## Flavorless physics with flavor

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#### Outline

- MFV
- MFV as a tool for EFT and flavor
- MFV instead of R-parity
- MFV for the spectrum
- Conclusions



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#### **EFTs**



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#### **EFTs**

The basic idea of EFTs is that we have NR operators. Consider the effective Fermi theory

$$a \; \frac{\bar{\mu}\nu_{\mu}\bar{\nu}_{e}e}{\Lambda^{2}} \quad \Longrightarrow \quad \mu \to e\nu\bar{\nu}$$

- If we know nothing, we assume  $a \sim 1$ , and get  $\Lambda \sim m_W$
- Unless we have more input, we assume that all dim 6 operators are roughly the same
- We could also have

$$a \; \frac{\bar{\mu} e \bar{e} e}{\Lambda^2} \quad \Longrightarrow \quad \mu \to e e e$$

Why is it not observed?

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### SM and Fermi theory

The Fermi theory is more than a simple EFT

- In the Fermi EFT we have an operator that gives  $\mu \rightarrow eee$
- Yet, the full theory is not generic, it is the SM
- We add the GIM mechanism as a guide for the size of FCNC operators
- Flavor is broken in a very specific way, and  $\mu \rightarrow eee$  is highly suppressed by small neutrino masses
- Can the full theory above the SM also be non-generic?

## The flavor problem

In fact, we know the full theory must be non-generic!

- Many reasons to expect NP at  $\Lambda \sim 1$  TeV
- Naive flavor bounds imply  $\Lambda > 10^4$  TeV
- This tension is called

The NP flavor problem

The full theory must have a non-trivial flavor structure





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## MFV

Minimal Flavor Violation (MFV)

- In this way of thinking let us define MFV as a guide to the full theory
- The NP flavor structure comes from the SM Yukawas
- We think about the Yukawas as small (other than top) so we have approximate flavor symmetries
- This idea can be realized in many NP models
- It basically solves the NP flavor problem: the NP operators are suppressed by small Yukawas

#### **MFV:** Definition

Without Yukawas the flavor group is

 $U(3)_Q \times U(3)_U \times U(3)_D$ 

- The Yukawas are spurions
- Recall  $U(3) = SU(3) \times U(1)$ . The diagonal U(1) is baryon number

Q(3,1,1) U(1,3,1) D(1,1,3)  $Y_D(3,1,\overline{3})$   $Y_U(3,\overline{3},1)$ 

- The idea of MFV is that the Yukawas are the only spurions that break the flavor symmetry
- Linear MFV is the idea that these Yukawas are small and it is justified to keep only leading order terms

#### MFV: Can it do more?

MFV is a nice idea to solve the NP flavor problem

Q: Can we use it to solve other problems or make other predictions?

A: Yes!



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# MFV instead of R-parity



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### Susy and R-parity

SUSY without imposing R-parity  $(R_p)$ 

It leads to rapid proton decay

 $W = \mu LH + \lambda LLE + \lambda' LQD + \lambda'' UDD$ 

• For example, using  $\lambda'$  and  $\lambda''$  we have

$$A(P^+ \to \pi^0 e^+) \propto \lambda' \lambda''$$

- Standard way out: impose  $R_P$  or just Baryon number
- Flavor symmetries could be used, but these ideas are not very elegant

## MFV instead of R-parity

Idea based on E. Nikolidakis and C. Smith (0710.3129, 0809.3152)

- We impose SU(3) MFV. The U(1) is baryon number, and it is "unfair" to impose it
- We can get the B violating term already with two Ys

 $(Y_u Y_d^{\dagger}) (UUD)$ 

• Some SU(3) algebra (Q, U, D)

 $(\bar{3},3,1) \times (3,1,\bar{3}) \times (1,3,1) \times (1,3,1) \times (1,1,3) \sim (1,1,1)$ 

This is not a huge suppression

#### MFV for L violation

The big effect comes from the leptons

• The MFV symmetry is  $SU(3)_L \times SU(3)_E$ 

 $L(3,1) \quad E(1,3) \quad Y_e(3,\bar{3})$ 

- Using powers of  $Y_e$  we cannot get the L violating terms!
- The L violating terms can be generated if we have neutrino masses  $m_{\nu}(6,1)$

 $m_{\nu} Y_e LLE$ 

Huge suppression!

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## MFV for proton decay

The small Yukawas and the need for neutrino mass save the proton

$$\mathcal{A} \sim Y_u^3 Y_d^2 Y_e^2 m_\nu$$

- Still, close to the current limit (depends on  $\tan \beta$ )
- The LSP decays hadronically
- May get "large"  $\Delta B = 2$  effects (*N*- $\overline{N}$  oscillations)
- We can understand it all in terms of symmetries. MFV implies that all L violating terms will be  $\Delta L = 3$  (and  $\Delta B = 1$ )

## MFV for the spectrum



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# Heavy quarks

Idea based on YG, Nir, Thaler, Volansky, Zupan, 0706.1845

- What can MFV tell us about new vector-like quarks,  $D_i$ ?
- We assume that these new quarks are heavy and decay into SM quarks
- MFV tells us there must be 3 of them!

$$Y_{ij}^{Dd}\bar{d}_i D_j H$$

- There are four different triplet flavor representations of the new quarks
- There are always three new quarks

## Spectrum

- Depending on the details we can have several spetra
  - Degenerate
  - **●** 2+1
  - Hierarchical
- MFV predicts that in the LHC we can only see degenerate heavy quarks!
- Observation of two non-degenerate heavy quarks can rule out MFV using a "flavorless" observable

## Conclusions



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#### MFV has a lot of flavor

- MFV is an effective idea that is an addition to the EFT description
- MFV can replace R-parity
- MFV can be probed by heavy states





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