

The NIKA2 Sunyaev Zel'dovich Large Program: cluster mass estimation

Miren Muñoz Echeverría, LPSC Grenoble,
on behalf of the NIKA2 collaboration



Cosmology with galaxy clusters

Galaxy clusters: largest gravitationally bound objects in the universe

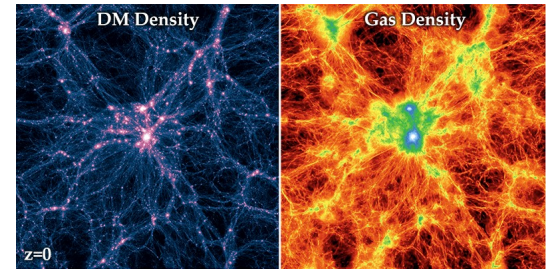
Created by gravitational collapse in the peaks of the matter density at the intersection of the filaments of the cosmic web

Mass $\sim 10^{13} - 10^{16} M_{\text{sun}}$

Temperature $\geq 10^7$ K (~ 1 keV)

Composition:

- 85% dark matter
- 12% hot ionised gas = the Intra-Cluster Medium (ICM)
- 3% galaxies



[Illustris Collaboration]

Clusters as cosmological probes

- tracers of the matter in the universe $f_{\text{gas}} \sim \frac{\Omega_{\text{b}}}{\Omega_{\text{m}}}$

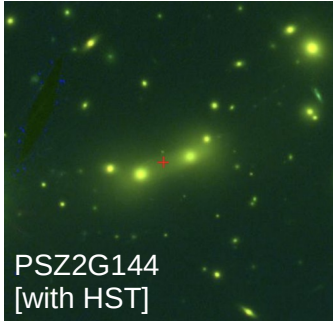
- cluster counts with respect to **mass and redshift distribution**

$$\frac{d^2 N}{dM dz} \sim \int d\Omega \frac{d^2 V}{dz d\Omega} \frac{dn}{dM}$$

Dependent of the evolution of the universe
and the structure formation rate

Cluster observables and mass estimates

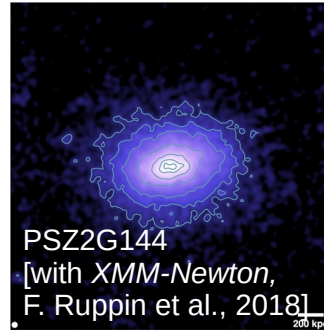
Optical/Infrared:
Stellar light of the galaxies



- Lensing of background galaxies
- Galaxy members richness and velocity dispersion

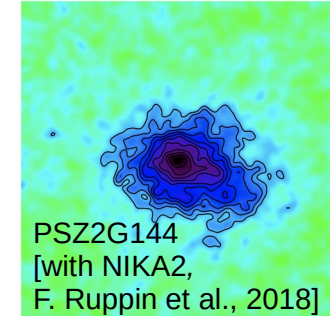
→ Different **mass estimates**

X-ray:
the hot ICM emission in X-ray



- Spectral line emission of metals
→ **Temperature**
- Bremsstrahlung of electrons
→ **Electron density**

Millimeter wavelengths:
Sunyaev-Zel'dovich (SZ) effect
from hot electrons in the ICM



→ **Thermal pressure of the ICM**

→ **Hydrostatic mass estimates** from X-ray and SZ

Different mass estimates affected by **different systematic** uncertainties. **Combining observables** may help building a consistent picture of the cluster physics to gain accuracy on the mass estimates

The Sunyaev-Zel'dovich effect

Thermal Sunyaev-Zel'dovich (**tSZ**) effect:

the inverse Compton scattering of CMB photons on the hot electrons of the ICM, which causes a spectral distortion of the CMB, independent of the redshift

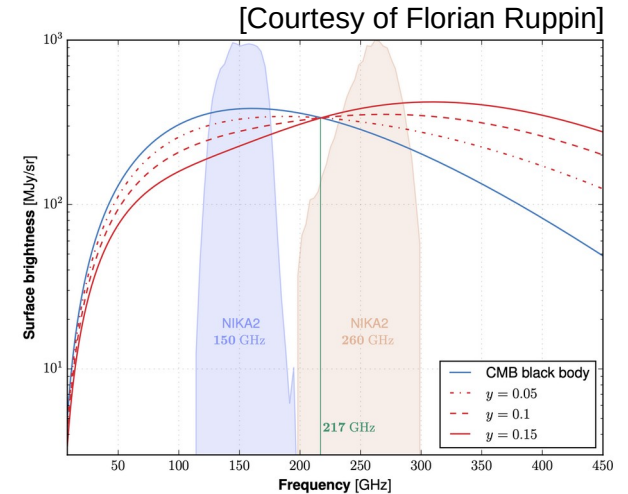
Relative variation in CMB intensity:

$$\left[\frac{\Delta I}{I_0} \right]_{\text{SZ}} = y \overset{\text{Spectral dependence of the SZ effect}}{f(\nu, T_e)}$$

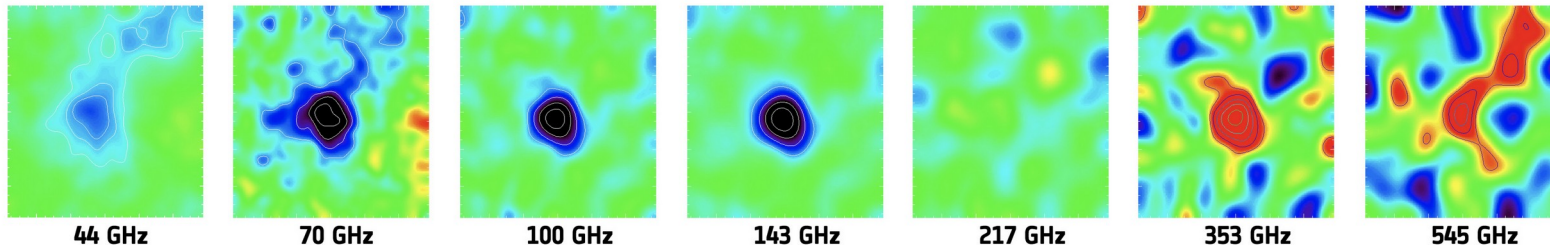
CMB specific intensity

$$y = \frac{\sigma_T}{m_e c^2} \int P_e dl$$

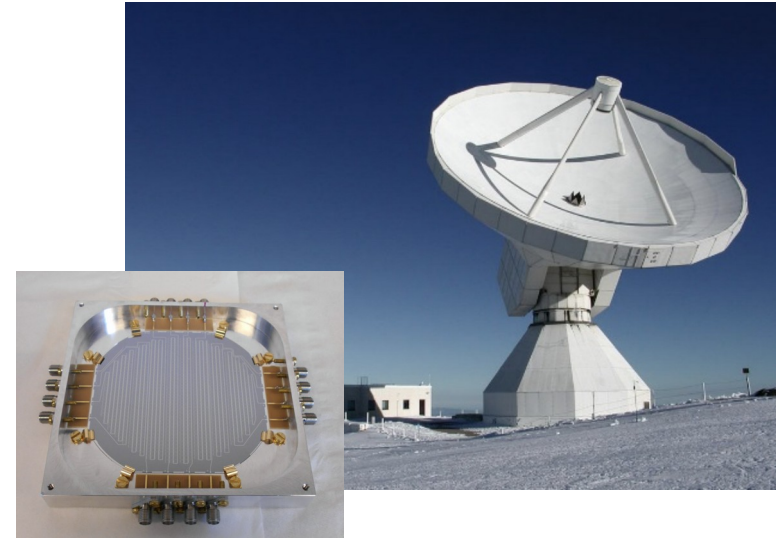
The amplitude of the tSZ effect or **Compton parameter: pressure** of the ICM integrated along the line-of-sight



Abel 2319 galaxy cluster [Planck]

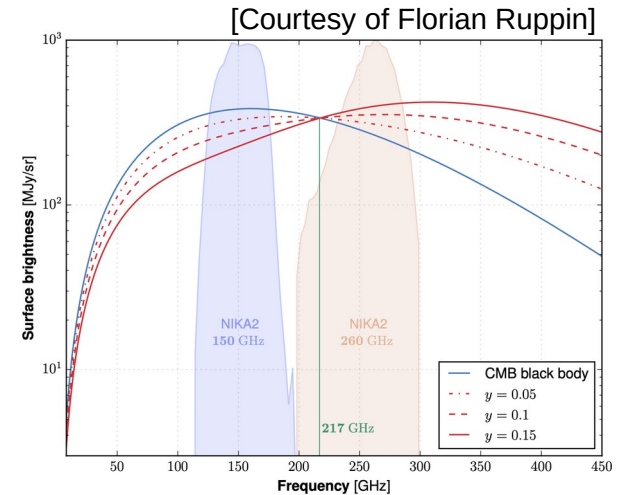


NIKA2 is a millimeter continuum camera of 2 900 Kinetic Inductance Detectors (KID) installed at the IRAM 30-meter telescope, and operating since 2017



Suited for observing the **SZ** effect:

- Operating at **150** and **260 GHz**: observe the distortion
- **High angular resolution** (18" and 12"): resolved maps of high redshift clusters
- 6.5' field of view: map the entire cluster
- High sensitivity: detect the weak signal of clusters



The NIKA2 Sunyaev-Zel'dovich Large Program (LPSZ)

Status before LPSZ:

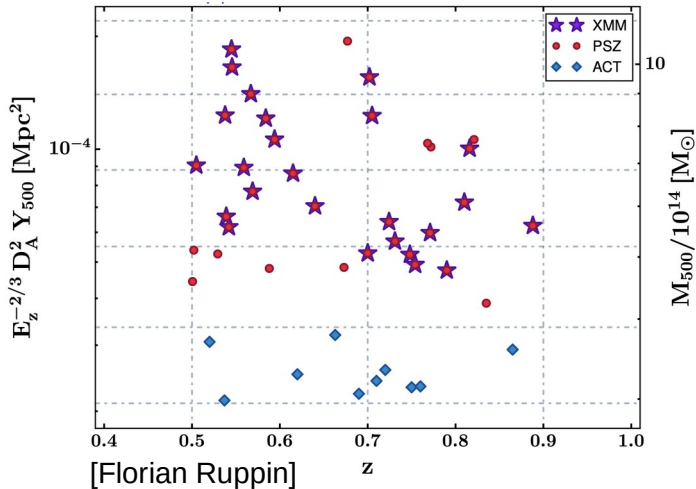
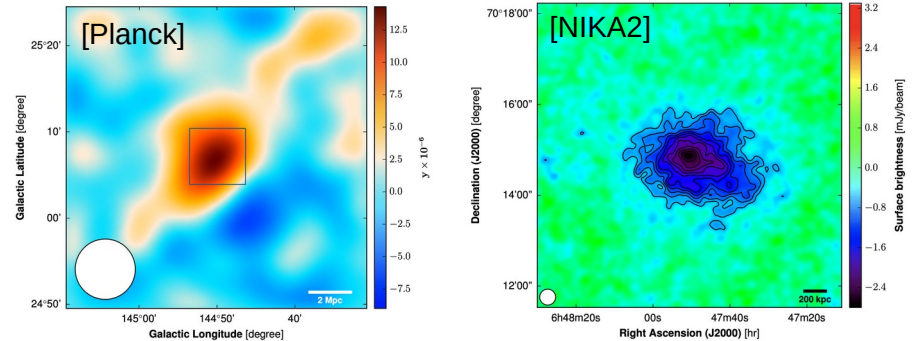
Catalogues of several thousands of clusters detected via the SZ effect with lower resolution instruments

→ need to resolve high redshift clusters to study their physical properties and masses

The **LPSZ**:

Study of 45 high redshift galaxy clusters selected in mass and redshift from Planck and ACT

PSZ2G144 [F. Ruppin et al., 2018]



Main goals:

- Thermodynamical **properties of the ICM**
- **Relation** between the **tSZ observable (Y_{500})** and the **mass (M_{500})**
- Probe the **low-mass and high-redshift clusters**

With:

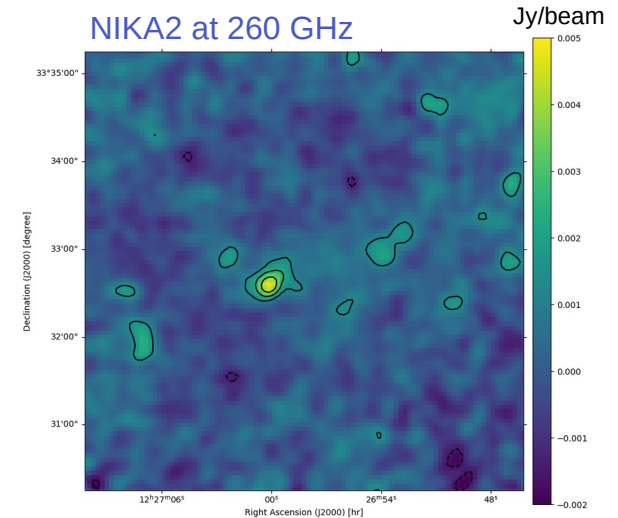
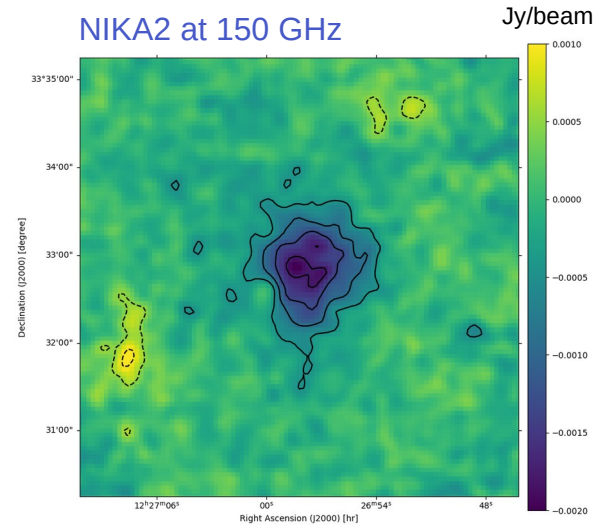
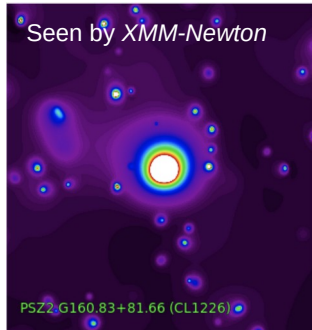
- tSZ measurements with **NIKA2** → $P_e(r)$
 - X-ray data from **XMM-Newton** → $n_e(r), T_e(r)$
- } $M_{HSE}(r)$

The galaxy cluster CL J1226.9+3332

Highest redshift cluster of the LPSZ: $z \sim 0.9$

Maps obtained with the NIKA2 collaboration's pipeline:

- Calibration
- Atmospheric and electronic contamination removal → Pipeline filtering
- Projection into sky map



5' maps, with a 10'' FWHM smooth. Contours are multiples of 3σ starting from 3σ
The SZ signal is contaminated by the positive signal of galaxies

Pressure profile of CL J1226.9+3332

We test the robustness of our estimate by:

- Fitting different pressure models

Radially binned

gNFW

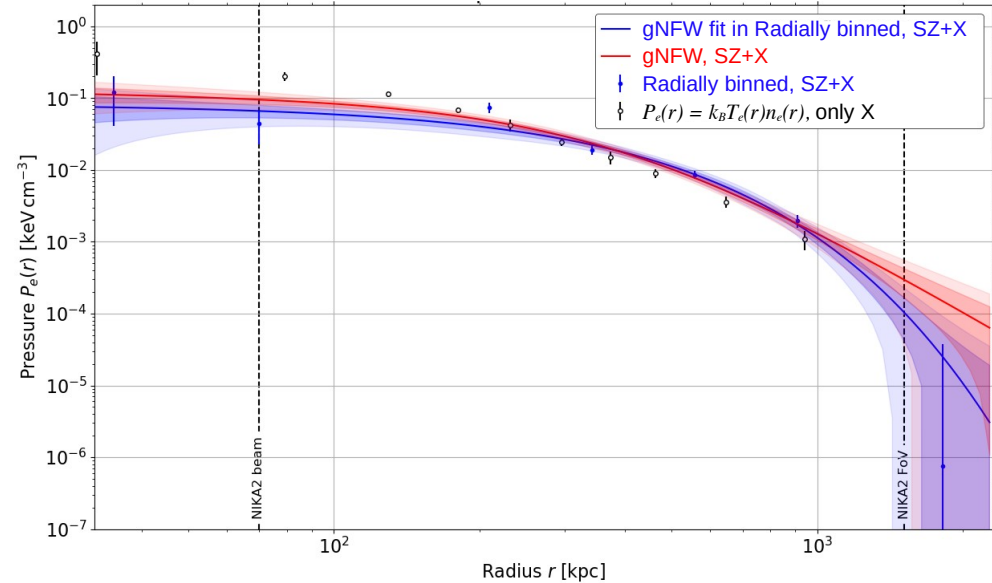
$$P_e(r) = P_i \left(\frac{r}{r_i} \right)^{-\alpha_i} \quad P_e(r) = \frac{P_0}{\left(\frac{r}{r_p} \right)^c \left(1 + \left(\frac{r}{r_p} \right)^a \right)^{\frac{b-c}{a}}}$$

- Defining the filtering induced by the data processing differently

2D: accounting for anisotropies in filtering

1D: assuming isotropic filtering

Reconstructed pressure profiles for CL J1226.9+3332



*Results assuming isotropic filtering (1D)

Hydrostatic mass and SZ signal

Hydrostatic equilibrium hypothesis for galaxy clusters: pressure compensates the gravity

The hydrostatic mass inside a sphere of radius r :

$$M_{\text{HSE}}(< r) = - \frac{r^2}{G \mu m_p n_e(r)} \frac{dP(r)}{dr}$$

Thermal pressure of the ICM from the SZ effect

Electron density from X-rays

We compare cluster masses and SZ signal inside a virialized sphere, in the literature very often at R_{500} :

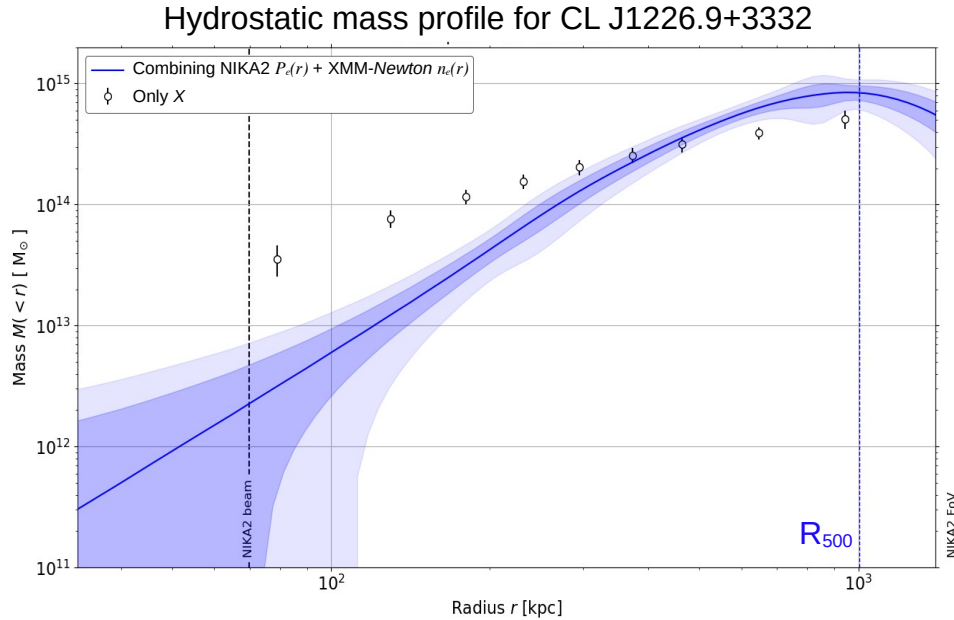
$$500 = \frac{\langle \rho(r < R_{500}) \rangle}{\rho_c(z)}$$

Hydrostatic mass and SZ signal at R_{500}

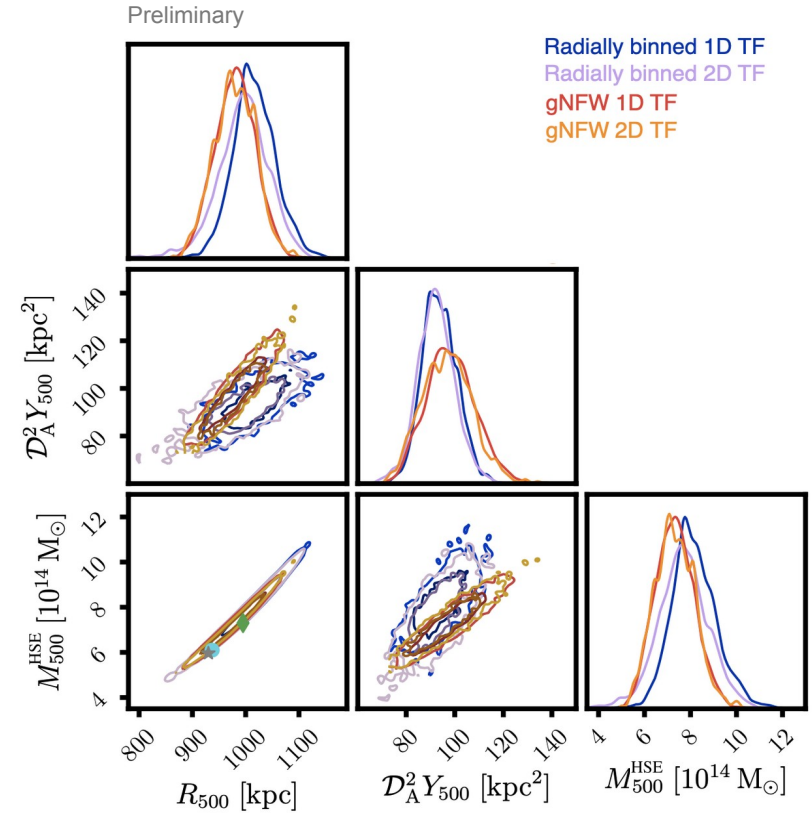
$$M_{500}^{\text{HSE}} = M^{\text{HSE}}(r < R_{500})$$

$$Y_{500} = \frac{\sigma_T}{m_e c^2 \mathcal{D}_A^2} \int_0^{R_{500}} P_e(r) 4\pi r^2 dr$$

Hydrostatic mass of CL J1226.9+3332



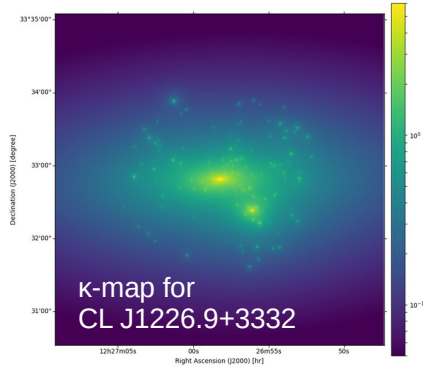
*Example from [Radially binned](#) pressure model assuming isotropic filtering (1D)



Tested systematic effects give robust hydrostatic mass estimates

Lensing mass of CL J1226.9+3332

The CLASH dataset provides lensing **convergence maps** that can be converted into **mass density maps**
 Convergence maps used here are the result of lensing raw data analysis described in Zitrin et al., 2015



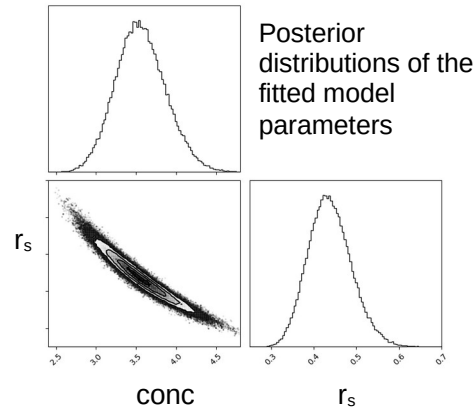
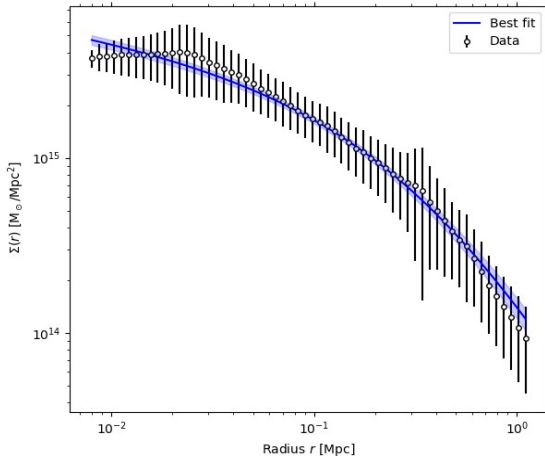
$$\kappa = \Sigma / \Sigma_{crit}$$

κ : convergence, Σ in critical density unities Σ_{crit}

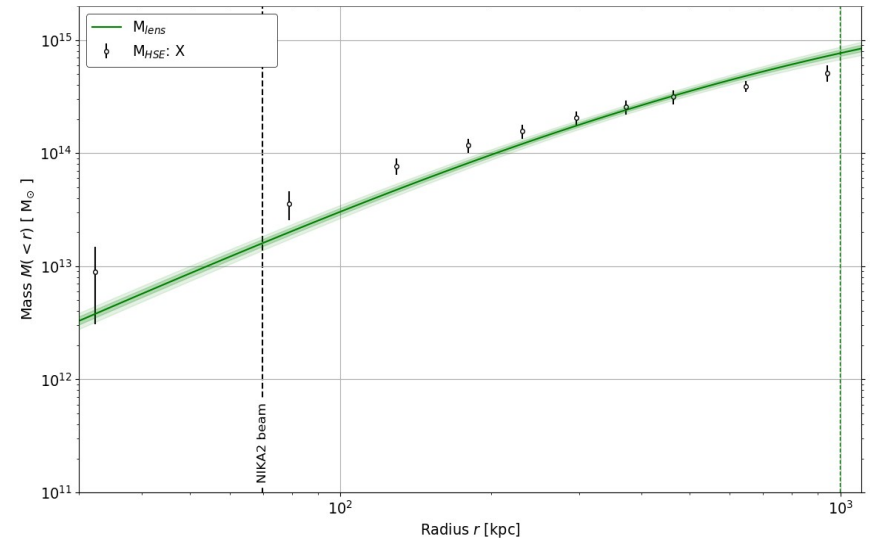
Σ : **surface mass density**

Σ_{crit} : the density needed for strong lensing to occur

Fit Σ to a projected density model using MCMC approach:



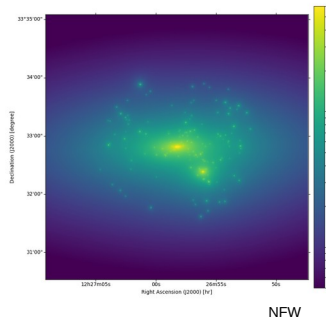
Lensing mass profile for CL J1226.9+3332



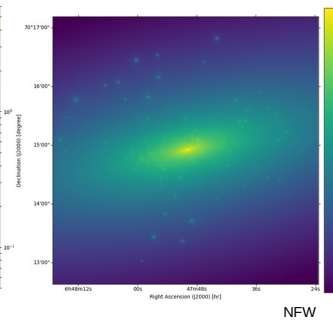
Hydrostatic and lensing mass of galaxy clusters

A combined NIKA-LPSZ-CLASH sample: convergence and SZ maps

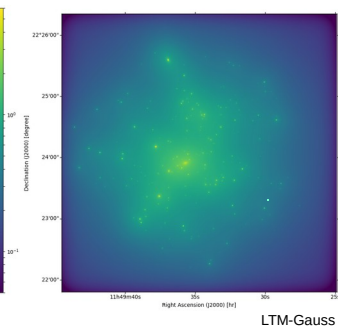
CL J1226.9+3332



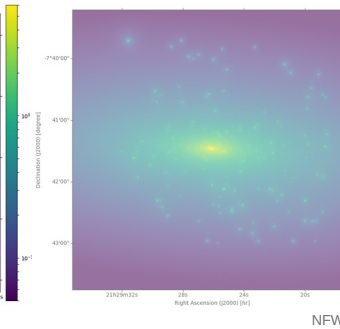
PSZ2G144



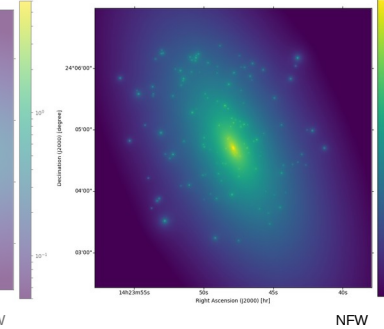
PSZ2G228



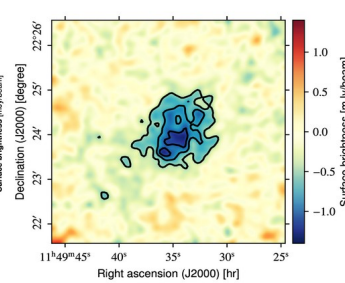
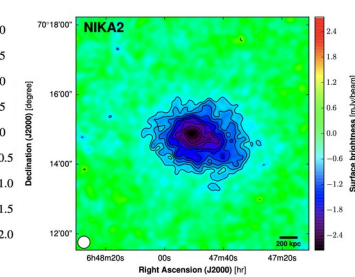
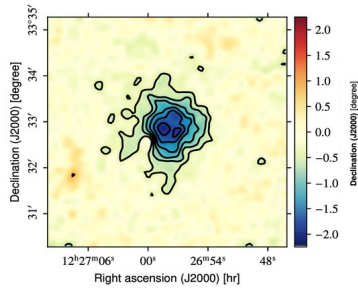
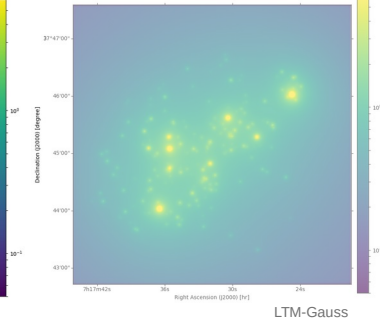
PSZ2G045



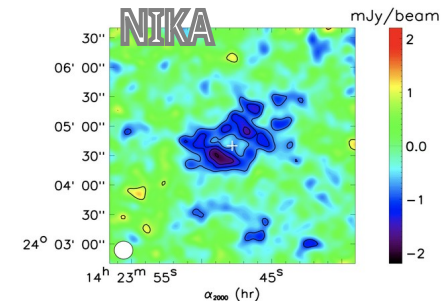
MACSJ1424



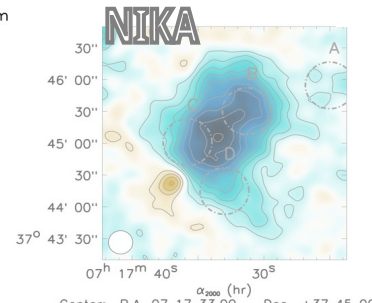
MACSJ0717



Ongoing analysis



High angular resolution Sunyaev-Zel'dovich observations of MACS J1423.8+2404 with NIKA: Multiwavelength analysis (Adam et al. 2018)



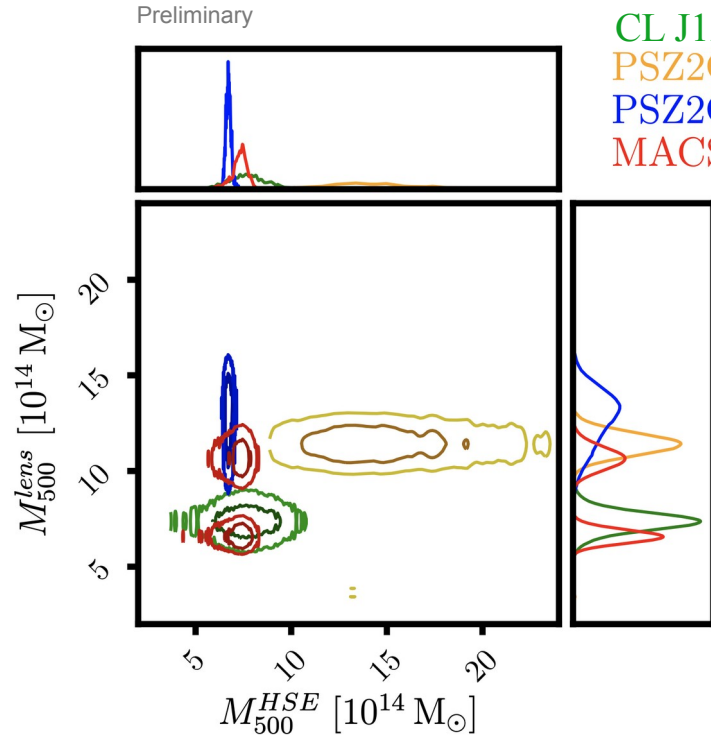
Mapping the kinetic Sunyaev-Zel'dovich effect toward MACS J0717.5+3745 with NIKA (Adam et al. 2016)

First Sunyaev-Zel'dovich mapping with the NIKA2 camera: Implication of cluster substructures for the pressure profile and mass estimate (Ruppin et al. 2018)

Hydrostatic and lensing mass of galaxy clusters

A combined NIKA-LPSZ-CLASH sample

For each cluster the marginalized M_{500} distributions for hydrostatic and lensing estimates



With this sample, we don't get a linear $M_{500}^{HSE} - M_{500}^{lens}$ relation

- PSZ2G228 has a very complex morphology to be fitted with spherical models
- PSZ2G144 has a very good M_{500}^{HSE} estimate compared to M_{500}^{lens}

Clusters properties affect differently the mass estimates:
essential to compare the lensing and hydrostatic mass estimates

Conclusions

The NIKA2 Sunyaev-Zel'dovich Large Program is well advanced:

- 30/45 clusters observed with high quality data
- Automatic analysis tools have been developed

This allows us to:

- Have good reconstructions of hydrostatic mass profiles combining NIKA2 and XMM-*Newton* data
- Characterize systematic effects on mass estimates

Combining LPSZ and CLASH data we can:

- Get the first hydrostatic to lensing mass bias in a combined sample at this redshift range ($0.5 < z < 0.9$)
- Study the mass dependence with clusters physical properties

Backup slides

Halo mass function

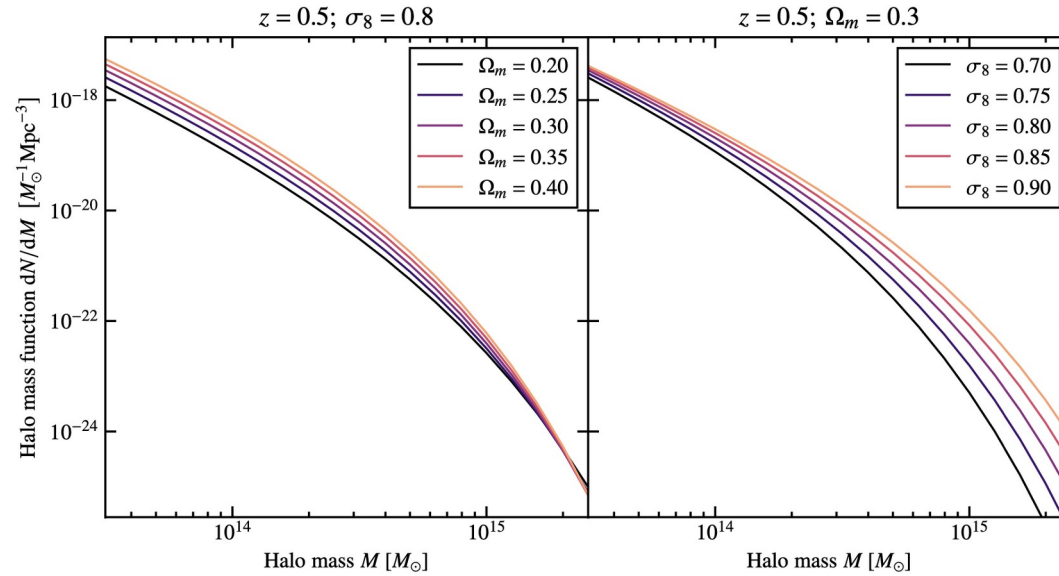


Figure 1.3 : Illustration de l'évolution de la fonction de masse avec les paramètres cosmologiques. La fonction de masse de Tinker *et al.* [32] est représentée pour différentes valeurs de Ω_m (gauche), et de σ_8 (droite). Le redshift est fixé à $z = 0.5$.

[Courtesy of Florian Kéruzoré]

Hydrostatic bias

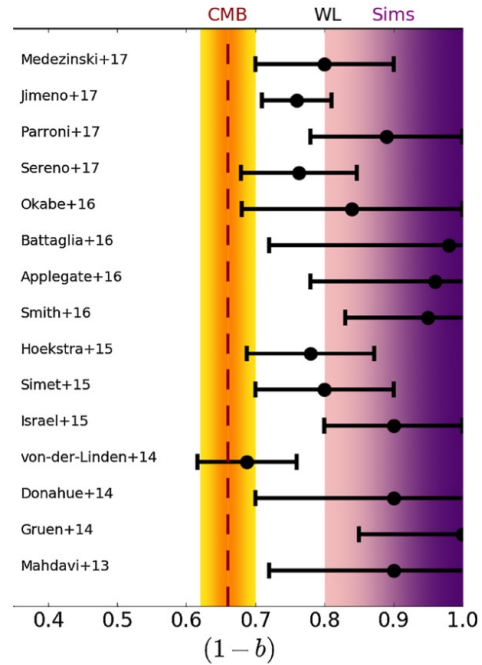


Figure 2.1 : Estimation du biais hydrostatique par comparaison de mesures de masse hydrostatique et par lentillage faible (voir section 2.2.1) pour différentes études (noir). La valeur nécessaire pour concilier les résultats du CMB et des amas de galaxies, mesurée par [51], est indiquée en orange. La région violette représente les valeurs de $(1-b)$ permises par plusieurs analyses d'amas simulés [52]. Figure extraite de [51].

[L. Salvati, et al. 2018]

Cosmo params

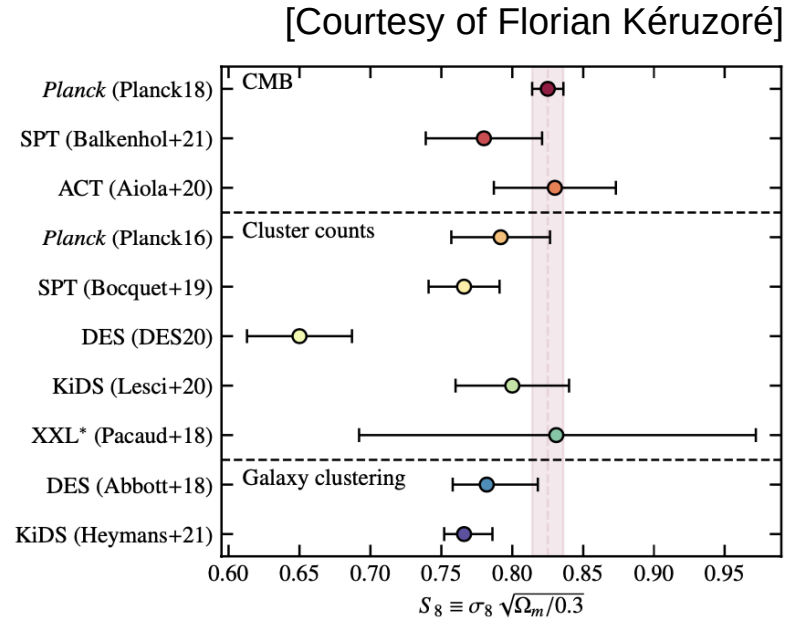


Figure 2.9 : Contrainte sur la combinaison de paramètres cosmologiques S_8 obtenues par différentes mesures du CMB (haut), de comptage d'amas (milieu), et de distribution de galaxies (bas). De haut en bas, les résultats sont issus de [13][119][120][100][99][103][107][102][121][122].

* : incertitudes estimées graphiquement.

Observed clusters in mass and redshift

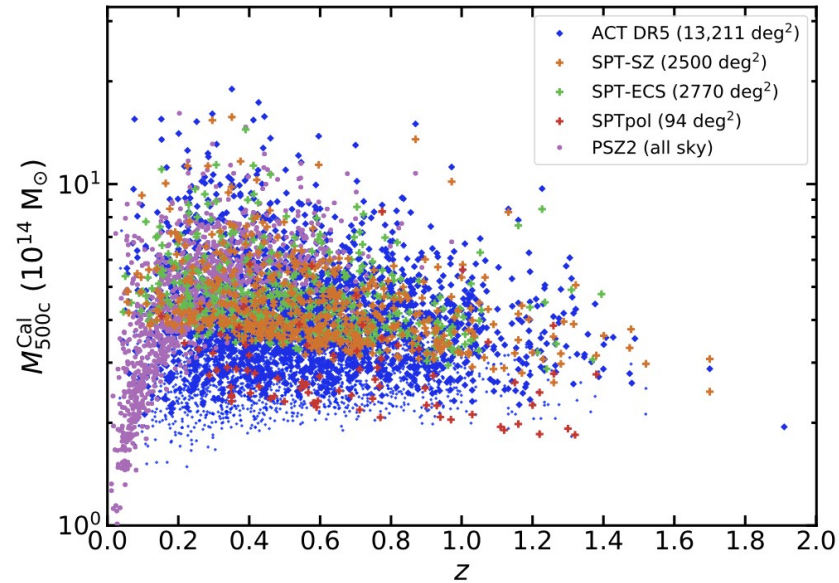


Figure 2.8 : Distribution dans le plan masse-redshift des amas détectés par les relevés ACT (bleu, [2]), SPT (orange [3], vert [4] et rouge [114]) et *Planck* (violet, [113]). Figure extraite de [2].

[M. Hilton et al., 2021]

SZ effect formalism

[Courtesy of Florian Kéruzoré]

La dépendance spectrale de la distorsion peut être calculée à partir de l'équation de Kompaneets [81, 82] : pour un bain d'électrons non-relativistes ($k_B T_e \ll m_e c^2$) traversé par des photons de fréquence réduite $x_e \equiv h\nu/k_B T_e$, on a

$$\frac{\partial n}{\partial y} = \frac{1}{x_e^2} \frac{\partial}{\partial x_e} \left[x_e^4 \left(\frac{\partial n}{\partial x_e} + n^2 + n \right) \right], \quad (2.8)$$

où n est l'indice d'occupation des photons, défini pour une intensité spécifique I_ν à une fréquence ν comme

$$n \equiv \frac{I_\nu c^2}{2h\nu^3}, \quad (2.9)$$

et y est le paramètre de Compton, défini comme

$$y \equiv \int \frac{k_B T_e}{m_e c^2} d\tau_e \quad (2.10)$$

où $\tau_e \equiv n_e \sigma_T l$ est la profondeur optique⁶ parcourue par les photons le long de la ligne de visée l ; n_e est la densité d'électrons, et σ_T la section efficace de diffusion Thompson.

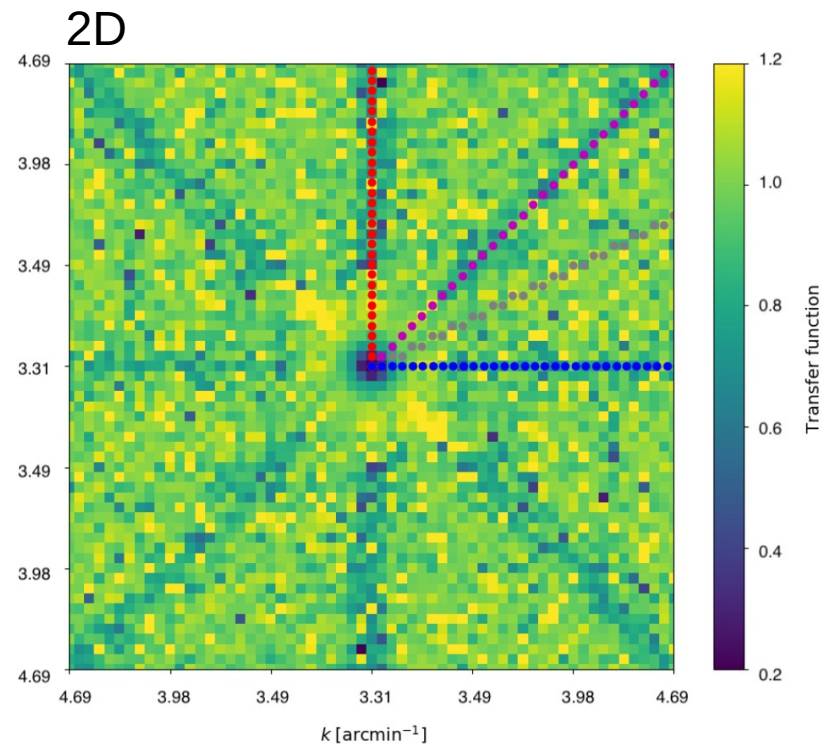
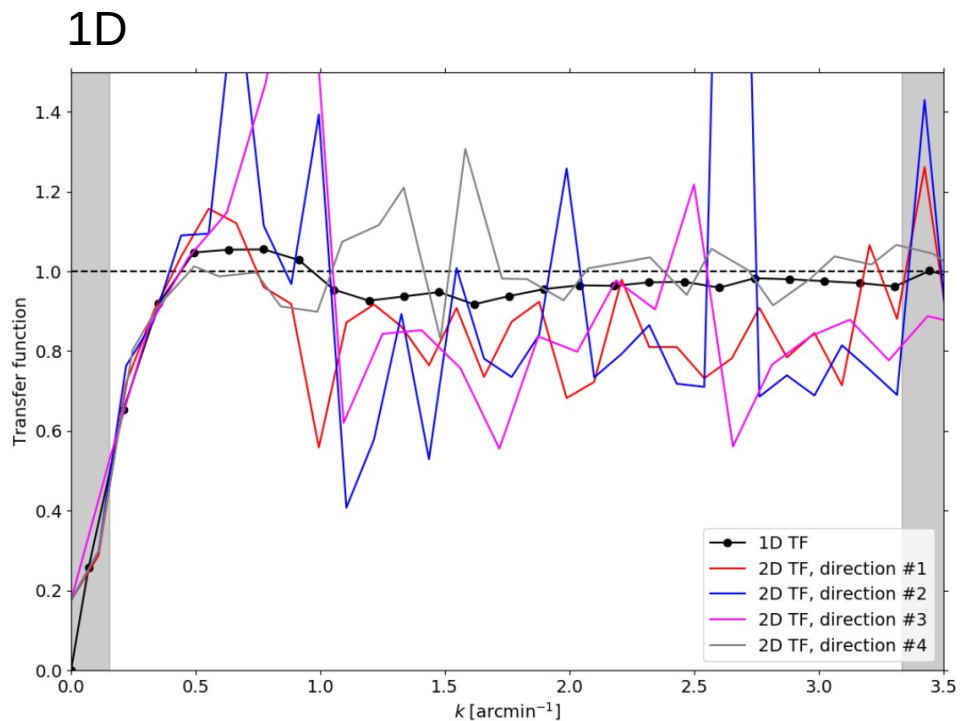
Cette équation peut alors aussi s'écrire

$$y = \int \frac{k_B T_e}{m_e c^2} n_e \sigma_T dl = \frac{\sigma_T}{m_e c^2} \int k_B T_e n_e dl. \quad (2.11)$$

Dans le cas de l'effet SZ, l'énergie des photons incidents est très faible par rapport à celle des électrons, soit $x_e \ll 1$ et donc $\partial n / \partial x_e \gg n, n^2$. L'équation (2.8) se simplifie alors, et en y injectant l'expression d'un spectre de corps noir pour les photons du CMB, on trouve l'expression de la variation d'intensité spécifique due à l'effet tSZ :

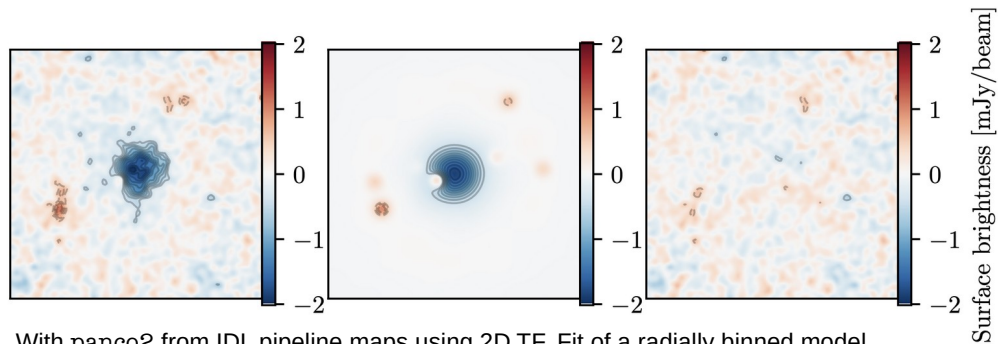
$$\frac{\Delta I_\nu^{\text{SZ}}}{I_0} = y \times \frac{x^4 e^x}{(e^x - 1)^2} [x \coth(x/2) - 4] [1 + \delta_r(x, T)], \quad (2.12)$$

Transfer function: raw data analysis-induced filtering



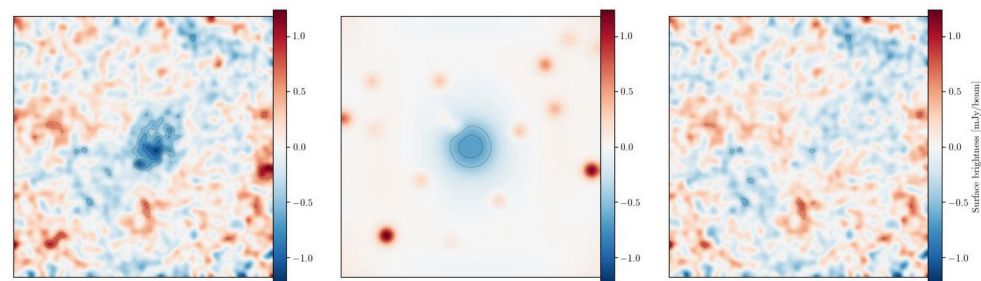
Fit of SZ maps to pressure models

CLJ1226.9+3332



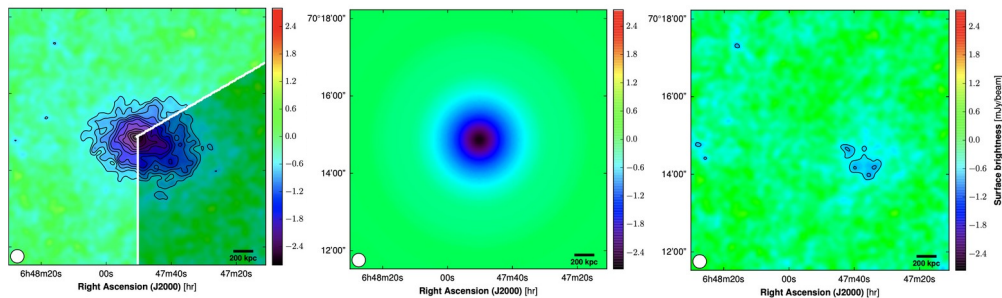
With `panco2` from IDL pipeline maps using 2D TF. Fit of a radially binned model.

PSZ2G228



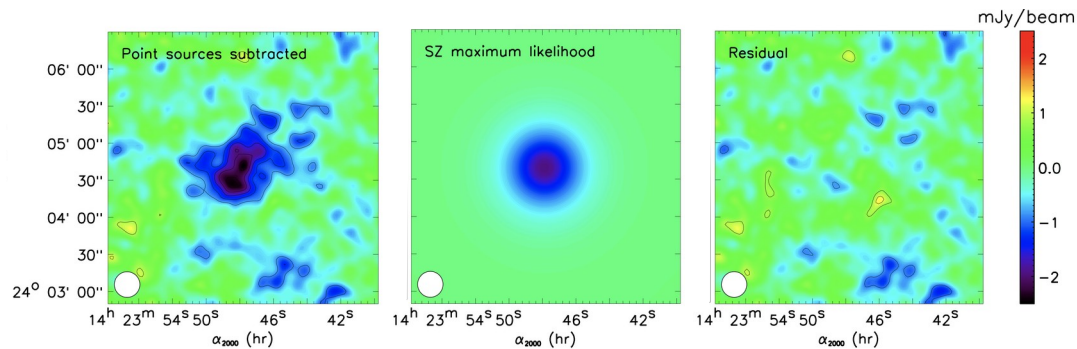
With `panco2` from IDL pipeline maps using 2D TF. Fit of a radially binned model. Big presence of contaminants and weaker cluster than expected.

PSZ2G144



(Ruppin et al. 2018)
Residual over-pressure from spherical fit.

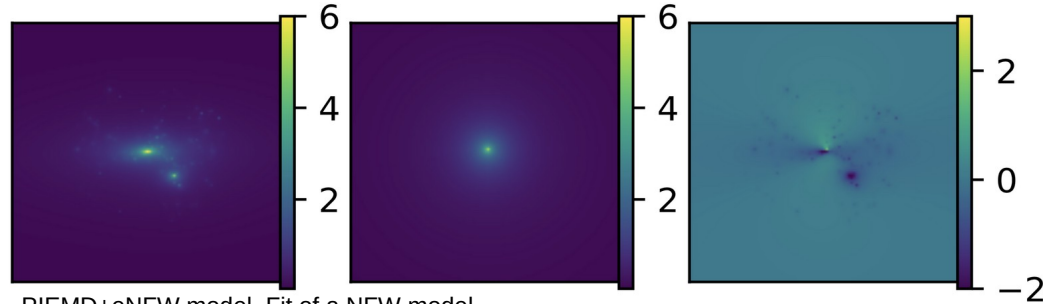
MACSJ1424



(Adam et al. 2018)
Central point source, subtracted for the fit of the pressure.

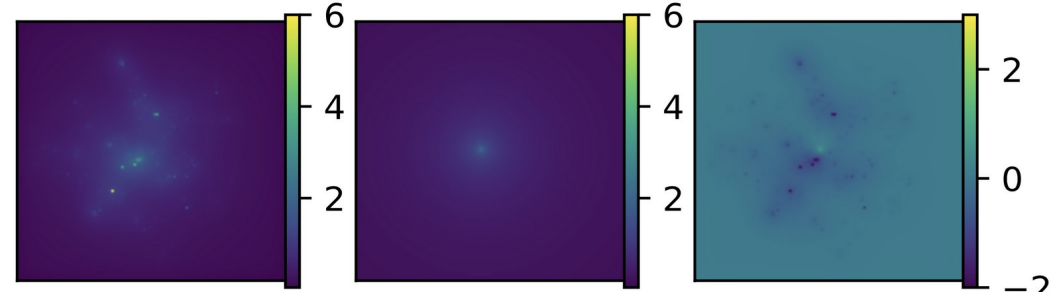
Fit of κ -maps to density models

CLJ1226.9+3332



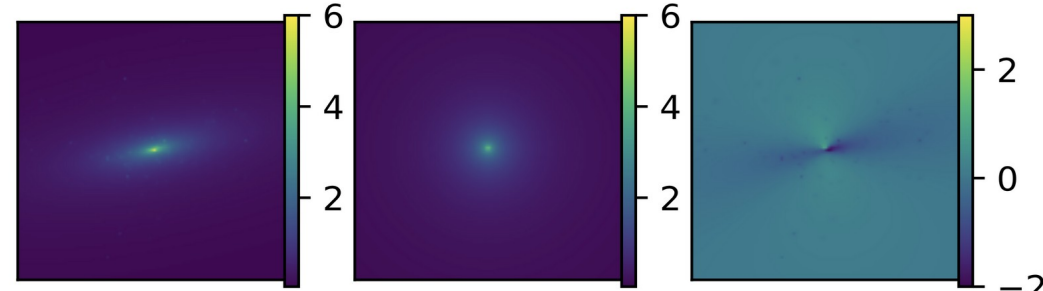
PIEMD+eNFW model. Fit of a NFW model.
Second clump present in projected density map

PSZ2G228



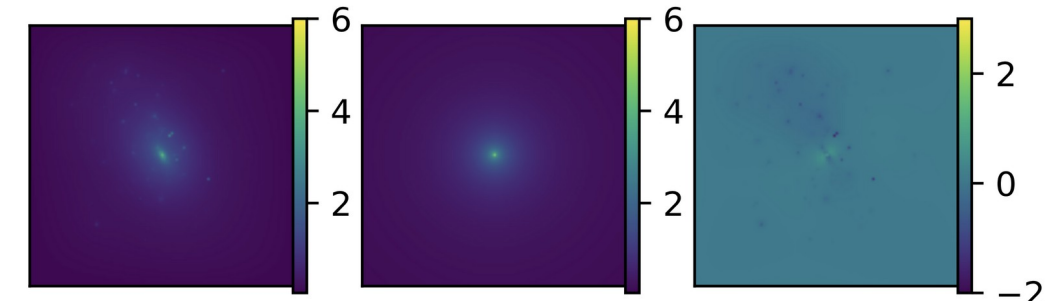
LTM model. Fit of a NFW model.
Complex morphology

PSZ2G144



PIEMD+eNFW model. Fit of a NFW model.
No hints of the over-pressure

MACSJ1424



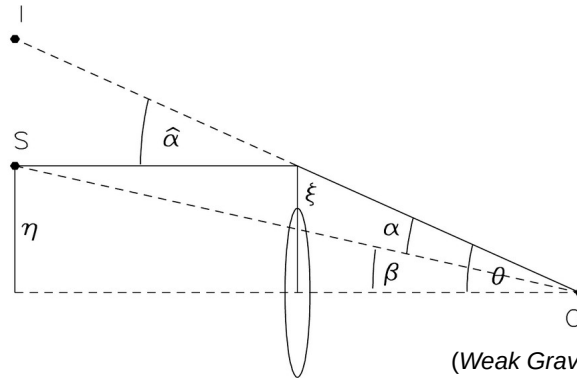
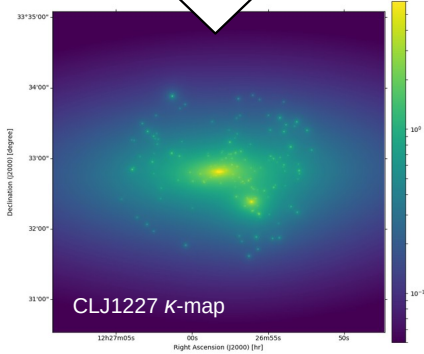
LTM model. Fit of a NFW model.

Deflection

A distribution of matter with surface mass density κ deflects the light an angle α



$$\nabla \cdot \vec{\alpha} = 2\kappa$$



(Weak Gravitational Lensing,
M.Bartelmann and P.Schneider,
2001)

$$\vec{\beta} = \vec{\theta} - \vec{\alpha}(\vec{\theta})$$

Real
position of
the source

Observed position
of the source