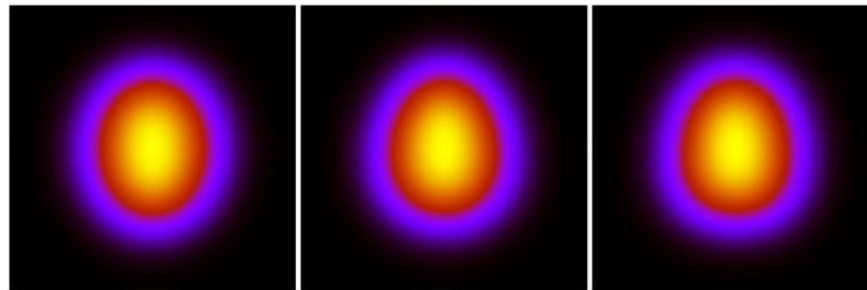


# „Multipole” exact solutions in relativistic hydrodynamics (& higher order flow coefficients)

Tamás Csörgő<sup>1</sup>, for Máté Csanád<sup>2</sup> & András Szabó<sup>2</sup>

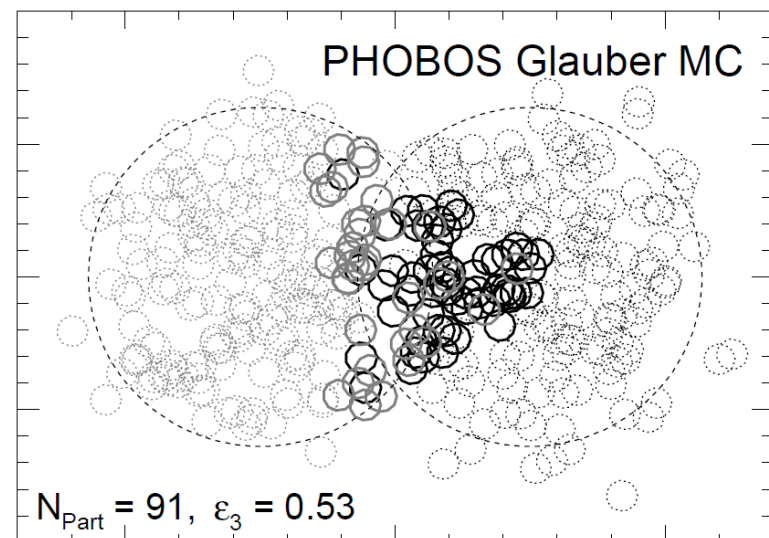
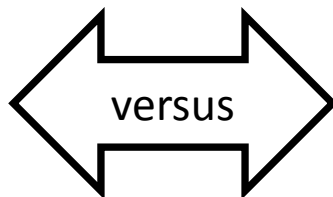
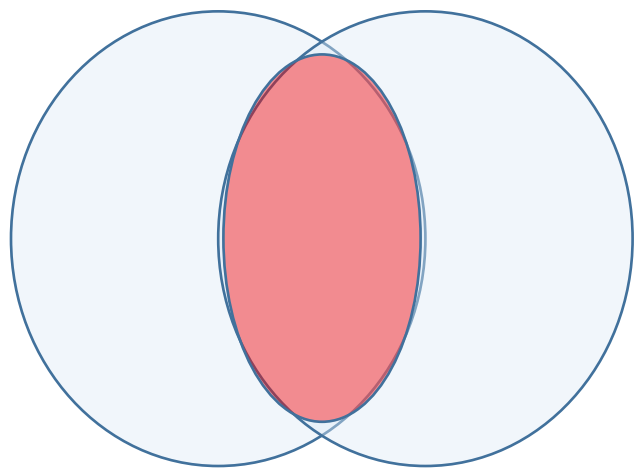
<sup>1</sup> Wigner Research Center for Physics, Nuclear & Particle Physics Inst.

<sup>2</sup> Eötvös University, Department of Atomic Physics



# Higher order flow and event-by-event hydro

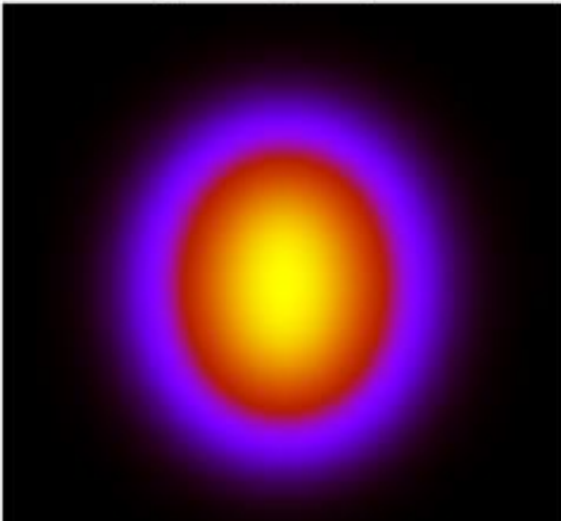
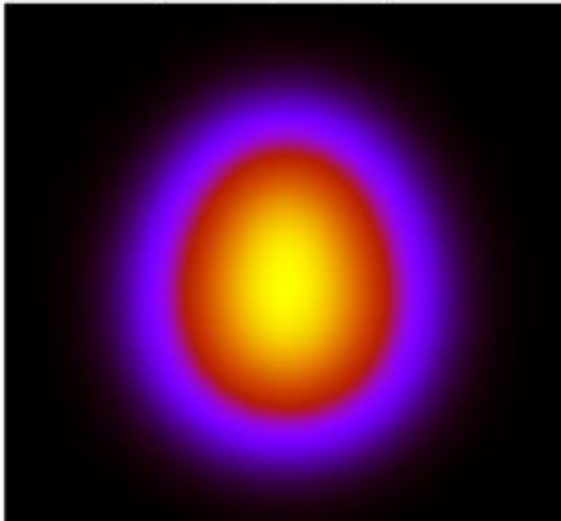
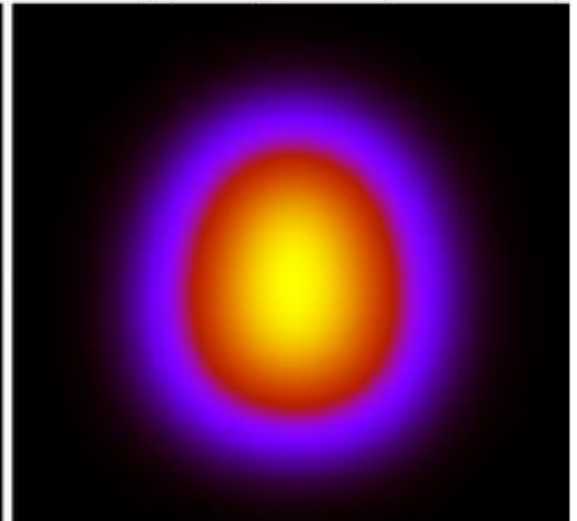
- Basic picture: elliptical region,  $v_2$  most important
- Finite number of nucleons: event-by-event fluctuating initial condition



- Higher order asymmetry (flow) coefficients arise from this effect!
- How to handle this in hydrodynamics?
- Numerically (see e.g. Broniowski et al.): implement fluctuating initial cond.
- Analytic hydro?

# New „multipole” solutions

- Based on Csörgő, Csernai, Hama, Kodama, Heavy Ion Phys. A21, 73 (2004)
- As  $\dot{R} = \text{const.}$ , Hubble-like velocity field:  $u^\mu = \frac{x^\mu}{\tau}$ , and  $\frac{\tau_0}{\tau} = \frac{\gamma R_0}{R}$
- At given  $\tau$ , temp. constant on scale variable defined surfaces, arbitrary profile
- Typical choice of temperature profile:  $\exp[-bs]$ , with  $b = \left\langle \frac{\Delta T}{T} \right\rangle_r$
- Several asymmetries superimposed, with various  $\epsilon_N$  coefficients
- Phase shifts define Nth order reaction plane, its angle is  $\psi_N$

 $\epsilon_2=0.6, \epsilon_3=0, \epsilon_4=0$ 

 $\epsilon_2=0.6, \epsilon_3=0.2, \epsilon_4=0$ 

 $\epsilon_2=0.6, \epsilon_3=0.2, \epsilon_4=0.1$ 


# Multipole symmetries in 2D

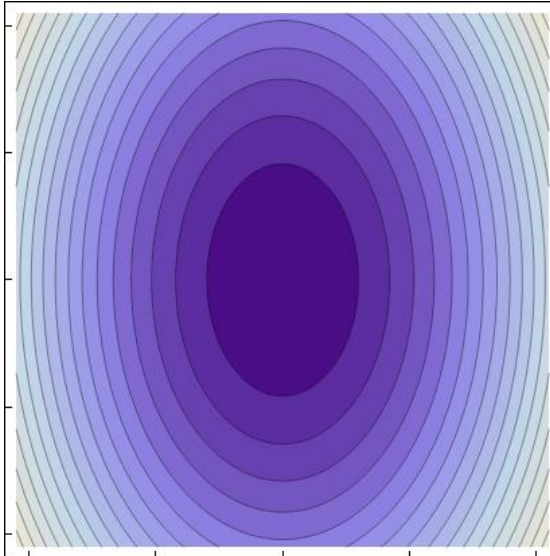
- Ellipsis in polar coordinates:

$$\frac{r_x^2}{x^2} + \frac{r_y^2}{y^2} = \frac{r^2}{R^2} \times [1 + \epsilon \cos(2\phi)], \text{ with } \frac{1}{x^2} + \frac{1}{y^2} = \frac{2}{R^2} \text{ and } \epsilon = \frac{x^2 - y^2}{x^2 + y^2}$$

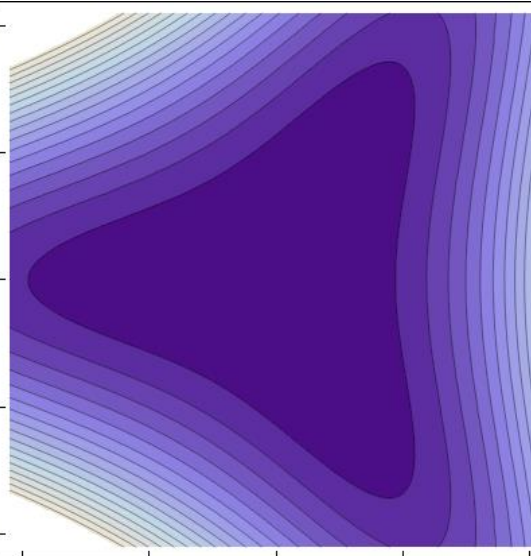
- Generalizing that to multipole symmetries:

$$\frac{r^N}{R^N} \times [1 + \epsilon \cos(N\phi)]$$

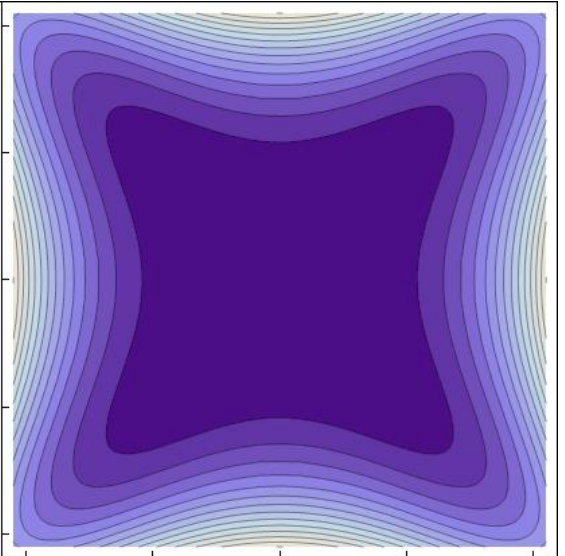
$N = 2$



$N = 3$



$N = 4$



- Are there hydro solutions that reflect these symmetries?

# Non-ellipsoidal solutions of relativistic hydro?

- Quite few non-spherical relativistic solutions
- Even those assume elliptical (ellipsoidal) symmetry
- Basics: thermodynamical quantities constant on ellipsoidal surfaces
- Defined by scaling variable:

$$s = \frac{r_x^2}{X(t)^2} + \frac{r_y^2}{Y(t)^2} + \frac{r_z^2}{Z(t)^2}$$

- How to modify this in the transverse plane?
- Ellipsis in polar coordinates:

$$\frac{r_x^2}{X(t)^2} + \frac{r_y^2}{Y(t)^2} = \frac{r^2}{R(t)^2} [1 + \epsilon_2(t) \cos(2\phi)]$$

- Generalization to „multipole” symmetries, with  $\Psi_N$  reaction plane angles):

$$s = \frac{r^N}{R(t)^N} [1 + \epsilon_N(t) \cos(N(\phi - \Psi_N))]$$

- Works even superimposing many of these

$$s = \sum_N \frac{r^N}{R(t)^N} [1 + \epsilon_N(t) \cos(N(\phi - \Psi_N))]$$

# Exact „multipole” hydro solutions

- A solution in axial coordinates:

$$s = \frac{r^N}{R^N} \left[ 1 + \epsilon \cos(N(\phi - \psi)) \right] + \frac{z^N}{R^N}$$

$$u^\mu = \gamma \left( 1, \frac{\dot{R}}{R} r \cos \phi, \frac{\dot{R}}{R} r \sin \phi, \frac{\dot{R}}{R} z \right) = \frac{x^\mu}{\tau}$$

$$T = T_0 \left( \frac{\tau_0}{\tau} \right)^{\frac{3}{\kappa}} f(s), \text{ with arbitrary } f \text{ function}$$

- Here  $R = u_t t$ , i.e.  $\dot{R} = u_t = \text{const.}$ , as well as  $\psi = \text{const.}$ ,  $\epsilon = \text{const.}$
- Constant EoS (but  $\kappa$  cancels from hadronic results)
- Works with other scaling variables as well, eg. (spherical coordinates):

$$s = \frac{r^N}{R^N} \left[ 1 + \epsilon \cos(N(\phi - \psi))(1 - \cos N\theta) + \chi \cos N\theta \right]$$

- Works combining many symmetries with different (constant)  $\epsilon$ 's and  $\psi$ 's

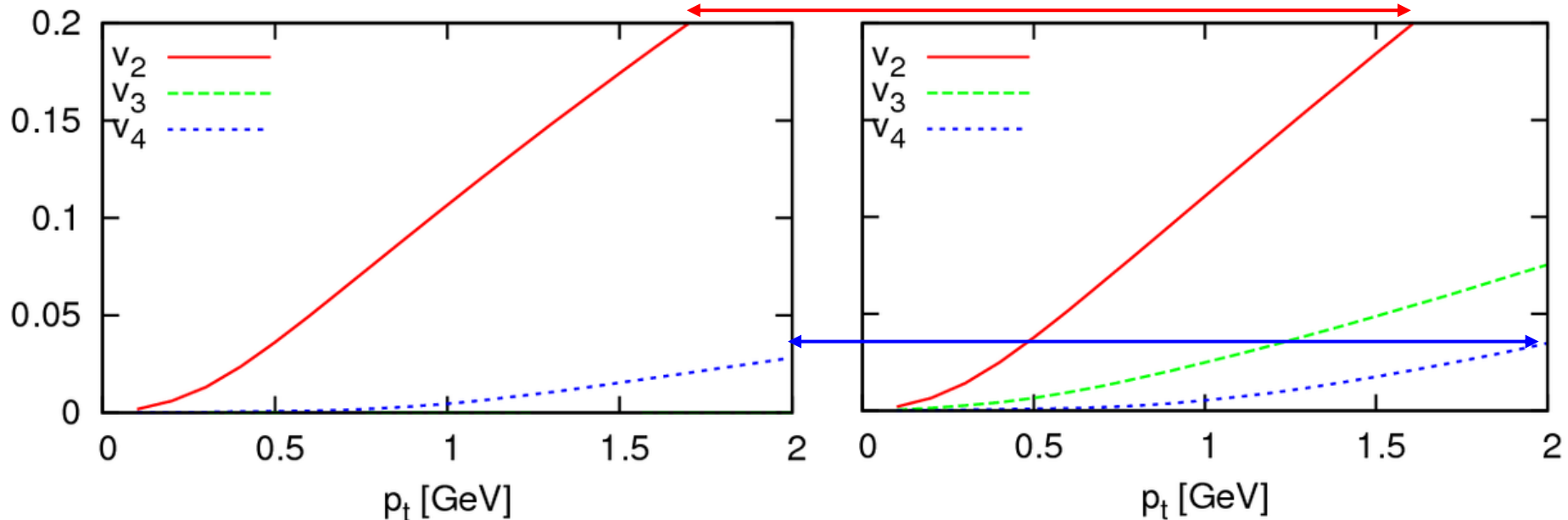
$$s = \sum_{N=2} \left[ \frac{r^N}{R(t)^N} (1 + \epsilon_N \cos(N(\phi - \psi_N))) + \frac{z^N}{R(t)^N} \right]$$

# Observables from the new solutions

- Define hadronic freeze-out, results independent of EoS ( $\kappa$ )
- Azimuthally integrated observables (spectra) are the same!
- Earlier results applicable here
  - Spectra, HBT and flow fits shown in Eur.Phys.J. A44 (2010) 473
- Odd  $v_N$  coefficients arise naturally here, even coefficients change slightly
- Relative orientation of  $\psi_N$  planes plays no role in the results
- Example:

$$\varepsilon_2=0.8, \varepsilon_3=0, \varepsilon_4=0$$

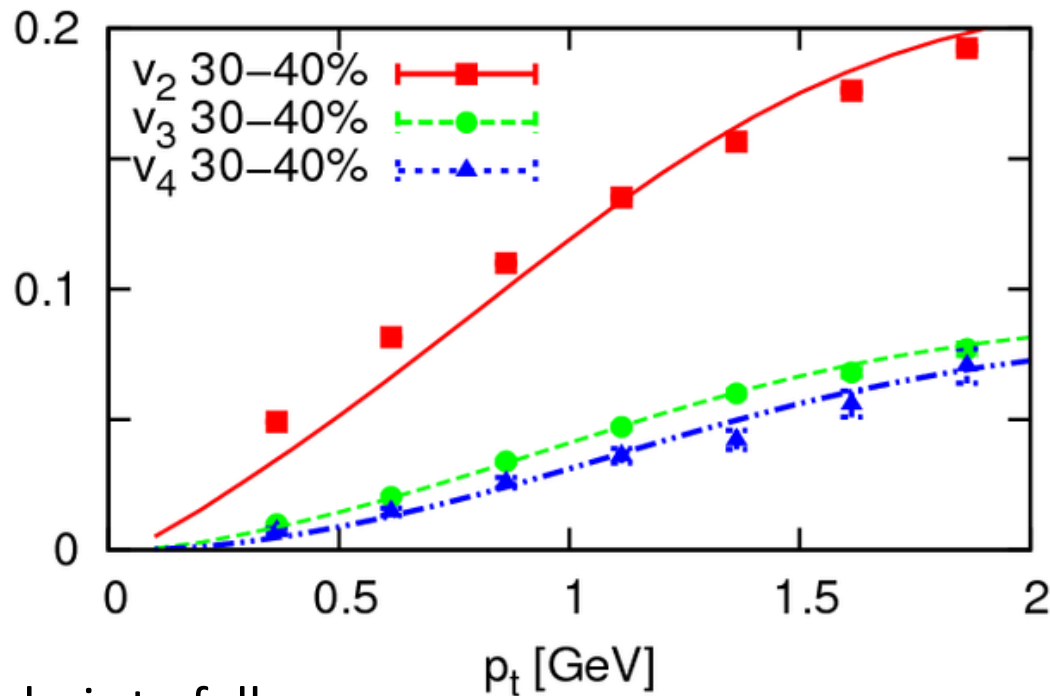
$$\varepsilon_2=0.8, \varepsilon_3=0.5, \varepsilon_4=0.1$$



# Comparison to data

- PHENIX data of higher order flow coefficients in several centrality bins:  
Phys. Rev. Lett. 107, 252301 (2011)
- Tune  $\epsilon_N$  asymmetry parameters to data (others from previous hadron fits):

$$\epsilon_2=0.47 \quad \epsilon_3=0.13 \quad \epsilon_4=0.06$$



- Detailed analysis to follow



# Summary

Higher order flow coefficients measure deviation from elliptical symmetry

Arise due to fluctuating initial conditions (c.f. wounded nucleon model)

Previously: predominantly in numerical hydro

This work: **exact analytic solutions for initial state fluctuations**

~ special case of T. Cs, Csernai, Hama, Kodama, Heavy Ion Phys. A21, 73 (2004)

New: Azimuthally integrated observables: nearly unchanged

Realistic higher order (odd) flow coefficients calculable

Detailed investigations to follow:

connection to other solutions, asHBT, higher order HBT, lattice QCD EOS)

# Thank you for your attention

And let me invite you to the 13th Zimanyi School

## ZIMÁNYI SCHOOL'13



M. M.: Women with Brushwood

13. Zimányi

**WINTER SCHOOL ON  
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YOUR INVITATION TO WPCF 2014

KRF, Gyöngyös, foot of the Mátra hills, Hungary

Last week of August 2014

T. Csörgő & T. Novák (Wigner RCP – KRF)