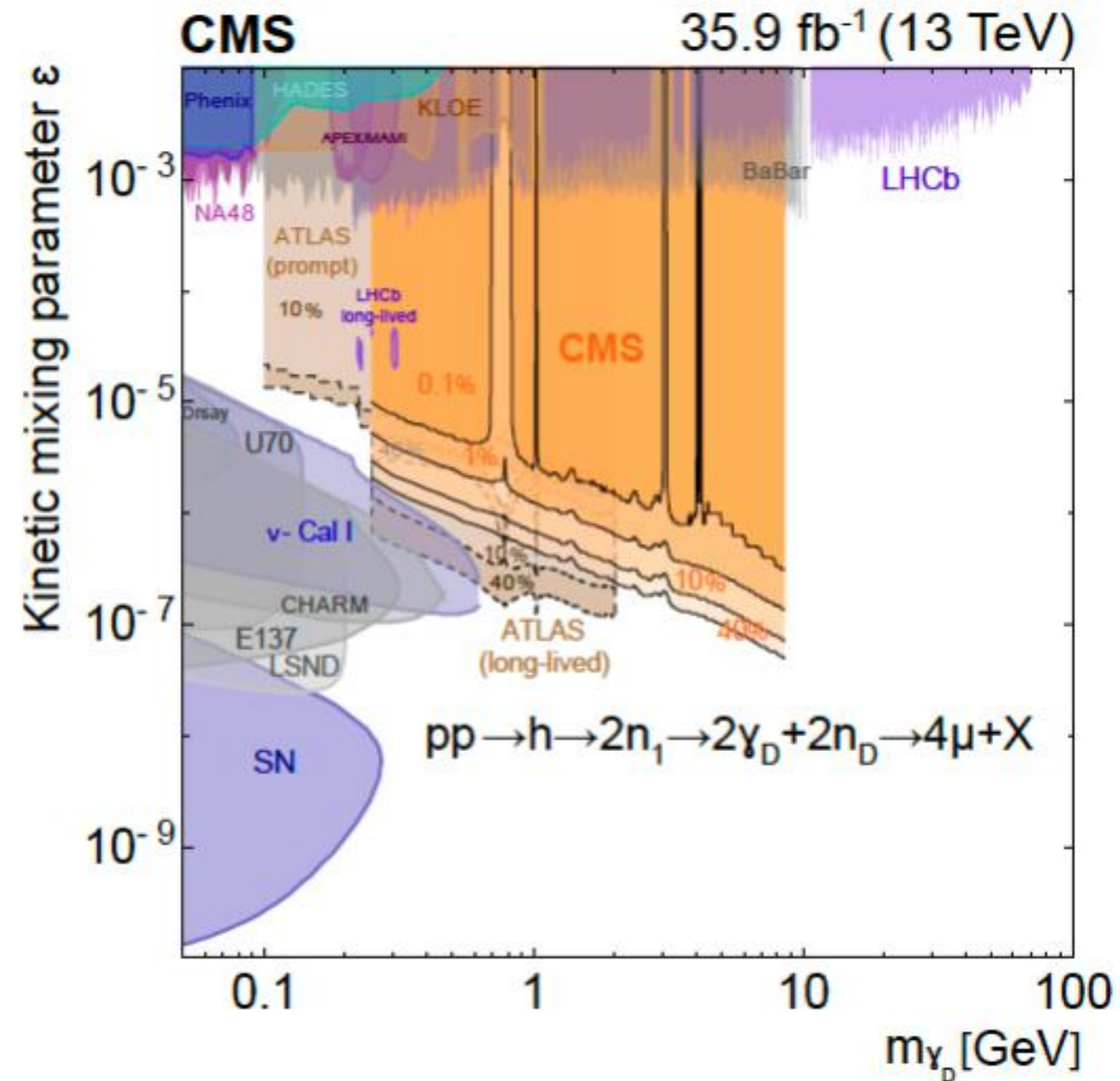
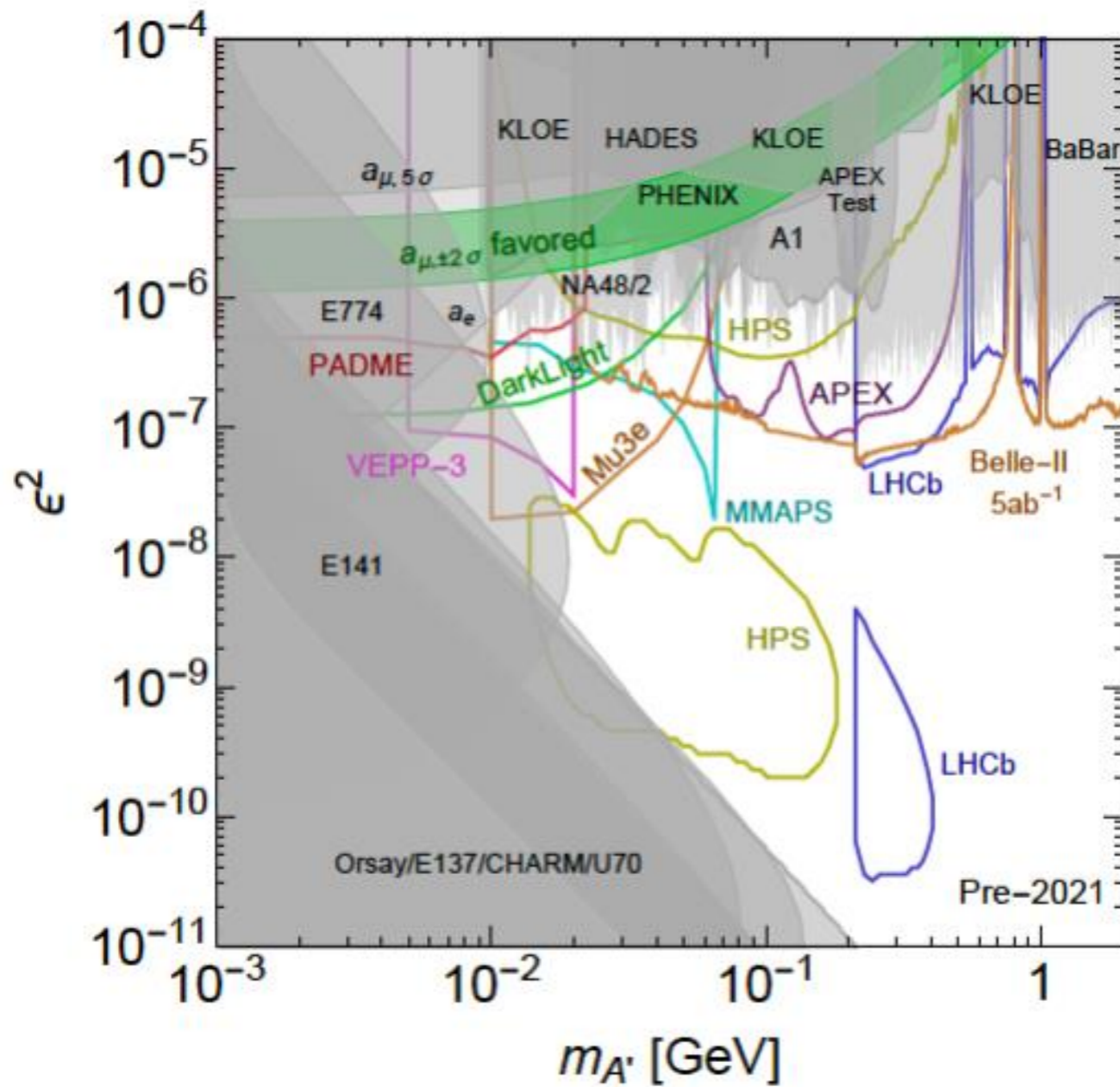
Axion portal
 [Weinberg, Wilczek, KSVZ, DFSZ]

Dark Photon

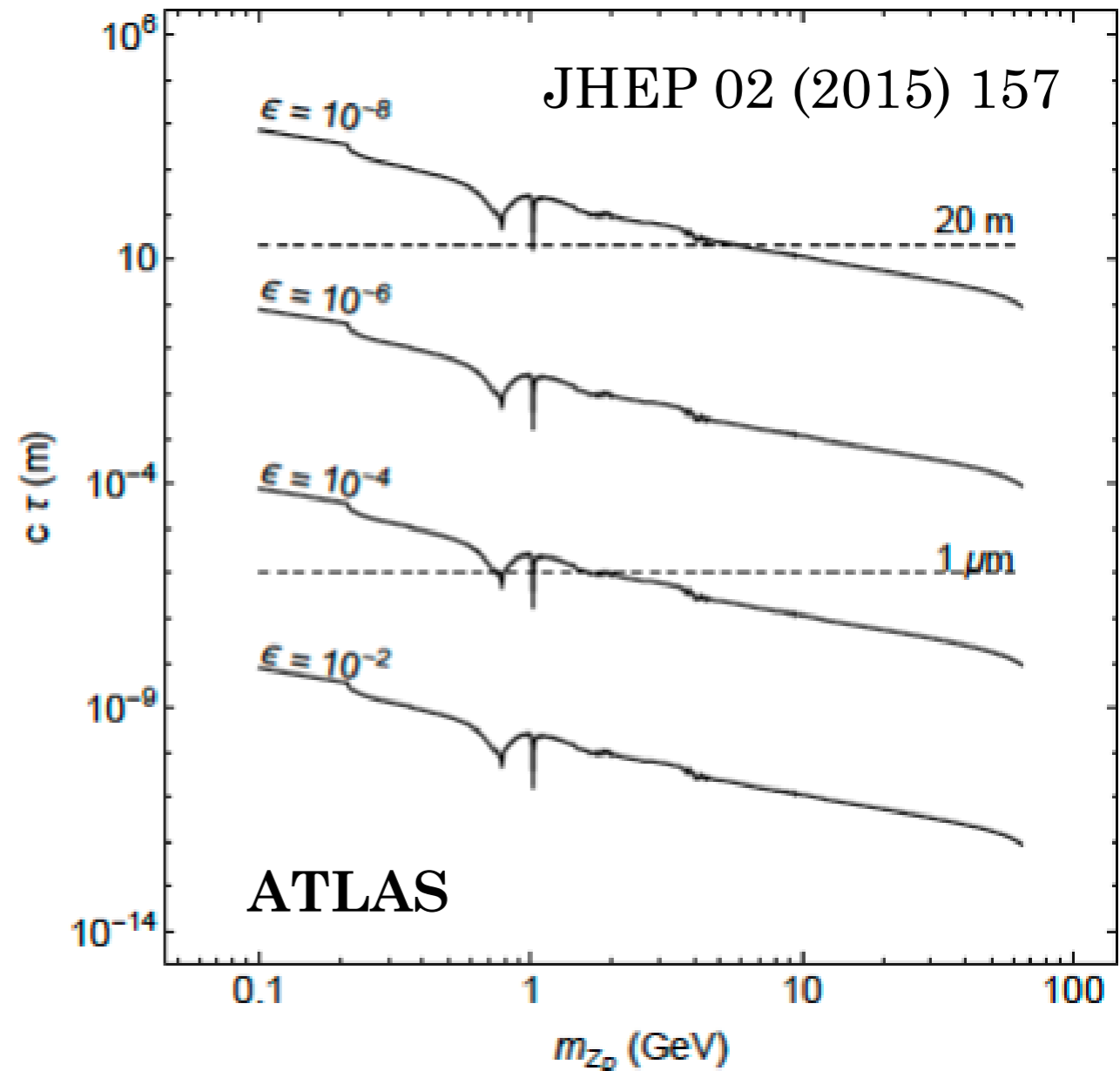
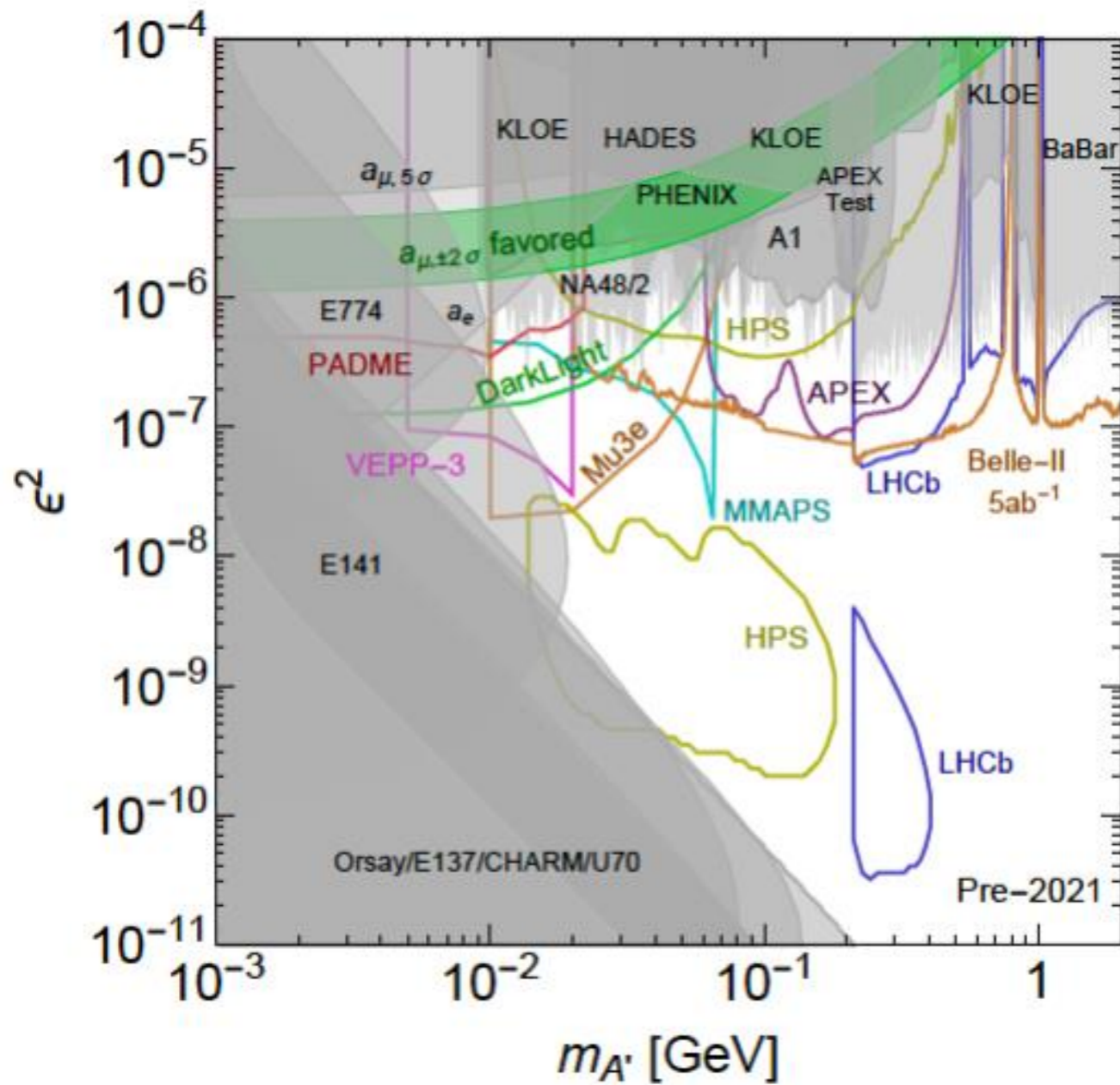
- Add a U(1)_D where massive dark gauge boson ($A'/Z_D/\gamma_D$) kinetically mix with SM photon
 - Parameters: kinetic mixing term, ϵ , and $m_{A'}$
- Dark photon decay can be:
 - Minimal: via same mixing as production
 - Generic: not ϵ -suppressed, dominant in cases where $m_\chi < m_{A'}$, with either χ stable & invisible, or χ decays to SM particles, $A' \rightarrow > 2$ charged particles.



Kinetic mixing in $A' \rightarrow ll$



Kinetic mixing in $A' \rightarrow ll$

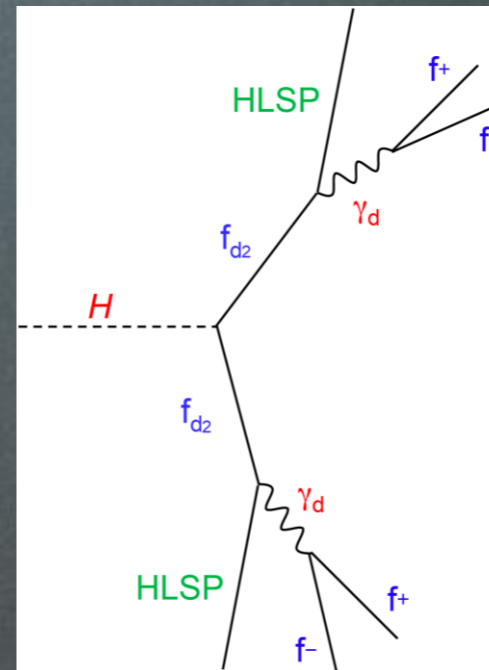
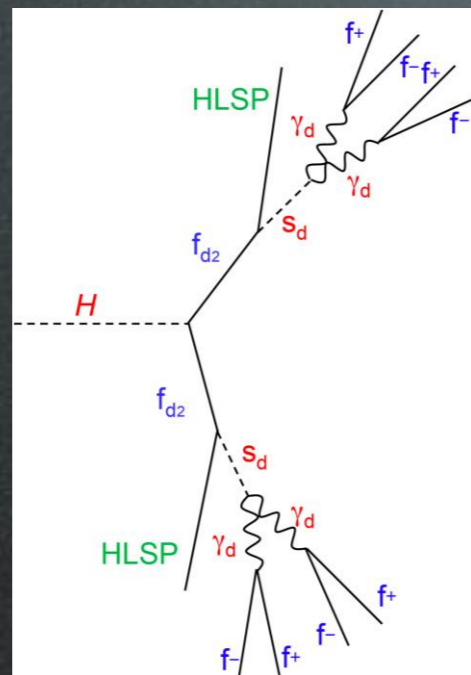


arXiv:1608.08632

Dark Photon may be prompt or long lived depending of ϵ and its mass

Lepton jets

- Search for collimated production of leptons: “lepton jets”.
- Dark fermions (f_{d2}) produced in H decays, which decay to γ_D (via dark scalar, s_d) and HLSP



HLSP = hidden lightest stable particle (fermion)

- Different experimental signatures depending on lifetime of dark photon.

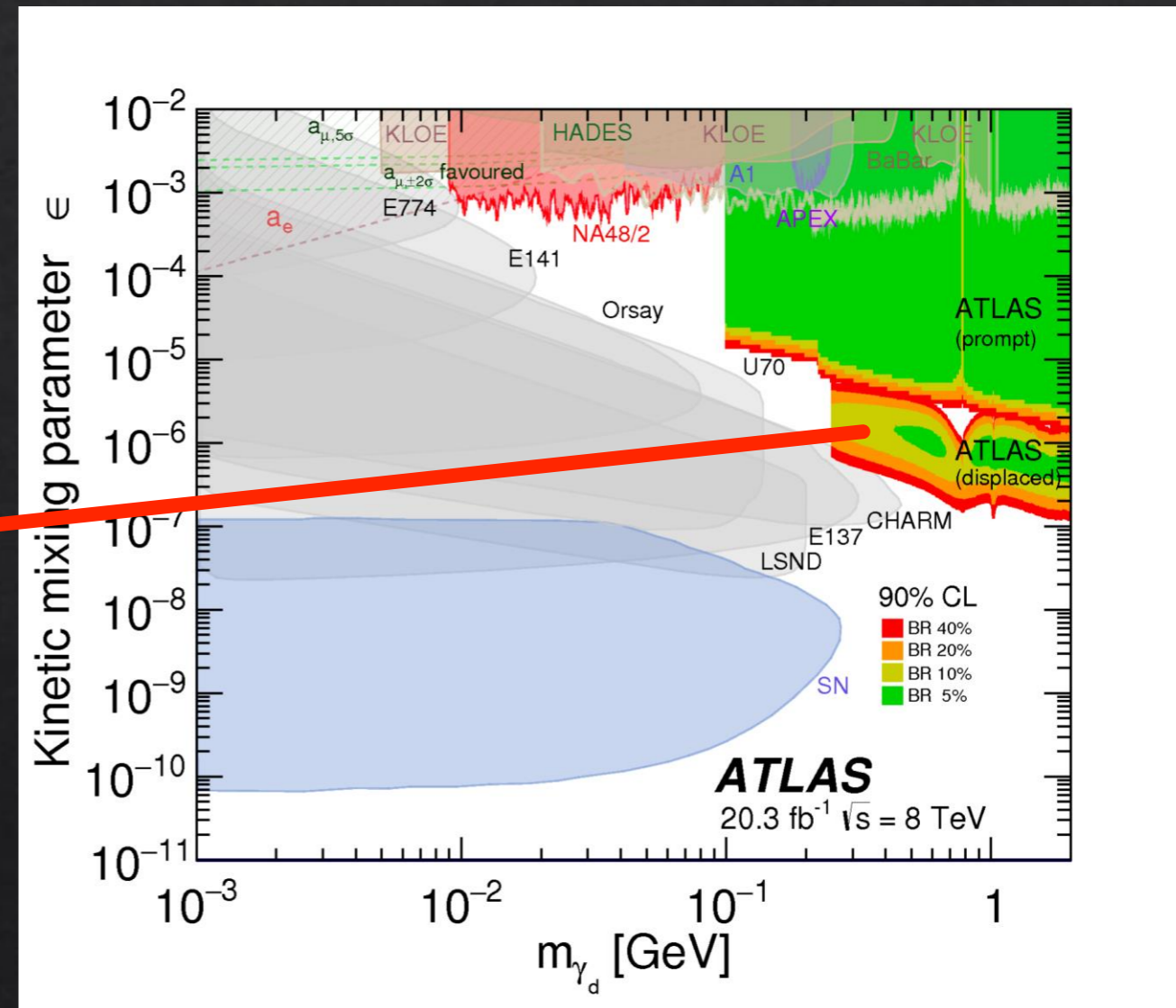
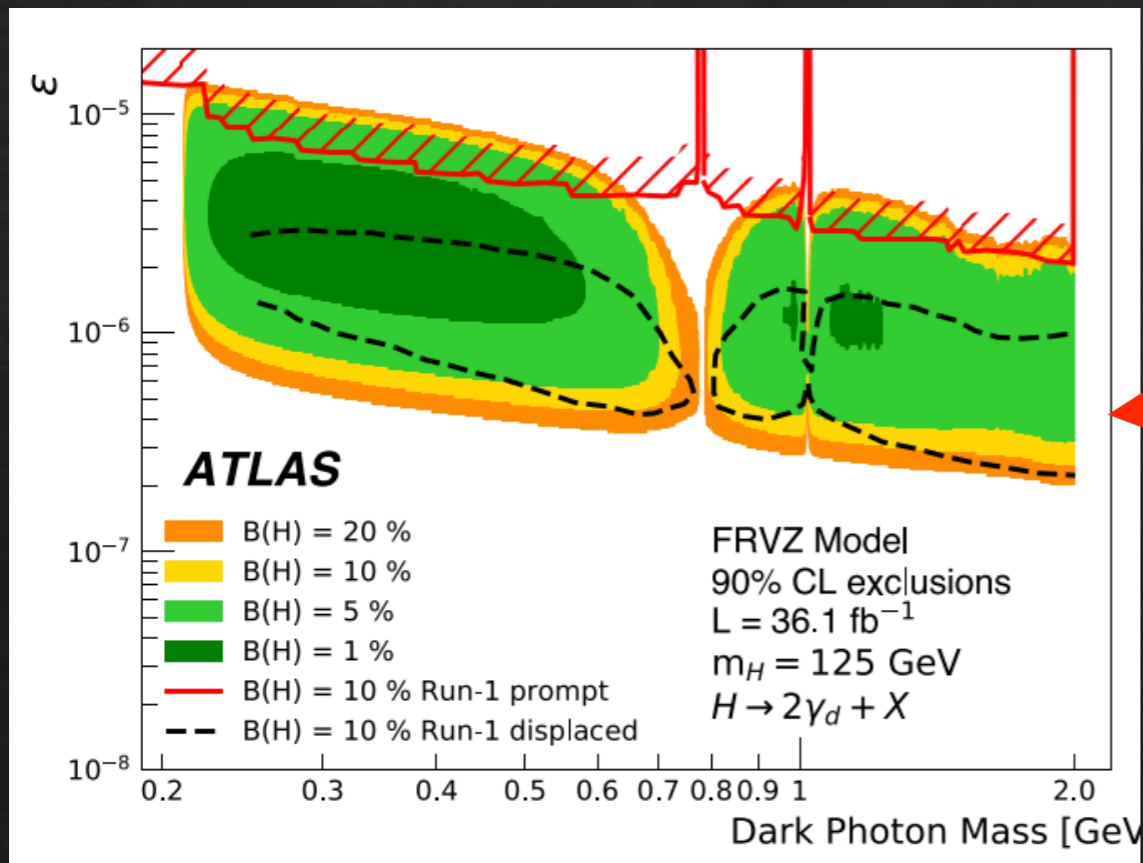
- Mean life time

$$\tau \propto \left(\frac{10^{-4}}{\epsilon} \right)^2 \left(\frac{100 \text{ MeV}}{m_{\gamma_d}} \right)$$

- Prompt and displaced lepton jet signatures

Lepton jets

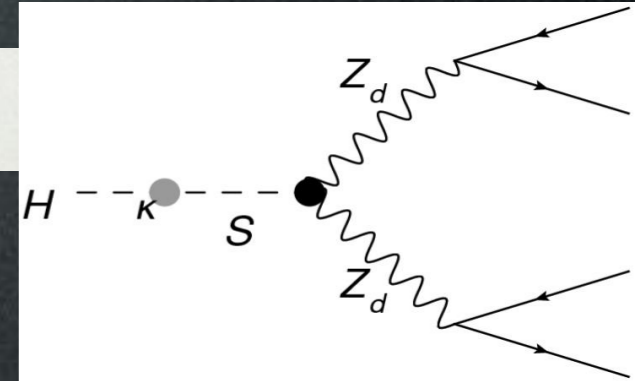
- Interpreted in terms of limits on the kinetic mixing parameter, ϵ , and m_A ,
- Limits are shown for $B(h \rightarrow 2\gamma_D + X)$ in range 1–20%.



Higgs Portal

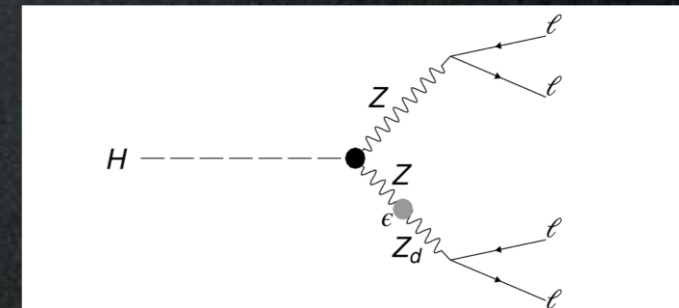
- There may be a dark Higgs boson; then there could also be a mixing between the SM Higgs boson (H) and the dark sector Higgs boson (S).

$$\epsilon_h |h|^2 |\phi|^2$$



- The mixing parameter κ between H and S, can be extracted from $H \rightarrow Z_d Z_d \rightarrow 4l$, a unique channel to access this parameter

- In addition to kinetic mixing, there could be also a mass mixing between SM Z and Z_d

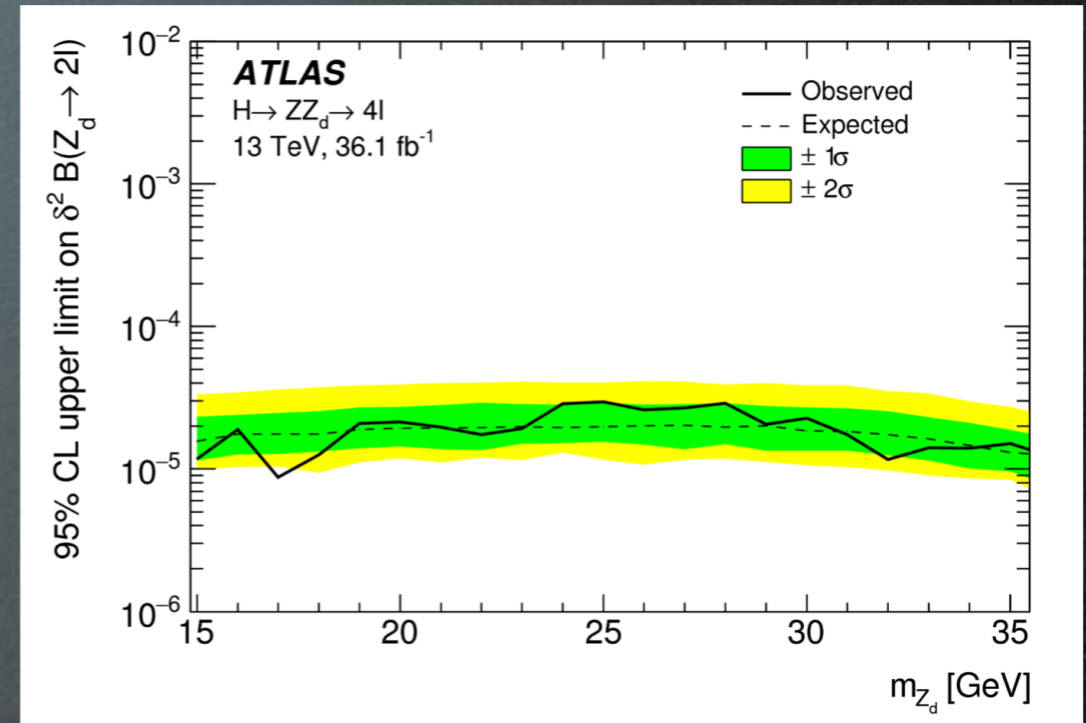
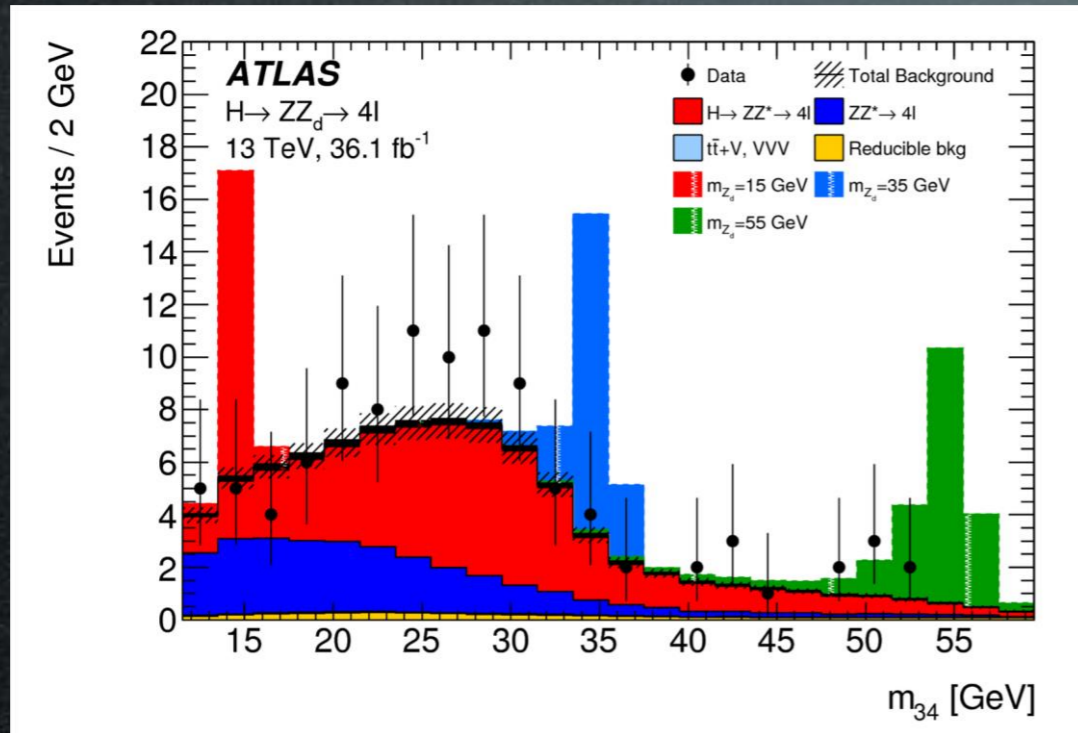


$$\epsilon_Z = \frac{m_{Z_d} \delta}{m_Z}$$

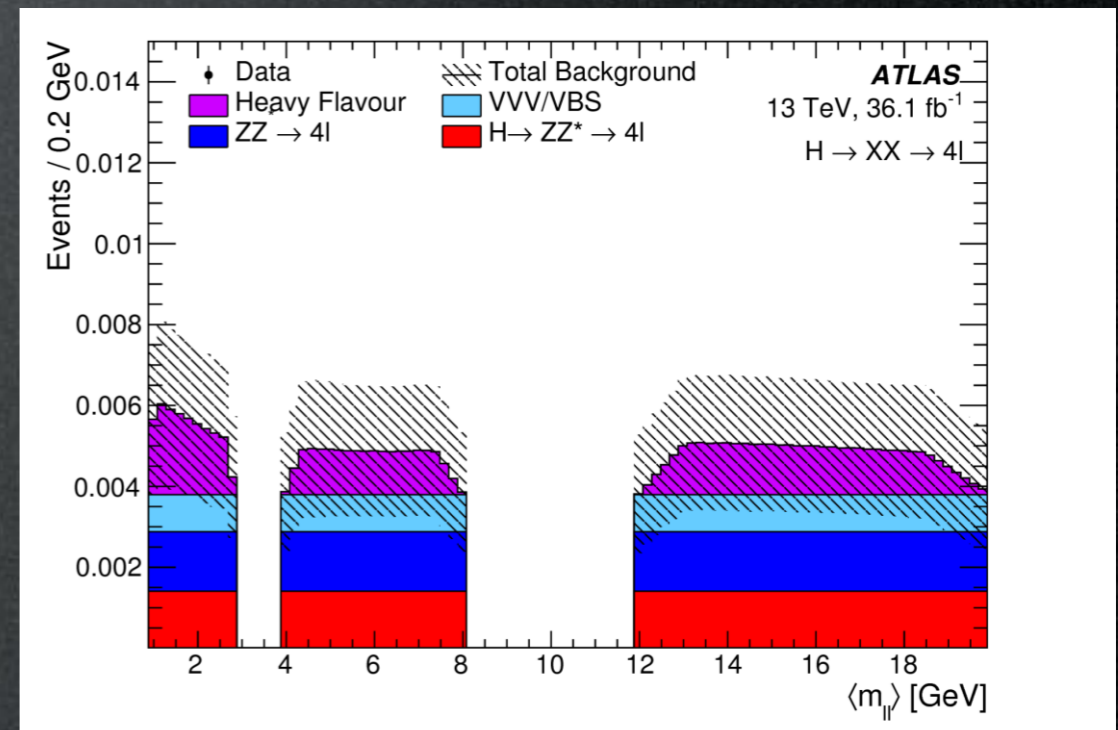
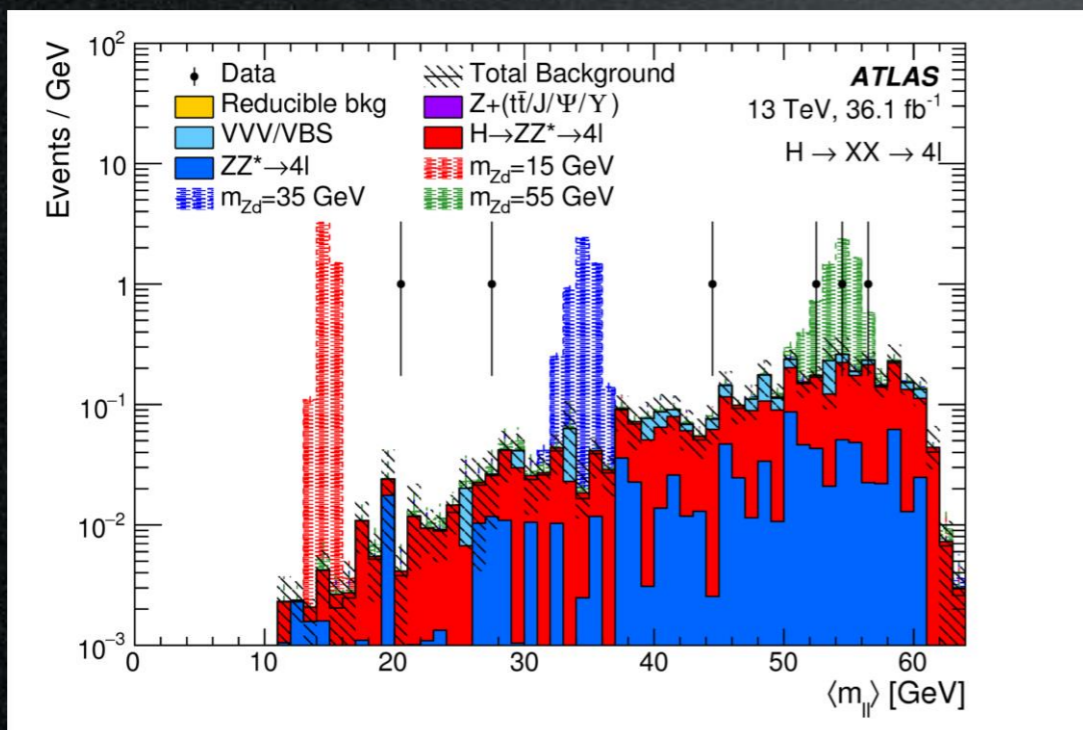
- $H \rightarrow Z Z_d \rightarrow 4l$ is sensitive to the kinetic mixing parameter ϵ , and to Z-Z_d mixing parameter δ . Unique channel to extract δ

Higgs Portal

- $H \rightarrow Z Z_d \rightarrow 4l$. No excess of events. Bound on ϵ (loose) and δ

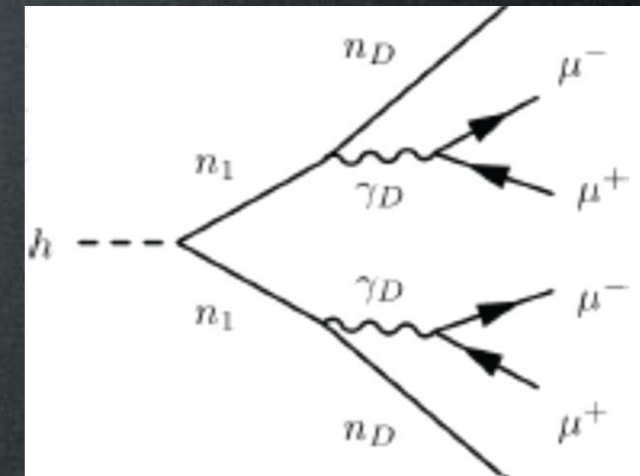


- $H \rightarrow Z_d Z_d \rightarrow 4l$. No excess of events. Searches for higher masses.



Higgs Portal

- Search for with dark SUSY benchmark model:
 - η_1 = lightest non-dark neutralino, and η_D = a dark neutralino that is undetected.

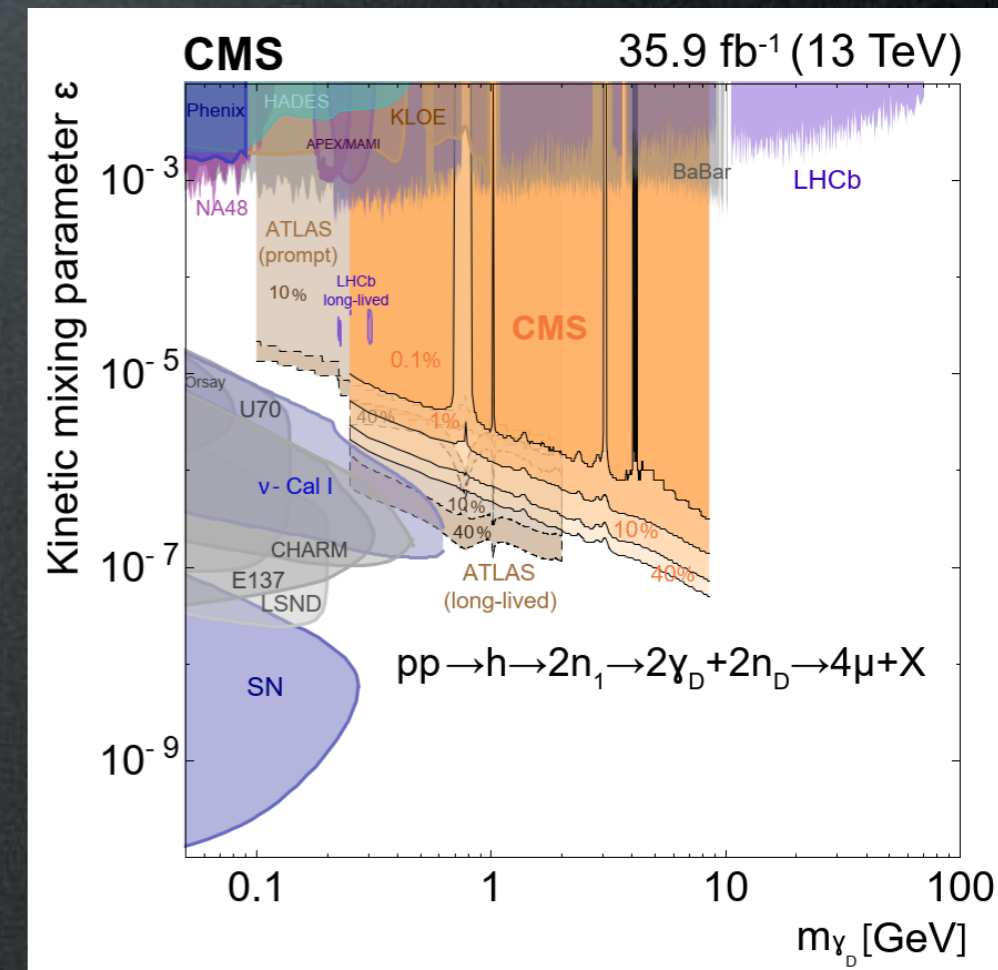


- Signal region defined by: $m(\mu\mu)_1 \sim m(\mu\mu)_2$
 $|m(\mu\mu)_1 - m(\mu\mu)_2| < 5\sigma(m_{\mu\mu})$

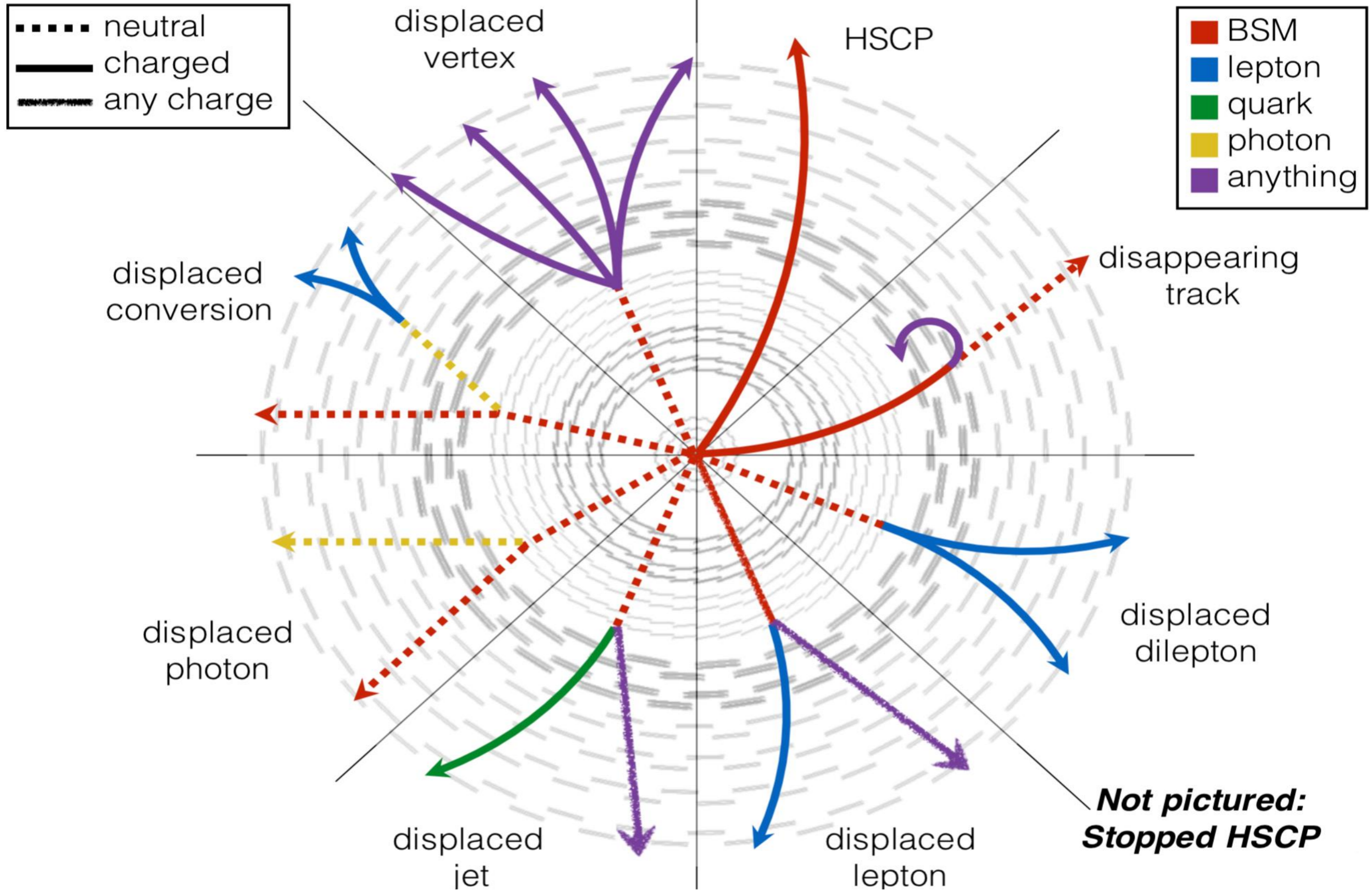
- Upper limits set on product of H production cross section and BR of Higgs boson (cascade) decay to a pair of dark photons:

$$\sigma(pp \rightarrow h \rightarrow 2\eta_1 \rightarrow 2\gamma_D + 2\eta_D) \times B(\gamma_D \rightarrow 2\mu)$$

- Limits on $B(h \rightarrow 2\gamma_D + X)$ in the range 0.1–40%

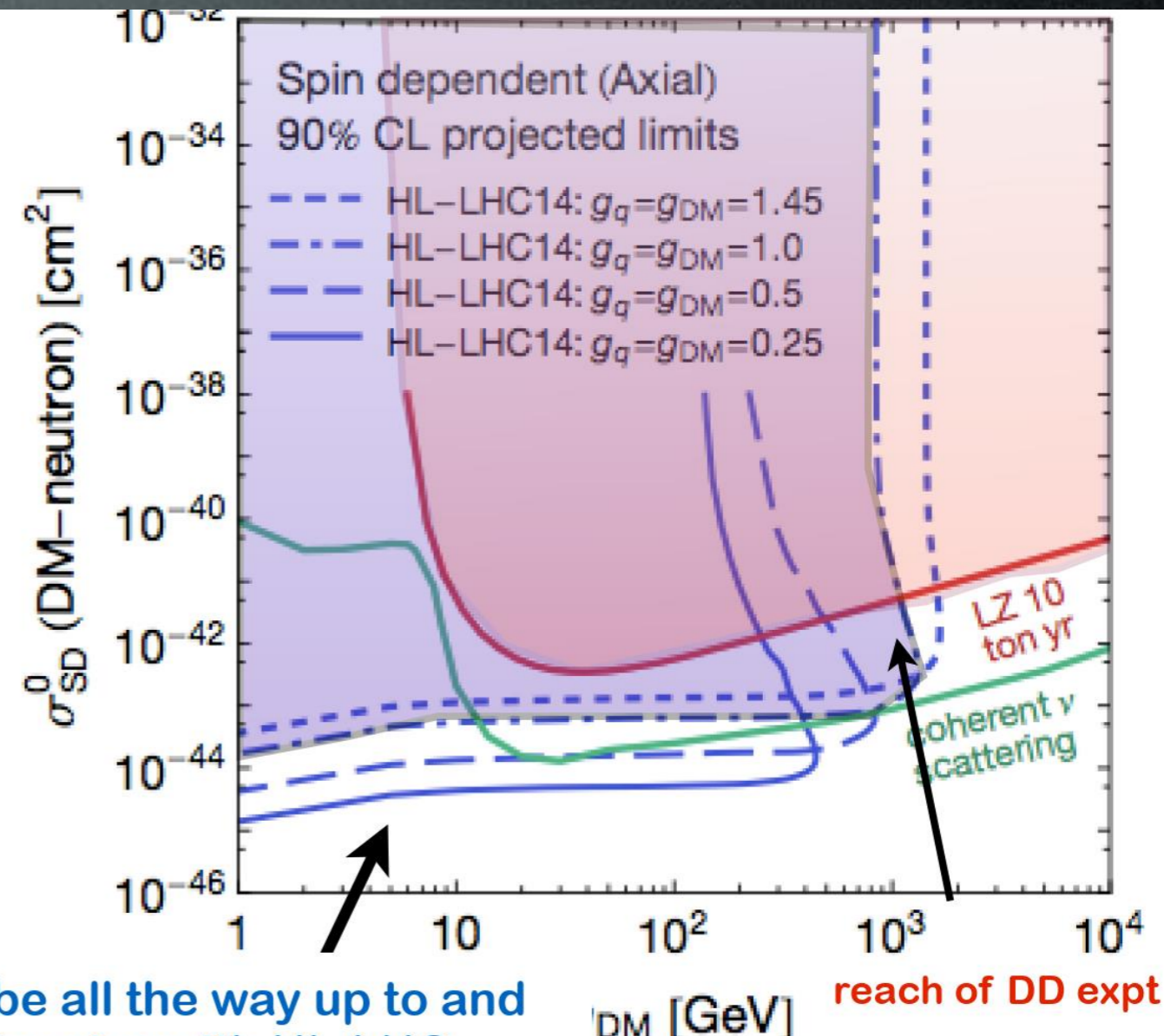
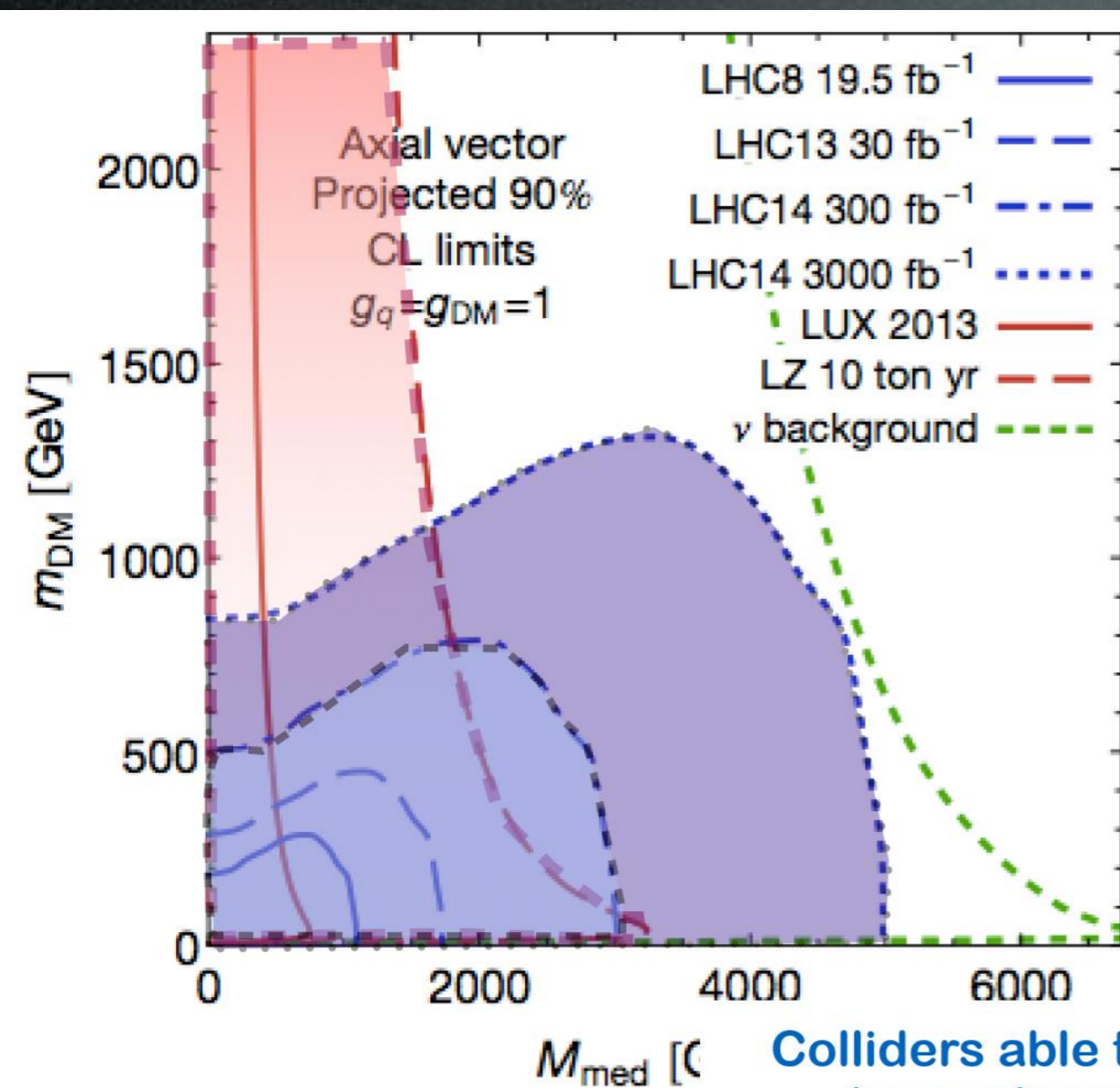


Beyond the WIMP



Future Prospects

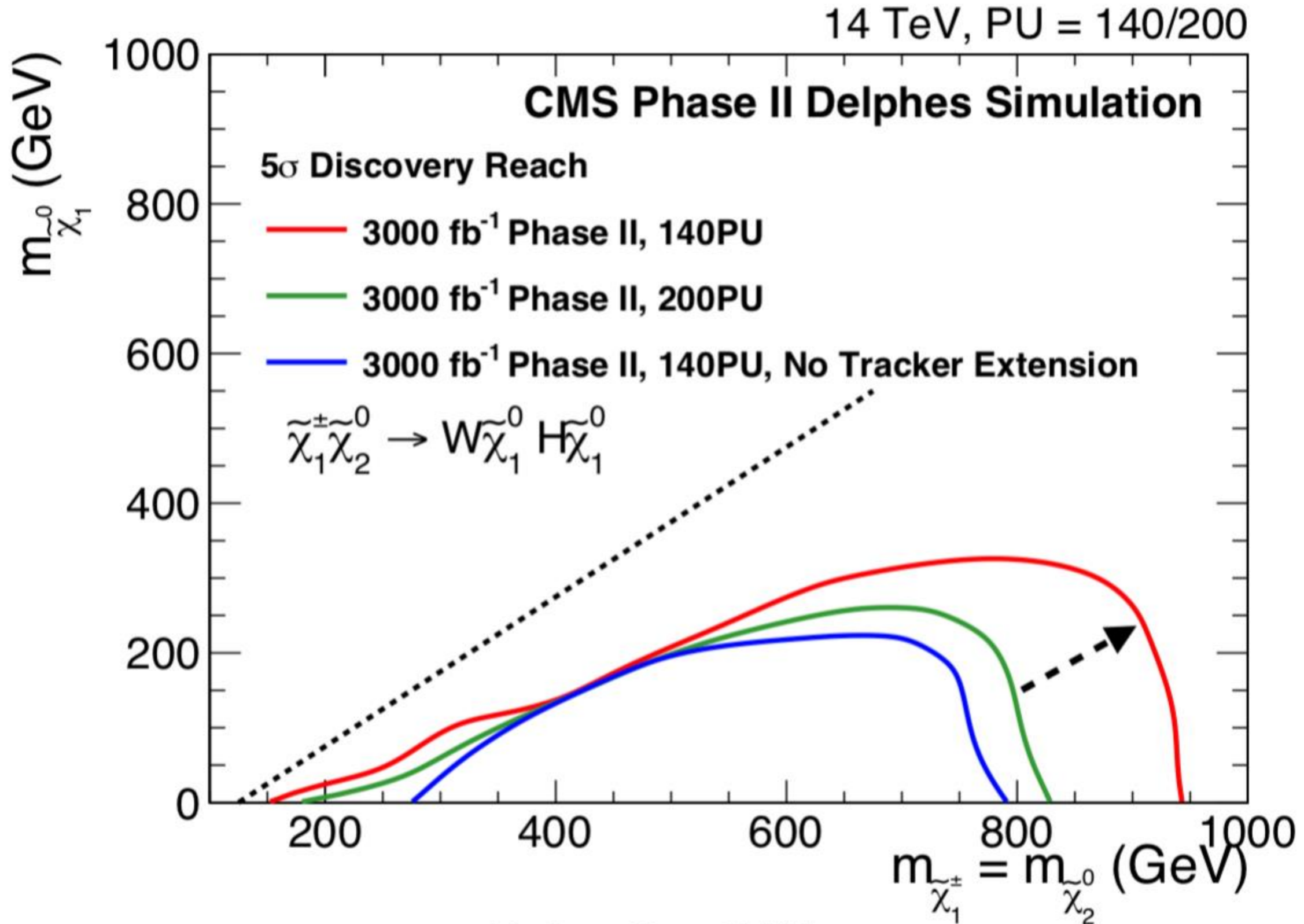
- Colliders can still do more with Run2 data. Reinterpretation, models..
- Future HL-LHC very challenging
 - High pile-up, Higher MET trigger threshold. But NEW and more performant detectors
 - Control of systematic uncertainties crucial
- Better interplay with DD experiments (Axion, Dark photons...)



Colliders able to probe all the way up to and beyond neutrino barrier with HL-LHC

Future Prospects

SUSY EW Searches, with conservative systematic uncertainties



Lindsey Gray, FNAL

Summary

- DM searches is the ideal inter-disciplinary field combining theoretical and experimental efforts from Cosmology, Astrophysics and Particle Physics.
- Interpretation (=benchmarks) important to understand the big picture
 - But the models we use inevitably influence motivations for searches
- Outlook for LHC DM searches:
 - expanding beyond WIMP simplified models
 - less simplified models (e.g. 2HDM), dark sectors
 - More complicated signatures (LLP), dark photon...
- Future LHC runs (Run3, HL-LHC) might open new era with higher reach in DM parameter space.
 - Significant improvement in detectors/Systematics to keep performance
- Only 1% of full LHC dataset analyzed so far
 - A long way to go to probe SM-DM interactions, whatever they are!