

Parton Shower Monte Carlo Event Generators

The Universit of Mancheste

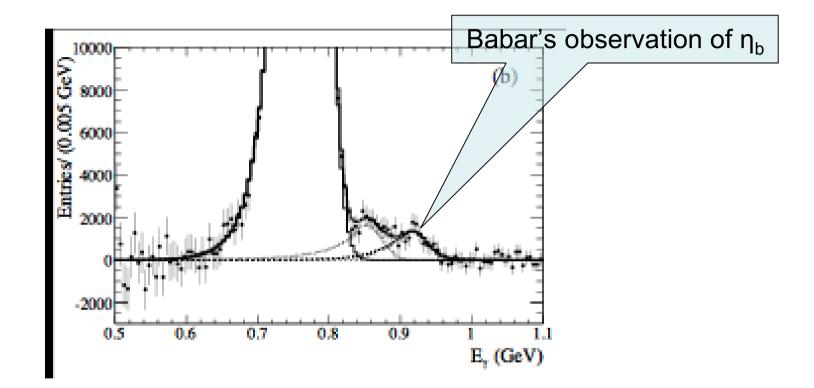
Mike Seymour University of Manchester & CERN

MC4LHC EU Networks' Training Event

May $4^{th} - 8^{th} 2009$

http://www.montecarlonet.org/

Overview and Motivation

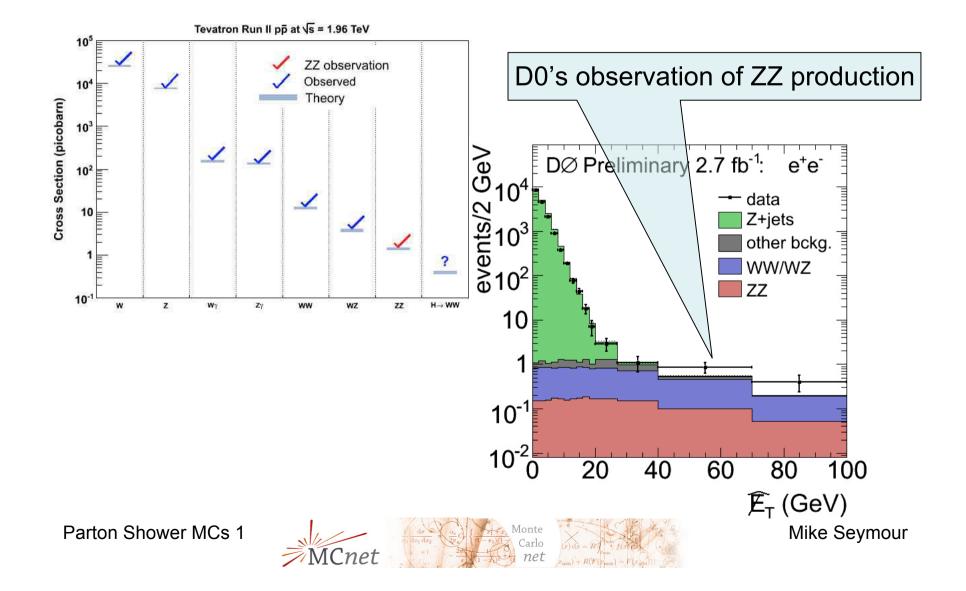




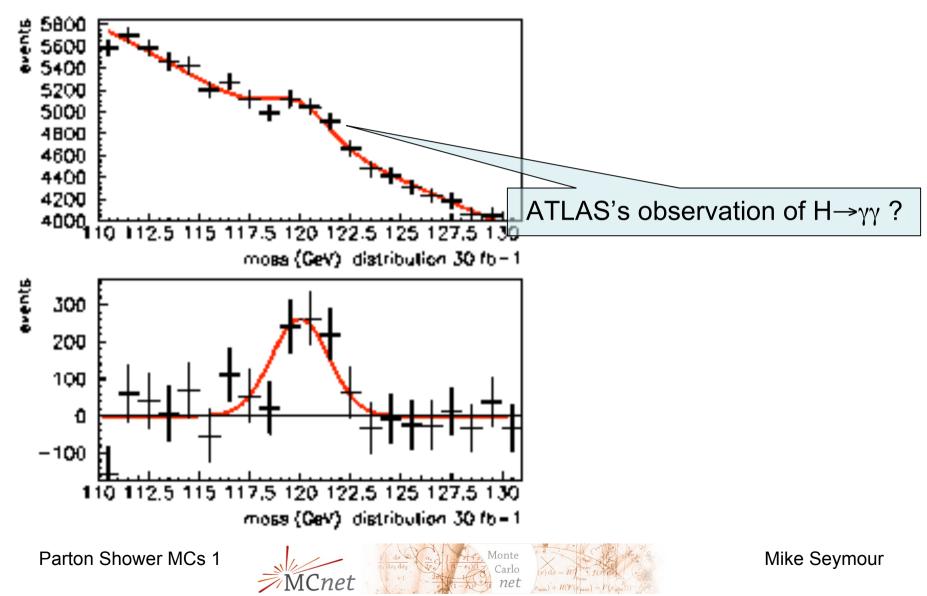
Mike Seymour

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Overview and Motivation



Overview and Motivation





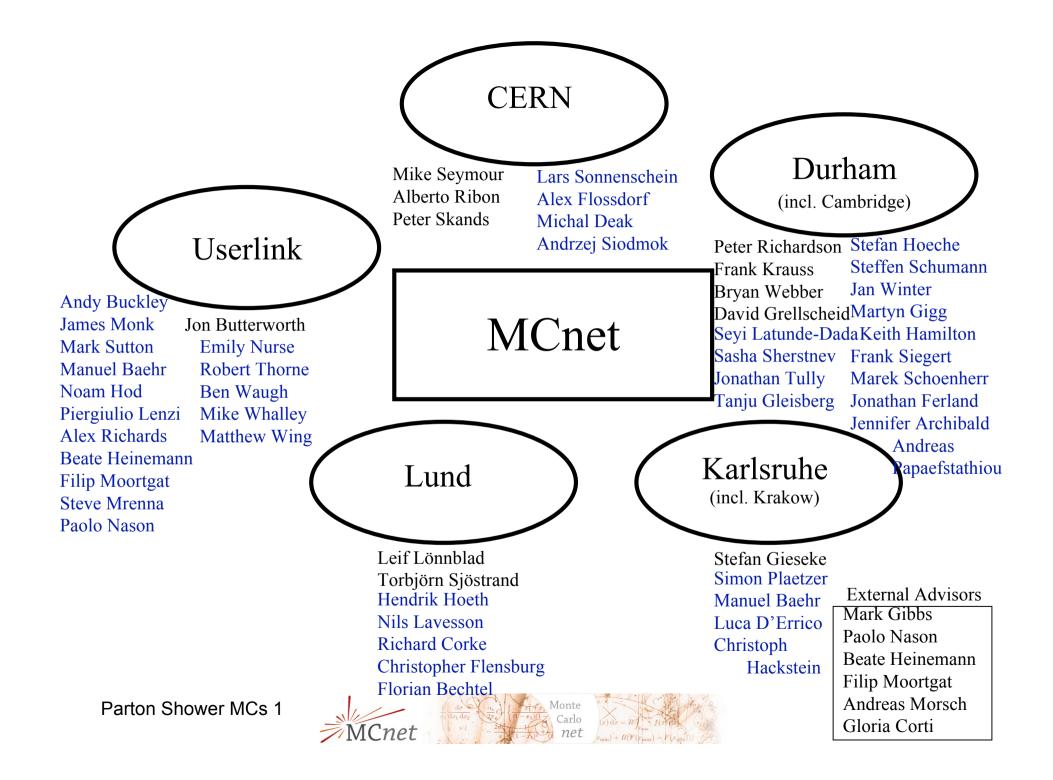


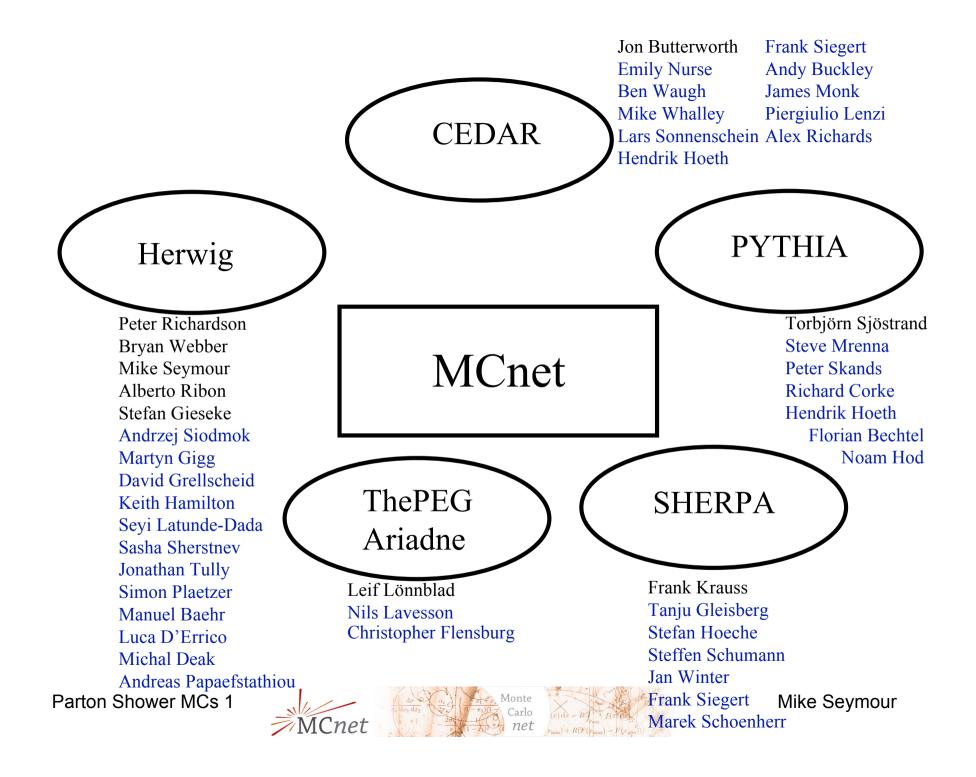
- Marie Curie Research Training Network
- for Monte Carlo event generator
 - development
 - validation and tuning
- Approved for four years from 1st Jan 2007



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MCnet objectives

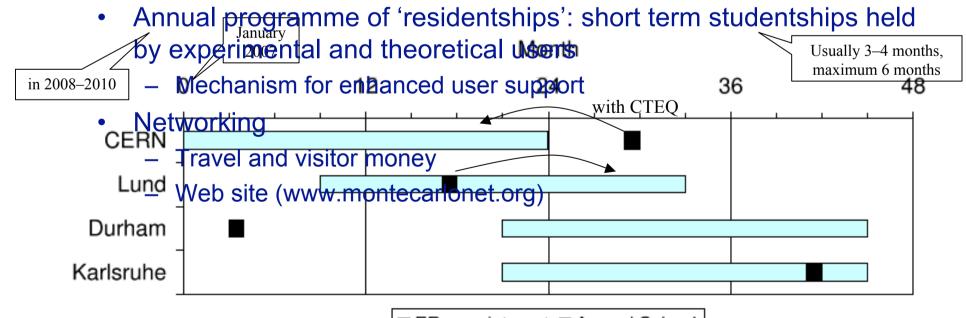
Training:

- To train a large section of the user base in the physics and techniques of event generators
- To train the next generation of event generator developers Through Research:
- To develop the next generation of event generators intended for use throughout the lifetimes of the LHC and ILC experiments
- To play a central role in the analysis of early LHC data and the discovery of new particles and interactions there
- To extract the maximum potential from existing data to constrain the modeling of the data from the LHC and other future experiments.



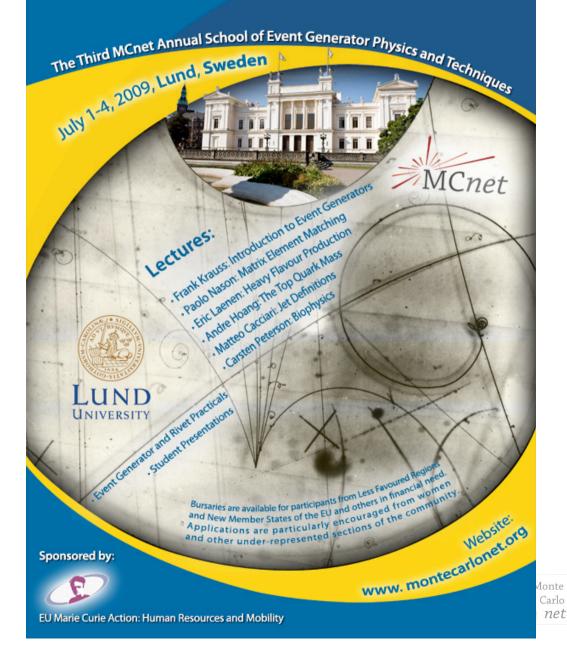
MCnet main activities

- Four postdoc positions
- Two joint studentships (Karlsruhe–Durham, Durham–UCL)
- Annual School
- Two annual meetings (Januarys @ CERN, summer with school)



ER appointment Annual School

2009 MCnet Summer School



MCnet opportunities 2009:

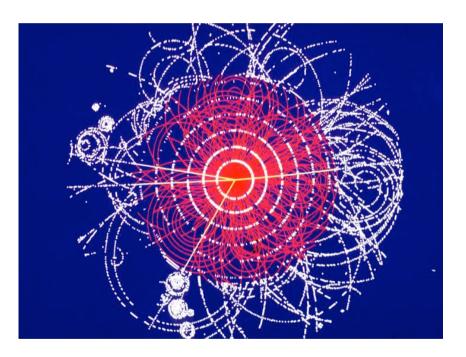
- MC4LHC (with Heptools and Artemis)
- MCnet school $(\rightarrow May 1^{st})$
- Short-term studentships: for th. and exptl. students to spend 3-6 months with MC authors in:
 - CERN
 - Durham/Cambridge
 - Karlsruhe
 - Lund
 - UCL

on a project of their choice

- Next closing dates:
 - May 4th
 - August 3rd

Introduction to Parton Shower Monte Carlo Event Generators

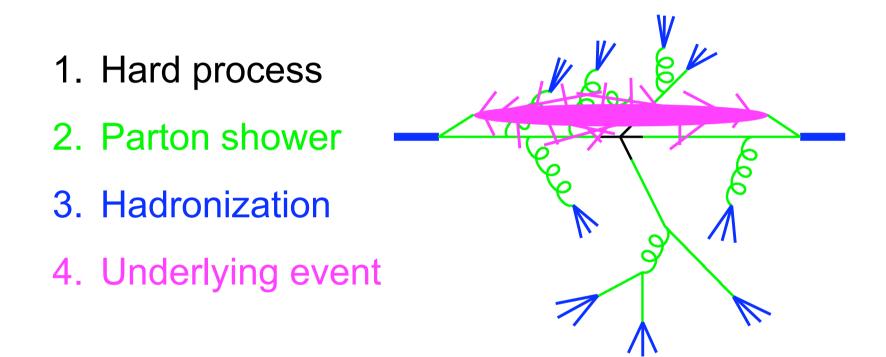
- Basic principles
- LHC event generation
- Parton showers
- Hadronization
- Underlying Events
- Practicalities



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Structure of LHC Event Simulations





Hard Process Simulation

Typically use fixed-order perturbative matrix elements Leading order can be largely automated...

- MADGRAPH/MADEVENT
- GRACE
- COMPHEP
- AMAGIC++ (SHERPA)
- ALPGEN
- HELAC

Matrix elements squared positive definite → simple Monte Carlo implementation

Next-to-leading order getting there...

- MCFM
- NLOJET++
- MC@NLO
-

Real and virtual contributions have equal and opposite divergences → naïve Monte Carlo fails

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Hard Process Simulation

Typically use fixed-order perturbative matrix elements Leading order can be largely automated...

- MADGRAPH/MADEVENT
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- AMAGIC++ (SHERPA)
- ALPGEN
- HELAC

But...

- Fixed parton/jet multiplicity
- No control of large logs
- Parton level

Matrix elements squared positive definite → simple Monte Carlo implementation

→ Need hadron level event generators



Intro to Monte Carlo Generators

- 1. Basic principles
- 2. Parton showers
- 3. Hadronization
- 4. Introduction to the MCnet Monte Carlo Event Generator projects



Parton Showers: Introduction

- QED: accelerated charges radiate.
- QCD identical: accelerated colours radiate.
- gluons also charged.
- \rightarrow cascade of partons.
- = parton shower.

- 1. e^+e^- annihilation to jets.
- 2. Universality of collinear emission.
- 3. Sudakov form factors.
- 4. Universality of soft emission.
- 5. Angular ordering.
- 6. Initial-state radiation.
- 7. Hard scattering.
- 8. Heavy quarks.
- 9. Dipole cascades.
- 10. Matrix element matching

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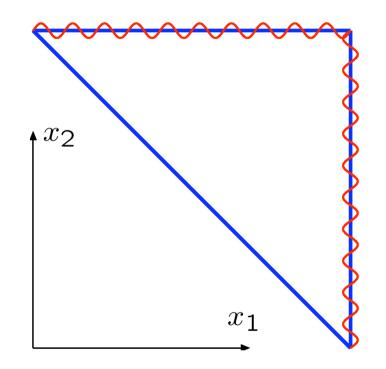
e^+e^- annihilation to jets

Three-jet cross section:

$$\frac{d\sigma}{dx_1 dx_2} = \sigma_0 C_F \frac{\alpha_s}{2\pi} \frac{x_1^2 + x_2^2}{(1 - x_1)(1 - x_2)}$$

singular as $x_{1,2} \to 1$

Rewrite in terms of quark-gluon opening angle θ and gluon energy fraction x_3 :



$$\frac{d\sigma}{d\cos\theta \, dx_3} = \sigma_0 \, C_F \frac{\alpha_s}{2\pi} \left\{ \frac{2}{\sin^2\theta} \, \frac{1 + (1 - x_3)^2}{x_3} - x_3 \right\}$$

Singular as $\sin \theta \rightarrow 0$ and $x_3 \rightarrow 0$.

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can separate into two independent jets:

$2 d\cos\theta$	_	$d\cos\theta$ _	$d\cos\theta$
$\sin^2\theta$	_	$1 - \cos \theta$	$\frac{1}{1+\cos\theta}$
		$d\cos\theta$	$d\cos\overline{ heta}$
	_	$1 - \cos \theta$	$1 - \cos ar{ heta}$
	\approx	$\frac{d\theta^2}{\theta^2} + \frac{d\overline{\theta}^2}{\overline{\theta}^2}$	

jets evolve independently

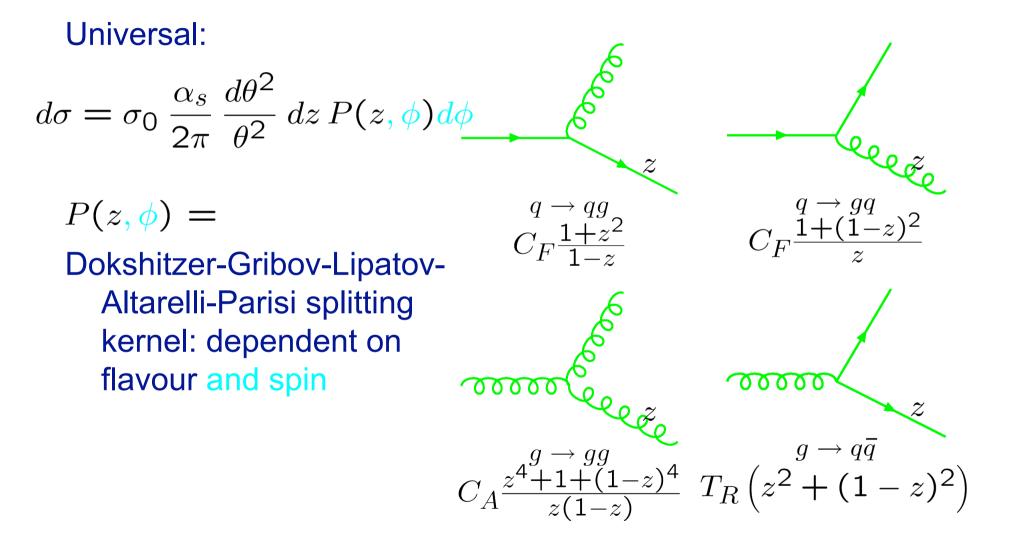
$$d\sigma = \sigma_0 \sum_{\text{jets}} C_F \frac{\alpha_s}{2\pi} \frac{d\theta^2}{\theta^2} dz \frac{1 + (1 - z)^2}{z}$$

Exactly same form for anything $\propto \theta^2$ eg transverse momentum: $k_{\perp}^2 = z^2(1-z)^2 \ \theta^2 \ E^2$ invariant mass: $q^2 = z(1-z) \ \theta^2 \ E^2$

$$\frac{d\theta^2}{\theta^2} = \frac{dk_{\perp}^2}{k_{\perp}^2} = \frac{dq^2}{q^2}$$

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Collinear Limit



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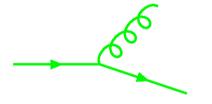
Resolvable partons

What is a parton?

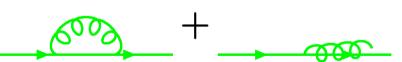
Collinear parton pair \longleftrightarrow single parton

Introduce resolution criterion, eg $k_{\perp} > Q_0$.

Virtual corrections must be combined with unresolvable real emission



Resolvable emission Finite



Virtual + Unresolvable emission Finite

Unitarity: P(resolved) + P(unresolved) = 1

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Sudakov form factor

Probability(emission between q^2 and $q^2 + dq^2$)

$$d\mathcal{P} = \frac{\alpha_s}{2\pi} \frac{dq^2}{q^2} \int_{Q_0^2/q^2}^{1-Q_0^2/q^2} dz \ P(z) \equiv \frac{dq^2}{q^2} \bar{P}(q^2).$$

Define probability(no emission between Q^2 and q^2) to be $\Delta(Q^2, q^2)$. Gives evolution equation

$$\frac{d\Delta(Q^2, q^2)}{dq^2} = \Delta(Q^2, q^2) \frac{d\mathcal{P}}{dq^2}$$
$$\Rightarrow \Delta(Q^2, q^2) = \exp - \int_{q^2}^{Q^2} \frac{dk^2}{k^2} \bar{P}(k^2).$$

c.f. radioactive decay

atom has probability λ per unit time to decay. Probability(no decay after time T) = $\exp - \int^T dt \lambda$



Sudakov form factor

Probability(emission between q^2 and $q^2 + dq^2$)

$$d\mathcal{P} = \frac{\alpha_s}{2\pi} \frac{dq^2}{q^2} \int_{Q_0^2/q^2}^{1-Q_0^2/q^2} dz \ P(z) \equiv \frac{dq^2}{q^2} \bar{P}(q^2).$$

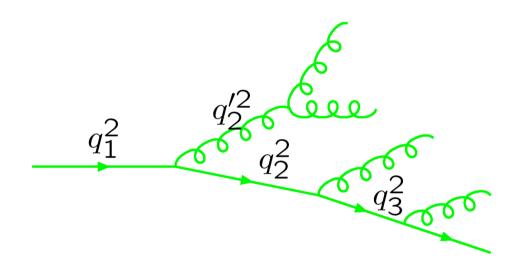
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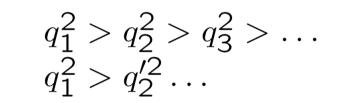
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 $\Delta(Q^2, Q_0^2) \equiv \Delta(Q^2)$ Sudakov form factor =Probability(emitting no resolvable radiation)

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$$\Delta_q(Q^2) \sim \exp{-C_F rac{lpha_s}{2\pi} \log^2 rac{Q^2}{Q_0^2}}$$

Multiple emission





But initial condition? $q_1^2 < ???$

Process dependent

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Monte Carlo implementation

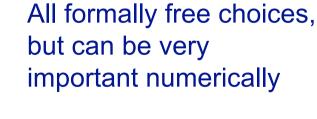
Can generate branching according to

$$d\mathcal{P} = \frac{dq^2}{q^2} \bar{P}(q^2) \,\Delta(Q^2, q^2)$$

By choosing $0 < \rho < 1$ uniformly: If $\rho < \Delta(Q^2)$ no resolvable radiation, evolution stops. Otherwise, solve $\rho = \Delta(Q^2, q^2)$ for q^2 =emission scale

Considerable freedom: Evolution scale: $q^2/k_{\perp}^2/\theta^2$? z: Energy? Light-cone momentum? Massless partons become massive. How? Upper limit for q^2 ?

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Running coupling

Effect of summing up higher orders:

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absorbed by replacing α_s by $\alpha_s(k_{\perp}^2)$. Much faster parton multiplication – phase space fills with soft gluons.

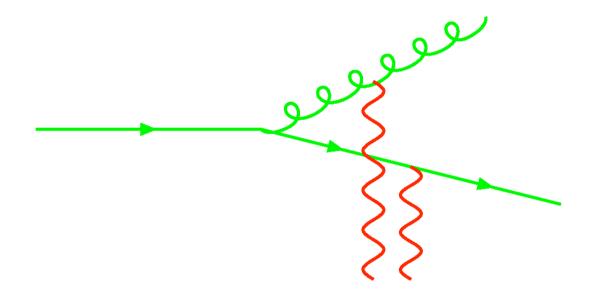
Must then avoid Landau pole: $k_{\perp}^2 \gg \Lambda^2$. Q_0 now becomes physical parameter!

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Soft limit

Also universal. But at amplitude level...



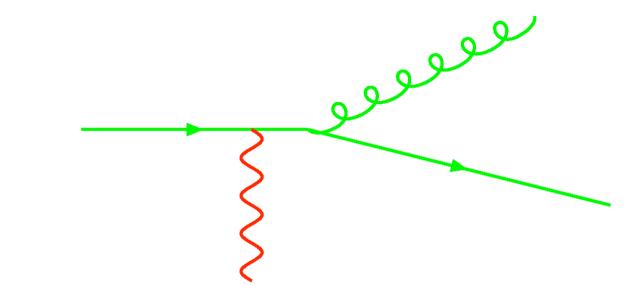
soft gluon comes from everywhere in event.

- \rightarrow Quantum interference.
- Spoils independent evolution picture?

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Angular ordering



outside angular ordered cones, soft gluons sum coherently: only see colour charge of whole jet.

Soft gluon effects fully incorporated by using θ^2 as evolution variable: angular ordering

First gluon not necessarily hardest!

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NO:

