TMD evolution and parton showers

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A. Bermúdez Martínez on behalf of the TMD PB collaboration



Why TMDs?

- Small transverse momentum phenomena
- Small-x phenomena
- DY, and semi-inclusive DIS
- Transverse momentum effects from intrinsic kt and evolution

Parton Branching (PB) method

- Evolution of TMDs (and collinear PDFs)
- Resummation of soft gluons at LL and NLL
- FH et al. [PLB 772 (2017) 446–451] FH et al. [JHEP 2018, 70 (2018)] ABM et al. [PRD 99, 074008 (2019)]

- Solution valid at LO, NLO and NNLO
- Determination of TMDs from the fully exclusive solution
- Backward evolution fully determines the TMD shower
 - consistently treats perturbative and non-perturbative transverse momentum effects

PB formulation of TMD evolution

[slide by M. van Kampen]

PB evolution equation for TMDs $\tilde{A}_a(x, k_t^2, \mu^2)$ can be solved iteratively with the Monte Carlo method:

$$\begin{split} \tilde{\mathcal{A}}_{a}(x,k_{t}^{2},\mu^{2}) &= \Delta_{a}(\mu^{2},\mu_{0}^{2})\tilde{\mathcal{A}}_{a}(x,k_{t,0}^{2},\mu_{0}^{2}) + \\ &+ \sum_{b} \left[\int \frac{d^{2}\mu'}{\pi\mu'^{2}} \int_{x}^{z_{M}(\mu')} dz \Theta(\mu^{2}-\mu'^{2})\Theta(\mu'^{2}-\mu_{0}^{2}) \right. \\ &\times \frac{\Delta_{a}(\mu^{2},\mu_{0}^{2})}{\Delta_{a}(\mu'^{2},\mu_{0}^{2})} P_{ab}^{(R)}(\alpha_{s}(q_{t}),z) \tilde{\mathcal{A}}_{b}\left(\frac{x}{z},\underbrace{k_{t,b}-q_{t,c}}_{k_{t,a}},\mu'^{2}\right) \right] \end{split}$$

JHEP 01 (2018) 070 [arXiv:1708.03279]



Kinematics in each branching governed by momentum conservation: $k_{t,b} = k_{t,a} + q_{t,c}$

 $P_{ab}^{(R)}(\alpha_s, z)$ real splitting function (resolvable branching probability), $\Delta_a(\mu^2, \mu_0^2)$ Sudakov (no branching probability) Angul

 $P_{ab}^{(R)}(\alpha_s, z) = \sum_{n=1}^{\infty} \left(\frac{\alpha_s}{2\pi}\right)^n P_{ab}^{(R)n-1}(z)$ $\Delta_a(\mu^2, \mu_0^2) = \exp\left(-\sum_b \int \frac{d\mu^2}{\mu^2} \int_0^{z_M} dz \ z \ P_{ab}^{(R)}(z, \alpha_s)\right)$



PB formulation of TMD evolution

[slide by M. van Kampen]

Backward evolution with PB method

The TMD evolution equation can be used to do a backward evolution:

$$\frac{\partial}{\partial \ln \mu^2} \left(\frac{\tilde{\mathcal{A}}_{a}(x, k_t, \mu)}{\Delta_{a}(\mu)} \right) = \sum_{b} \int_{x}^{z_{M}} dz P_{ab}^{(R)} \frac{\tilde{\mathcal{A}}_{b}(x/z, k'_t, \mu)}{\Delta_{a}(\mu)}$$

normalize to $\frac{\tilde{\mathcal{A}}_{a}(x,k_{t},\mu)}{\Delta_{a}(\mu)}$ and integrate over μ' from μ_{i} down to μ_{i-1}

$$\Delta_{bw}(x, k_t, \mu_i, \mu_{i-1}) = \exp\left\{-\sum_b \int_{\mu_{i-1}^2}^{\mu_i^2} \frac{d\mu'^2}{\mu'^2} \int_x^{z_M} dz P_{ab}^{(R)} \frac{\tilde{\mathcal{A}}_b(x/z, k_t', \mu')}{\tilde{\mathcal{A}}_a(x, k_t, \mu')}\right\}$$

In each splitting

This Sudakov is used as the no-branching probability in the TMD parton shower.



$$k_{t,b} = k_{t,a} + q_{t,c}$$

$$= \mathbf{k}_{t,a} + (1-z)\mu$$

Total transverse momentum:

$$k_t = k_{t,0} + \sum_c q_{t,c}$$



Implemented in the CASCADE event generator

S. Baranov et al. [Eur. Phys. J. C 81 (2021) 425]

PB method recap

 $A_{a}^{(1)}(x, \mathbf{k}_{t}; Q^{2}) = \Delta_{a}(Q^{2}, Q_{0}^{2})A_{a}(x, \mathbf{k}_{t}; Q_{0}^{2}) + \sum_{b} \int_{Q_{0}^{2}}^{Q^{2}} \frac{d^{2}\mathbf{Q}'}{\pi Q'^{2}} \frac{\Delta_{a}(Q^{2}, Q_{0}^{2})}{\Delta_{a}(Q'^{2}, Q_{0}^{2})} \int_{x}^{z_{M}} dz P_{ab}^{(R)}\left(z, \alpha_{s}(Q'^{2})\right) \Delta_{b}(Q'^{2}, Q_{0}^{2})A_{b}\left(\frac{x}{z}, \mathbf{k}_{t} + (1-z)\mathbf{Q}'; Q_{0}^{2}\right)$



- kinematics of the splittings is known
- physics mapping of evolution variables to splitting kinematics
- TMD from cumulative kt of the branchings in forward PB evolution
- Automatically includes resummation at NLL
- Initial-state shower fully determined by TMD and its backward PB evolution
- Parton shower exactly matches the evolution of the TMD

Pert. and non-pert. PB TMD contributions



ABM et al. [PRD 99, 074008 (2019)]

- \circ DIS measurements from HERA I+II
- \circ fitting procedure in a nutshell:
 - parametrize the integrated distribution at Q_0
 - with the PB method evolve the TMD to $Q>Q_0$

(implemented in xFitter)

- fit the measurements and extract the initial

parametrization

- store the TMD in a grid for later use

(TMDlib, complementary slides)

N. A. Abdulov et al. [Eur. Phys. J. C 81 (2021) 752]

Consider the integrated distribution above the jet pT scale:

$$a_j(x,m{k},\mu^2) = \int rac{d^2m{k}'}{\pi} \; \mathcal{A}_j(x,m{k}',\mu^2) \; \Theta(m{k}'^2 - m{k}^2)$$

- e.g. probability of 0.3 that the gluon develops a kt larger than 20 GeV, for μ = 100 GeV

- TMD evolution effects crucial at describing jet production



A. Bermúdez Martínez

Z pT in a wide range of DY mass

Application to high mass DY production

DY pt spectrum

- Combined with MC@NLO
- Excellent description of DY pT spectrum

ABM et al. [PRD 100, 074027 (2019)] ABM et al. [EPJC 80, 598 (2020)]

- Non-perturbative TMD effects not significant at high pT
- Missing contributions at high pT



Application to low mass DY production

DY pt spectrum

- Combined with MC@NLO
- Excellent description of DY pT spectrum

ABM et al. [PRD 100, 074027 (2019)] ABM et al. [EPJC 80, 598 (2020)]

- First simultaneous description of both low and high-mass DY pT spectrum
- No more low pT crisis Bacchetta et al. [PRD 100 (2019) 014018]; ABM et al. [EPJC 80, 598 (2020)]



Jet production

Z + jets production

Jet multiplicity and third jet pT spectrum

- PB shower combined with Z + j and Z + jj MC@NLO
- Suffers at higher multiplicities
- Good description of third jet pT spectrum given by higher orders



Z + jets production

Jet multiplicity and third jet pT spectrum

- Significant contribution from perturbative part of TMD
- Final state emissions important at higher multiplicities



Trying TMD shower with higher orders

- Important deficit at high pT with Z at NLO
- Potentially large corrections by higher orders
- Try combining high pT TMD effects with multiple higher orders



What we want:

- Treat perturbative and non-perturbative TMD effects
- Include soft gluon resummation
- Include corrections from higher-order fixed-order calculations

Develop a method to combine PB-TMDs with multijet calculations

Multi-jet merging

- Make higher-order ME exclusive by Sudakov suppression
- Avoid double counting between PS and ME
- Improvement of hard, wide-angle emissions
- Description of high-pT phenomena



DESY.

TMD merging method

ABM et al. [arXiv:2107.01224]

- Evaluate the ME for n-jet cross sections
- Reweight the strong coupling according to shower history
- Evolve the ME using the TMD PB evolution
- Shower the events using the backward PB evolution for ISR
- Apply the MLM^[1] prescription between the PB-evolved ME and the showered events

NB: The method could also be applied to merging criteria other than $\ensuremath{\mathsf{MLM}}$

TMD $k_t > 0$ DY $k_t > 0$ TMD

New merging procedure applicable to TMDs!

[1] M. L. Mangano [NPB 632 (2002) 343-362]

ABM et al. [paper in preparation] What about the original MLM applied to TMD events?

• very strong dependence on Rclus

• at large scales ME accuracy lost!





d(n,n+1): square of scale at which (n+1)-jet configuration becomes n-jet

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d(n,n+1): square of scale at which (n+1)-jet configuration becomes n-jet



DESY.

ABM et al. [paper in preparation]

TMD merging

- little dependence on Rclus
- at large scales ME accuracy recovered!





ABM et al. [paper in preparation]

d(n,n+1): scale at which (n+1)-jet configuration becomes n-jet



- Smoothness merging follows shower Sudakov suppression
- Merging scale divides phase space for different jet multiplicities avoiding double counting

DY pt spectrum

• TMD evolution with multi-jet merging achieved at LO

New! ABM et al. [arXiv:2107.01224]

- Low as well as high-pt now nicely described
- Consistent with MCatNLO PB-NLO at low pT



Exclusive jet multiplicity in Z events

ATLAS data: [Eur. Phys. J. C77 (2017) 361]



• Not only the overall recoil but also the number of jets are described

DESY.

Jets pt spectrum

• Not only overall recoil but also jet pT

New! ABM et al. [arXiv:2107.01224]

• The description of jet pT improves at high multiplicities



Systematics

Multi-jet cross section in Z production

Merging scale [GeV]	σ[tot] [pb]	σ[≥ 1 jet] [pb]	σ[≥ 2 jet] [pb]	σ[≥ 3 jet] [pb]	σ[≥ 4 jet] [pb]
23	572.98	87.26	20.27	4.84	1.18
33	563.04	86.15	20.48	4.86	1.19

- 10 GeV variation gives < 2% change in jets cross sections
- Standard merging algorithms can give over 10 % change for the same variation of the merging scale CF: J. Alwall et al. [EPJC 53, 473–500 (2008)]
 - Dependence on merging scale reduced by treating transverse momentum in the initial state

Systematics

Z pt and phi*



Less than 10% effect localized around the merging scale



^{dσ}/dd] [pb/GeV]

Ratio

A. Bermúdez Martínez

ptj3 [GeV] 27



Conclusions

- PB TMD evolution provides excellent description of DY pt spectrum in a wide range of DY mass
- Parton shower from PB TMD evolution have significant contribution to jet multiplicity and jet pt spectra
- Higher fix-order contributions to PB TMD evolution are significant
- First combination of TMD evolution effects with multi-jet merging for Z pt and jet spectra
- Dependence of the results on the merging scale are smaller than that of standard algorithms



PB framework Phys. Lett. B 772:446451 (2017) JHEP 01:070 (2018) Phys. Rev. D 99, 074008 (2019)

- TMD determined, no extra parameters
- Full access to splitting kinematics
- TMD evolution implemented in xFitter



anti-up, x = 0.01, µ = 100 GeV

Where to find them:

arXiv:2103.09741 (accepted for publication in EPJC)

- TMDlib: library of parametrization of TMDs and uPDFs
- TMDplotter: TMD plotting tool

