

Analytical Study of Planar Flow

Lorenzo Mannelli

Weizmann Institute of Science



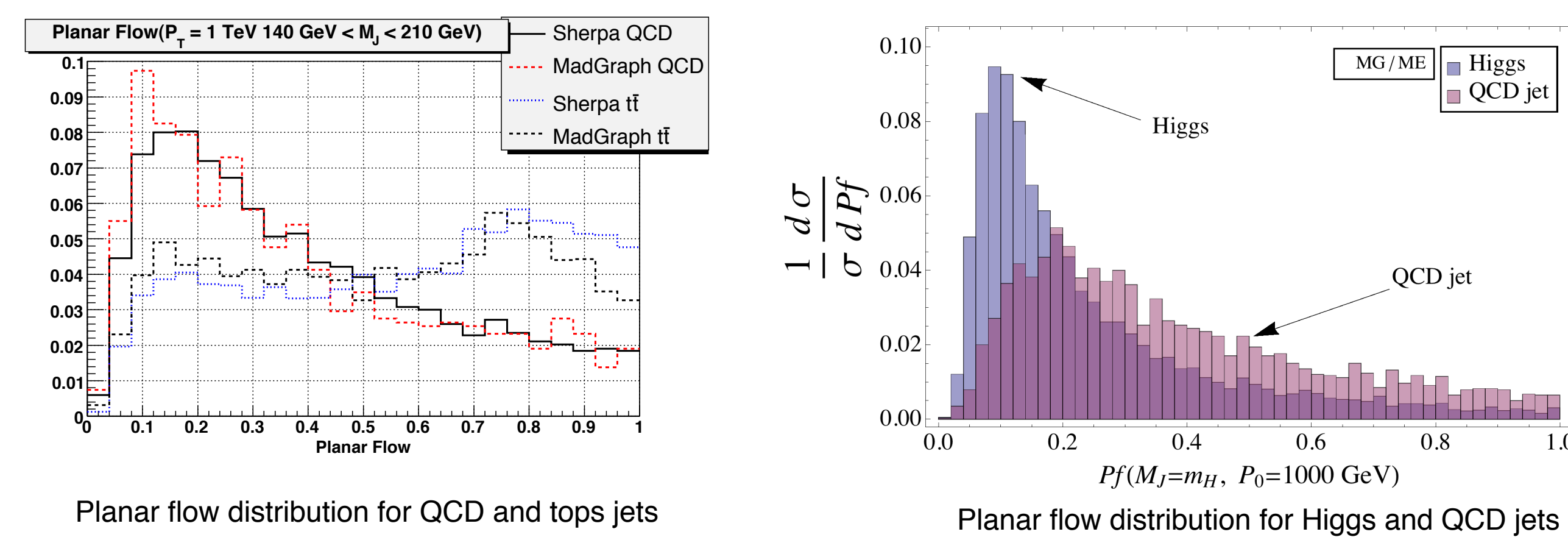
Introduction

At the LHC heavy SM particles will be frequently produced with a transverse momentum greatly exceeding their rest mass. In a typical Beyond the Standard Model (BSM) scenario an unknown heavy resonance X_{heavy} decays to highly boosted particles with intermediate masses that then decay to light quarks

$$X_{heavy} \rightarrow Y_{interm.} \rightarrow \text{jets}$$

where $Y_{interm.}$ may be a SM particle (W , Z , top) or a BSM particle. The reconstruction of hadronic decays of boosted objects $Y_{interm.}$ is particularly challenging and require and analysis of the substructure of the produced jet.

Planar flow [1] represent a valuable tool in the analysis of the jet substructure and the separation of the QCD background .



In this work we provide and study an analytical expression for the jet cross section at fixed jet mass m and planar flow P_f .

Planar Flow

In order to define the planar flow observable we introduce for a given jet the matrix I_w

$$I_w^{kl} = \frac{1}{m} \sum_i w_i \frac{p_{i,k}}{w_i} \frac{p_{i,l}}{w_i}$$

where m is the jet mass, w_i the energy of the particle i and $p_{i,k}$ is the projection of its momentum on the plane perpendicular to the jet axis. Given I_w planar flow is defined as

$$P_f = \frac{4 \det(I_w)}{\text{tr}(I_w)^2} = \frac{4\lambda_1 \lambda_2}{(\lambda_1 + \lambda_2)^2}$$

where $\lambda_{1,2}$ are the eigenvalues of I_w

Analytic Expression

In the collinear limit the amplitude matrix element factorize as

$$|M(2 \rightarrow n)|^2 \simeq \frac{4}{s_{123}^2} (4\pi\alpha)^2 |M(2 \rightarrow n-2)|^2 \langle P_{1 \rightarrow 3} \rangle$$

where $\langle P_{1 \rightarrow 3} \rangle$ is the 1 to 3 splitting function [2]. The reason to use the 1 to 3 rather than the more familiar 1 to 2 splitting function is that it allow to go beyond the angular ordered approximation $\theta_i \ll \theta_j$, approximation that break down in the region of large planar flow.

In the collinear limit the jet cross section at fixed jet mass and planar flow is

$$\frac{1}{\sigma_0} \frac{d\sigma}{dP_f dm^2}(P_f, m^2) = \frac{1}{m^4} \frac{\alpha^2}{\pi^3} \int d\theta_1 d\theta_2 d\phi \frac{1}{z_3} \bar{p}_1 \bar{p}_2 \theta_1 \theta_2 J(\bar{p}_1, \bar{p}_2, \theta_1, \theta_2) \cdot \langle P_{1 \rightarrow 3} \rangle(\bar{p}_1, \bar{p}_2, \theta_1, \theta_2)$$

where J is a Jacobian factor, θ_1, θ_2 are angles between the partons momenta and the jet axis and \bar{p}_1, \bar{p}_2 are functions that parametrize the jet mass and planar flow constraints.

The previous equation is well described by the following fitting formula

$$A \left(\frac{C_A^2 \alpha^2}{\pi^3} \right) \frac{q^2}{m^4 P_f} \log \left[B \frac{2qR}{m} \frac{1}{P_f} \right]$$

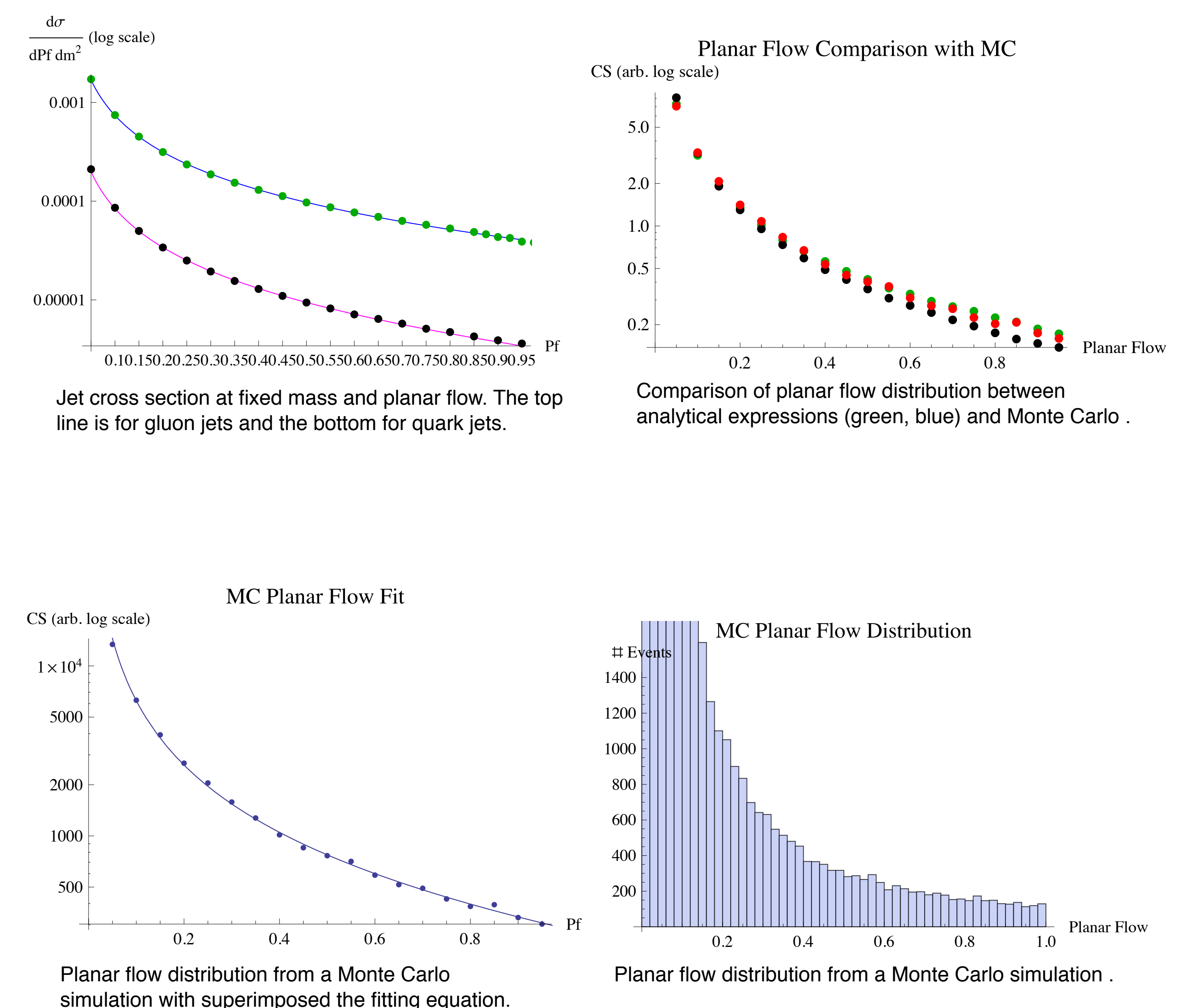
where $C_A = N_c = 3$ is a Casimir. The fit parameters $A \approx 0.4$, $B \approx 1.4$ and $A \approx 0.06$, $B \approx 0.4$ describe well gluon and quarks jets respectively.

Conclusion

Highly boosted particles are important signals of new Physics, planar flow turn out to be a valuable observable for the study of jet substructures produced by boosted objects.

We have derived an analytical description of planar flow.

The fitting formula and a comparisons with Monte Carlo simulation is shown in the following figures.



References

- [1] L. G. Almeida, S. J. Lee, G. Perez *et al.*, *Substructure of high-pT Jets at the LHC*, Phys. Rev. **D79**, 074017 (2009) [SPIRES].
- [2] S. Catani and M. Grazzini, *Collinear Factorization and Splitting Functions for Next-to-next-to-leading Order QCD Calculations*, Phys. Lett. B **446**, 143 (1999) [hep-ph/9810389];